

Title: TBA - Particle Physics Seminar

Speakers: Dhong Yeon Cheong

Collection/Series: Particle Physics

Subject: Particle Physics

Date: May 05, 2026 - 5:00 PM

URL: <https://pirsa.org/26050037>

Kinetic Isocurvature Perturbations

w / K.J. Bae, J.-O. Gong, K. Harigaya, C.S. Shin

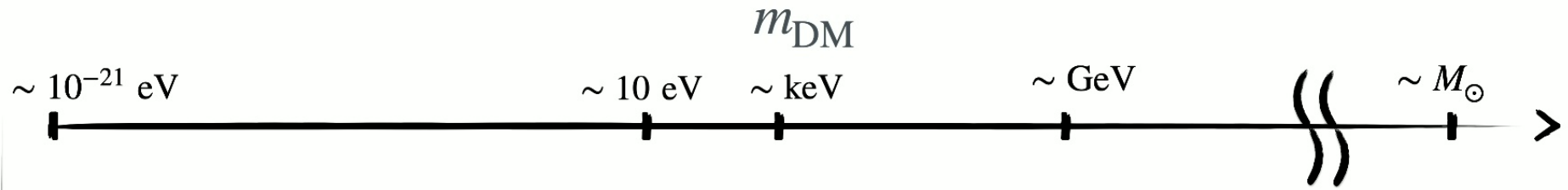
[2603.22394]

Dhong Yeon (D.Y.) Cheong

@ Perimeter Institute Particle Physics Seminar



Motivation



$$\Omega_{\text{DM}} \neq 0$$

Many possibilities : origin, properties, signatures

Important to identify unique (cosmological) signatures

Motivation

Q : A new type of isocurvature perturbation?

- Kinetic Isocurvature Perturbations
- Adiabatic Number + Non-Adiabatic Momentum

Q : Any unique observational probe?

- Modulation of the Free Streaming Length
- Modulation of the Matter Power Spectrum

DM Energy Density Perturbations

The energy density of DM

$$\rho_{\text{DM}} = E_{\text{DM}} n_{\text{DM}}, \quad E_{\text{DM}} \equiv \sqrt{m_{\text{DM}}^2 + p_{\text{DM}}^2}$$

The energy density perturbation decomposes to

$$\frac{\delta \rho_{\text{DM}}}{\rho_{\text{DM}}} = \frac{\delta n_{\text{DM}}}{n_{\text{DM}}} + \frac{p_{\text{DM}}^2}{E_{\text{DM}}^2} \frac{\delta p_{\text{DM}}}{p_{\text{DM}}}$$

DM Energy Density Perturbations

The energy density of DM

$$\rho_{\text{DM}} = E_{\text{DM}} n_{\text{DM}}, \quad E_{\text{DM}} \equiv \sqrt{m_{\text{DM}}^2 + p_{\text{DM}}^2}$$

The energy density perturbation decomposes to

$$\frac{\delta \rho_{\text{DM}}}{\rho_{\text{DM}}} = \frac{\delta n_{\text{DM}}}{n_{\text{DM}}} + \frac{p_{\text{DM}}^2}{E_{\text{DM}}^2} \frac{\delta p_{\text{DM}}}{p_{\text{DM}}}$$



Normally the dominant source

DM Energy Density Perturbations

The energy density of DM

$$\rho_{\text{DM}} = E_{\text{DM}} n_{\text{DM}}, \quad E_{\text{DM}} \equiv \sqrt{m_{\text{DM}}^2 + p_{\text{DM}}^2}$$

The energy density perturbation decomposes to

$$\frac{\delta \rho_{\text{DM}}}{\rho_{\text{DM}}} = \frac{\cancel{\delta n_{\text{DM}}}}{\cancel{n_{\text{DM}}}} + \frac{p_{\text{DM}}^2}{E_{\text{DM}}^2} \frac{\delta p_{\text{DM}}}{p_{\text{DM}}}$$

↓
0!

DM Energy Density Perturbations

The energy density of DM

$$\rho_{\text{DM}} = E_{\text{DM}} n_{\text{DM}}, \quad E_{\text{DM}} \equiv \sqrt{m_{\text{DM}}^2 + p_{\text{DM}}^2}$$

The energy density perturbation decomposes to

$$\frac{\delta \rho_{\text{DM}}}{\rho_{\text{DM}}} = \frac{\cancel{\delta n_{\text{DM}}}}{\cancel{n_{\text{DM}}}} + \frac{p_{\text{DM}}^2}{E_{\text{DM}}^2} \frac{\delta p_{\text{DM}}}{p_{\text{DM}}}$$



Kinetic Isocurvature Perturbation

Fluctuations imprinted at the production of DM, inheriting p modulation.

Kinetic Isocurvature Perturbation

$$\frac{\delta\rho_{\text{DM}}}{\rho_{\text{DM}}} \simeq \frac{p_{\text{DM}}^2}{E_{\text{DM}}^2} \frac{\delta p_{\text{DM}}}{p_{\text{DM}}}$$

For (warm) DM, as DM becomes non-relativistic

$$v_{\text{DM}} = \frac{p_{\text{DM}}}{E_{\text{DM}}} \propto \frac{1}{a} \qquad \frac{\delta\rho_{\text{DM}}}{\rho_{\text{DM}}} \propto \frac{1}{a^2}$$

Free-streaming length perturbation governed by initial p distribution

$$\delta\lambda \equiv \frac{\delta\lambda_{\text{FS}}}{\lambda_{\text{FS}}} \propto \frac{\delta p_{\text{DM}}}{p_{\text{DM}}} \qquad \lambda_{\text{FS}} \equiv \int_{t_i}^{t_{\text{eq}}} dt \frac{v(t)}{a(t)} = \int_{a_i}^{a_{\text{eq}}} \frac{da}{a^2 H(a)} \frac{p_{\text{DM}}(a)}{E_{\text{DM}}(a)}$$

Kinetic Isocurvature Perturbation

17:00 17:00
17:00 17:00

Setup to realize kinetic isocurvature perturbations?

