

**Title:** Lecture - Quantum Fields & Strings, PHYS 77

**Speakers:** Pedro Vieira

**Collection/Series:** Quantum Fields & Strings (Elective), March 30 - May 1, 2026

**Subject:** Quantum Fields and Strings

**Date:** April 20, 2026 - 1:00 PM

**URL:** <https://pirsa.org/26040055>

$$Z_{SG} = \int \mathcal{D}\phi \exp\left(-\int \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{\alpha}{\beta^2} \cos \beta \phi + \gamma\right)$$

$$Z_{\text{Thirring}} = \int \mathcal{D}\Psi \exp\left(-\int \bar{\Psi} \not{\partial} \Psi - m \bar{\Psi} \Psi - \frac{g}{2} (\bar{\Psi} \gamma^\mu \Psi)\right)$$

$\begin{pmatrix} \Psi_R \\ \Psi_L \end{pmatrix}$

Winding  
current


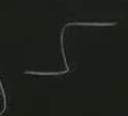
Dictionary:

$$\beta^2 = \frac{4\pi}{1 + g/\pi}$$

$$\frac{\alpha_0}{\beta^2} \cos \beta \phi \longleftrightarrow -m_0 \bar{\Psi} \Psi$$

winding current  $\equiv -\frac{\beta}{2\pi} \epsilon^{\mu\nu} \partial_\mu \phi \longleftrightarrow \bar{\Psi} \gamma^\mu \Psi \equiv$  fermion current

$\begin{matrix} \text{---} \\ \text{---} \end{matrix} \begin{matrix} = \text{soliton} \\ = \delta\phi \end{matrix} \longleftrightarrow \begin{matrix} \Psi \\ \Psi \bar{\Psi} \end{matrix} \text{ bound state}$

- \* Lumps & anti-lumps attract for  $\vec{p}^2 < 4\pi$  ←
  - \* They form bound-states. Lightest bound-state =  $\delta\phi$  naive particle 
  - \* Free @  $4\pi$
  - \* Repel for  $\vec{p}^2 > 4\pi$  (only  $\int$  exist!) 
  - \* Theory exists for  $\vec{p}^2 < 8\pi$
- nice in  $g$  ( $g > 0, g < 0$ )

$\int_{-\infty}^{\infty} \frac{d^2 p}{p^2} \sim \text{finite}$ . Only

$\int_{-\infty}^{\infty} \frac{d^3 p}{p^3} \sim \text{divergent}$

Normal Order Works!

$$\mathcal{H} = \frac{1}{2} \hat{\pi}^2 + \frac{1}{2} (\partial_x \hat{\phi})^2 + \frac{\alpha}{\beta^2} \cos \beta \hat{\phi} + \gamma$$

$$\hat{\phi} = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \frac{1}{\sqrt{2\omega_k}} \left( \hat{a}_k e^{-ikx} + \hat{a}_k^\dagger e^{+ikx} \right), \quad \omega_k = \sqrt{k^2 + m^2}$$

ren quantities

$\hat{a}^\dagger \hat{a}$

freedom!

$$\begin{array}{c} \circ \\ \circ \end{array} \gamma \begin{array}{c} \circ \\ \circ \end{array} \mathcal{H} \begin{array}{c} \circ \\ \circ \end{array} m = \begin{array}{c} \circ \\ \circ \end{array} \gamma \begin{array}{c} \circ \\ \circ \end{array} \mathcal{H} \begin{array}{c} \circ \\ \circ \end{array} \mu \quad \left| \quad \alpha \rightarrow \alpha \left( \frac{\mu^2}{m^2} \right) \right. \begin{array}{l} \beta^2 / 8\pi \\ \text{finite} \end{array} \\
 \gamma \rightarrow \gamma + \frac{\mu^2 - m^2}{8\pi}
 \end{array}$$

Warm-up:

free massive scalar

$$\mathcal{H} = \int \frac{d^3x}{2} \left[ \dot{\hat{\phi}}^2 + (\nabla_x \hat{\phi})^2 + \frac{M^2}{2} \hat{\phi}^2 \right] + \int \frac{d^3x}{2} m^2 \hat{\phi}^2$$

$$= \int \frac{d^3x}{8\pi} \left[ \mu^2 - m^2 - M^2 \log \frac{\mu^2}{m^2} \right]$$

$M, \mu, m$  all finite!

