

Title: Lecture - AdS/CFT, PHYS 777

Speakers: David Kubiznak

Collection/Series: AdS/CFT (Elective), PHYS 777, March 31 - May 2, 2025

Subject: Quantum Fields and Strings, Quantum Gravity

Date: April 24, 2025 - 9:00 AM

URL: <https://pirsa.org/25040033>

WILSON LOOPS = NON-LOCAL GAUGE INVARIANT

OPERATORS \simeq PARALLEL TRANSPORT OF QUARK ALONG
CLOSED LOOP C .

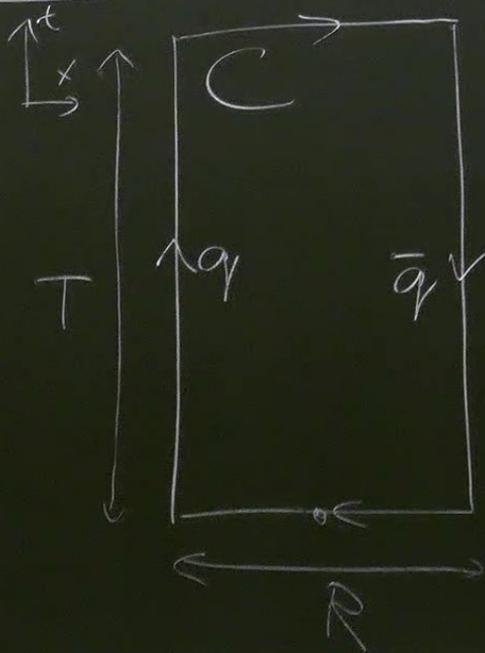
$$W(C) = \frac{1}{N_c} \text{Tr} \left(\text{P exp} \left(i \oint_C A_\mu dx^\mu \right) \right)$$

$$A_\mu = (V, \vec{\Phi}) \quad dx^\mu = (dt, dx^i)$$

$\leftrightarrow iT$

\downarrow

SPECIFICALLY CONSIDER C : " $q\bar{q}$ PAIR CREATION & ANNIHILATION"

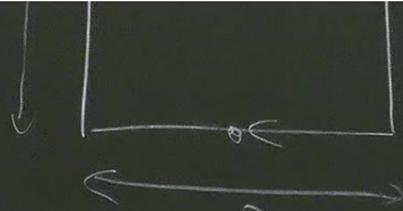


$\langle W(C) \rangle \propto e^{-TV(R)}$

EUCL. QUARK POTENTIAL

FIELD THEORY WISDOM:

$$V(R) = \begin{cases} \sim R & \text{CONFINED PHASE} \\ \sim \frac{1}{R} & \text{UNCONFINED PHASE} \end{cases}$$



$$V(R) = \begin{cases} \sim R & \\ \sim \frac{1}{R} & \end{cases}$$

CONFINED PHASE

UNCONFINED PHASE

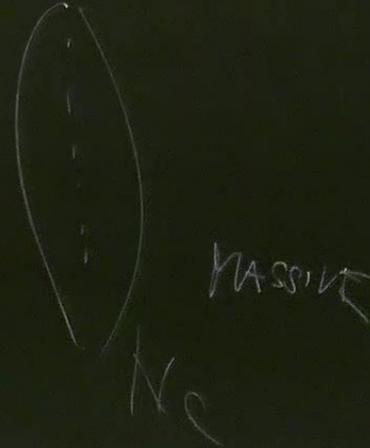
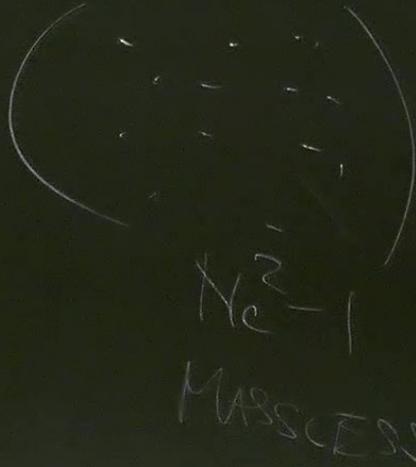
② CAN WE DERIVE THIS FROM AdS/CFT ③

