

Title: Long-term stable non-linear evolutions of ultracompact black hole mimickers

Speakers: Seppe Staelens

Collection/Series: Strong Gravity

Subject: Strong Gravity

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Abstract:

Ultracompact black hole mimickers formed through gravitational collapse under reasonable assumptions obtain light rings in pairs, where one is unstable and the other one is not. Stable light rings are believed to be a potential source for dynamical instability due to the trapping of massless perturbations, as their decay is relatively slow.

We study the stability of ultracompact boson stars admitting light rings combining a perturbative analysis with 3+1 numerical-relativity simulations with and without symmetry assumptions. We observe excellent agreement between all perturbative and numerical results which uniformly support the hypothesis that this family of black-hole mimickers is separated into stable and unstable branches by extremal-mass configurations. This separation includes, in particular, thin-shell boson stars with light rings located on the stable branch which we conclude to represent long-term stable black-hole mimickers. Our simulations suggest that the proposed mechanism may not be efficient after all to effectively destroy ultracompact black hole mimickers.

Stable evolutions of ultracompact black hole mimickers

(2504.17775)

Seppe J. Staelens

- with G. Marks, T. Evstafyeva, U. Sperhake -

Strong Gravity Seminar, Perimeter Institute for Theoretical Physics



UNIVERSITY OF
CAMBRIDGE



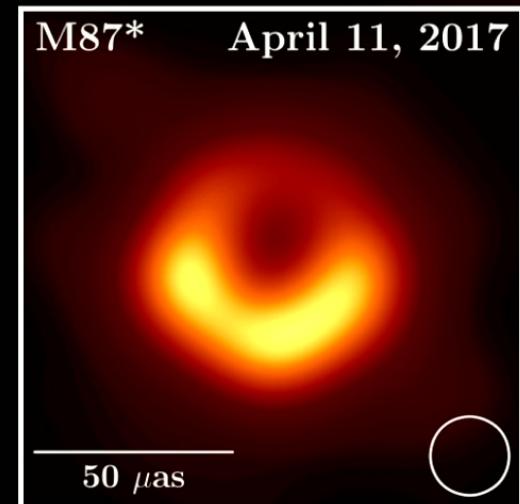
St. Edmund's College

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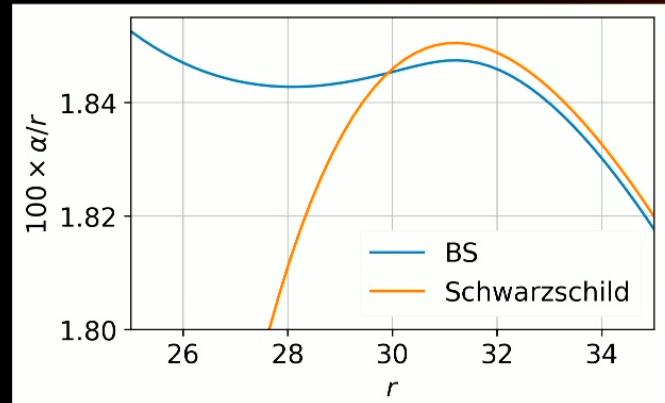
Introduction

- Black hole mimickers
 - Horizonless / regular
 - Dark matter ?
- Light ring vs. Event horizon
 - Ringdown
 - VLBI images
- Ultracompact objects



Introduction

- Light rings (LRs) come in pairs (Cunha+, '17)
 - Stable + unstable
- Stable LR
 - Slow decay of linear waves (Keir, '16)
 - Trapping well
- LR instability is still debated (Cardoso+, '14) (Cunha+, '23)
 - (Guo+, '24) (Bonomio+, '24)
 - (Siemonsen, '24) (Redondo-Yuste+, '25)



Contents

- Boson stars
- Radial stability
- Numerical simulations
- Effective potential & light rings
- Teaser: VLBI images

Boson stars

Boson stars

- GR + complex scalar field

$$S = \int \left(\frac{R}{16\pi G} + \mathcal{L}_\varphi \right) \sqrt{-g} \, d^4x, \quad \mathcal{L}_\varphi = -\frac{1}{2} \left(\nabla_\alpha \bar{\varphi} \nabla^\alpha \varphi + V(\varphi) \right)$$

- Einstein-Klein-Gordon equations

$$G_{\alpha\beta} = 8\pi G T_{\alpha\beta}, \quad \nabla_\mu \nabla^\mu \varphi = \frac{dV}{d\bar{\varphi}}$$

Boson stars

(Liebling+, '23)

- GR + complex scalar field

$$S = \int \left(\frac{R}{16\pi G} + \mathcal{L}_\varphi \right) \sqrt{-g} \, d^4x, \quad \mathcal{L}_\varphi = -\frac{1}{2} \left(\nabla_\alpha \bar{\varphi} \nabla^\alpha \varphi + V(\varphi) \right)$$

- Einstein-Klein-Gordon equations

- Scalar potential

- Mini

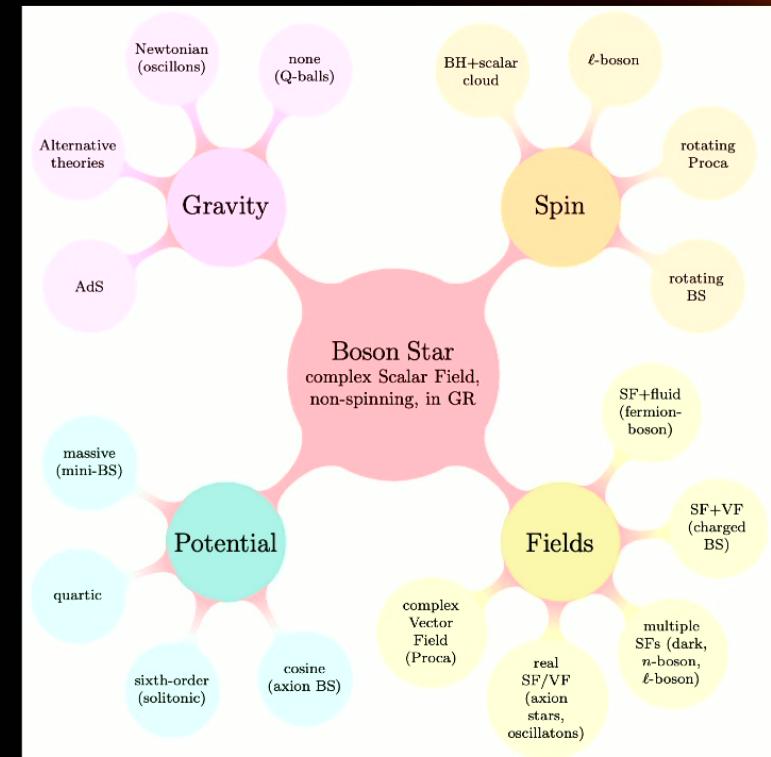
$$V(\varphi) = \mu^2 |\varphi|^2$$

- Solitonic

$$V(\varphi) = \mu^2 |\varphi|^2 \left(1 - 2 \frac{|\varphi|^2}{\sigma^2} \right)^2$$

- Axion (real scalar)

$$V(\varphi) = \mu^2 f^2 \left| 1 - \cos \frac{\varphi}{f} \right|$$



Why boson stars?

- BH mimickers
 - Black if no coupling to EM
 - Regular and no horizon
 - Varying compactness
- Only one additional complex scalar
- ‘Easy’ to simulate
- Good first-order approximation of non-BH?

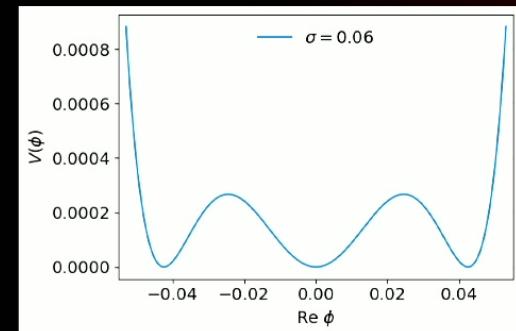
Solitonic boson stars

- Solitonic potential
 - Parameter σ

$$V(\varphi) = \mu^2 |\varphi|^2 \left(1 - 2 \frac{|\varphi|^2}{\sigma^2}\right)^2$$

- Spherically symmetric, harmonic Ansatz

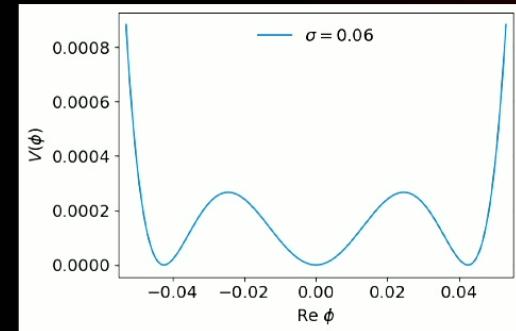
$$\varphi = A(r) e^{i\omega t}$$



Solitonic boson stars

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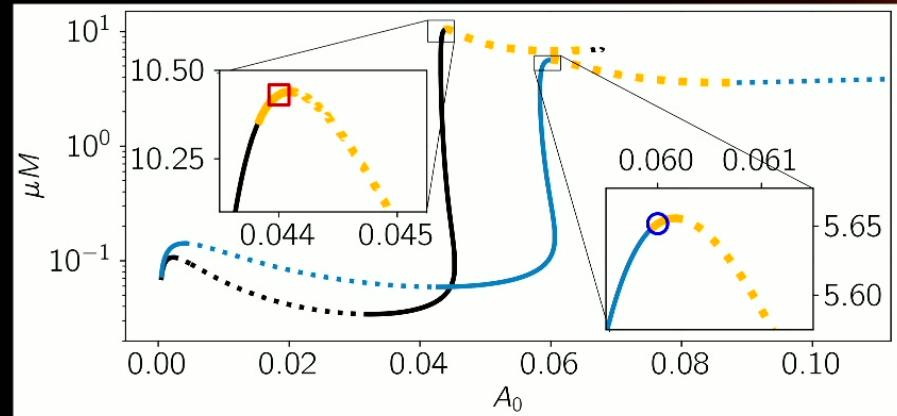
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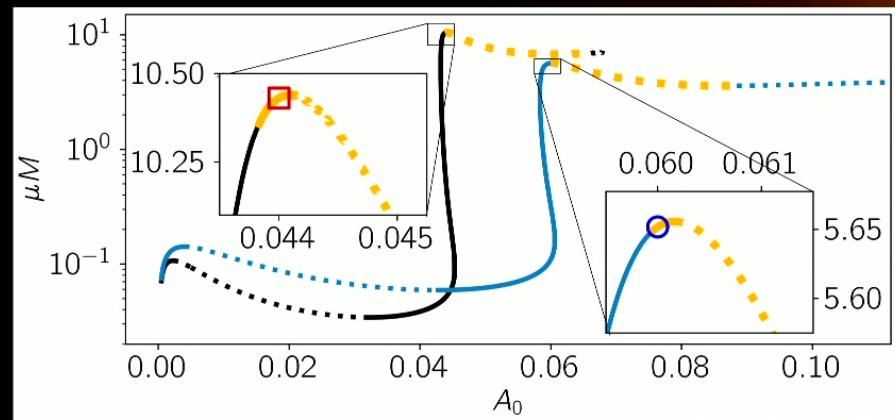
$$\varphi = A(r) e^{i\omega t}$$

