

**Title:** Building and (hints of) seeing gravitational statistical mechanics

**Speakers:** Seth Major

**Collection/Series:** Lee's Fest: Quantum Gravity and the Nature of Time

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**Abstract:**

A statistical mechanics of geometry based on the quasi-local energy of constantly accelerating observers outside black holes is proposed. By assuming particle statistics for the geometric quanta, the geometry condenses in the large area limit, providing a quantum gravity cutoff for the entropy of the thermal atmosphere of geometric excitations. These excitations model the fluctuating ``shape of the black hole" and form an effectively two dimensional quantum thermal geometric atmosphere. The entropy from these ``Debye" excitations is proportional to the area. Prospects of for the observation of the fluctuations are briefly discussed.

# Building and (hints of) seeing gravitational statistical mechanics

Seth Major

3 June 2025

Lee's Fest @ PI

*Thank you Lee!*



• What is the gravitational statistical mechanics for observers maneuvering to ensure that their proper acceleration is constant?

- For high acceleration observers around black holes

there is simple form of the quasi-local energy for such observers

$$E_g \simeq \frac{A}{8\pi\ell} \simeq \frac{gA}{8\pi}$$

with  $\ell$  proper distance to horizon  
and  $g$  the proper acceleration

Frodden, Ghosh, Perez, Phys. Rev. D 87 (2013) 121503(R)

See also “Black Holes: The Membrane Paradigm”

S. Major Phys Rev D 105 (2022) 104050

# Statistical Mechanics Model

- Two-stage statistical model in high  $g$ , large area limit in Schwarzschild spacetime

I - Use quasi-local energy for constantly accelerating observers

- Assume quantization of area (and thus energy)

- Assume **quantum statistics** for the atoms of geometry

- Model gas of indistinguishable, non-interacting atoms of geometry

- Condenses in large area limit to the  $j = 1$  state

- Entropy of condensed gas is low

II (Sketch of work in progress) `quasi-normal' phonons - Debye model

- The phonon frequency is cutoff by the condensed state area

- Entropy of phonons is proportional to area

- Thoughts on possible observations

# Statistical Model - I

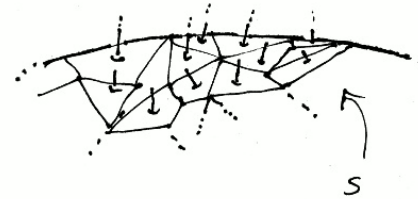
Schwarzschild, large  $g$ ,

- With spin network states observers surface has area

$$\hat{A}_S | \vec{j} \rangle = 8\pi l_{eff}^2 \sum_{i \in S \cap G} \sqrt{j_i(j_i + 1)} | \vec{j} \rangle$$

where  $l_{eff} = \sqrt{\hbar\gamma} *$  ( $G = c = 1$ )

$$\hat{H}_g | \vec{j} \rangle = \frac{g\hat{A}_S}{8\pi} | \vec{j} \rangle = gl_{eff}^2 \sum_{i \in S \cap G} \sqrt{j_i(j_i + 1)} | \vec{j} \rangle$$



- A gas of indistinguishable, non-interacting quantum geometric particles

Each geometric particle has degeneracy

$$d_j = 2j + 1$$

(**No** 'holographic' degeneracy  $\exp(\lambda A/l_p^2)$ )

Asin et. al., Ghosh et. al.

Partition function with non-vanishing chemical potentials  $\mu_{\pm}$

$$Z_{\pm} = \prod_j^{\infty} \left( 1 \pm e^{-\beta(\epsilon\sqrt{j(j+1)} - \mu_{\pm})} \right)^{\pm(2j+1)} \quad \text{with } \epsilon = gl_{eff}^2$$