$AK (V, \theta) dx^M = (dt, dx^i)$ (2) CAN WE DERIVE

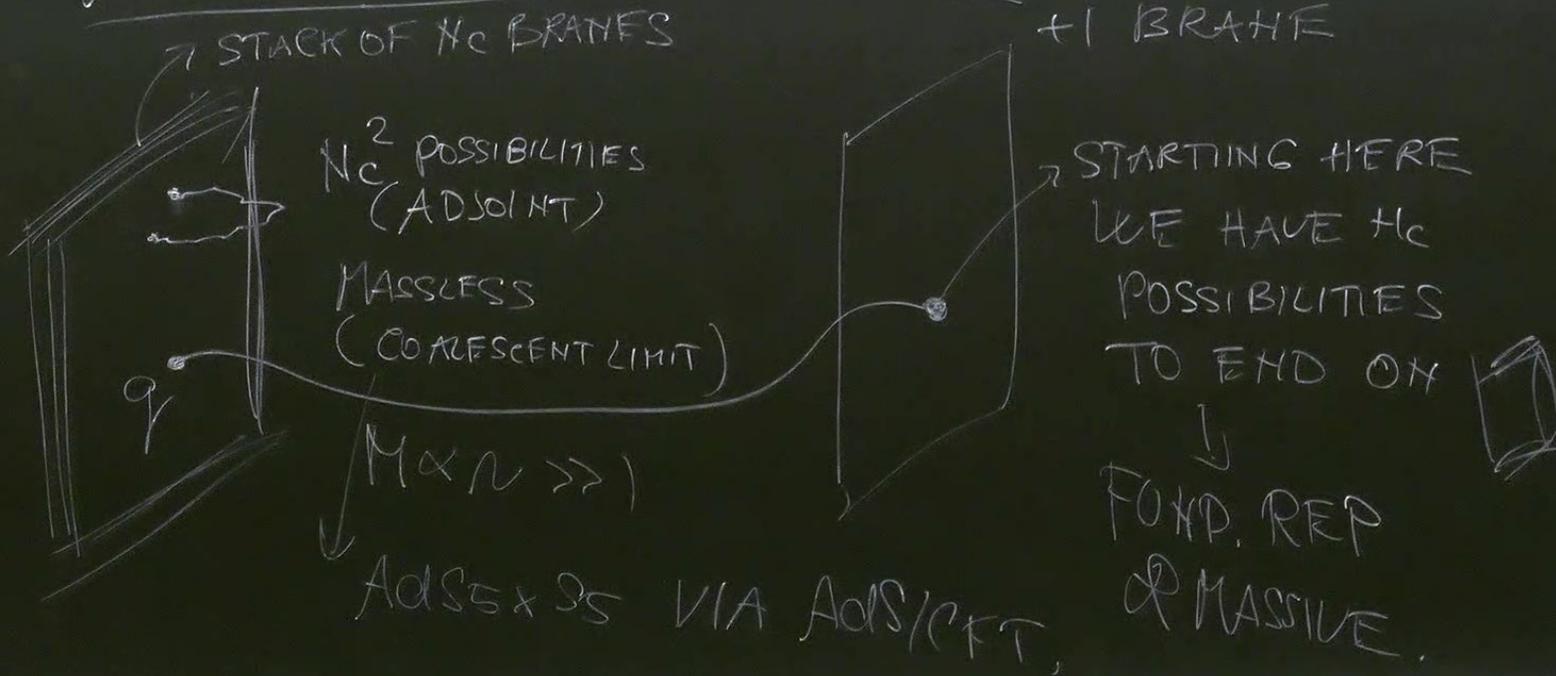
• HERE q IS MASSIVE & IN FUNDAMENTAL REP OF $SO(N)$
 (C.F. MASSLESS & ADJOINT REP. OF $SO(N) \sim SU(N)$)

ADS

FOND



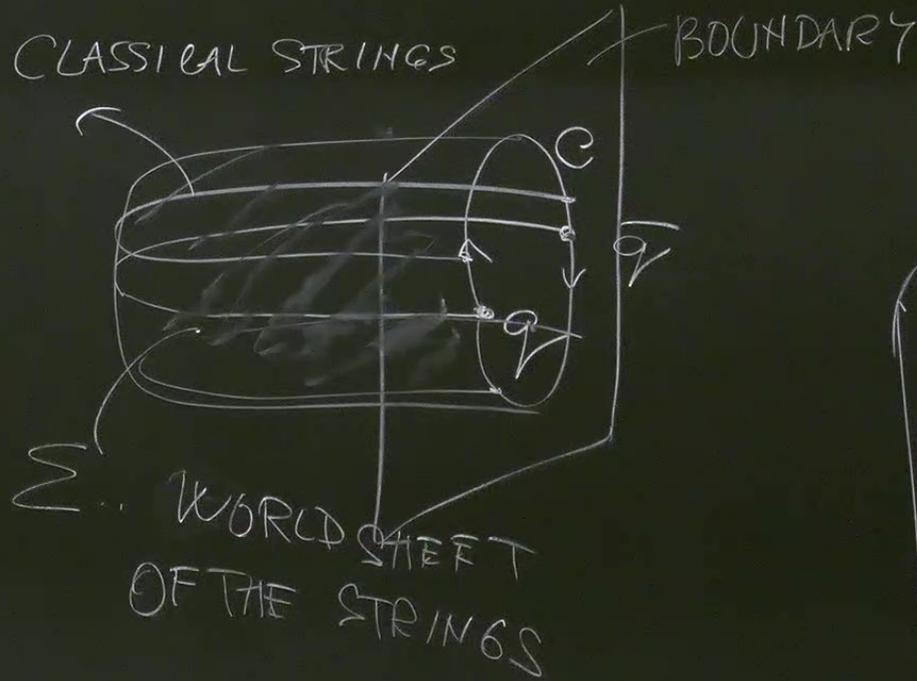
WHAT IS THE ADS PRESCRIPTION ?