$\mu^2 - m^2 - M^2$  all finite.

bound on GS energy using trial state

$$| \text{trial} \rangle = | 0, \mu \rangle$$

vacuum for a scalar of mass  $\mu$ ,  $a_{\mathbf{k}, \mu} | 0, \mu \rangle = 0$

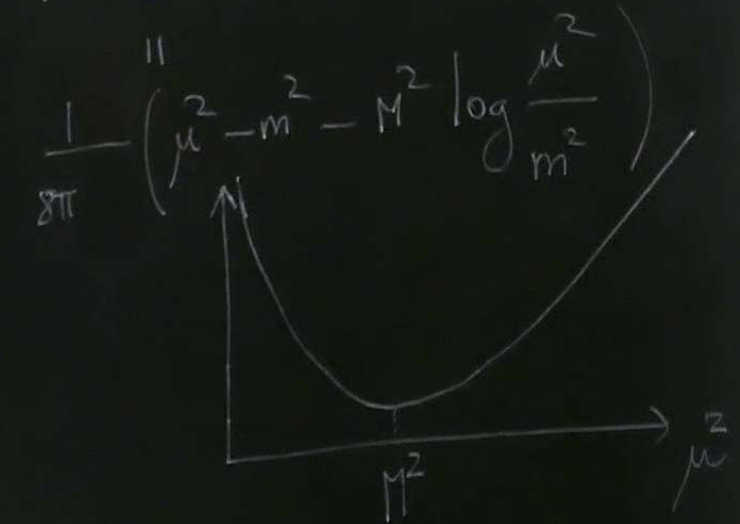
$M, \mu, m$  all finite!

$$+ \frac{M^2}{2} \phi^2 + \dots + m^2$$

$$\left( \mu^2 - m^2 - M^2 \log \frac{\mu^2}{m^2} \right)$$

trial state  
all finite.

$$\langle \text{trial} | \mathcal{H} | \text{trial} \rangle$$

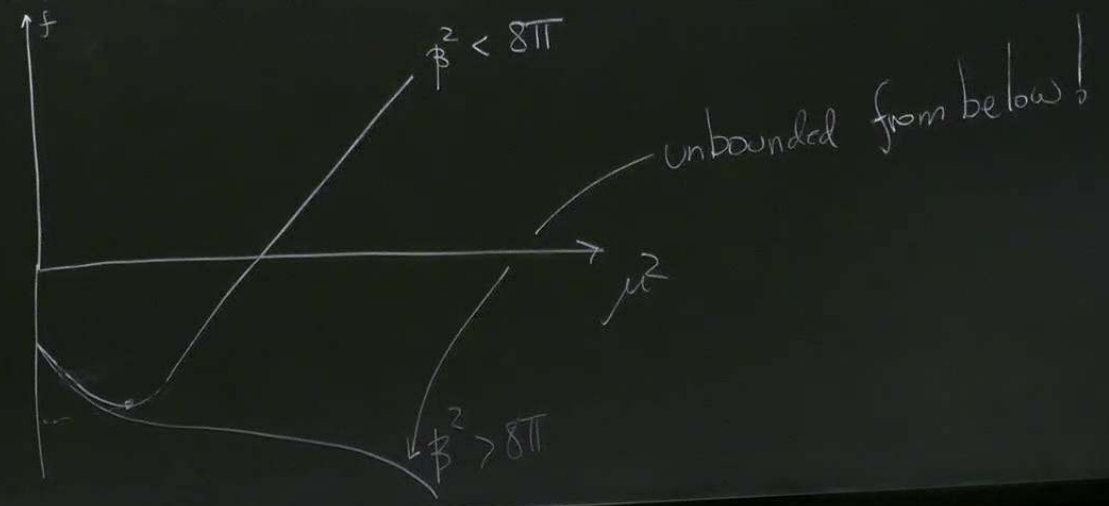


Scalar mass  $\mu$ ,  $a_{K,\mu} |0,\mu\rangle = 0$

↑ value for  $\alpha$

Same for SG:

$$\langle 0, \mu | : \mathcal{H}_m : | 0, \mu \rangle = - \frac{\alpha}{\beta^2} \left( \frac{\mu^2}{m^2} \right)^{\beta^2/8\pi} + \frac{\mu^2 - m^2}{8\pi} - \gamma$$



from pt of view of SG = CFT + pert.