2. Setup

Setup

Key : adiabatic number perturbations + modulated kinetic distribution

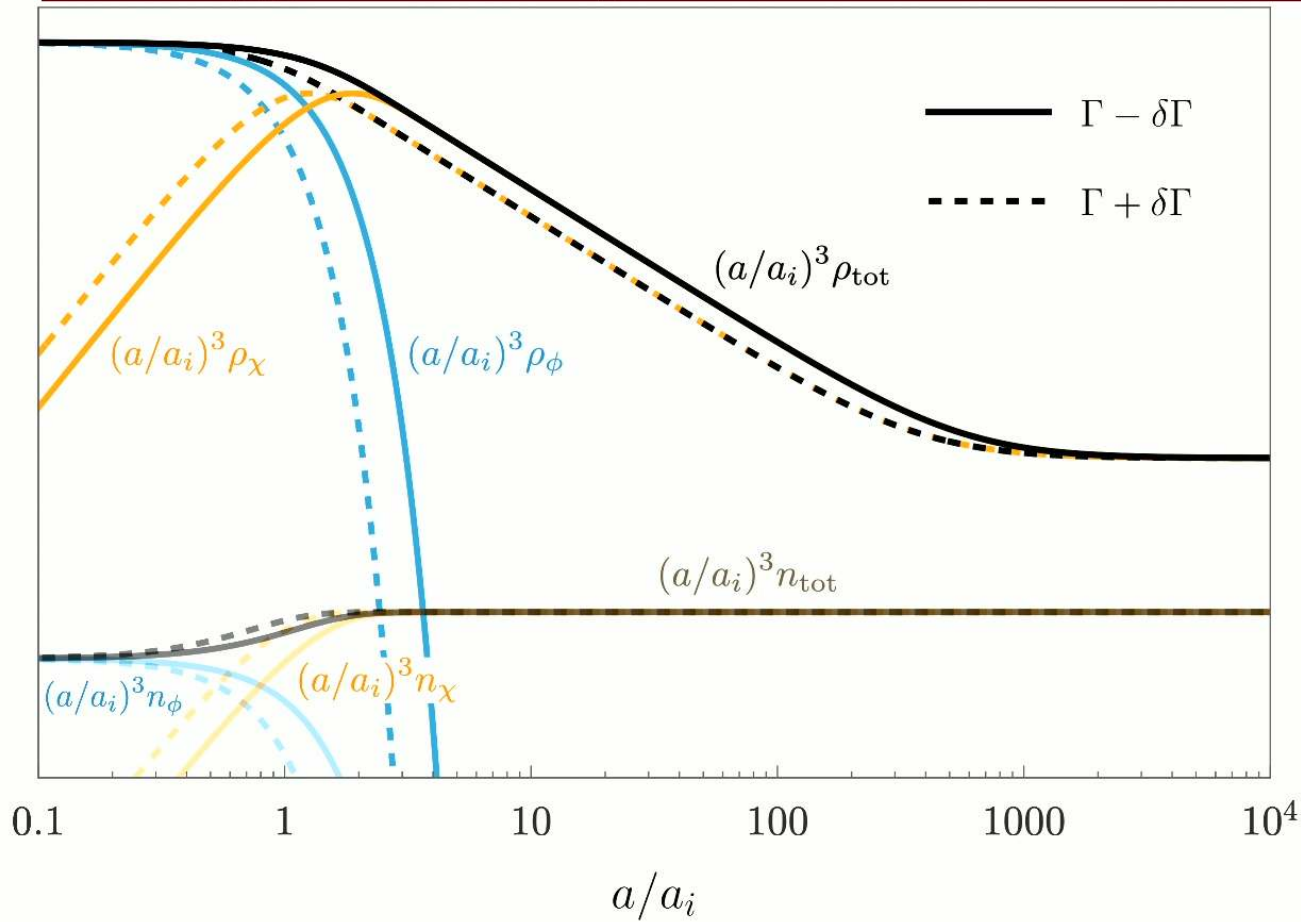
Decaying massive scalar field : ϕ (subdominant component)