HAVE THE FOLLOWING PRESCRIPTION

AdS₅

BULK



STRINGS GO

→ SHAPE C

COO

$$\langle W(C) \rangle \propto \Rightarrow V(R) =$$

... VIA ADS/CFT, ... MASSIVE.

GOVERNED BY CLASSICAL HG ACTION

OF Σ EXTREMIZING HG SUBJECT TO BOUNDARY

CONDITION: $\partial \Sigma = \Sigma / \partial \text{ADS} = C$

$-S_{NG, \text{EXT}} \times \mathcal{L} = \frac{S_{NG}}{T}$

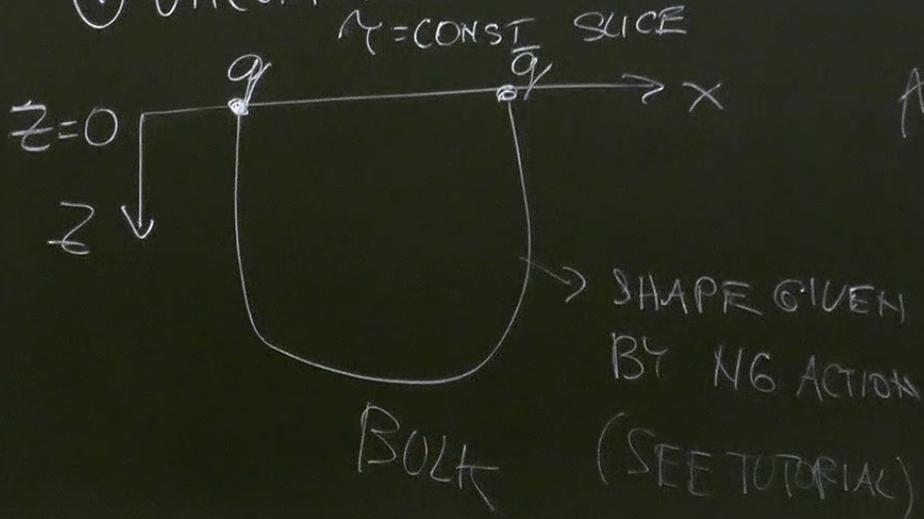
$$S_{NG} = \frac{1}{2\pi\alpha'} \int_{\Sigma} d^2\sigma \sqrt{\det g_{AB}}$$

