

Title: Lecture - Quantum Matter, PHYS 777

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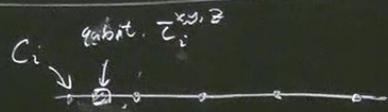
Collection/Series: Quantum Matter (Elective), PHYS 777, March 31 - May 2, 2025

Subject: Condensed Matter

Date: April 04, 2025 - 9:00 AM

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\mathbb{Z}_2 gauge theory, Majorana & Ising



Impose gauge symmetry

$$C_i \rightarrow S_i C_i, \quad \tau_i^x \rightarrow S_i S_{i+1} \tau_i^x$$

$$S_i = \pm 1$$

↓ gauge connection.

$$u \exp \left(i \int_i^i \vec{A} \cdot d\vec{x} \right)$$

Hamiltonian

$$(D \times A)^2$$

$$C_i = \gamma_{2i-1} + i \gamma_{2i}$$

$$H = -J \sum_i \tau_i^x \tau_{i+1}^x - g \sum_i \gamma_{2i-1} \gamma_{2i}$$

We can pick a "smart" gauge.

$$C_i \rightarrow C_i, \quad C_{i+1} \rightarrow \left(\prod_{j \leq i} \tau_j^x \right) C_{i+1}$$

Hamiltonian

$$(\nabla \times A)^2$$

$$C_i = \delta_{2i-1} + i \delta_{2i}$$

$$\prod_{\square} \tau^x \sim \text{e}$$

$$H = -J \sum_i i \delta_{2i} \delta_{2i+1} \tau_i^x - g \sum_i i \delta_{2i} - \delta_{2i}$$

We can pick a "smart" gauge

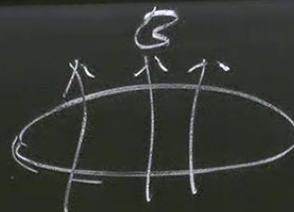
$$C_i \rightarrow C_i, \quad C_{i+1} \rightarrow \left(\prod_{j < i} \tau_j^x \right) C_{i+1}$$

$$\Rightarrow H = -J \sum_{1 \leq i \leq L-1} i \delta_{2i} \delta_{2i+1} - g \sum_i i \delta_{2i} - \delta_{2i} - J i \delta_{2L} \delta_1 \tau_L^x$$

In 1d \mathbb{Z}_2 gauge field, only gauge-inv. def. is the Wilson loop made of τ_L^x

$\gamma_{zi} - \gamma_{zi}$

$$\prod_{\square} \tau^x \sim e^{i \oint \vec{A} \cdot d\vec{\ell}}$$

gauge flux

Constraint from gauge invariance:

Total gauge charge fixed in \mathcal{H} .

Two choices:

① $\prod_{i=1}^L i \gamma_{2i-1} \gamma_{2i} = 1$

② $\left(\prod_{i=1}^L i \gamma_{2i-1} \gamma_{2i} \right) \left(\prod_{i=1}^L \tau_x \right) = 1$

$\xrightarrow{\text{JW transformed}}$
 Ising model.
 $(-1)^F = \tau_L^x = \prod_i X_i$

$\gamma_{zi} = \prod_{j=1}^i \tau_L^x$

made of τ_L^x

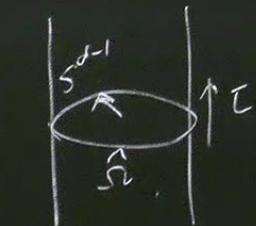
gauge-inv. def. is the Wilson loop

" Ising \leftrightarrow Majorana / \mathbb{Z}_2 "

Now: $\mathbb{Z}_2 \leftrightarrow \underbrace{\tau_{2i-1}^z \tau_{i-1}^z}_{\text{flux-changing operator}}$ a.k.a. "instanton"

How to compute $\Delta_{\mathbb{Z}}$

ory : State-operator correspondence



Cylinder: $S^{d-1} \times \mathbb{R}$.

$$r = e^z$$



$$z = -\infty \leftrightarrow r = 0.$$

Eigenstate on $S^{d-1} \leftrightarrow$ Operator $\mathcal{O}(r=0)$
 $|\Omega\rangle$

$$\text{G.S. } |\Omega\rangle \leftrightarrow 1$$

Time-translation $z \rightarrow z + \lambda \leftrightarrow r \rightarrow r \cdot e^\lambda = \text{scaling}$