Light DM particle : χ / Modulating field σ controlling the decay

$$\mathcal{L} \supset \begin{cases} A\phi\chi^2 & \chi : \text{scalar} \\ y\phi\bar{\chi}\chi & \chi : \text{fermion.} \end{cases} \quad \Gamma = \begin{cases} \frac{A^2}{16\pi m_\phi} \left(1 - \frac{4m_\chi^2}{m_\phi^2}\right)^{1/2} \\ \frac{y^2 m_\phi}{8\pi} \left(1 - \frac{4m_\chi^2}{m_\phi^2}\right)^{3/2} \end{cases}$$

$$\delta\sigma \rightarrow (\delta A, \delta y) \rightarrow \delta\Gamma$$

Energy density modulation



Decay yield

$$n_{\chi}(a_i) = 2n_{\phi}(a_i).$$

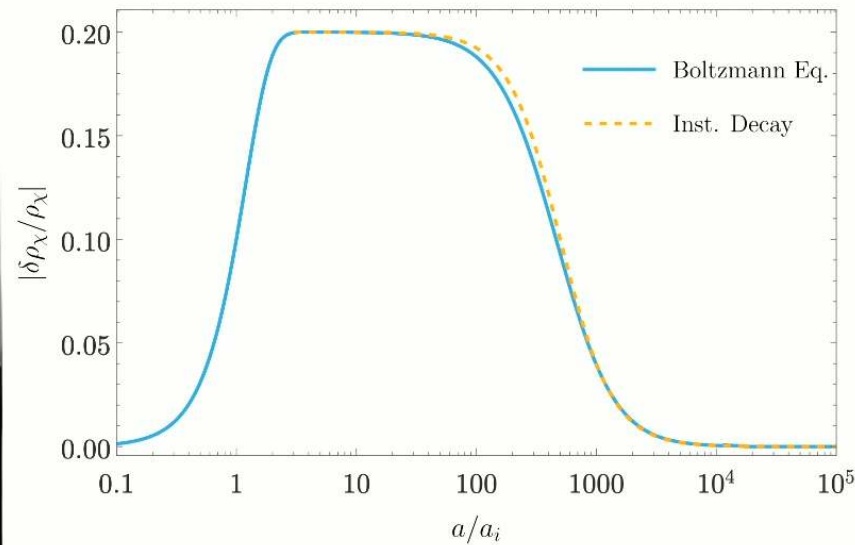
Fixed, no δn

Different decay times t_i

Momentum redshift
 $p(a) = p_i(a_i/a)$

Induced $\delta\rho_{\chi}$

Relation between $\delta\rho$ and $\delta\lambda$



Analytic relation w/ instantaneous decay

$$\frac{\delta\rho_\chi}{\rho_\chi} \simeq -\frac{1}{2} \left[\frac{\left(m_\phi^2/4\right) \left(a_i/a\right)^2}{m_\chi^2 + \frac{m_\phi^2}{4} \left(\frac{a_i}{a}\right)^2} \right] \frac{\delta\Gamma}{\Gamma}.$$

$$\delta\lambda \simeq -\frac{1}{2} \frac{\delta\Gamma}{\Gamma}$$

$$\frac{\delta\rho_\chi}{\rho_\chi} \simeq \left[\frac{\left(m_\phi^2/4\right) \left(a_i/a\right)^2}{m_\chi^2 + \frac{m_\phi^2}{4} \left(\frac{a_i}{a}\right)^2} \right] \delta\lambda$$

Relation between $\delta\rho$ and $\delta\lambda$

Within our region of interest,

$$u_i = p_i/m_\chi$$

$$\lambda_{\text{FS}} \simeq 1\text{Mpc} \left(\frac{u_i^2 t_i}{10^6 \text{sec}} \right)^{1/2} \left[1 - 0.07 \ln \left(\frac{u_i^2 t_i}{10^6 \text{sec}} \right) \right]$$

$$|\delta_\lambda| \simeq \left| \frac{\delta\rho_\chi}{\rho_\chi} \right| \left[1 + 1.56 g_*^{-1/2} \left(\frac{1\text{Mpc}}{\lambda_{\text{FS}}} \right)^2 \left(1 + 0.27 \ln \left(\frac{1\text{Mpc}}{\lambda_{\text{FS}}} \right) \right) \left(\frac{6.7\text{Mpc}^{-1}}{k} \right)^2 \right]$$

A Concrete Model

Massive scalar field : ϕ (zero mode), scalar DM χ , modulating scalar σ .

$$\mathcal{L}_{\text{int}} = g\sigma\phi\chi^2, \quad \Gamma_\phi \simeq \frac{1}{8\pi} \frac{g^2\sigma_i^2}{m_\phi}$$

Temperature scales : $T_{\text{dec}}, T_{\text{osc}}, T_{\text{NR}}$, assuming RD at T_{dec} and T_{osc} .

$$T_{\text{dec}} \simeq \sqrt{\Gamma M_P}, \quad T_{\text{osc}} \simeq \left(\frac{90}{\pi^2 g_*} \right)^{1/4} \sqrt{m_\phi M_P}, \quad T_{\text{NR}} \simeq T_{\text{dec}} \frac{m_\chi}{m_\phi}$$

$$\frac{\rho_\chi}{s} \simeq \frac{m_\chi n_\chi}{s} \simeq m_\chi \frac{n_\phi}{s} \sim \frac{m_\chi m_\phi \phi_i^2}{g_* (m_\phi M_P)^{3/2}}$$

$$\phi_i \sim \left(\frac{g_* T_{\text{dec}} T_{\text{eq}}}{T_{\text{NR}}} \right)^{1/2} \frac{M_P^{3/4}}{m_\phi^{1/4}} \sim 10^{15} \text{GeV} \times \left(\frac{T_{\text{dec}}}{10^9 \text{GeV}} \right)^{1/2} \left(\frac{10 \text{keV}}{T_{\text{NR}}} \right)^{1/2} \left(\frac{10^8 \text{GeV}}{m_\phi} \right)^{1/4}$$