$\Sigma: \partial \Sigma = C$

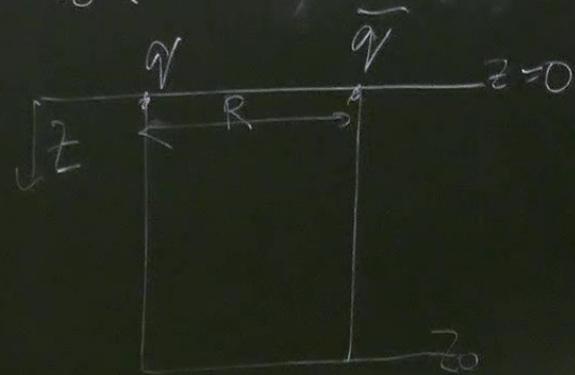
QUANTUM POTENTIAL

• QUALITATIVE ESTIMATES

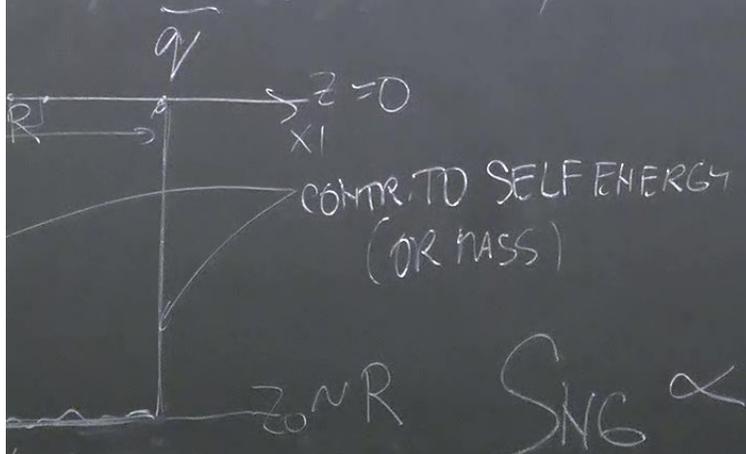
① UNCONFINED PHASE: USE POINCARÉ $ds^2 = \frac{r^2}{z^2} (d\tau^2 + dz^2 + dx_1^2 + dx_2^2 + dx_3^2)$



APPROX \approx



$$+ dz^2 + dx_1^2 + dx_2^2 + dx_3^2$$



• INDUCED METRIC

$$d\eta^2 = \frac{l^2}{z^2} (d\tau^2 + dx_i^2)$$

$$\sqrt{\det \eta} = \frac{l^2}{z^2} = \frac{l^2}{z_0^2}$$

$$S_{NG} \propto \frac{1}{l_s^2} \frac{l^2}{z_0^2} \int d\tau dx_i = \left(\frac{l^2}{l_s^2} \right) \frac{TR}{R^2} \sim \frac{\sqrt{X} T}{R}$$

$\Rightarrow V(R) \sim \frac{\sqrt{X}}{R}$

 STRONG COUPLING TR

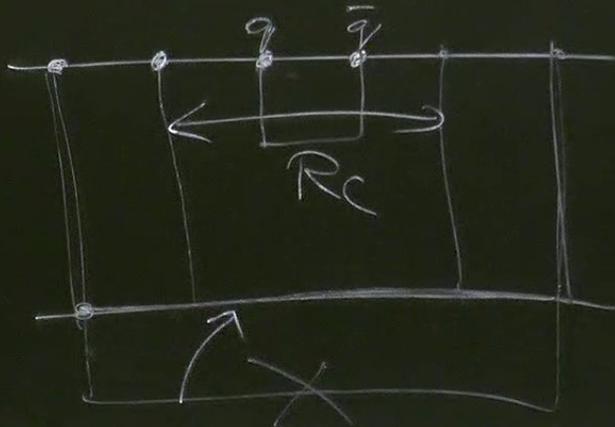
 UNCONFINED PHASE

② CONFINEMENT / DE-CONFINEMENT PT

$R < R_c$

NEED TO INTRODUCE A NEW SCALE (WHICH TELLS US WHERE THE PT HAPPENS)

... HOW DEEP WE CAN GO IN AdS:



$z=0$

$z=z_c$... NEW SCALE

• $R < R_c$... SITUATION AS BEFORE ... UNCONFINED PHASE $V \sim \frac{1}{R}$

• $R > R_c$... DEEPEST WE CAN GO IS $z = z_c$.

$$\sqrt{\det g} = \frac{l^2}{z_c^2}$$

$$S_{NG} \propto \frac{1}{l_s^2} \frac{l^2}{z_c^2} \int dr dx_1 \sim \frac{l^2}{l_s^2} \frac{RT}{z_c^2} \sim \sqrt{\frac{RT}{z_c^2}}$$

$$\boxed{V \sim \sqrt{\frac{RT}{z_c^2}}} \dots \text{CONFINED PHASE!}$$

(3) FINITE TEMP: PLASMA PHASE.

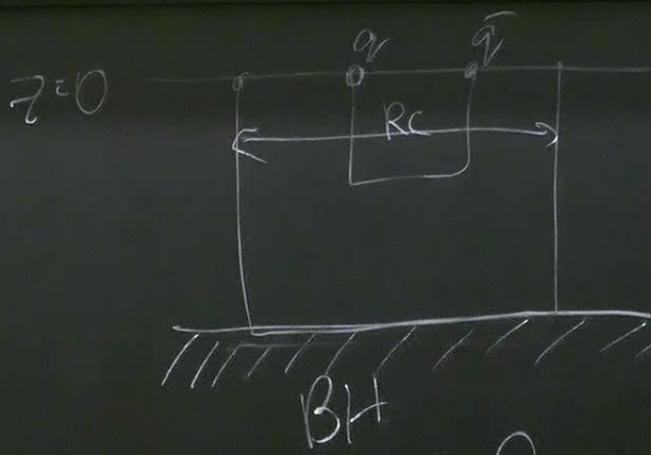
BH IN THE BULK GIVES CFT A FINITE
TEMPERATURE $T = T_H$