$$\left\langle e^{i\beta\phi(x_1)} e^{-i\beta\phi(x_2)} \right\rangle_{\text{CFT}} =$$

$$\frac{1}{|x_1 - x_2|^{\frac{\beta^2}{2\pi}}}$$

$$2\Delta_0 = e^{i\beta\phi}$$

$d = D + 1$   
irrelevant

$$\Delta_0 = \frac{\beta^2}{4\pi}$$

$> 2$  if  $\beta^2 > 8\pi$

$< 2$  if  $\beta^2 < 8\pi$

↑ relevant



comment

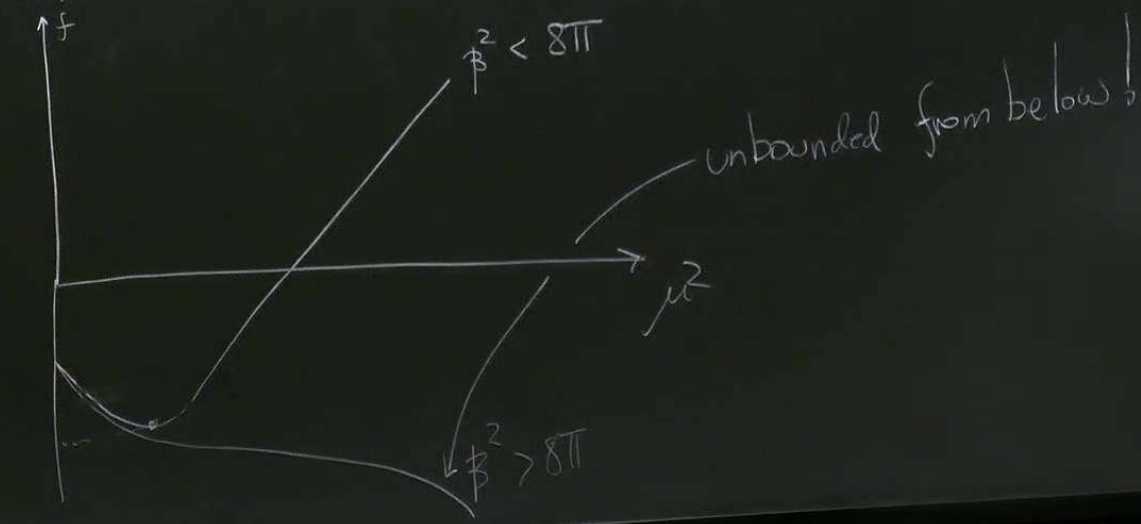
$$\langle e^{i\beta\phi(x_1)} e^{-i\beta\phi(x_2)} \rangle = \frac{1}{|x_1 - x_2|^{\beta^2/2\pi}} \left( 1 + \# \int \frac{dx_3 dx_4}{|x_3 - x_4|^{\beta^2/2\pi}} \right)$$

CFT  
 + part  
 ↑  
 does nothing if  $\beta^2 > 8\pi$

↑  
 get this at  $\beta^2 > 8\pi$   
 the part was irrelevant!

↓  
 converges for  $\beta^2 > 8\pi$   
 if  $\beta^2 < 8\pi$

$$|0, \mu\rangle = - \frac{\alpha}{\beta^2} \left( \frac{\mu^2}{m^2} \right)^{\beta^2/8\pi} + \frac{\mu^2 - m^2}{8\pi} - \gamma$$



from pt of view of

$$\langle e^{i\beta\phi(x_1)} e^{-i\beta\phi(x_2)} \rangle$$

$\begin{matrix} \text{---} \\ \text{---} \end{matrix} = \text{soliton} \longleftrightarrow \Psi$   
 $\begin{matrix} \text{---} \\ \text{---} \end{matrix} = \text{sg} \longleftrightarrow \Psi \Psi \text{ bound state}$

