

Title: Lecture - Quantum Matter, PHYS 777

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Collection/Series: Quantum Matter (Elective), PHYS 777, March 31 - May 2, 2025

Subject: Condensed Matter

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$U(1)$: Insulator, Superfluid, Superconductor.

Ginzburg-Landau theory

Effective field theory, in term $\phi(x, t)$ smoothly varying in spacetime. (e

$U(1)$ symmetry: $\phi(x, t) \rightarrow e^{i\alpha} \phi(x, t)$

$$Z[A] = \int D\phi D\phi^* e^{-\int d^4x d\tau \mathcal{L}[\phi, A_\mu]} \quad A_\mu: \text{background/probe gauge field.}$$

$$= e^{-F[A]}$$

$$\mathcal{L}[\phi, A] = i\phi^*(\partial_t - iA_t)\phi + \frac{1}{2}|\nabla - iA|\phi|^2 + \frac{r}{2}|\phi|^2 + \frac{\lambda}{4}|\phi|^4$$

acetime. (energy/momentum cut off. $\sim \Lambda \ll \text{lattice scale}$)

elcl.

$$\frac{r}{2} |\phi|^2 + \frac{\lambda}{4} |\phi|^4 + \dots$$

↗ higher order in ϕ, ∂ .

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$$\mathcal{L}[\phi, A] = \underbrace{i\phi^* (\partial_t - iA_t)\phi}_{\text{Sometimes extra sym condition } |(\partial_t - iA_t)\phi|^2} + \frac{1}{2} |\vec{\nabla} - i\vec{A})\phi|^2 + \frac{r}{2} |\phi|^2 + \frac{\lambda}{4} |\phi|^4$$

Consider $|r| \gg 0$: mean field.

$r \gg 0$. Vacuum (g.s.) fluctuates around $\langle \phi \rangle = 0$. $U(1)$ not spontaneously broken.

Correlation length $\xi \sim 1/\sqrt{r}$. $\langle \phi^*(x,t) \phi(0,t) \rangle \sim e^{-x/\xi}$.

Energy gap $\Delta \sim r$ ($\sqrt{r}$ if relativistic).

$F(A)$ should be local theory at length $\gg \xi$, energy $\ll \Delta$.

Maxwell.

$$F(A) = \int c_1 (\partial_t \vec{A} - \vec{\partial} A_t)^2 + c_2 (\partial_i A_j - \partial_j A_i)^2 + c_3 \partial_\mu A^\mu$$

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Energy gap $\Delta \sim r$ ($\sqrt{r}$ if relativistic)

$F(A)$ should be local theory, at length $\gg \xi$, energy $\ll \Delta$.

Maxwell

$$F(A) = \int c (\partial_t \vec{A} - \vec{\partial} A_t)^2 + C (\partial_i A_j - \partial_j A_i)^2 + o(c \partial_\mu A)$$

Dielectric response.

ntaneously broken.
 $\sim e^{-x/\xi}$

Some response theory:
 Current density $\vec{j}_i = - \frac{\delta F(A)}{\delta A_i} \stackrel{B=0}{=} C \cdot \partial_t \vec{E} = G \cdot \vec{E}$

$$\Rightarrow G(\omega) = i C \omega.$$

① Non-dissipative.

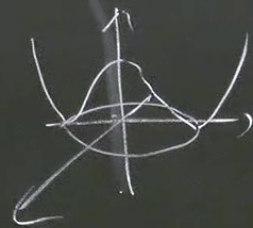
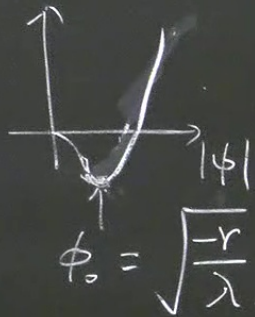
② Vanishes as $\omega \rightarrow 0$ (DC limit)

$\Rightarrow$ Insulator.

$\Rightarrow$ Gapped system are insulators



$r < 0$ minimize $V(\phi)$



SSB of $U(1)$

Expand around ϕ_0 , $\phi = (\phi_0 + \delta\rho) e^{i\theta}$ ^{real}

Amplitude mode $\delta\rho$ is gapped. $\rightarrow$ int

$$\Rightarrow \mathcal{L}[\theta] = \frac{K}{2} [(\partial_t \theta)^2 + (\nabla \theta)^2] \quad \theta =$$

$$U(1): \theta(x) \rightarrow \theta(x) + \alpha \quad \cos \theta, \sin \theta$$

Compactness: $\theta \sim \theta + 2\pi$, θ, θ^*

Superfluidity:

$$(\overset{\text{real}}{\phi_0} + \delta\phi) e^{i\theta}$$

apped. $\rightarrow$ integrate it out!