- Shooting method
 - Fix central amplitude A_0
 - Guess ω
 - Integrate EKG equations outwards
 - Repeat for different ω



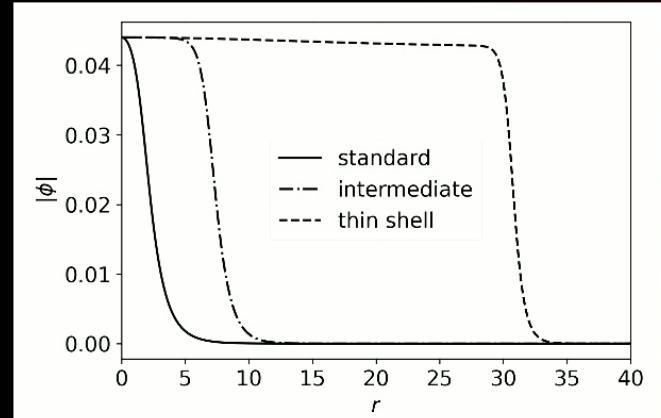
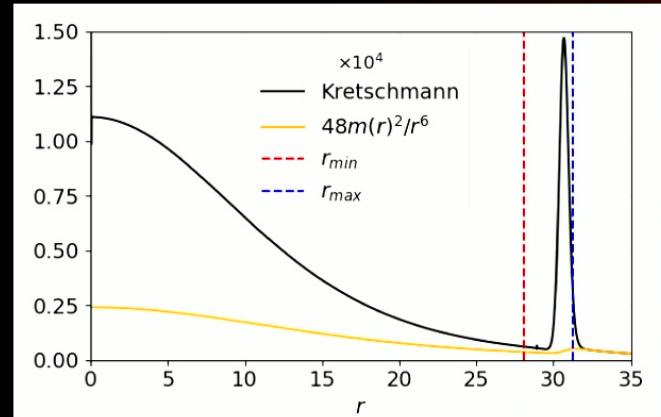
Thin shell models

- For small σ
 - Multivalued
 - Ultracompact models (orange)



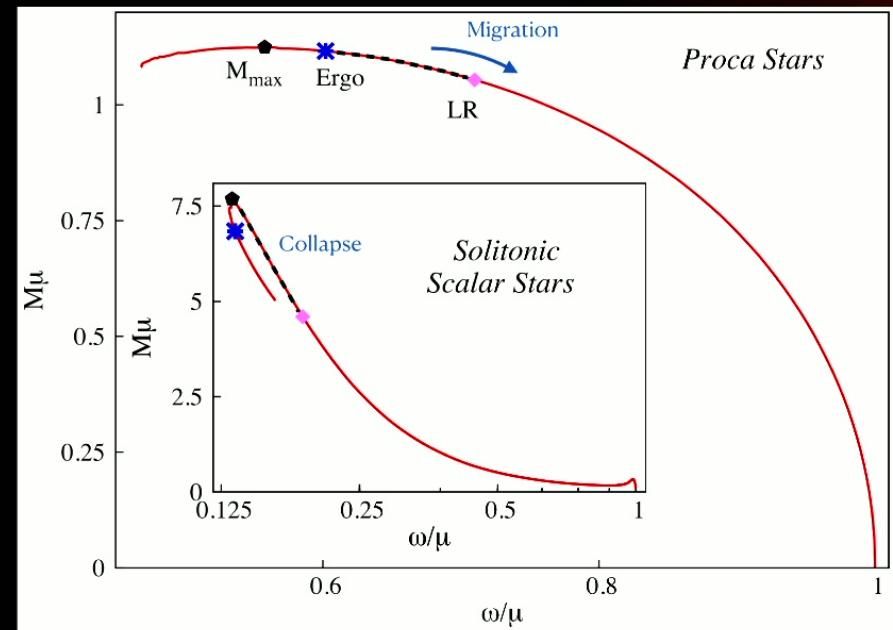
Thin shell models

- For small σ
 - Multivalued
 - Ultracompact models (orange)
- Thin shell models
 - Center of BS close to non-zero vacuum
 - Energy concentrated in a shell
 - \approx Schwarzschild outside of the shell



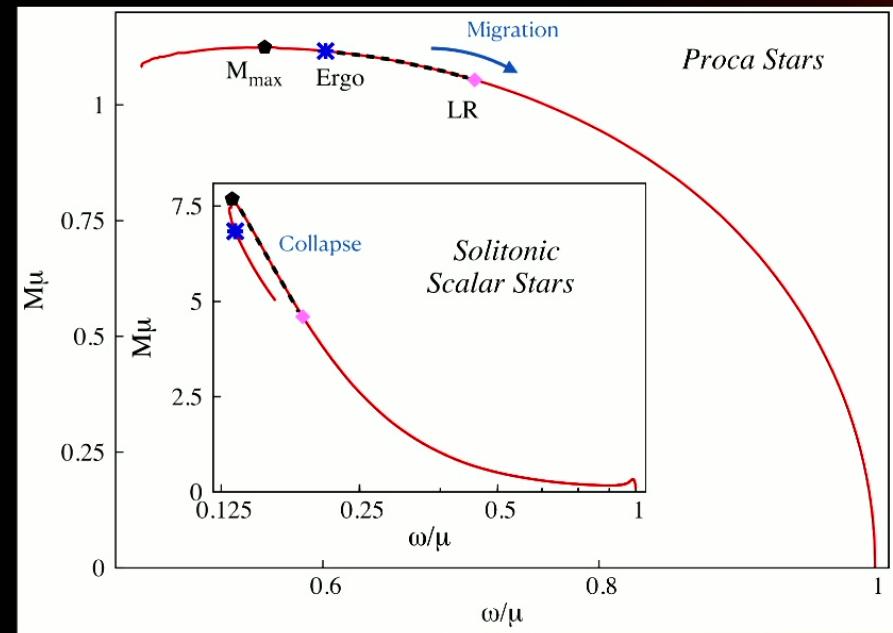
Instability: past work

- (Cunha+, '23)
- BS and Proca stars
- Find instability onset at LR regime
- Difference: rotating models



Instability: past work

- (Cunha+, '23)
- BS and Proca stars
- Find instability onset at LR regime
- Difference: rotating models
- Our motivation
 - Simplify by neglecting rotation
 - If unstable: non-spherical instability?
 - Analytical arguments for / against



Radial Perturbative Stability (Gareth's work)

Theory

- Linear stability against radial perturbations
- Expand EKG equations around background BS

$$ds^2 = -\alpha^2(r, t) dt^2 + X^2(r, t) dr^2 + r^2 d\Omega^2$$

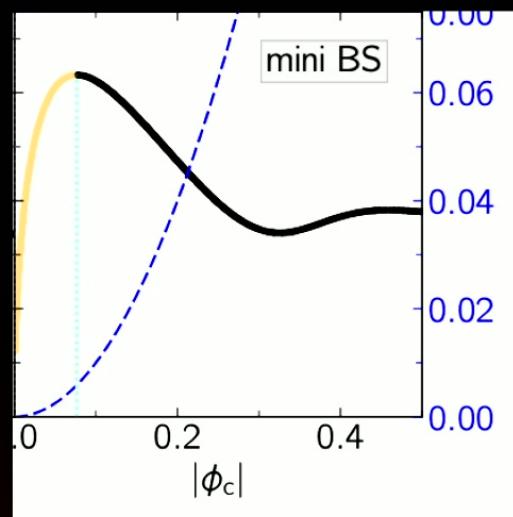
$$\begin{aligned}\varphi &= e^{i\omega t} (\psi_1 + i\psi_2), \quad \psi_1 = A (1 + \delta\psi_1), \quad \psi_2 = A\delta\psi_2, \\ \alpha &= \alpha_0 \left(1 + \frac{1}{2}\delta\nu \right), \quad X = X_0 \left(1 + \frac{1}{2}\delta\lambda \right)\end{aligned}$$

Theory

- Assume perturbations of the form $\delta f(r, t) = e^{i\chi t}g(r)$
- Result: system of two coupled ODEs
- Positive (negative) χ^2 corresponds to linear radial (in)stability

Theory

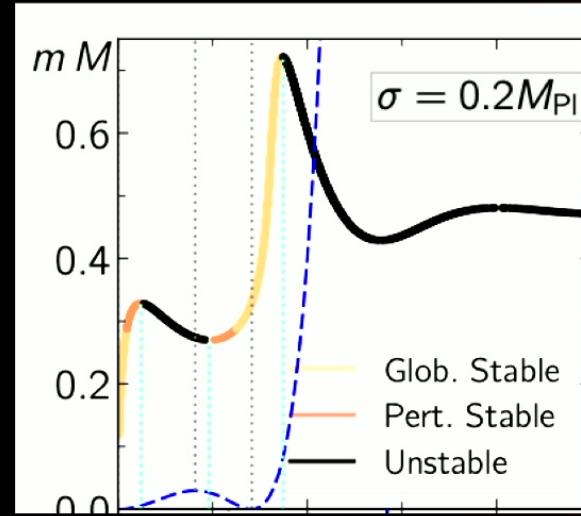
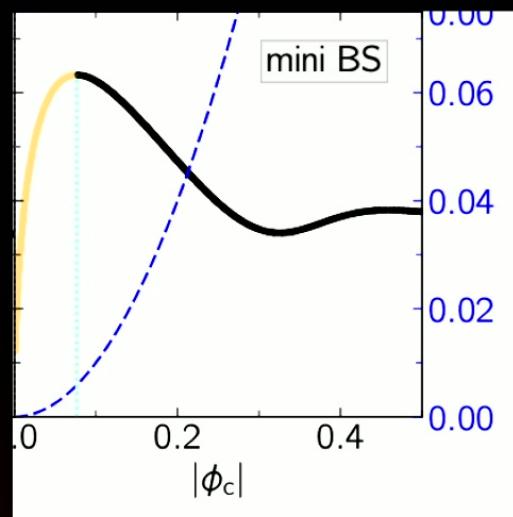
- (Un)stable branches are separated by extrema of $M(A_0)$



(Ge+, 2410.23839)