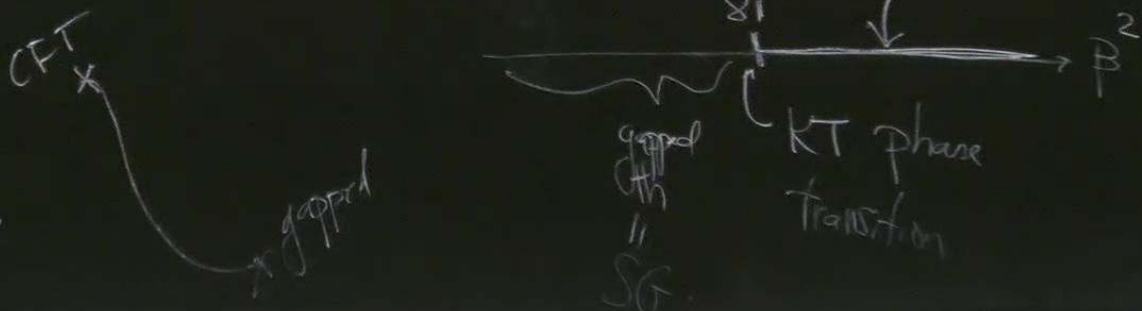
$$\left( \frac{(X_1 - X_3)^2 (X_2 - X_4)^2}{(X_1 - X_4)^2 (X_2 - X_3)^2} \right)^{\beta/4\pi} + \dots = \frac{N}{|X_1 - X_2|} \left( \frac{\beta}{8\pi} \right)^{\beta/8\pi}$$

at  $|X_1 - X_2| \gg 1$  set to 1

$\langle e^{i\beta\phi_1} e^{-i\beta\phi_2} e^{i\beta\phi_3} e^{-i\beta\phi_4} \rangle$

Physical  $\leftarrow \beta^2 > 8\pi$   
 Physical!  $\leftarrow \beta/8\pi$

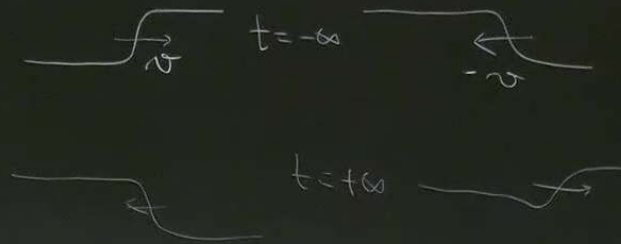
$\beta^2 > 8\pi$   
 $< 8\pi$ , need RG



irrelevant!

2 Lumps

$$\phi(x,t) = \frac{4}{\beta} \operatorname{arctg} \left( \frac{\eta \sin \omega t}{\cosh \eta x} \right)$$



$\sqrt{1-v^2}$  like  $\operatorname{erfc} < 1$   
also  $v = i v$

irrelevant!

If  $\beta^2 < 8\pi$ , need RG too  $\rightarrow$

gapped

actg  $\left( \frac{\eta \sin \omega t}{\cosh \eta \omega} \right)$  ← oscillates w/ frequency  $\omega$

Where  $\eta = \frac{\sqrt{\alpha^2 - \omega^2}}{\omega}$ ,  $\omega \in [0, \alpha]$

$\sqrt{1-v^2}$  likes  $\beta < 1$   
also  $v = i v$

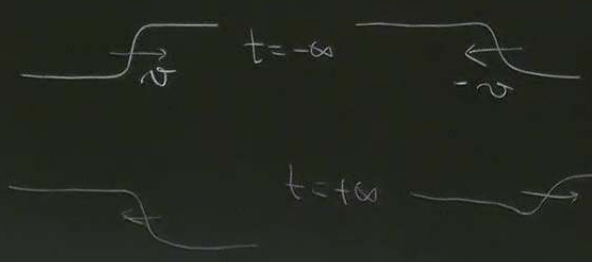
irrelevant!

