

Title: Lecture - Quantum Matter, PHYS 777

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Subject: Condensed Matter

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Ginzburg-Landau $\mathcal{L}[\Phi, A_\mu] \xrightarrow{\langle \Phi \rangle \neq 0}$ Goldstone-Higgs $\mathcal{L}[\theta, A_\mu] = \frac{K}{2} (\partial_\mu \theta)^2$

Comment: in a real superconductor $\Phi \in \mathbb{C}$ two electrons a.k.a. Cooper pair
 on
 Universality. Why this is preferred $\leftarrow$ "Pairing mechanism"

$$|\partial \phi|^2 + \gamma |\phi|^2 + \lambda |\phi|^4$$

Good metals, understood. electron-phonon
 "Unconventional S.C." "High- T_c Superconductors"

$\rightarrow$ Goldstone-Higgs $\mathcal{L}[\theta, A_\mu] = \frac{K}{2} (\partial\theta - A)^2$.

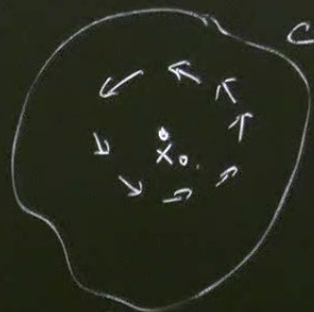
for $\Phi \in \mathbb{C}$ two electrons a.k.a. Cooper pair.

Why this is preferred $\leftarrow$ 'Pairing mechanism'

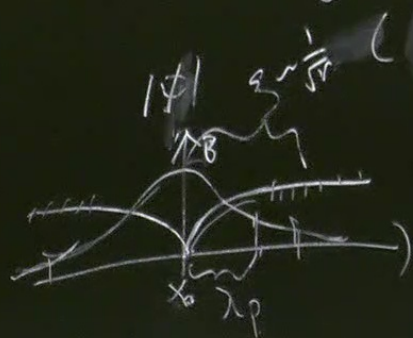
Good metals, understood. electron-phonon coupling $\leftarrow$ BCS. theory.

'Unconventional S.C.' 'High- T_c Superconductor'

Vortex: θ can have a topological defect.



2d.



$$\oint_C \vec{\nabla} \theta \cdot d\vec{l} = 2\pi n, \quad n \in \mathbb{Z} \Rightarrow \text{Vortex cannot be created/destroyed}$$

$$\vec{\nabla} \times \vec{\nabla} \theta = 2\pi n \delta^{(2)}(\mathbf{x} - \mathbf{x}_0)$$

θ is not defined at x_0

($|\phi|e^{i\theta} \rightarrow 0$ as $x \rightarrow x_0$)

$$|\phi| \sim \frac{1}{\sqrt{r}}$$

In a S.C., vortex costs energy

For a dynamical gauge field

For large C ; $A \rightarrow A + \partial\theta$

$$A = \partial\theta$$

$$\int_B d\vec{x} = \int A \cdot d\vec{l} = \int \partial\theta \cdot d\vec{l} = 2\pi n$$

$$\frac{K}{2} (\partial\theta - A)^2 + (\partial \times A)^2$$

topological defect.

$n, n \in \mathbb{Z} \Rightarrow$ Vortex cannot be created/destroyed locally.

In a S.C., vortex costs energy, but is stable once formed.

For a dynamical gauge field A .

For large C , $A \rightarrow A + \partial\theta$ illegal with vortex

$$A = \partial\theta$$

$$\oint B dx = \oint A \cdot dl = \oint \partial\theta \cdot dl = 2\pi n.$$

$$\frac{K}{2} (\partial\theta - A)^2 + (\partial \times A)^2$$

minimize if $A = \partial\theta$

Put back Maxwell

$$B = \nabla \times A = \nabla \times \partial\theta = 2\pi n \delta(x - x_0)$$

Charge-Vortex duality. (Refs: David Tong's notes on Gauge Theory. Chp. 8).