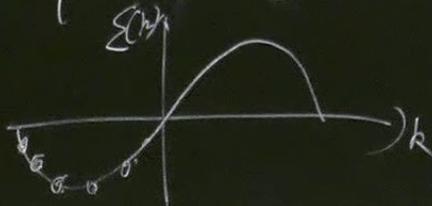
$$\text{Energy } E \leftrightarrow \Delta_{\mathcal{O}}$$

" Ising \leftrightarrow Majorana / Z_2 "

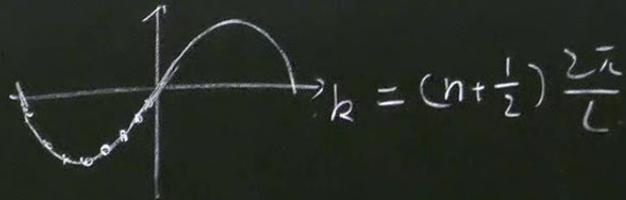
Now: $Z_i \leftrightarrow \underbrace{\tau_{i-1}^z}_{\tau_{i-1}^z}$: flux-changing operator
a.k.a. "instanton"

How to compute Δ_Z ?

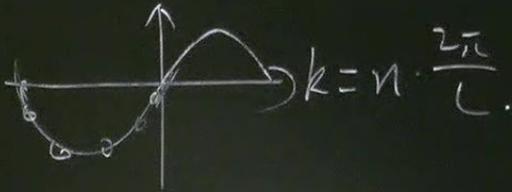
Just compute E of some states.



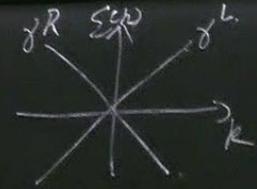
G.S. : anti-p.b.c. $\Leftrightarrow \prod_i X_i = 1$



Lowest excitation w/ $\prod_i X_i = -1 \Leftrightarrow$ p.b.c.



Massless Majorana.



Trick: 2x Majorana = Dirac.

$$\bar{\gamma}_1 \not{\partial} \gamma_1 + \bar{\gamma}_2 \not{\partial} \gamma_2 \longleftrightarrow \bar{\psi} \not{\partial} \psi$$

$\psi = \gamma_1 + i\gamma_2$

$$E_{\text{Majorana}} = \frac{1}{2} E_{\text{Dirac}}$$

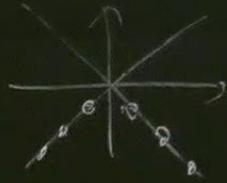
$$E_{\text{Majorana}} = \frac{1}{2} E_{\text{Dirac}}$$

$$= \frac{1}{2} \sum_{k: \epsilon(k) < 0} \epsilon(k)$$

$$= - \sum_{k > 0} |k|$$

$$\text{Anti-pbc} = -\frac{1}{2} (1+3+5+7+\dots)$$

$$\text{Pbc} = -(1+2+3+4+\dots)$$



$$\epsilon(k) = -|k|$$

$$k = \frac{n+\frac{1}{2}}{n} \cdot \frac{2\pi}{L}$$

Deep facts:

$$1+2+3+4+\dots = -\frac{1}{12}$$

$$1+3+5+7+\dots = \frac{1}{12}$$

$$\Rightarrow \Delta_2 = E_{\text{pbc}} - E_{\text{Apbc}} = \frac{1}{12} + \frac{1}{24} = \frac{1}{8}$$

Deep facts:

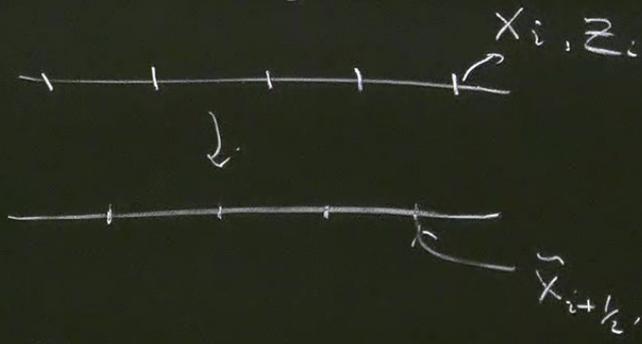
$$1+2+3+4+\dots = -\frac{1}{12}$$

$$1+3+5+7+\dots = \frac{1}{12}$$

$$\Rightarrow \Delta_Z = E_{pbc} - E_{Apbc} = \frac{1}{12} + \frac{1}{24} = \frac{1}{8}$$

$$\langle Z_i Z_j \rangle \sim \frac{1}{|i-j|^{1/4}}$$

Another duality for Ising = Kramers-Wannier (KW)



\tilde{Z} : domain-wall creation.