$4P^2 < 8$

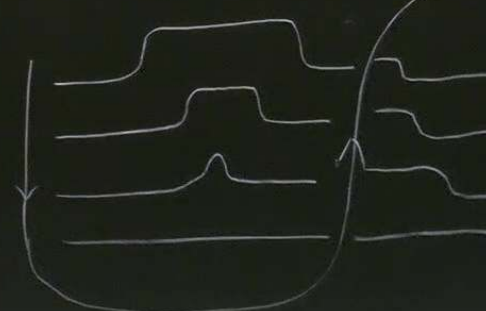
2 Lumps

$$\phi(x,t) = \frac{4}{\beta} \operatorname{arctg} \left( \frac{\eta \sin \omega t}{\cosh \eta x} \right)$$

Where  $\omega$  oscillates w/ frequency



$\sqrt{1-v^2}$  like  $\operatorname{erfc} < 1$   
 also  $v = i v$



$\leftarrow 8\pi$ , need RG soon  $\rightarrow$

gapped

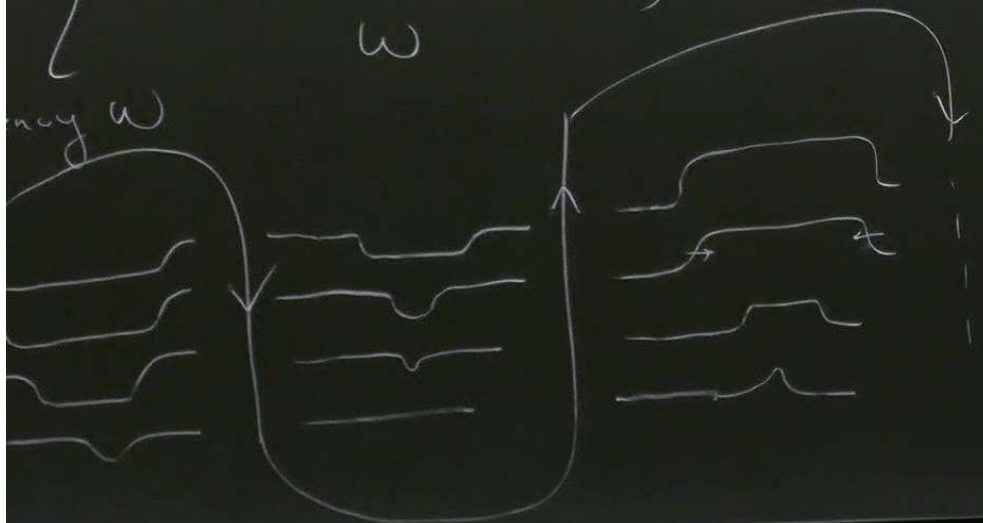
SG

transition

$$\eta = \frac{\sqrt{\alpha^2 - \omega^2}}{\omega}$$

may  $\omega$

$$\omega \in [0, \alpha]$$



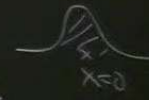
breather



$$E = \frac{1}{2} \int dx \dot{\phi}^2(x, 0)$$

$$= \frac{8\eta^2 \omega^2}{\beta^2} \int \frac{dx}{\cosh^2(\eta \omega x)}$$

$$E = \frac{16\eta \omega}{\beta^2}$$



\*  $\beta^2 = 4\pi$   
no breather!

\*  $\beta^2 > 4\pi$   
 $E = m n$  for  $n=1$

$\beta \ll 1$   
- #  $\beta^2 (n^3 - n)$   
 $> 0$

$n=2$ : second excited state

$n=3$ : third excited state

= band states

$$E(\omega) = 2 \underbrace{\left( \frac{8\sqrt{\alpha}}{\beta^2} \right)}_{M_{\text{Lump}}} \left[ 1 - \frac{\omega^2}{\alpha} \right]^{\frac{1}{2}} < 2M_{\text{Lump}}$$

B.S. :  $\int dE_n = \omega(E) dn$  Exact!

$$\sin\left(\frac{\beta^2 n}{16}\right)$$

$$E_n = 2M_{\text{Lump}}$$

$n = 0, 1, 2, 3, \dots$   
↑ Lightest breather

BS  
"WKB" DN ← prime