# Statistical Model - I

Schwarzschild, large  $g$ , large  $A$  limit

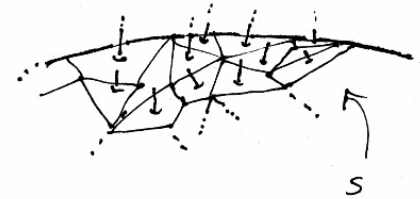
- Observers are at the Unruh temperature

$$\beta \simeq \left( \frac{2\pi}{g \hbar} \right)$$

At Unruh temperature the Boltzmann weight scales with

$$\beta \epsilon = 2\pi\gamma$$

The Barbero-Immirzi parameter acts like an inverse temperature



- Bosonic area

(Fermionic excitations sub-dominant.)

$$\frac{\langle A \rangle_b}{8\pi l_{eff}^2} = \sum_{j=1}^{\infty} \frac{(2j+1)\sqrt{j(j+1)}}{e^{2\pi\gamma(\sqrt{j(j+1)} - \tilde{\mu}_b)} - 1}$$

Positivity requires maximum chemical potential with

$$\tilde{\mu}_{max} = \frac{\mu_b}{\epsilon} = \sqrt{2} - \delta \quad \langle A \rangle_b \simeq \frac{12\sqrt{2}l_P^2}{\delta}$$



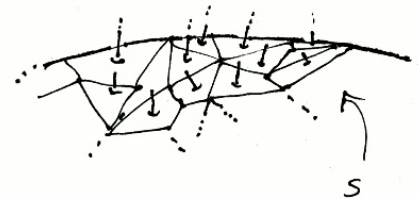
Large area requires condensation  $\delta \rightarrow 0$  (or at least suitably small)

# Statistical Model - I

Schwarzschild, large  $g$ , large  $A$  limit

The free energy is  $\Omega = -\frac{1}{\beta} \ln Z$

Entropy:  $S = \beta^2 \partial_{\beta} \Omega = 3 \ln \left( \frac{\langle A \rangle_b}{24\sqrt{2}\pi l_{eff}^2} \right)$



In condensed state the entropy is low, sub-dominant

**Story so far:** a low entropy geometric 'condensate' with area quanta  $a_1$  and

$$a_1 = 8\pi\sqrt{2}l_{eff}^2$$
$$A \simeq Na_1$$

See also Asin et. al.

See also Major and Setter

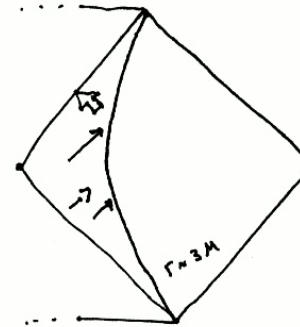
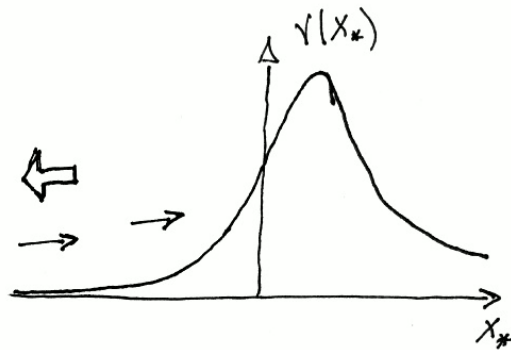
Dreyer PRL 2003

# Statistical Model - II

Schwarzschild, large  $g$ , large  $A$  limit

- Thermal or quantum atmosphere via Hawking process

York, 't Hooft, Frolov and Novikov, Jacobson and Parentani, ...



- These 'phonon modes' contribute entropy in near horizon geometry.