Claim: $U(1)$ Goldstone theory (superfluid) is dual to a gauge theory

$$\begin{aligned}
 (2+1)d \\
 Z = \int D\theta e^{-\int d^3x \frac{K}{2} (\partial\theta)^2} &= \int D\theta DA_\mu Df_{\mu\nu} e^{-\int \frac{K}{2} A_\mu^2} e^{i \int \frac{E_{\mu\nu\lambda}}{2\pi} f_{\mu\nu} (A_\lambda - \partial_\lambda \theta)} \\
 &\quad \text{U(1) Maxwell} \\
 &\quad \text{integrate out } A_\mu \\
 &= \int Df_{\mu\nu} D\theta e^{-\int \frac{1}{2} \frac{1}{4\pi^2 K} f_{\mu\nu}^2} e^{i \int \frac{E_{\mu\nu\lambda}}{2\pi} \partial_\lambda \theta f_{\mu\nu}} \\
 &\quad \int D\theta \Rightarrow E_{\mu\nu\lambda} \partial_\lambda f_{\mu\nu} = 0
 \end{aligned}$$

in Gauge Theory (Chp. 8) "T-duality"

a gauge theory

1611 Maxwell

$$\int \frac{G_{\mu\nu\lambda}}{2\pi} f_{\mu\nu} (A_\lambda - \partial_\lambda \theta)$$

$$\mathcal{Q} \text{ Compact} \Rightarrow \oint_{C_1} f_{\mu\nu} = 2\pi n$$

$$\int \mathcal{D}\theta \left(e^{-\int \frac{1}{4e^2} (\partial_\mu a_\nu - \partial_\nu a_\mu)^2} \mathcal{D}a_\mu \right)$$

$$e^{i \int \frac{e_{\mu\nu\lambda}}{2\pi} \theta_\lambda \partial f_{\mu\nu}}$$

$$\int \mathcal{D}\theta \Rightarrow G_{\mu\nu\lambda} \partial_\lambda f_{\mu\nu} = 0$$

$$f_{\mu\nu} = \partial_\mu a_\nu - \partial_\nu a_\mu$$

$(2+1)d$ only has 1 polarization $\equiv$ 1 mass scalar.

$(2+1)d$ photon $\equiv$ Goldstone mode.

Conserved U(1) current $j_\mu = iK \partial_\mu \theta \longleftrightarrow$ Maxwell, $j_\mu = \frac{\epsilon_{\mu\nu\lambda}}{2\pi} \partial_\nu a_\lambda$, $\partial_\mu j_\mu = 0$.

Quantized charge. $\longleftrightarrow$ Quantized. $j_0 = \frac{b}{2\pi}$, $\vec{j} = \frac{\vec{e} \times \vec{e}}{2\pi}$ "compact"
 $\int d^2x \frac{b}{2\pi} = n \in \mathbb{Z}$ U(1) gauge field.

or, $\hat{=}$ 1 mass scalar.

$\hat{=}$ Goldstone mode.

$\langle \partial_\mu \theta \rangle \longleftrightarrow$ Maxwell, $j_\mu = \frac{\epsilon_{\mu\nu\lambda}}{2\pi} \partial_\nu a_\lambda$, $\partial_\mu j_\mu = 0$.

$\longleftrightarrow$ Quantized. $j_0 = \frac{b}{2\pi}$, $\vec{j} = \frac{\vec{E} \times \vec{E}}{2\pi}$ "Compact"
 $\int d^3x \frac{b}{2\pi} = n \in \mathbb{Z}$ U(1) gauge field.

Dirac quantization.

$\Rightarrow$ 1 mass scalar.

$\Rightarrow$ Goldstone mode.

$\leftrightarrow$ Maxwell, $j_\mu = \frac{\epsilon_{\mu\nu\lambda}}{2\pi} \partial_\nu a_\lambda$, $\partial_\mu j_\mu = 0$.

$\rightarrow$ Quantized, $j_0 = \frac{b}{2\pi}$, $\vec{j} = \frac{\vec{E} \times \vec{E}}{2\pi}$ "Compact"
 $\int d^3x \frac{b}{2\pi} = n \in \mathbb{Z}$ U(1) gauge field.

Dirac quantization.

$\leftrightarrow$ An operator changes $\int d^3x \frac{b}{2\pi}$: (Space-time) monopole (instanton)



$\Rightarrow$ (2+1)d Maxwell theory is a condensate (superfluid) of monopoles.

photon in (2+1)d only has 1 polarization $\hat{=}$ 1 mass scalar.

(2+1)d photon $\hat{=}$ Goldstone mode.