Theory

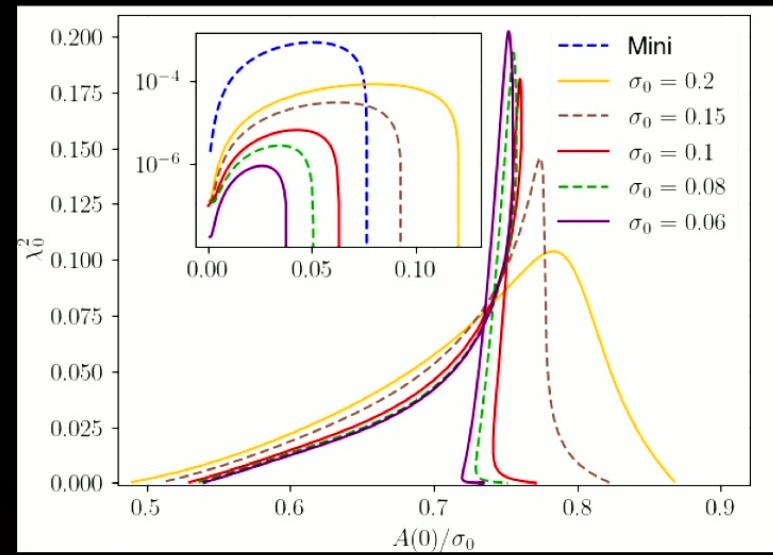
- (Un)stable branches are separated by extrema of $M(A_0)$



(Ge+, 2410.23839)

Gareth's results

- Calculated perturbation frequencies for solitonic families that include thin-shell models.
- Multi-stage shooting code with quadruple precision



Numerical Evolutions

Numerical Relativity

- 3+1 metric

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij} (dx^i + \beta^i dt) (dx^j + \beta^j dt)$$

- 6 spatial evolution equations

$$\partial_t \gamma_{ij} = -2\alpha K_{ij} + \dots, \quad \partial_t K_{ij} = \dots$$

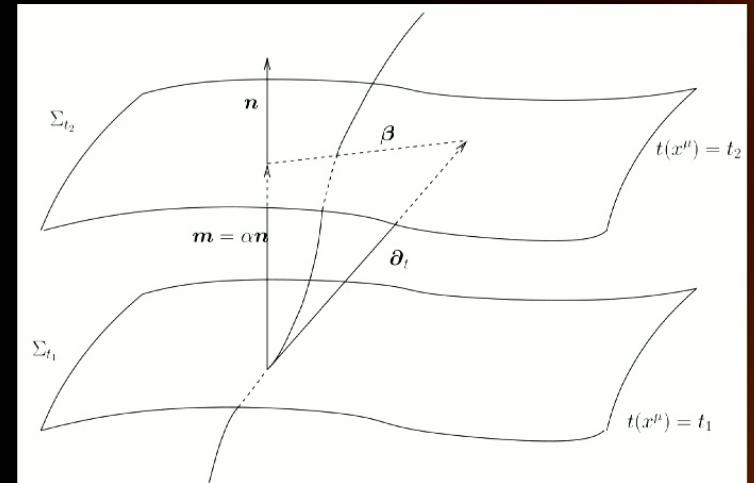
- 4 constraint equations

$$\mathcal{H} = 0, \quad \mathcal{M}_i = 0$$

- 4 evolution equations for freely chosen α, β^i

- In practice: BSSN / CCZ4 formulation

- Latter: Constraint damping κ_i



Numerical Relativity

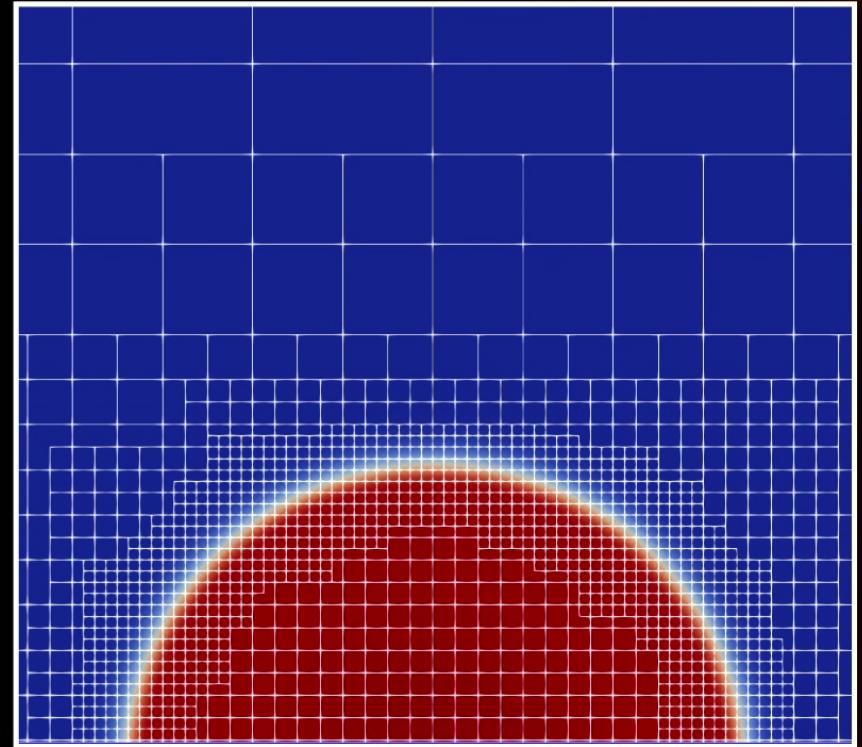
- Initial data: FORTRAN for quadruple precision
- 4 codes
 - ▶ Cartoon method
 - ▶ Spherical symmetry (GAM)
 - ▶ Axial symmetry: ExoZvezda (TE)
 - ▶ Full NR*
 - ▶ LEAN (US)
 - ▶ ExoZvezda (SJS)



3+1 simulations

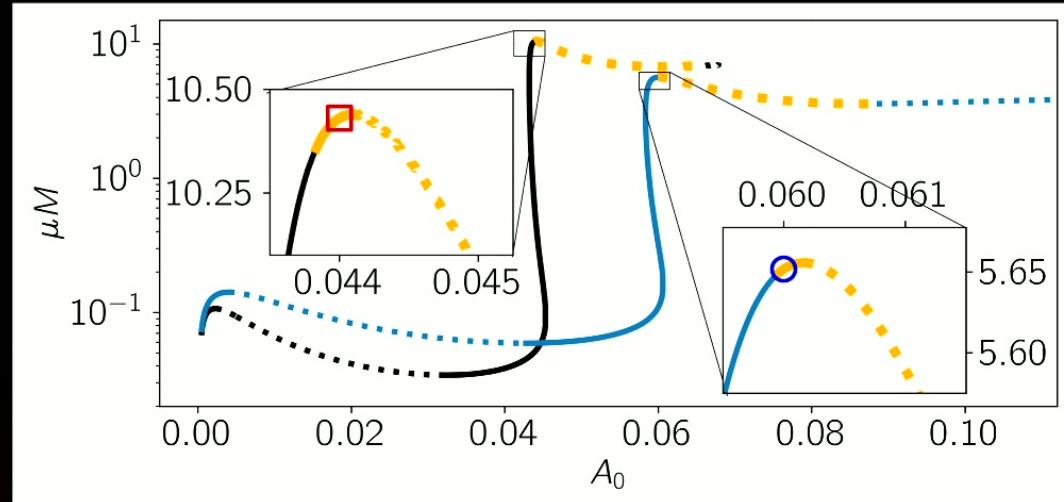


- **ExoZvezda**: Numerical Relativity with Adaptive Mesh Refinement (AMR)
- Challenging simulations
 - Initial data (quadruple precision)
 - Sensitive and slow
 - Separation of scales



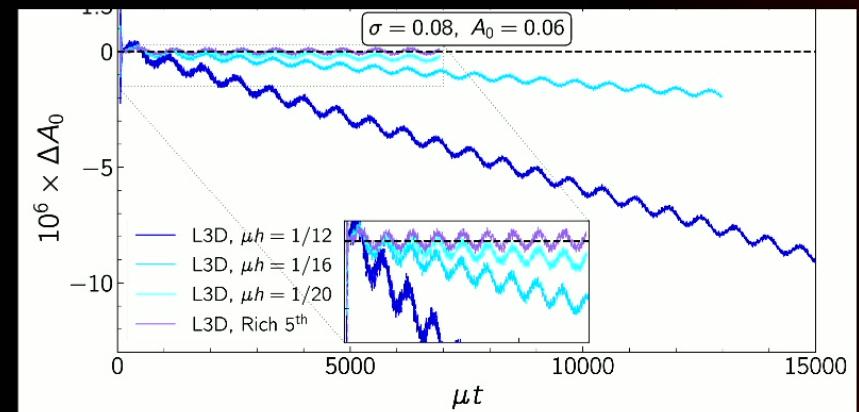
Reminder

- Two ultracompact thin-shell boson stars



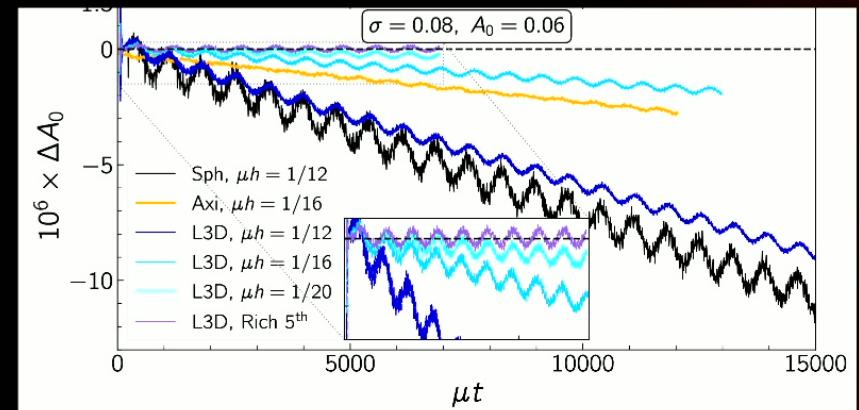
Central amplitude

- $\sigma = 0.08$: Evolution dominated
 - ▶ Observed drift converges away



Central amplitude

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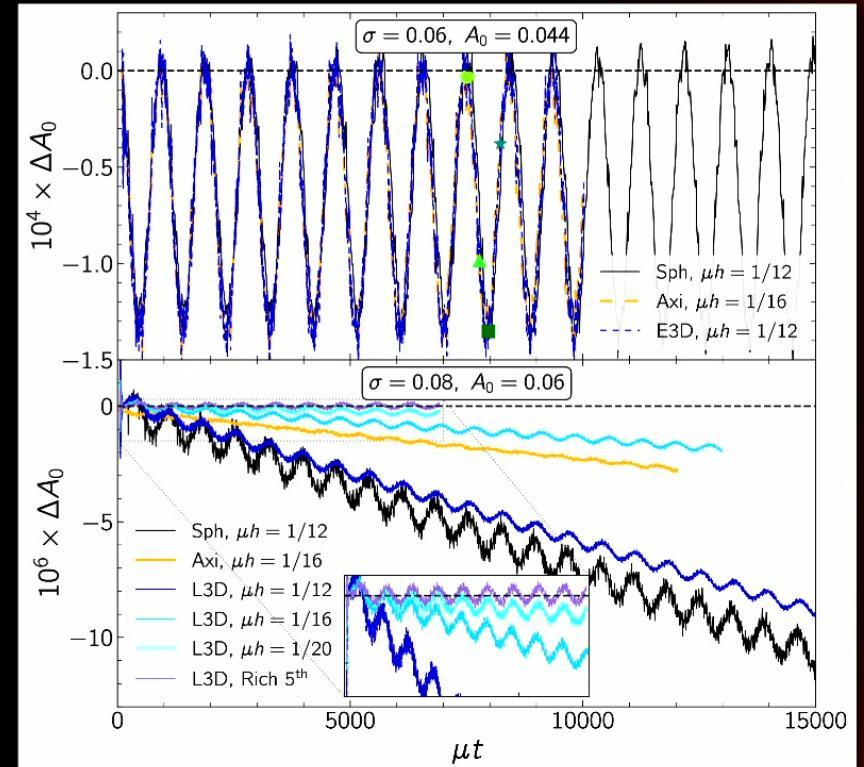