A Concrete Model

Requirements on m_σ

In order for large Γ modulation, σ should start oscillating at $T < T_{\text{dec}}$

$$m_\sigma < \frac{T_{\text{dec}}^2}{M_P} = 1\text{GeV} \left(\frac{T_{\text{dec}}}{10^9\text{GeV}} \right)^2$$

In order for $\delta m_\sigma^{1\text{-loop}} \sim \frac{g}{4\pi} \phi < \frac{T_{\text{dec}}^2}{M_P}$, $g \lesssim 10^{-8} \left(\frac{T_{\text{NR}}}{10\text{keV}} \right)^{1/2} \left(\frac{m_\phi}{10^8\text{GeV}} \right)$

A Concrete Model

σ eventually starts oscillations and dissipated at T_{dis} , needs suppressed

$$\mathcal{R}_\sigma \sim \frac{\delta\sigma}{\sigma_i} \frac{\sigma_i^2}{M_P^2} \frac{\sqrt{m_\sigma M_P}}{T_{\text{dis}}}$$

$$\sim 10^{-6} \times \frac{\delta\sigma}{\sigma_i} \frac{m_\phi}{10^8 \text{GeV}} \left(\frac{T_{\text{dec}}}{10^9 \text{GeV}} \right)^2 \left(\frac{m_\sigma}{1 \text{GeV}} \right)^{1/2} \frac{10^3 \text{GeV}}{T_{\text{dis}}} \left(\frac{10^{-8}}{g} \right)^2$$

\mathcal{R}_σ suppressed for sufficiently large T_{dis}

e.g. If $\mathcal{L} \supset A_H \sigma |H|^2$, σ is dissipated with a rate $\sim \frac{A_H^2}{8\pi T}$

$$T_{\text{dis}} \sim 10^3 \text{GeV} \left(\frac{A_H}{10^{-3} \text{GeV}} \right)^{2/3}$$

$m_\sigma \gg \text{MeV}$ for suppressing the ΔN_{eff}

Relation between $\delta\rho$ and $\delta\lambda$

Within our region of interest,

$$u_i = p_i/m_\chi$$

$$\lambda_{\text{FS}} \simeq 1\text{Mpc} \left(\frac{u_i^2 t_i}{10^6 \text{sec}} \right)^{1/2} \left[1 - 0.07 \ln \left(\frac{u_i^2 t_i}{10^6 \text{sec}} \right) \right]$$

$$|\delta\lambda| \simeq \left| \frac{\delta\rho_\chi}{\rho_\chi} \right| \left[1 + 1.56 g_*^{-1/2} \left(\frac{1\text{Mpc}}{\lambda_{\text{FS}}} \right)^2 \left(1 + 0.27 \ln \left(\frac{1\text{Mpc}}{\lambda_{\text{FS}}} \right) \right) \left(\frac{6.7\text{Mpc}^{-1}}{k} \right)^2 \right]$$

The $\propto k^{-2}$ dependence gives overall enhancement on $\delta\lambda$ associated with a suppressed $\delta\rho_\chi/\rho_\chi$. Also a λ_{FS} dependence.

Any constraints on $\delta\rho_\chi/\rho_\chi \leftrightarrow \delta\lambda$?

3. Isocurvature Constraints

Kinetic Isocurvature Constraints

We consider kinetic isocurvature perturbations entering the horizon with

$$\langle \delta_\lambda(k) \delta_\lambda(k') \rangle = (2\pi)^3 \delta(k + k') \frac{2\pi^2}{k^3} P_{\delta_\lambda}(k)$$

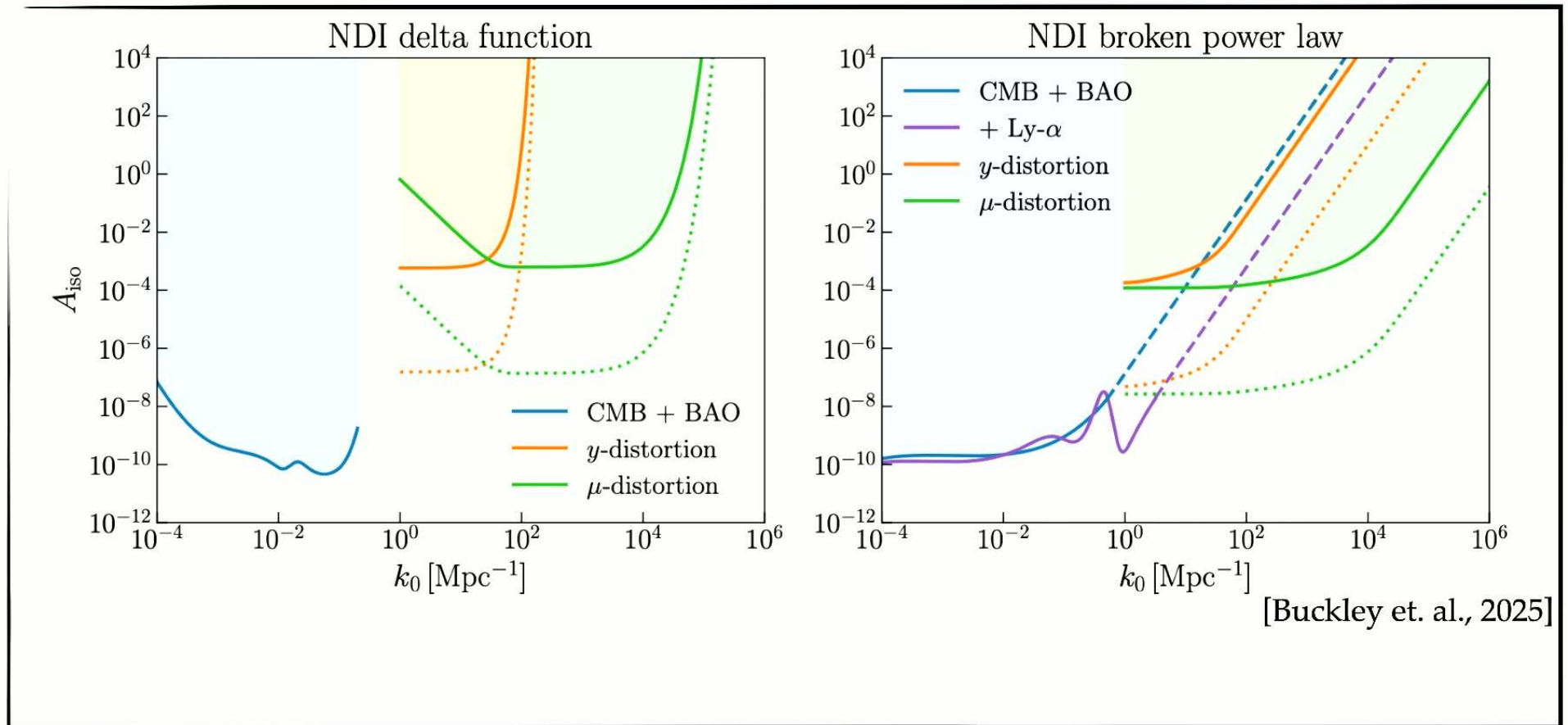
$$P_{\text{kin}}(k) \simeq 2 \times 10^{-4} g_* (\lambda_{\text{FS}} k)^4 P_{\delta_\lambda}(k).$$