θ : gapless mode, $E \sim |k|$

"Goldstone mode".

Goldstone theorem: such gapless mode always exist for continuous SSB.

Superfluidity

Put back gauge field $\mathcal{L} = \frac{k}{2} (\partial_\mu \theta - A_\mu)^2$.

If I turn on $A_x = \text{const}$, $\mathcal{L}$ minimized by $\partial_x \theta = A_x$.

$\Rightarrow$ G.S. has current $j_x = k \partial_x \theta = k A_x$ per

field $\mathcal{L} = \frac{k}{2} (\partial_\mu \theta - A_\mu)^2$

$$\int dt \mathbf{j} \cdot \mathbf{E} = 0$$

$A_x = \text{const}$, $\mathcal{L}$ minimized by $\partial_x \theta = A_x$.

$\Rightarrow$ G.S. has current $\mathbf{j}_x = k \partial_x \theta = k A_x$ persistent (non-dissipative)

$A_x = \frac{E_x}{i\omega} \Rightarrow G(\omega, \vec{q}=0) = \frac{k}{i\omega}$ ① Non-dissipative
② Diverges as $\omega \rightarrow 0$

Superconductor

Now upgrade A_μ to a dynamical gauge field i.e., "Gauging" $-F(A) = -\int_{\text{Maxwell}} [A]$

$$Z = \int D\phi DA_\mu e^{-S[\phi, A_\mu]} e^{-\frac{1}{4e^2} \int (\partial_\mu A_\nu - \partial_\nu A_\mu)^2} = \int DA_\mu e$$


$$A_x = \frac{\alpha}{L} = \partial_x \left(\frac{\alpha}{L} x \right)$$

$$\phi \rightarrow \phi e^{i\alpha}$$

field ϕ , "Gauging"

$$= \int_{D\theta} DA_\mu e^{-\frac{K}{2}(\partial_\mu \theta - A_\mu)^2 - \frac{1}{4e^2}(\partial_\mu A_\nu - \partial_\nu A_\mu)^2}$$

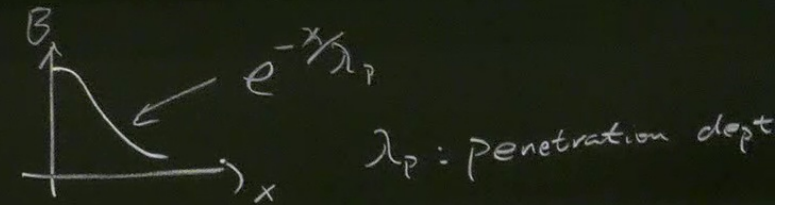
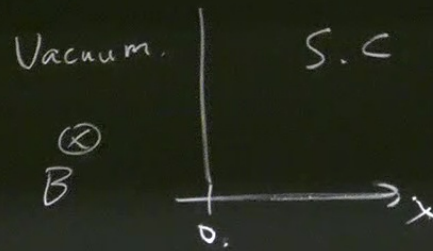
If θ is smooth, choose gauge $A \rightarrow A_\mu + \partial_\mu \theta$

$$\mathcal{L} = \frac{K}{2} A^2 + \frac{1}{4e^2} (\partial \times A)^2 \quad \text{massive vector field, } m = e\sqrt{K}.$$

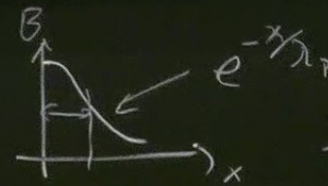
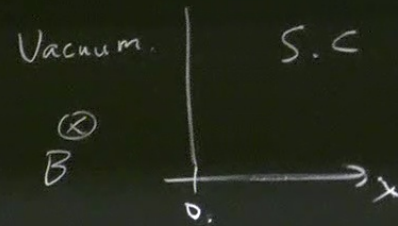
$A - \partial\theta$: gapped. Gauge field & Goldstone gap out each other.

Anderson-Higgs mechanism

Meissner effect:



Meissner effect:



λ_p : penetration depth $\sim \frac{1}{e\sqrt{K}}$