- Asymptotic QNM spectrum is  $\omega_n \simeq \frac{1}{8\pi M} \left( \ln 3 + 2\pi i \left( n + \frac{1}{2} \right) \right)$  Motl, Neitzke

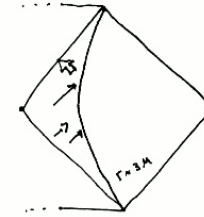
# Statistical Model - II

Schwarzschild, large  $g$ , large  $A$  limit

- Energy of quasi-normal-like modes at energy  $\Delta E = \hbar \frac{\ln 3}{8\pi M} \rightarrow \Delta E_{local} = \hbar \frac{g \ln 3}{2\pi}$  Hod PRL 1998

Fixes upper bound on phonon modes

2D Debye model has a Debye (max) frequency of  $\omega_D = \frac{2}{L} \sqrt{\frac{\pi A}{a_1}} = \frac{g \ln 3}{2\pi}$



Geometric 'phonon' partition function  $Z_P = \prod_{n=0}^{n_{max}} \frac{1}{e^{\beta \hbar \omega_n} - 1}$

Helmholtz free energy  $F = -\frac{1}{\beta} \ln Z_P$

- The phonon entropy is

$$S_P = \beta^2 \partial_{\beta} F \simeq \left( \frac{4A}{a_1} \right) \left( \frac{1}{x_D^2} \right) I(x_D) \quad \text{with } x_D = \beta \hbar \omega_D = \ln 3$$

and 
$$I(x_D) = \int_0^{x_D} dx \left[ \frac{x^2}{e^x - 1} - x \ln(1 - e^{-x}) \right]$$

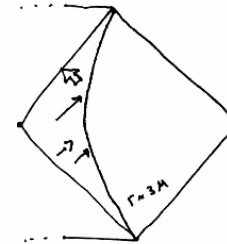
## Statistical Model - II

Schwarzschild, large  $g$ , large  $A$  limit

Observers experience thermal bath at Unruh temperature thus

$$x_D = \beta_U \hbar \omega_D = \ln 3 \text{ and } I(x_D) \simeq 0.863$$

$$\Rightarrow S_P \simeq \frac{4 \cdot 0.863 A}{(\ln 3)^3 a_1} \simeq 0.294 \frac{A}{l_P^2} *$$

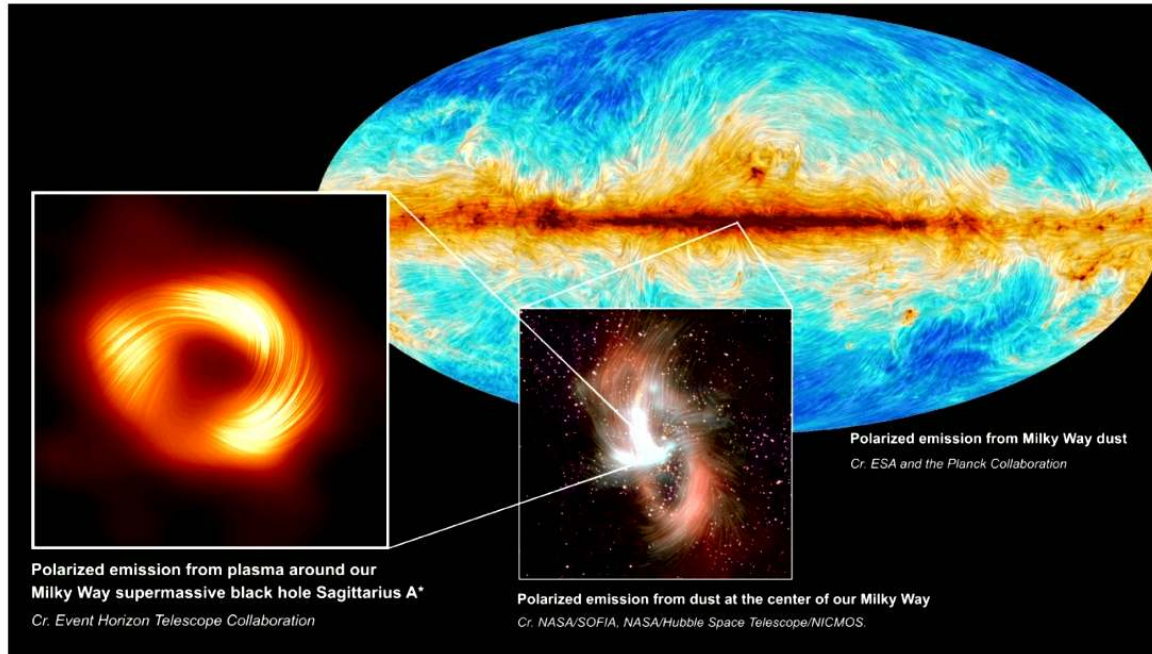


Not surprising from earlier studies by Frolov and Novikov, 't Hooft, and others but still nice :-)

These fluctuations in geometry are **outside** the black hole. They may be observable.