Central amplitude

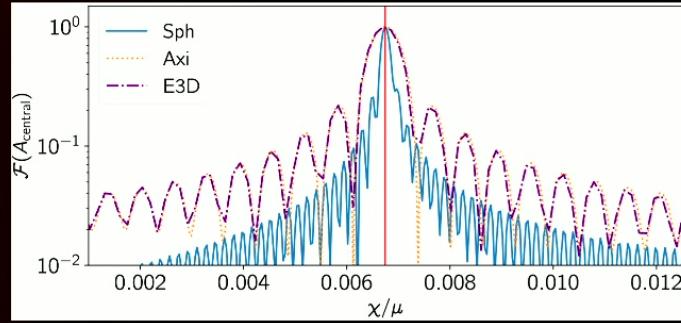
- $\sigma = 0.06$: Initial data dominated

BOTTOM LINE: Stability

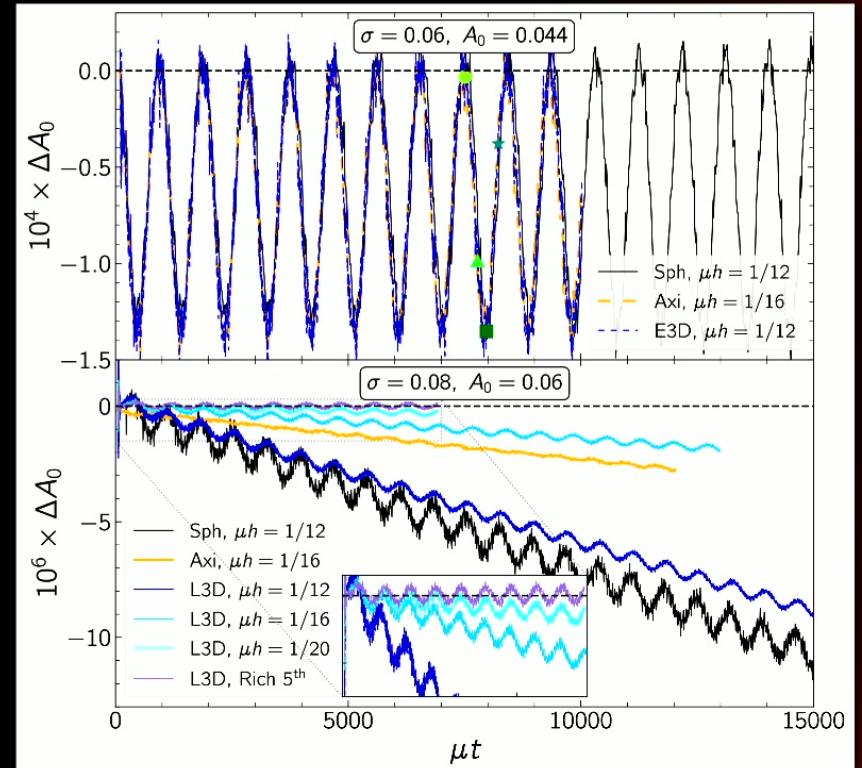
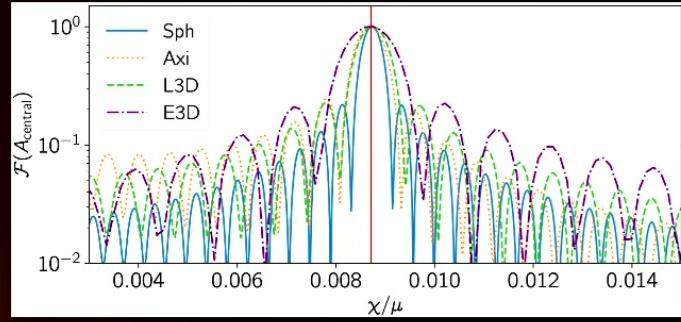
- $\sigma = 0.08$: Evolution dominated
 - Observed drift converges away



Oscillation Frequencies



Consistent with perturbation theory



But: is the BS still ultracompact?

Effective potential

- Stationary and spherical symmetric spacetime

$$ds^2 = -\alpha^2(r) dt^2 + X^2(r) dr^2 + r^2 d\Omega^2$$

- **Light rings:** circular null geodesics $\dot{r} = \ddot{r} = 0$

Effective potential

- Stationary and spherical symmetric spacetime

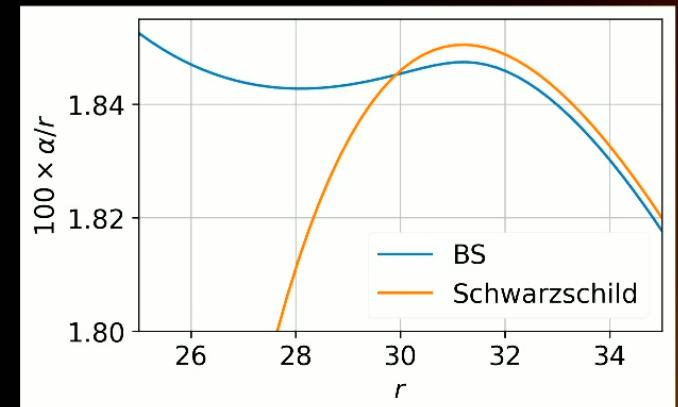
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- **Light rings:** circular null geodesics $\dot{r} = \ddot{r} = 0$

- Effective potential

$$V_{\text{eff}}(r) = \sqrt{\frac{-g_{tt}}{g_{\phi\phi}}} = \frac{\alpha}{r}, \quad \partial_\mu V_{\text{eff}} = 0$$

- (in)stability: second derivative



Adiabatic effective potential

- Numerical evolutions: departure from assumed symmetries
 - Assume slow evolution and small departures
- **Adiabatic Effective Potential (AEP)** (Cunha+, '23)

Adiabatic effective potential

- Numerical evolutions: departure from assumed symmetries
 - Assume slow evolution and small departures
- **Adiabatic Effective Potential (AEP)** (Cunha+, '23)
 - Average over coordinate spheres

$$V_{\text{eff}}(r) = \frac{\alpha}{r} \quad \longrightarrow$$

$$H_{\text{eff}}(r) = \sqrt{\frac{\langle \alpha^2 \rangle - \langle \gamma_{ij} \beta^i \beta^j \rangle}{(\oint \sqrt{\det q})/4\pi}}$$

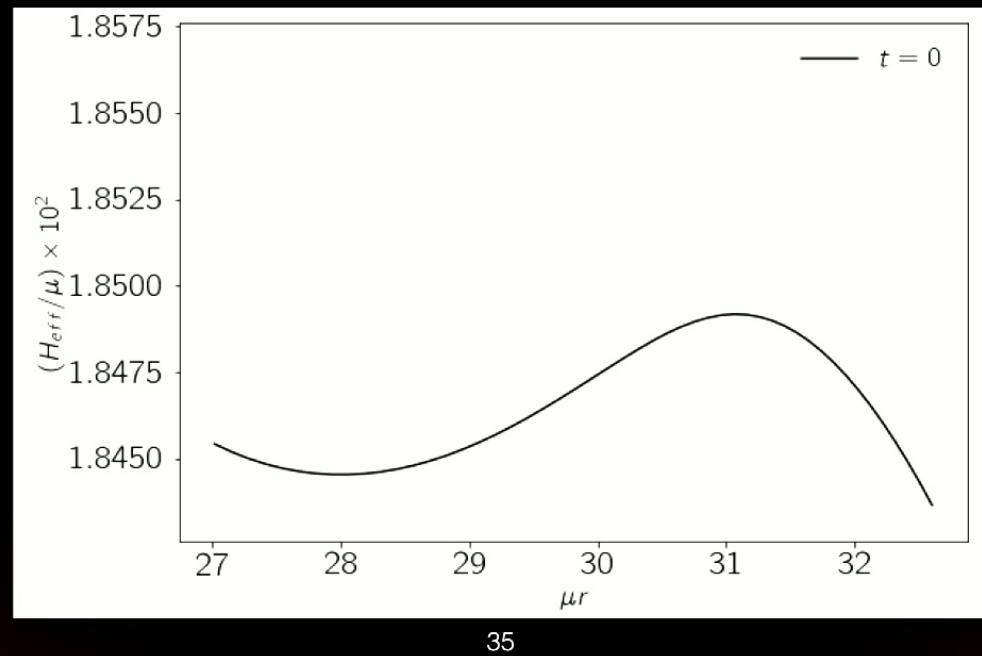
The diagram illustrates the decomposition of the Adiabatic Effective Potential $H_{\text{eff}}(r)$ into four components. The potential is represented by a vector pointing towards the right. Four green arrows originate from the tip of this vector and point towards the right, each labeled with a component name:

- Lapse
- Spatial metric
- Shift
- Induced metric on sphere

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Adiabatic effective potential