Two representative profiles :

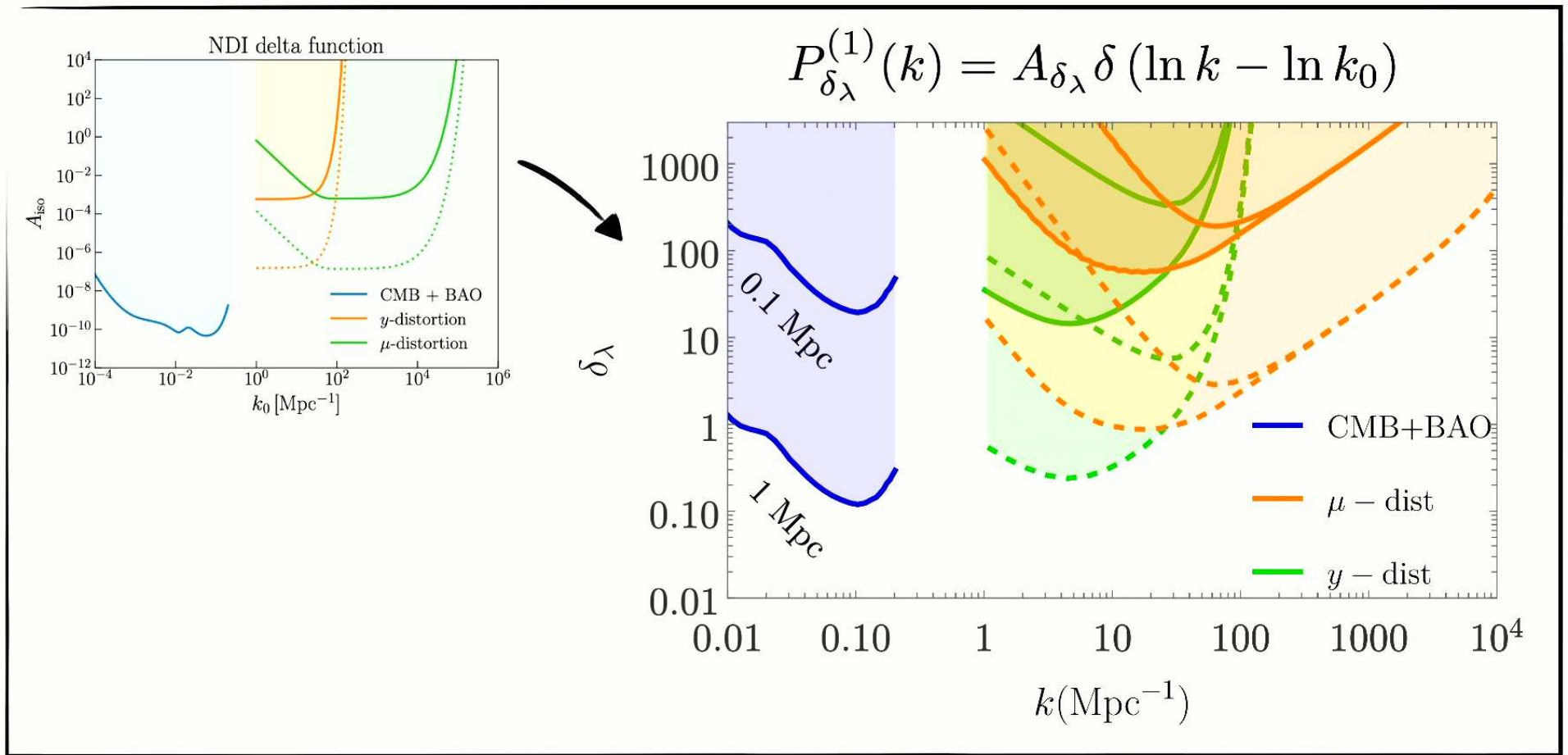
$$P_{\delta_\lambda}^{(1)}(k) = A_{\delta_\lambda} \delta(\ln k - \ln k_0)$$