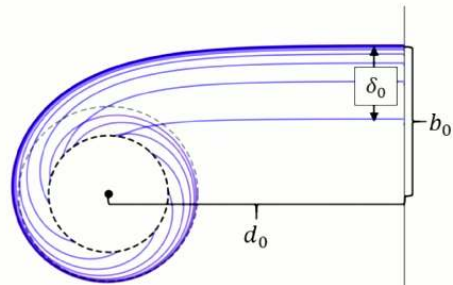
# 'Seeing' black hole entropy

- Event Horizon Telescope images (and movies) of Sgr A\*



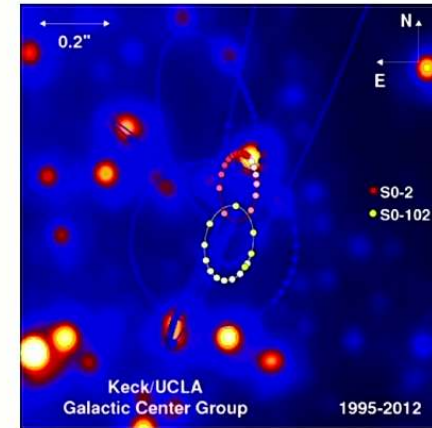
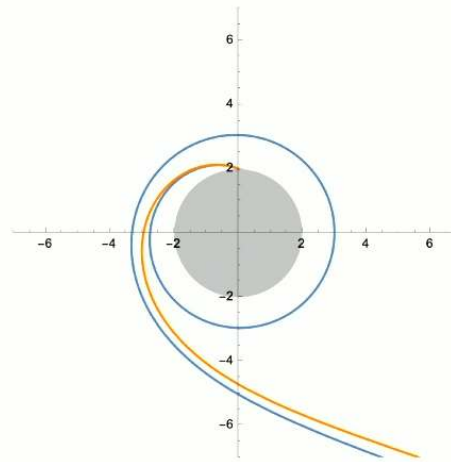
# 'Seeing' black hole entropy

- Flare from accreting matter around black holes e.g. Sgr A\*



Sneppen 2107.04044

Characteristic fluctuation  $\omega_\infty \sim \frac{1}{M}$



- Photon ring sub-structure via very long baseline interferometry

Johnson et al Sci. Adv. 2020

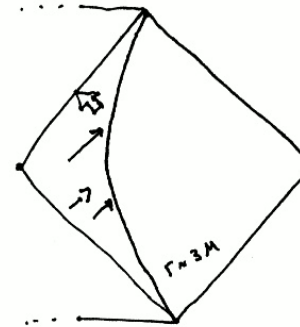
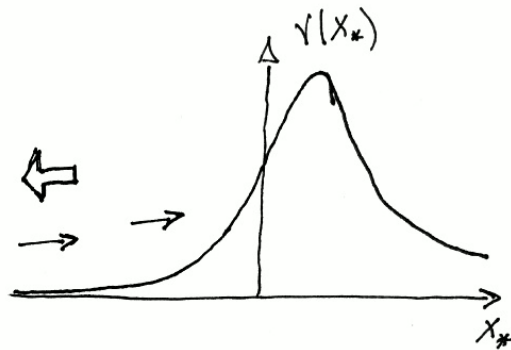
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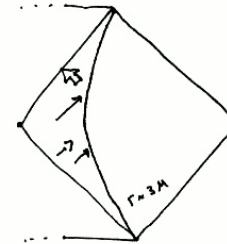
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