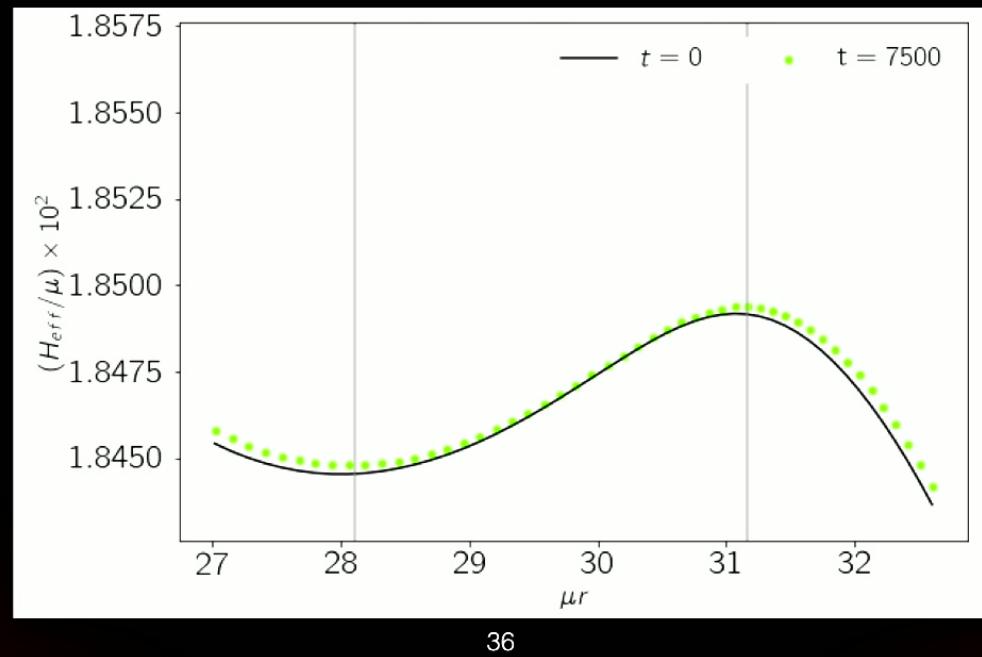
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$$H_{\text{eff}}(r) = \sqrt{\frac{\langle \alpha^2 \rangle - \langle \gamma_{ij} \beta^i \beta^j \rangle}{(\phi \sqrt{\det q})/4\pi}}$$

Adiabatic effective potential

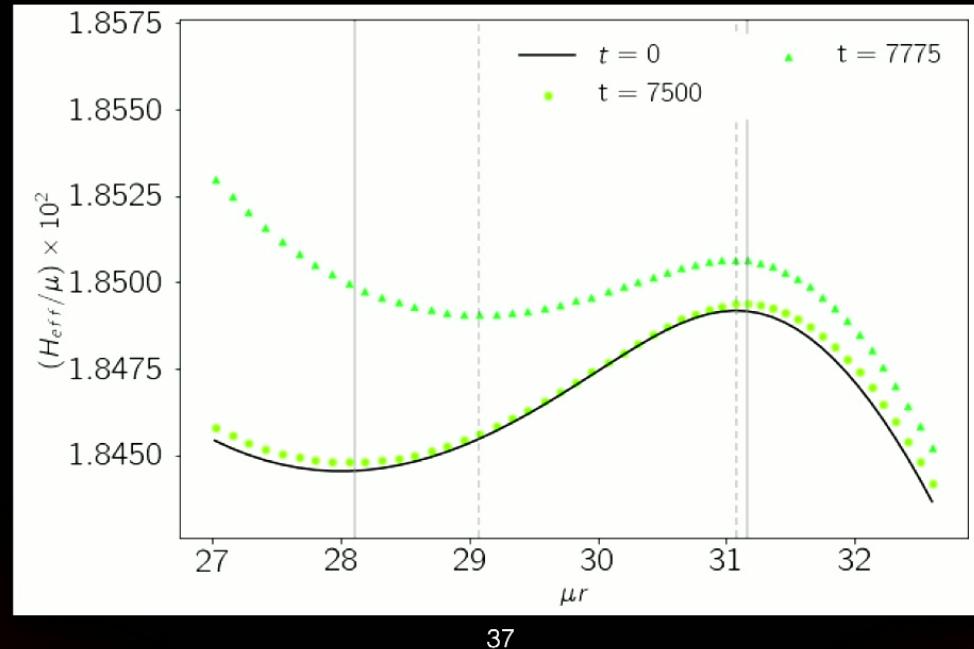
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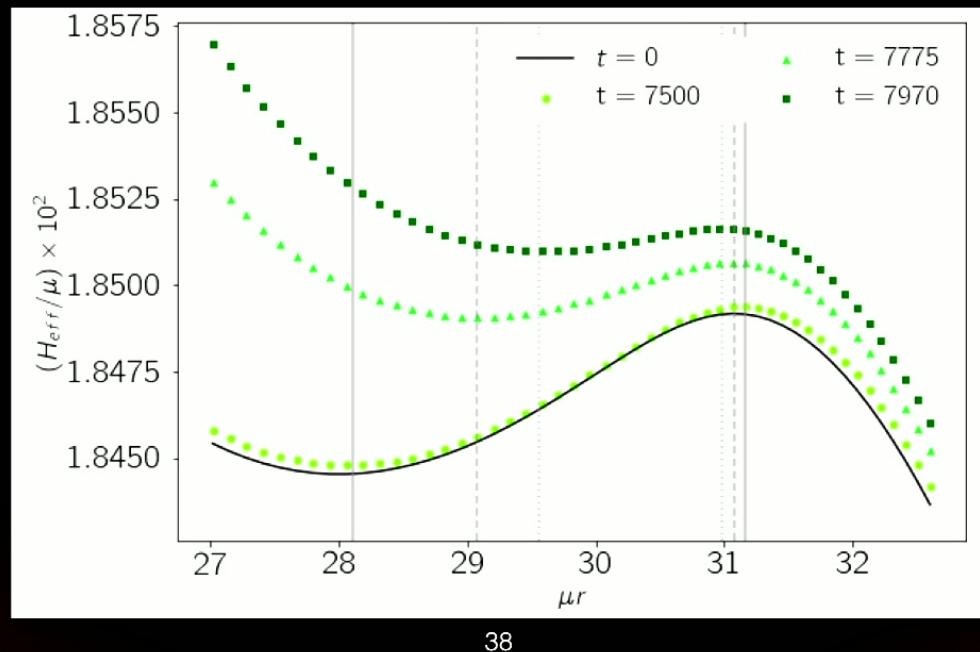
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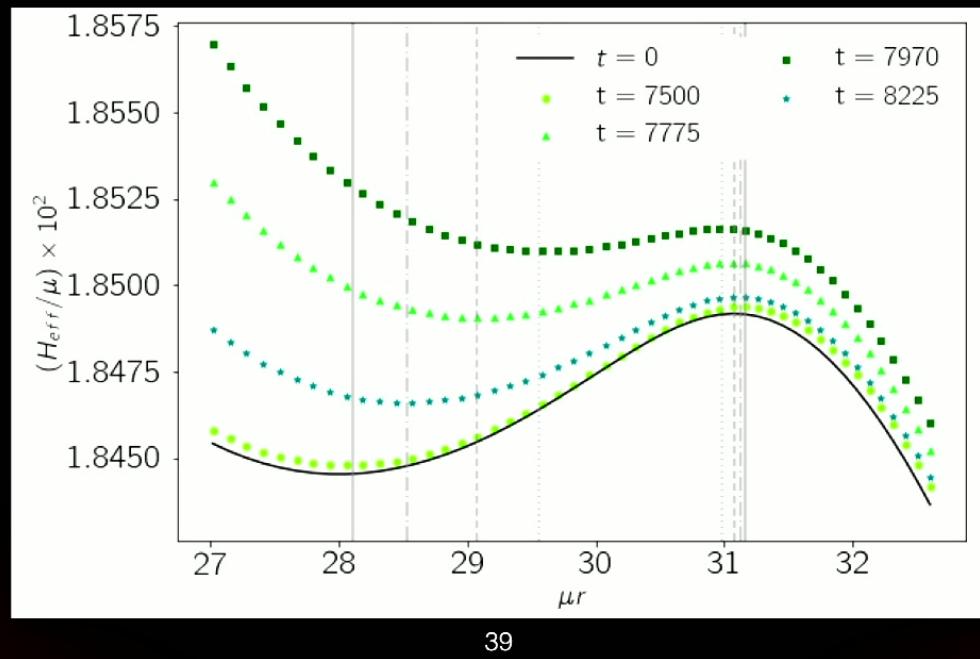
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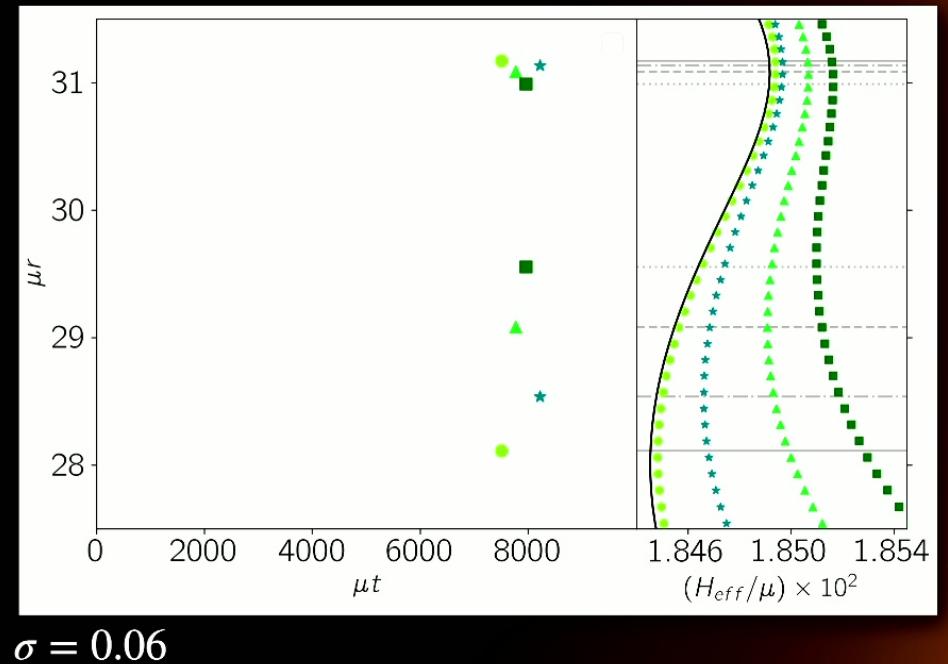
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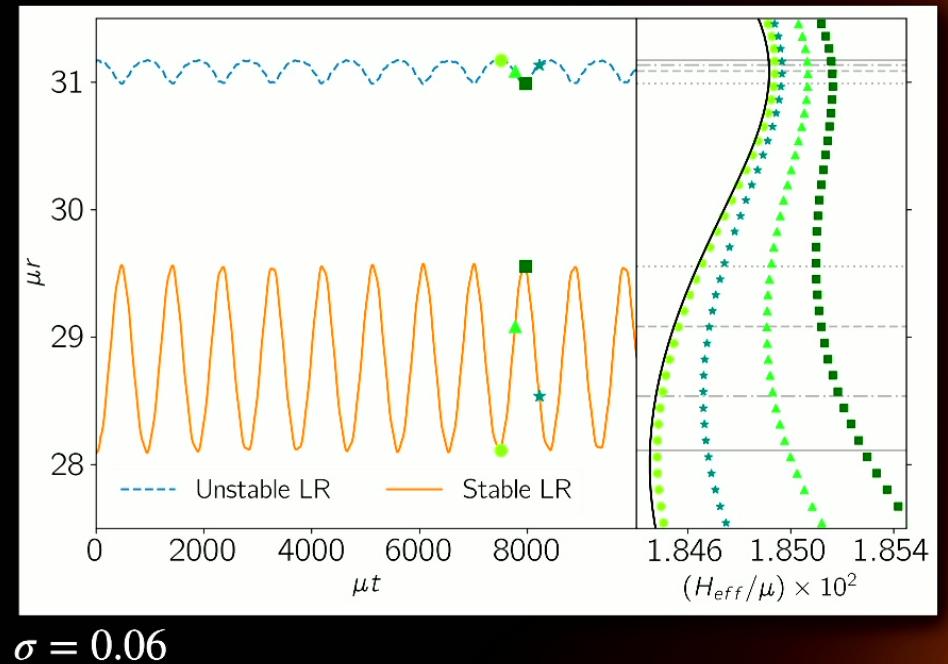
Light ring evolution

