

Title: Lecture - Quantum Matter, PHYS 777

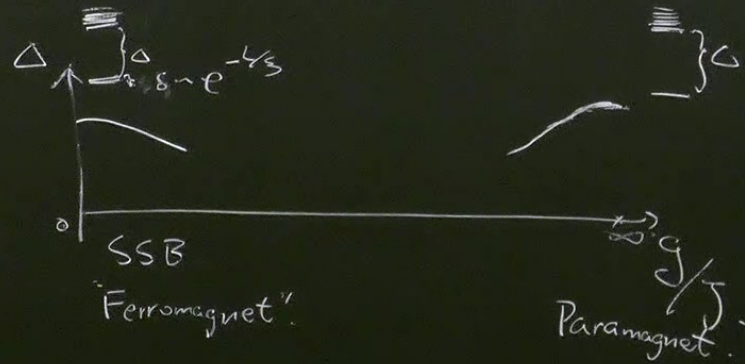
Speakers: Chong Wang

Collection/Series: Quantum Matter (Elective), PHYS 777, March 31 - May 2, 2025

Subject: Condensed Matter

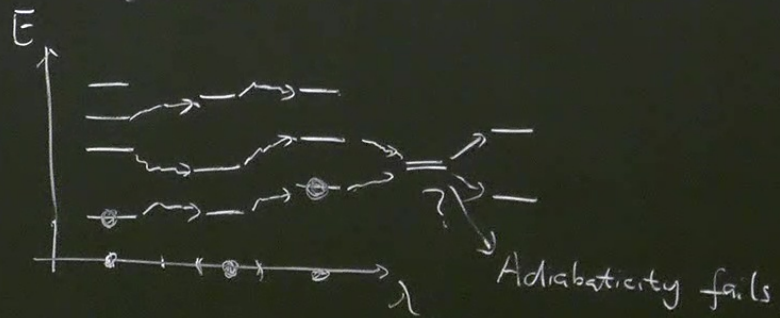
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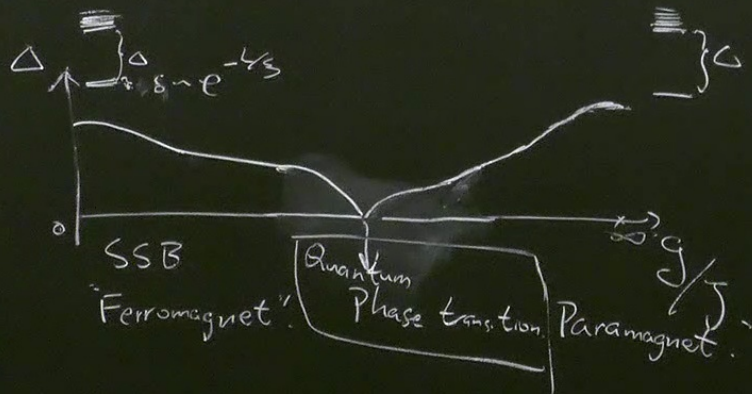
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$$H = -J \sum_i z_i z_{i+1} - g \sum_i x_i$$

Principle of Adiabatic continuation





$$H = -J \sum_i z_i z_{i+1} - g \sum_i x_i$$

Solving 1d TFIM. w/ Jordan-Wigner (JW)

1d spin chain \longrightarrow fermion

Use Majorana basis. $C_i = \frac{\gamma_{2i-1} + i\gamma_{2i}}{2}$

γ_{2i-1}
 γ_{2i}

1d. TFIM. w/ Jordan-Wigner (JW)

Ising chain

→ "fermion chain"

C_i . $\{C_i\}$

X_i, Y_i, Z_i

Use Majorana basis.

$$C_i = \frac{\gamma_{2i-1} + i\gamma_{2i}}{2}$$

$$C_i^+ = \frac{\gamma_{2i} - i\gamma_{2i-1}}{2}$$

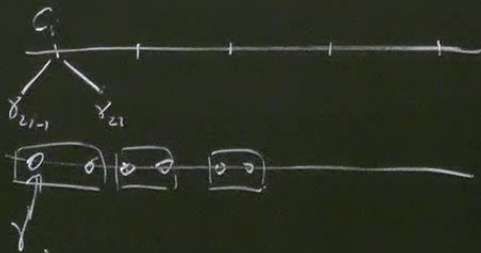
or,

$$C_i + C_i^+, \quad \gamma_{2i} = \frac{C_i - C_i^+}{i}$$

$$\gamma_j = \gamma_j^+$$

$$\{\gamma_j, \gamma_k\} = 2\delta_{jk} \Rightarrow$$

$$2\gamma_{2i-1}\gamma_{2i}$$



→ "fermion chain" c_i . $\{c_i, c_j^\dagger\} = \delta_{ij}$

or, $\gamma_{2i-1} = c_i + c_i^\dagger$, $\gamma_{2i} = \frac{c_i - c_i^\dagger}{i}$

$$\gamma_j^2 = 1$$

↓

Majorana fermion.

$$i \gamma_{2i-1} \gamma_{2i} = 2 \frac{c_i^\dagger c_i - c_i c_i^\dagger}{i} = (-1)^{n_i}$$

JW transform.

$$X_i = i \gamma_{2i-1} \gamma_{2i}$$

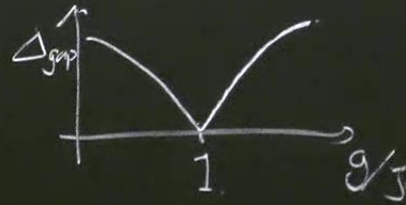
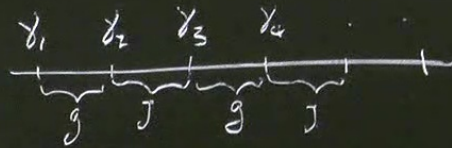
$$Z_i = \left[\prod_{1 \leq j < i} (i \gamma_{2j-1} \gamma_{2j}) \right] \gamma_{2i-1}$$

Check: $X_i^2 = Z_i^2 = 1$

$$X_i Z_j = (-1)^{\delta_{ij}} Z_j X_i$$

$$H = -J \sum_i \psi_{2i} \psi_{2i-1} - g \sum_i \psi_{2i-1} \psi_{2i}$$

free fermion model!