$$ds^2_{\text{AdS}_5 \text{ BH}} = \frac{l^2}{z^2} \left(f dr^2 + \frac{dz^2}{f} + dx_1^2 + dx_2^2 - dx_3^2 \right)$$

$$f = 1 - \left(\frac{z}{z_0} \right)^4$$

BLACKENING
FACTOR

$$\left(\frac{f}{f_{\text{IR}}} = 0 \right)$$



$$d\eta^2 = \frac{l^2}{z^2} (f dr^2 + dx_1^2)$$

$z = z_0$ HORIZONTAL

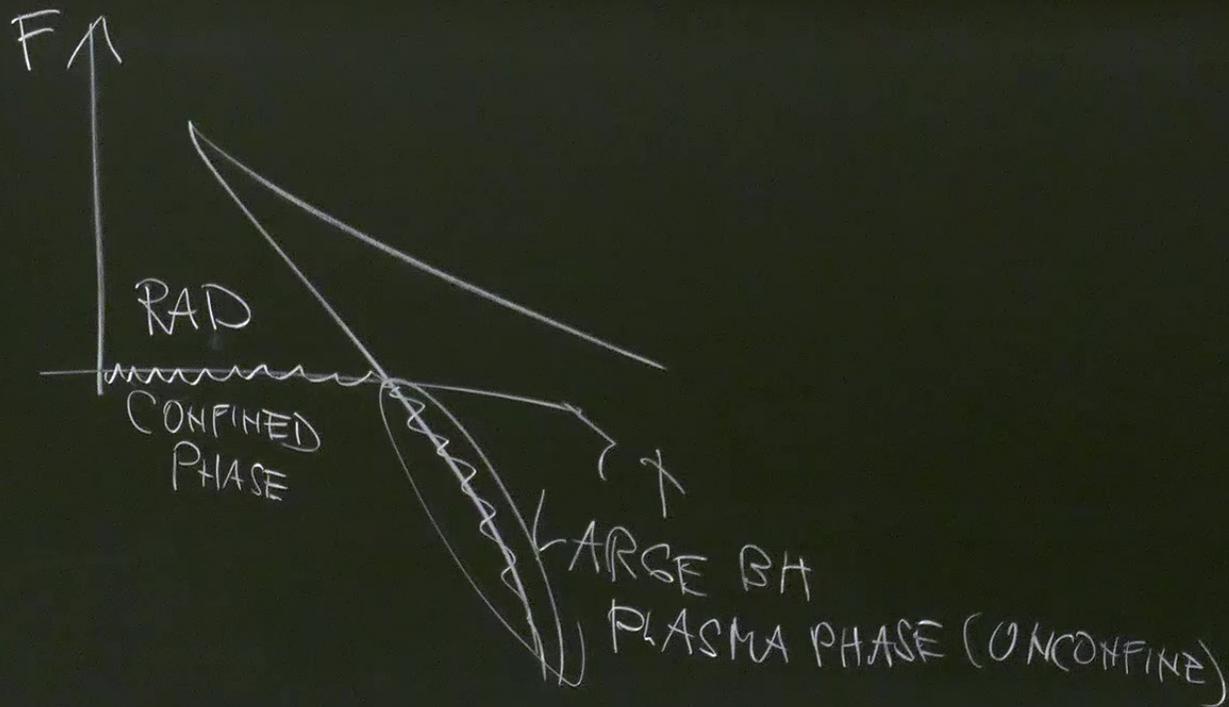
$$\sqrt{\det \eta^2} = \frac{l^2}{z^2} \sqrt{f} = 0 \quad \uparrow \text{ON HORIZONTAL}$$

BH $\Rightarrow S_{HG} = 0 = V(R)$

DEBYE SCREENING ... STANDARD IN PLASMA

$\frac{1}{R} = 0$

CONNECT THIS TO H-P PHASE TRANSITION, SPHERICAL BH



③ FINITE TEMP: PLASMA PHASE.

BH IN THE BULK GIVES CFT A FINITE TEMPERATURE $T = T_H$

$$ds^2_{\text{AdS}_5 \text{ BH}} = \frac{l^2}{z^2} \left(f dz^2 + \frac{dz^2}{f} + dx_1^2 + dx_2^2 + dx_3^2 \right)$$

$$f = 1 - \left(\frac{z}{z_0} \right)^4 \leftarrow \text{BLACKENING FACTOR} \quad \left(\frac{f}{f_{\text{HOR}}} = 0 \right)$$

$z=0$