$$P_{\delta_\lambda}^{(2)}(k) = A_{\delta_\lambda}$$

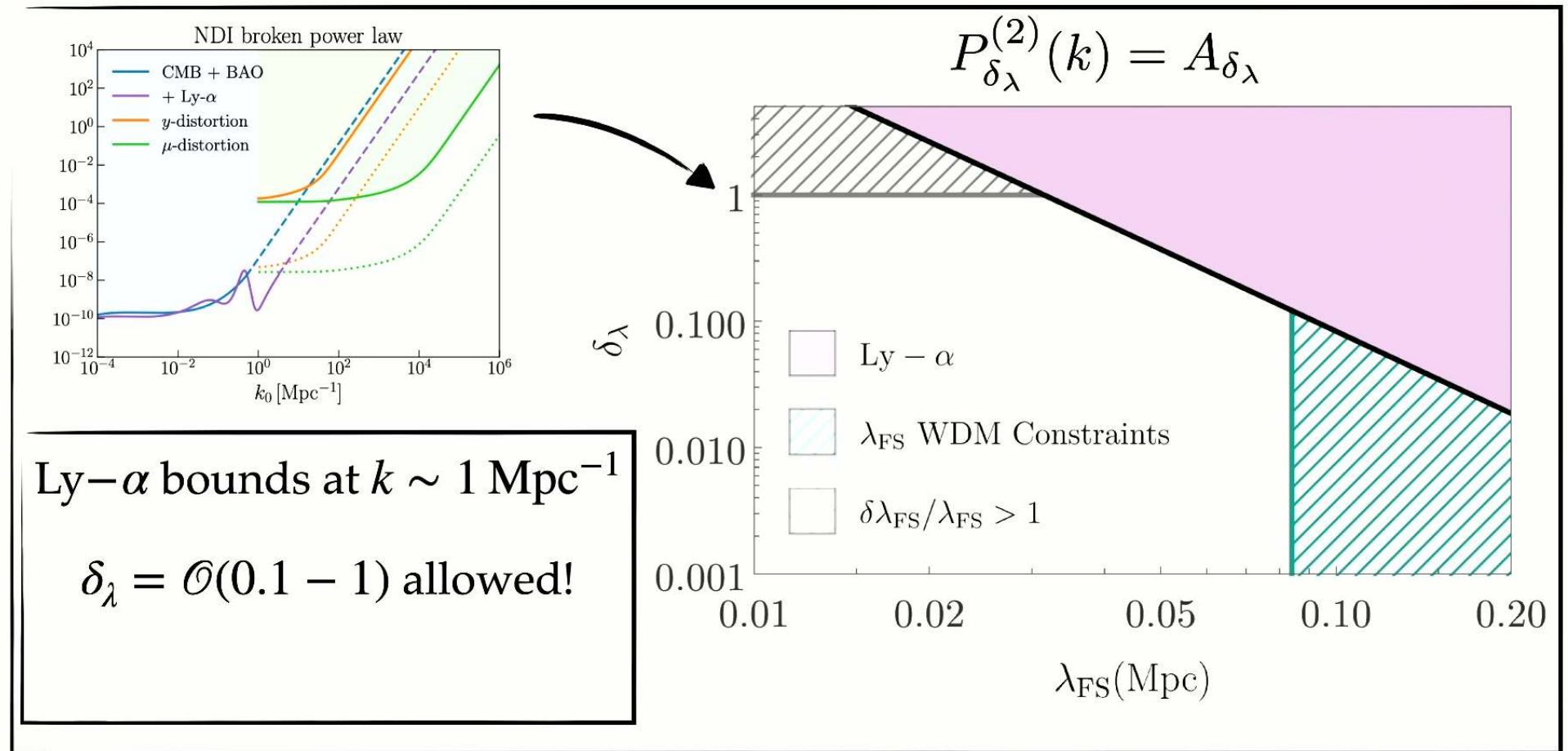
Kinetic Isocurvature Constraints



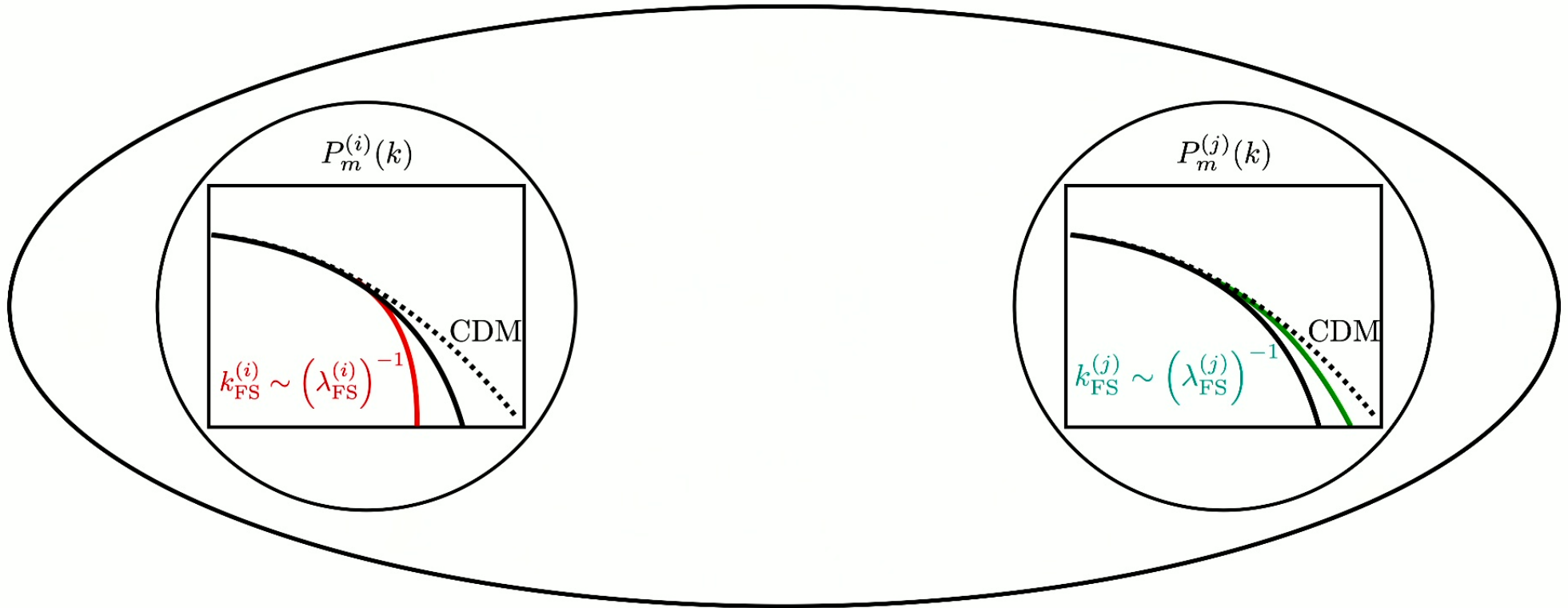
Kinetic Isocurvature Constraints



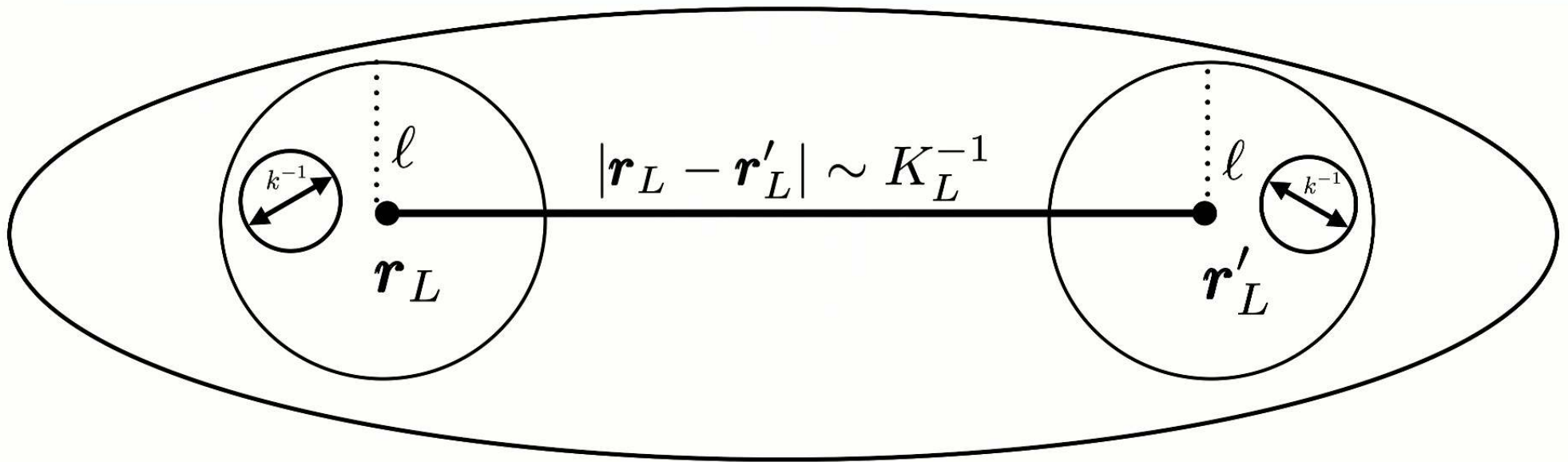
Kinetic Isocurvature Constraints



Modulating Free Streaming Length



Position Dependent Power Spectrum



[Takada & Wu (2013), Chiang et.al. (2014)]

Position dep. power spectrum
within sub volume V_s

$$P_\delta(\mathbf{k}, \mathbf{r}_L) \equiv \frac{1}{V_s} |\delta(\mathbf{k}, \mathbf{r}_L)|^2,$$

$$\delta(\mathbf{k}, \mathbf{r}_L) \equiv \int d^3r \delta(\mathbf{r}) W_s(\mathbf{r} - \mathbf{r}_L) e^{-i\mathbf{k} \cdot \mathbf{r}} = \int \frac{d^3q}{(2\pi)^3} \delta^{(3)}(\mathbf{k} - \mathbf{q}) W_s(\mathbf{q}) e^{-i\mathbf{q} \cdot \mathbf{r}_L}$$