Conserved (U(1)) current $j_\mu = k \partial_\mu \theta \longleftrightarrow$ Maxwell, $j_\mu = \frac{\epsilon_{\mu\nu\alpha} \partial_\nu a_\alpha}{2\pi}$, $\partial_\mu j_\mu = 0$

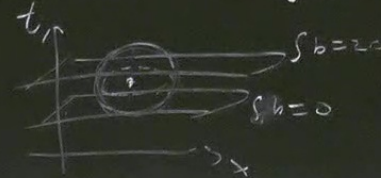
Quantized charge. $\longleftrightarrow$ Quantized. $j_0 = \frac{b}{2\pi}$, $\vec{j} = \frac{\vec{e} \times \vec{e}}{2\pi}$ "Compact"
 $\int d^2x \frac{b}{2\pi} = n \in \mathbb{Z}$. U(1) ga

Charged operator. $e^{i\theta}$

$\longleftrightarrow$ An operator changes $\int d^2x \frac{b}{2\pi}$: (space-time)

$$\int D a_\mu \text{Max} e^{-S[a]} = \int D a_\mu e^{-S[a]}$$

$\odot \oint \partial_\mu \vec{a} = 2\pi$



$\Rightarrow$ (2+1)d Maxwell

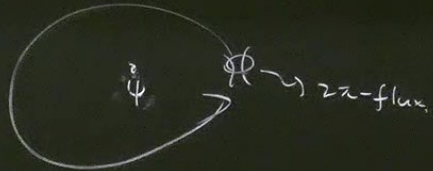
$$\int D a_\mu \text{Max} \mathcal{L}$$

$$= \int D a_\mu e^{-i \int \partial_\mu \vec{a} = 2\pi}$$

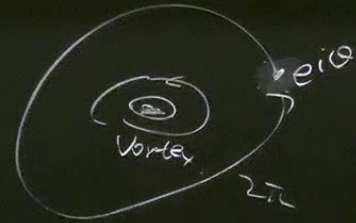


Gauge charge ψ Maxwell phase ψ gapped

Superfluid = vortex is gapped

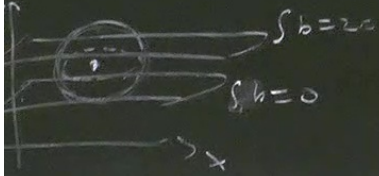


AB phase of 2π



$$\frac{1}{2} |(\partial - i a) \psi|^2 + r |\psi|^2 + |\psi|^4 + \frac{1}{4e^2} f_{\mu\nu}^2$$

$$\leftrightarrow \frac{1}{2} |\partial \phi|^2 - r |\phi|^2 + |\phi|^4$$



$\Rightarrow$ (2+1)d Maxwell theory is a condensate (superfluid) of monopoles.

is gapped.

$r < 0$. Insulator phase, trivially gapped.

Gauge theory $\langle \psi \rangle \neq 0$, gapped vacuum (Higgs).

$\Rightarrow$ Insulator is a condensate of vortices.

Topological
Defect

$(2+1)d$ charge-vortex duality.

Vortex, $\pi_1(U(1)) = \mathbb{Z}$.

$(1+1)d$

Domain

$\mathbb{Z}$

Dual theory is
a gauge theory

Maxwell,

Defect = gauge charge.



$SSB \leftrightarrow$ Gapped, symmetric
phase of defect.



Symmetric phase $\leftrightarrow$ Higgs phase.



(1+1)d Kramers-Wannier duality (Ising)

Domain wall $\Pi_0(\mathbb{Z}_2) = \mathbb{Z}_2$

$\mathbb{Z}_2$ gauge theory

✓

✓

✓



Symmetric phase $\leftrightarrow$ Higgs phase. ✓

Charged operator

$\leftrightarrow$ Instantons.

Monopole

Boundary-condition changing op.

Bonus: Maxwell $\rightarrow [M(x) M^*(x)]$

Photon becomes gapped
due to instanton

$\Rightarrow$ Gauge charge confined.

Polyakov
mechanism.

$$\leftrightarrow (\partial\theta)^2 = \lambda (e^{i\theta} + e^{-i\theta}) \sim \theta^2$$

Gapped

$\leftrightarrow$

$$\leftrightarrow \frac{2\pi}{r} \left(\frac{2\pi}{r} \right) \quad E \sim O(1/r)$$