- Extract light rings at extrema



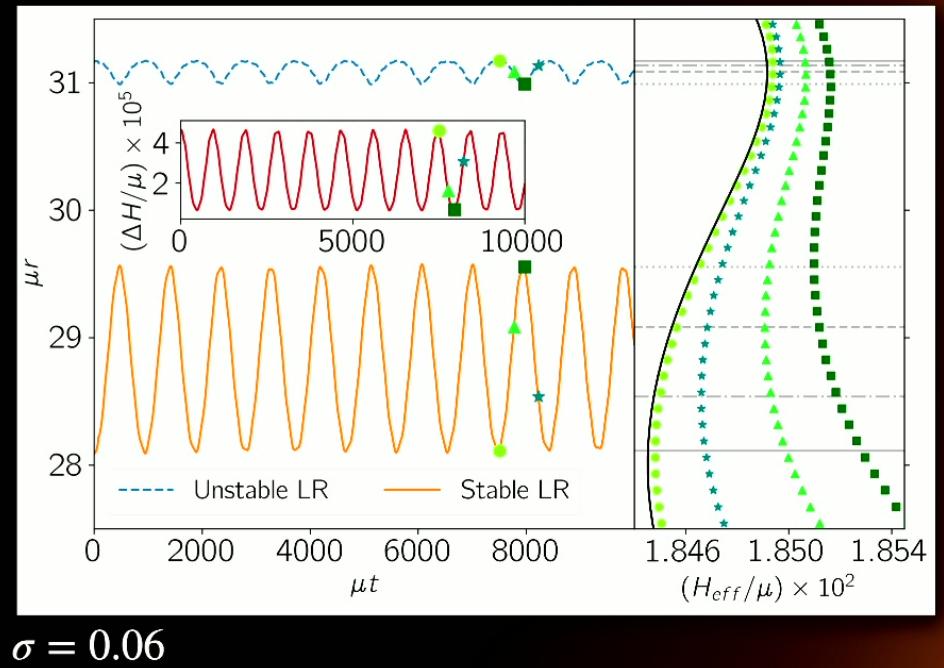
Light ring evolution

- Extract light rings at extrema
- Light ring structure preserved



Light ring evolution

- Extract light rings at extrema
- Light ring structure preserved



Robustness of AEP

- Variances over the extraction sphere

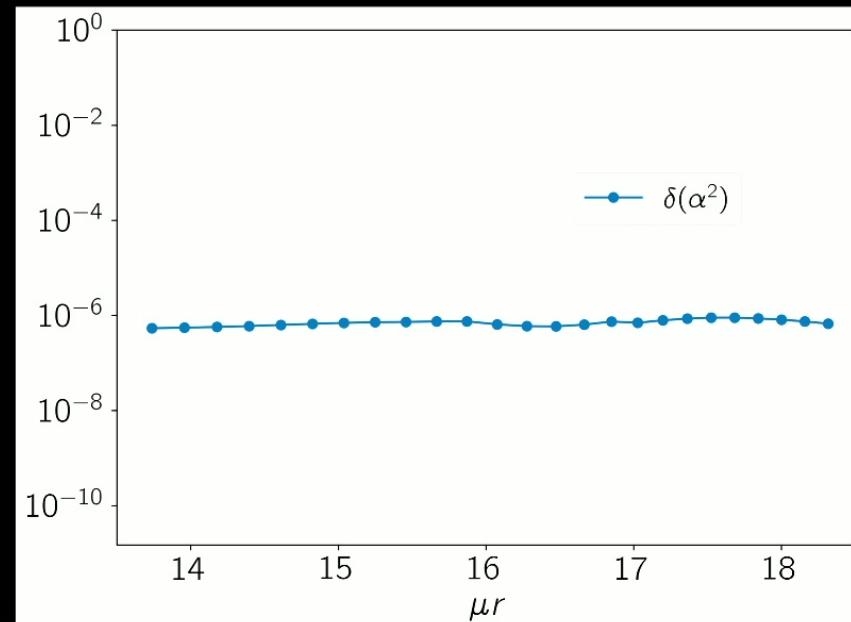
$$H_{\text{eff}}(r) = \sqrt{\frac{\langle \alpha^2 \rangle - \langle \gamma_{ij} \beta^i \beta^j \rangle}{\left(\oint \sqrt{\det q}\right)/4\pi}}$$
$$\sigma(X), \langle X \rangle$$
$$\delta(X) = \frac{\sigma(X)}{\langle X \rangle}$$

Robustness of AEP

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$\sigma = 0.08$

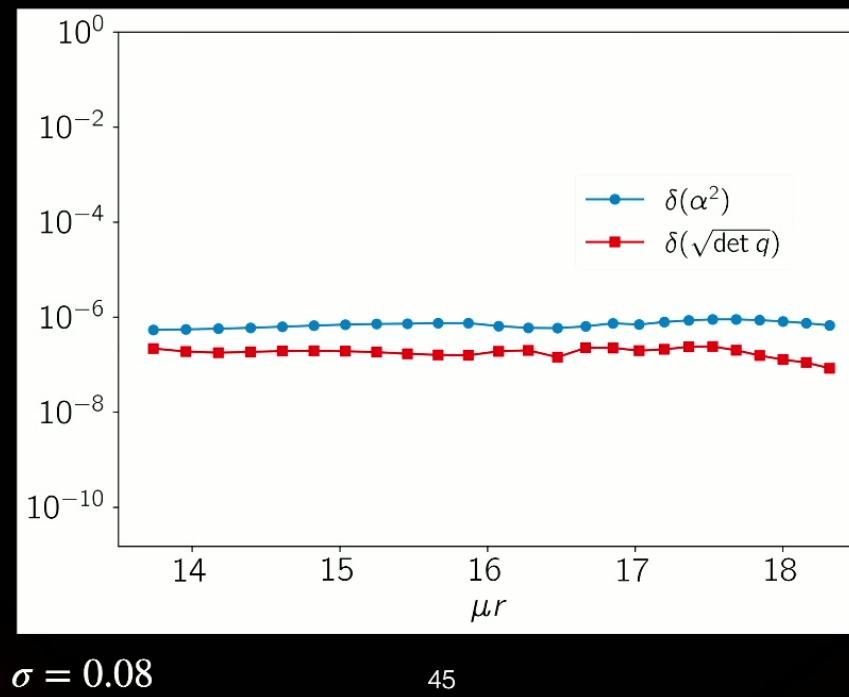
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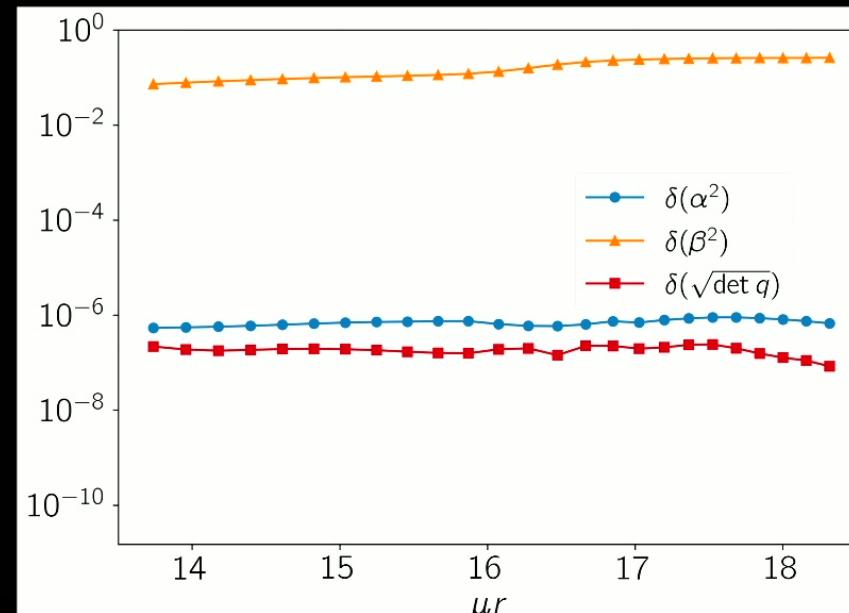
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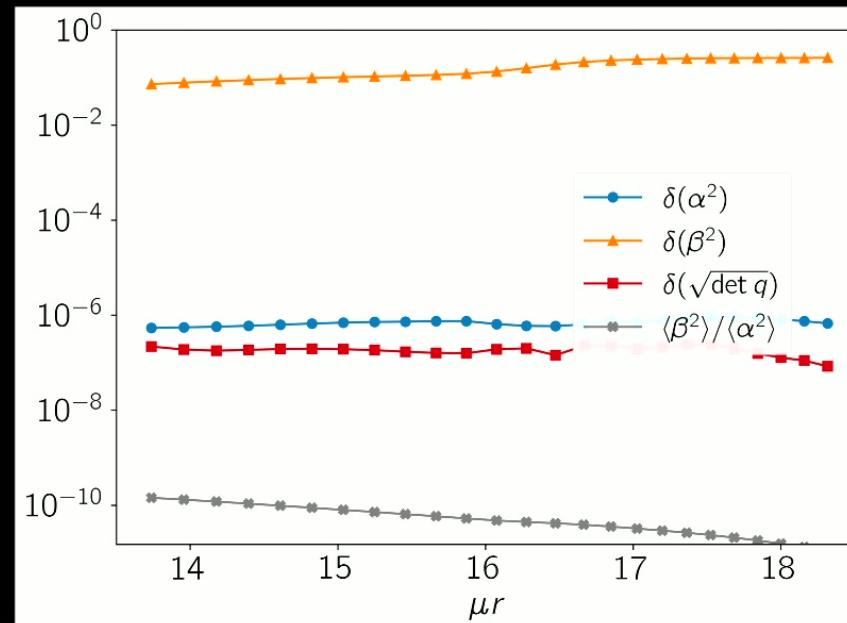
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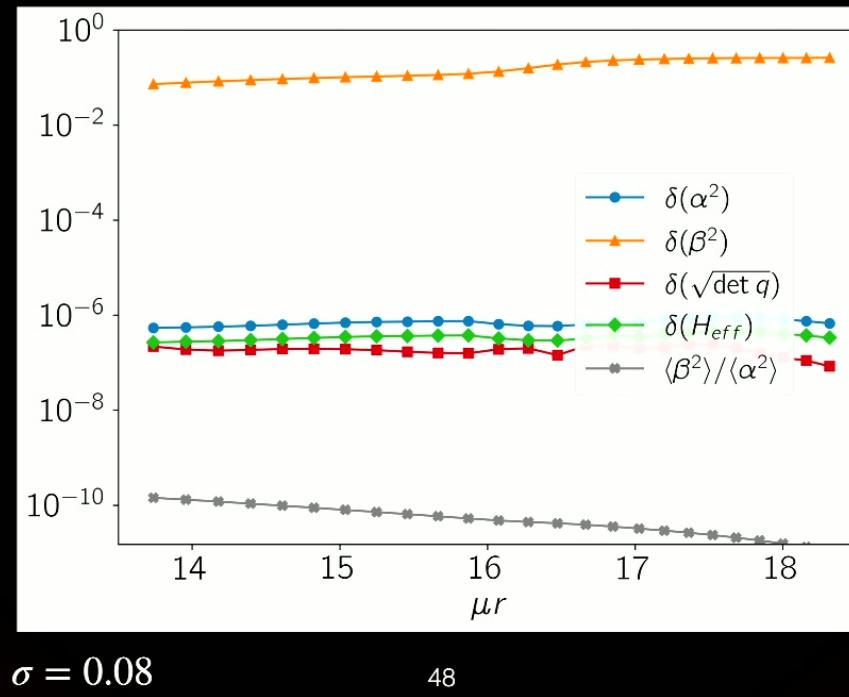
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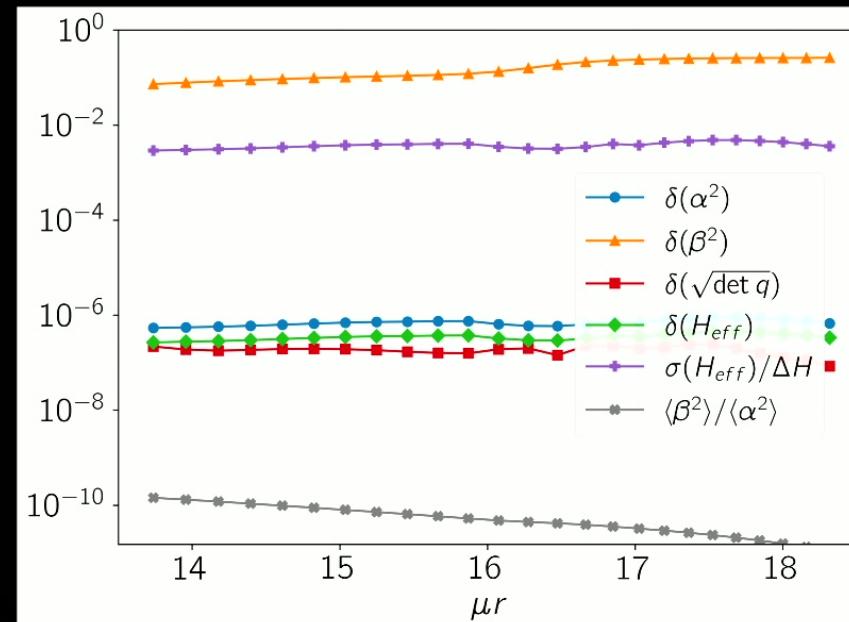