D.Y. Cheong | Kinetic Isocurvature Perturbations

26

Position Dependent Power Spectrum

Modulation in λ_{FS}

$$\lambda_{\text{FS}}(\mathbf{x}) = \bar{\lambda}_{\text{FS}} + \delta\lambda_{\text{FS}}(\mathbf{x}).$$

Separate universe approach : λ_{FS} constant within a V_s patch.

$$\Delta P_m(\mathbf{k}, \lambda_{\text{FS}}(\mathbf{r}_L)) \simeq \lambda_{\text{FS}} \left. \frac{dP_m}{d\lambda_{\text{FS}}} \right|_{\lambda_{\text{FS}} = \bar{\lambda}_{\text{FS}}} \delta\lambda(\mathbf{r}_L)$$

$$\Delta P_m(\mathbf{k}, \lambda_{\text{FS}}(\mathbf{r}_L)) \equiv P_m(\mathbf{k}, \lambda_{\text{FS}}(\mathbf{r}_L)) - P_m(\mathbf{k}, \bar{\lambda}_{\text{FS}})$$

Position Dependent Power Spectrum

Correlation of ΔP_m governed by P_{δ_λ}

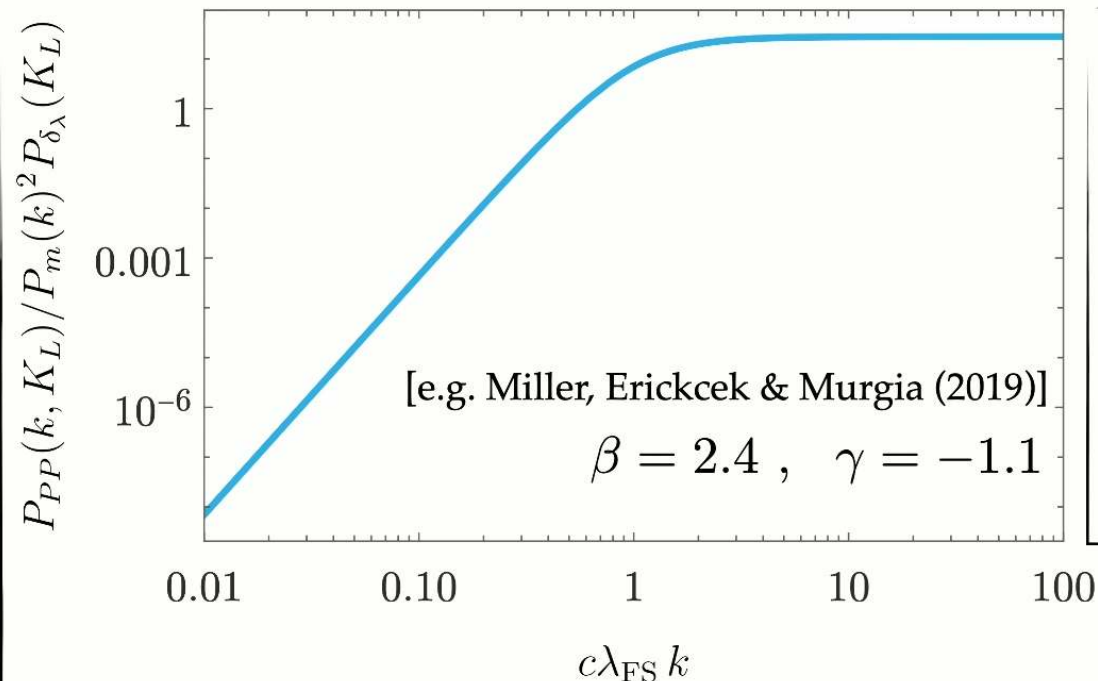
$$P_{PP}(\mathbf{k}, \mathbf{K}_L) = \lambda_{\text{FS}}^2 \left(\frac{dP_m}{d\lambda_{\text{FS}}} \right)^2 P_{\delta_\lambda}(\mathbf{K}_L)$$

$$\langle \Delta P_m(\mathbf{k}, \lambda_{\text{FS}}; \mathbf{K}_L) \Delta P_m(\mathbf{k}, \lambda_{\text{FS}}; \mathbf{K}'_L) \rangle \equiv (2\pi)^3 \delta^{(3)}(\mathbf{K}_L + \mathbf{K}'_L) P_{PP}(\mathbf{k}, \mathbf{K}_L)$$

Size of effect governed by P_m dependence on λ_{FS} , through the T_{FS} !

Modulating Power Spectrum

$$P_m(k) = T_{\text{FS}}(k)^2 P_{\mathcal{R}}(k) |W_s(k)|^2 \quad \text{Representative form } T_{\text{FS}}(k, \lambda_{\text{FS}}) = \left[1 + (c\lambda_{\text{FS}}k)^\beta \right]^\gamma$$



$$\frac{P_{PP}(k_{\text{FS}}, K_L)}{P_m(k_{\text{FS}})^2} \simeq 10 P_{\delta_\lambda}(K_L)$$

$\mathcal{O}(1)$ fluctuations in normalized correlation

Modulations in galaxy surveys, small scale observations?

Conclusions and Discussions

- Kinetic Isocurvature Perturbations as a new class of isocurvature!
- Governed by p fluctuation, $O(1)$ fluctuations compatible with isocurvature constraints.
- $O(1)$ modulations in the patch-by-patch free streaming length.

Open questions :

UV complete model, optimized observables, observational test?