Z : charge creation.

$$\tilde{X}_{i+1/2} = Z_i Z_{i+1}$$

$$\tilde{Z}_{i+1/2} = \prod_{1 \leq j \leq i} X_j$$

Consider a ferromagnet

$|\uparrow\uparrow\uparrow\uparrow\rangle$

$\tilde{Z}_{i+1/2} |\uparrow\uparrow\uparrow\uparrow\rangle = |\downarrow\downarrow\downarrow\uparrow\rangle$

Let Z_2 gauge field, only gauge-inv. def. is the Wilson loop $\prod_{i=1}^L (1 + b_{2i-1} \tau_i) \prod_{i=1}^L (1 + b_x) = (-1)^F$

w) Side note:
 JW transform: $\gamma_i = \tilde{Z}_{i-1/2} Z_i$. Domain wall + charge \Rightarrow fermion.
 ("Flux-attachment")

$$-J \sum Z_i Z_{i+1} - g \sum X_i \Leftrightarrow -g \sum \tilde{Z}_{i-1/2} \tilde{Z}_{i+1/2} - J \sum \tilde{X}_{i+1/2}$$

KW: $g \leftrightarrow J$ Symmetric phase = Defect (domain wall) condensate.
 Para \leftrightarrow Ferro
 Ferro \leftrightarrow Para.

Kramers-Wannier (KW)

\tilde{Z} : domain-wall creation.

Z : charge creation.

$\uparrow\uparrow\uparrow \downarrow\downarrow\downarrow$

ferrimagnet $\pi_0(Z_2) = Z_2$

$\uparrow\uparrow\uparrow \rightarrow$

$\uparrow\uparrow\uparrow\uparrow\uparrow \rightarrow \downarrow\downarrow\downarrow\downarrow\downarrow \uparrow\uparrow\uparrow\uparrow \rightarrow$

Side note:

JW transform: $\gamma_i = \tilde{Z}_{i-1/2} Z_i$. Domain wall \Rightarrow ("Flux-attach")

$$-J \sum Z_i Z_{i+1} - g \sum X_i \Leftrightarrow -g \sum \tilde{Z}_{i-1/2} \tilde{Z}_{i+1/2}$$

KW: $g \leftrightarrow J$

Para \leftrightarrow Ferro

Ferro \leftrightarrow Para.

Symmetric phase = Def

Let's learn something about gauge theory

$g \gg J$. Ising: unique g.s. $|\rightarrow \rightarrow \rightarrow \dots \rightarrow$

KW: $\tilde{\Sigma}$ ferromagnet,

$|\uparrow \uparrow \uparrow \dots \uparrow + \downarrow \downarrow \downarrow \dots \downarrow \rightarrow \checkmark$
&

$|\uparrow \uparrow \uparrow \dots \uparrow - \downarrow \downarrow \downarrow \dots \downarrow \rightarrow$: killed by gauge field ($\Rightarrow \vec{\pi} \cdot \vec{x} = 1$)

Low energy states due to SSB, Killed by gauge field.

\simeq "Higgs Mechanism".

$g \ll J$ Ising: G.S: $|\uparrow\uparrow\uparrow\uparrow \dots \rightarrow \pm 1 \downarrow\downarrow\downarrow \dots \rangle$

KW: $\sum \vec{S}$ paramagnet, at low energy \Rightarrow pure \mathbb{Z}_2 gauge theory

\Rightarrow Pure \mathbb{Z}_2 gauge theory has 2-fold degeneracy w./ dif

As long as \vec{S} perturbation is not allowed

But if I turn on $H_0 - \delta \sum_i Z_i$

$\Rightarrow |\uparrow\uparrow\uparrow\uparrow \rangle$ unique G.S. ($\Delta E_{\downarrow\downarrow} \propto \delta \cdot L$)

Domain wall $|\uparrow\uparrow\uparrow \underbrace{\downarrow\downarrow\downarrow\downarrow\downarrow}_{\vec{r}} \uparrow\uparrow\uparrow \rangle$ Energy cost $\begin{cases} \delta=0: \Delta E \sim 0(L) \\ \delta \neq 0: \Delta E \sim 0(L \cdot r) \end{cases}$

"Instanton \Rightarrow confinement"

v./ different Z_2 flux

$E \sim O(1)$
 $E \sim O(16 \cdot r)$

kW

Z_2 gauge charge "confined"