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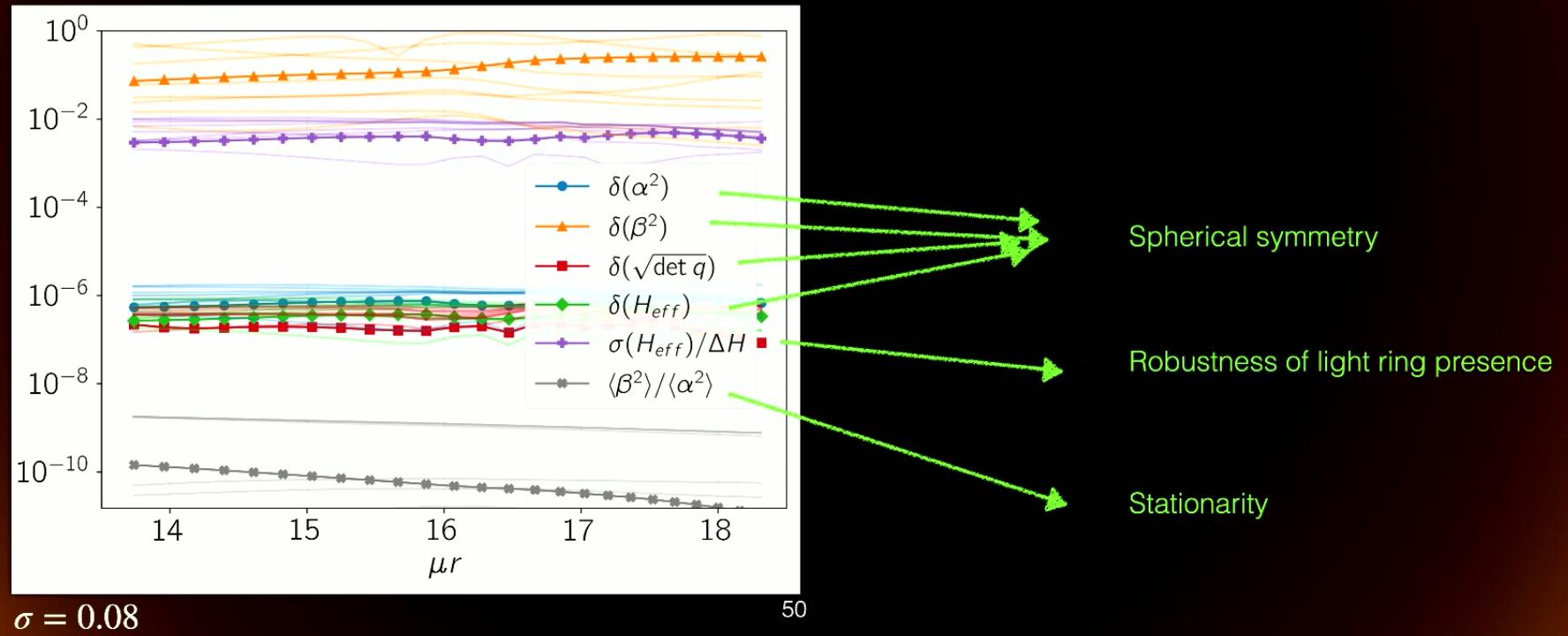
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Robustness of AEP

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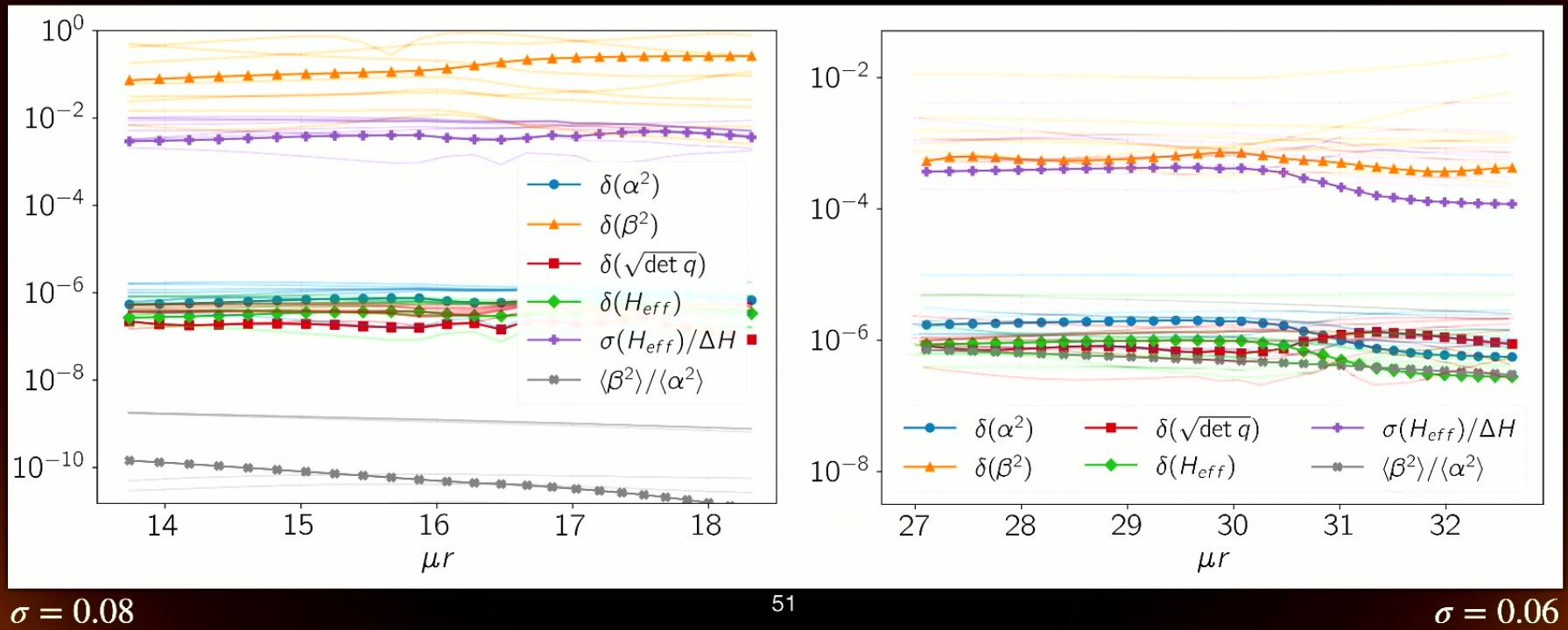
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Robustness of AEP

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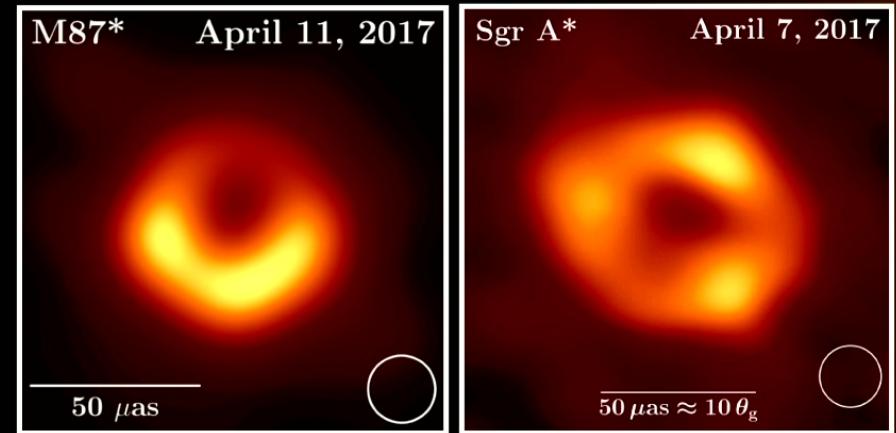


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Teaser: VLBI images

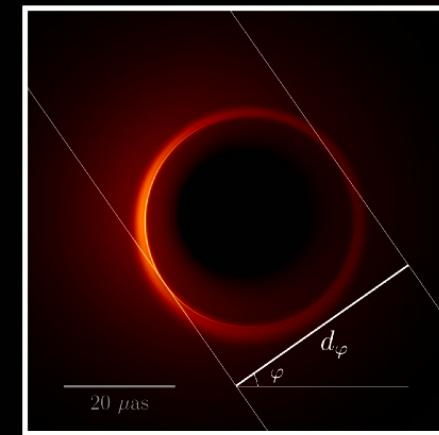
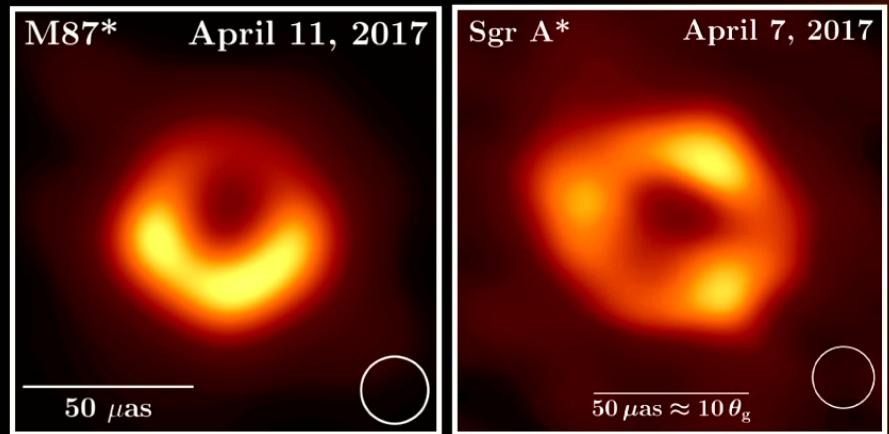
VLBI images

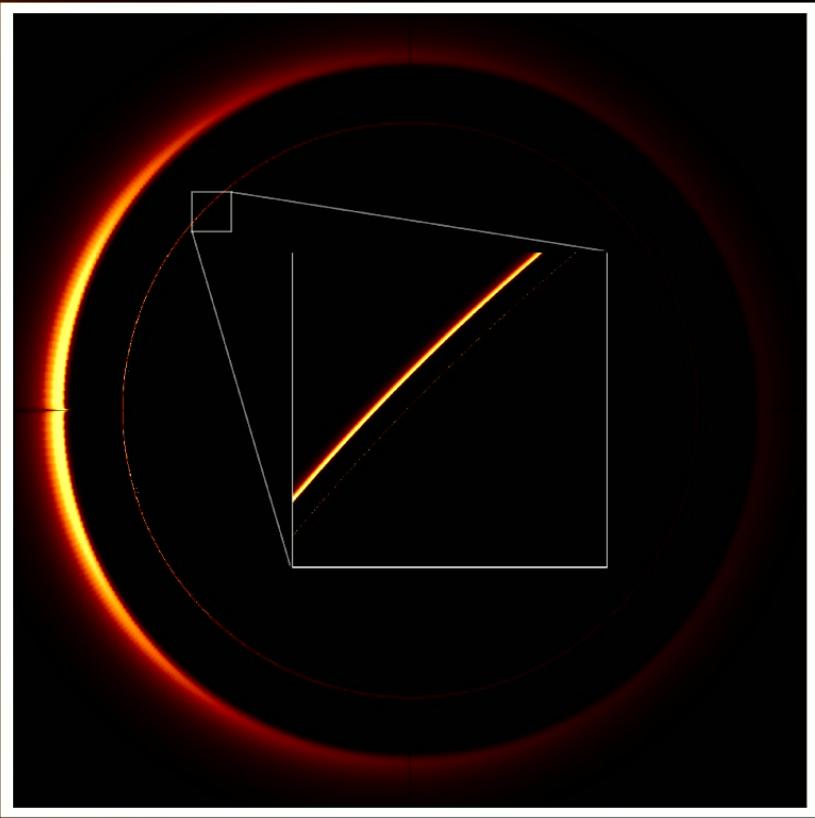
- Event Horizon Telescope Collaboration
- Test the BH paradigm



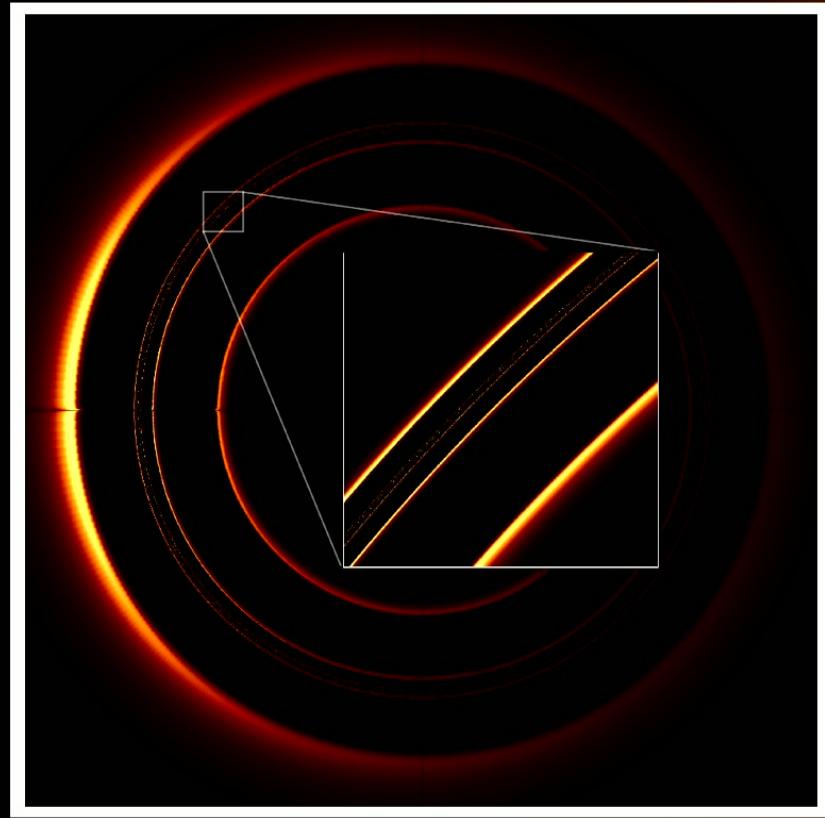
VLBI images

- Event Horizon Telescope Collaboration
- Test the BH paradigm
- Photon rings of special interest (Johnson+, '20) (Gralla+, '20)





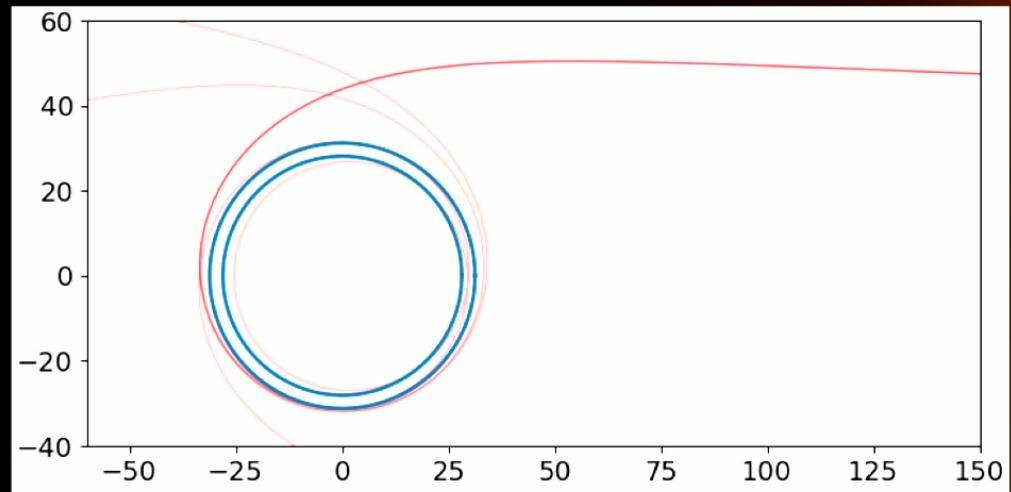
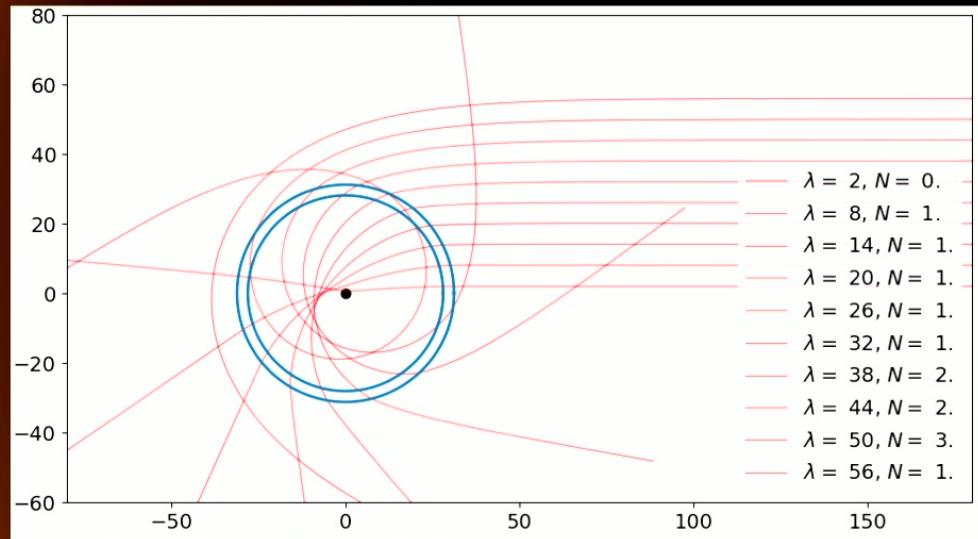
Schwarzschild



Boson Star

(Images produced with FOORT (Mayerson+ (incl. SS), in prep.))

Ray tracing



Conclusions & Outlook

- Stable evolutions of ultracompact solitonic BSs
- Light ring structure is preserved throughout
- Excellent agreement with radial perturbation analysis

Conclusions & Outlook

- Stable evolutions of ultracompact solitonic BSs
- Light ring structure is preserved throughout
- Excellent agreement with radial perturbation analysis

- Extend to rotating BSs
- Simulations of mergers, formation, non-spherical perturbations