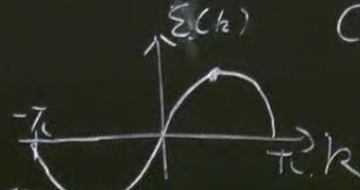
Take $g=J$. $H = -iJ \sum_{j=1}^{2L} \psi_j \psi_{j+1}$

Try $\psi_k := \frac{1}{\sqrt{4L}} \sum_j e^{-ikj} \psi_j$, k

Can check: if $k=0, \pi$, $\sqrt{2}\psi_k$ is for other k , $\psi_k^\dagger = \psi_{-k}$

$\Rightarrow d_R = \psi_k$

$\Rightarrow H = 2J \sum_k \sin k \psi_{-k} \psi_k$



$S, N_k = d_k^\dagger d_k = \psi_{-k} \psi_k = \begin{cases} 1, & \pi < k < 0 \\ 0, & 0 < k < \pi \end{cases}$ Same condition.

free fermion model!

$$\text{Take } g=j, \quad |-\rangle = -i \sum_{j=1}^{2L} \psi_j \psi_{j+1}$$

$$\text{Try } \gamma_k := \frac{1}{\sqrt{4L}} \sum_j e^{-ikj} \psi_j, \quad k = \frac{2\pi}{2L} \cdot n, \quad n=0, 1, \dots, 2L-1$$

Can check: if $k=0, \pi$, $\sqrt{2}\gamma_k$ is a Majorana.

for other k , $\gamma_k^\dagger = \gamma_{-k}$, $\{\gamma_k, \gamma_{-k'}\} = \delta_{k,k'}$.

$$\Rightarrow d_R = \gamma_k, \quad d_k^\dagger = \gamma_{-k}$$

Low energy ($E \ll J$): Expand near $k=0, \pi$

Take a "cutoff" momentum Λ , $\frac{2\pi}{L} \ll \Lambda \ll \pi$

For $\delta k \ll \Lambda$, define

$$\gamma_{\delta k}^L := \gamma_{\delta k}, \quad \gamma_{\delta k}^R = \gamma_{\delta k + \pi}$$

$$\Rightarrow H \approx \sum_{|k| < \Lambda} 2J \left[k \gamma_{-k}^L \gamma_k - k \gamma_{-k}^R \gamma_k \right]$$

Define "continuum" fermion

$$\psi^{L,R}(x) := \sqrt{\frac{\pi}{L\Lambda}} \sum_k e^{ikx} \gamma_k^{L,R}$$

↓
Hermitian.

$$\Rightarrow H \approx 2J \int_0^L dx$$

$$\Rightarrow H \approx 25 \int_0^L dx \left[\psi^L (-i\partial_x) \psi^L - \psi^R (-i\partial_x) \psi^R \right]$$

$\begin{pmatrix} \psi^L \\ \psi^R \end{pmatrix}$: Relativistic, Massless Majorana fermion in $(1+1)d$.

Lorentz inv. emerges at low energy.

In Lagrangian form:

$$S = \int dt dx \left[\psi^L (i\partial_t - i\partial_x) \psi^L + \psi^R (i\partial_t + i\partial_x) \psi^R \right]$$

For $\delta k \ll \Lambda$, define.

$$\gamma_{sk}^L := \gamma_{sk} \quad \gamma_{sk}^R = \gamma_{sk+\pi}$$

$$\Rightarrow H \approx \sum_{|k| < \Lambda} 2\epsilon_k \left[k \gamma_{-k}^L \gamma_k - k \gamma_{-k}^R \gamma_k^R \right]$$

Define "continuum" fermion

$$\psi^{L,R}(x) := \sqrt{\frac{\pi}{L\Lambda}} \sum_k e^{ikx} \gamma_k^{L,R}$$

↓
Hermitian.

$(\bar{\psi} \psi)$: Relativistic, Massless Majorana fermion in (1+1)d.

Lorentz inv. emerges at low energy.

In Lagrangian form:

$$S = \int dt dx \left[\bar{\psi}^L (i\partial_t - i\partial_x) \psi^L + \bar{\psi}^R (i\partial_t + i\partial_x) \psi^R \right] \quad (*)$$

Relevant op.

only one: $i\partial_x \psi_R$ = mass. $\Delta_m = 1 < 2 \Leftrightarrow$

Scaling invariant: $x \rightarrow bx, t \rightarrow bt, \psi^{L,R} \rightarrow b^{-1/2} \psi^{L,R} \Rightarrow$ Renormalization group (RG) fixed point.

Recall in RG: scaling operator $O \rightarrow b^{-\Delta_O} O, \Delta_O$: scaling dim.

As a perturbation $S + \lambda \int dt dx O, \lambda \rightarrow b^{2-\Delta_O} \lambda,$

Δ_O also controls many other things,

e.g. $\langle \mathcal{O}^+(x) \mathcal{O}^-(y) \rangle \sim \frac{1}{|x-y|^{2\Delta_O}}$

$\Delta_O < 2$: relevant

$\Delta_O > 2$: irrelevant

$\Delta_O = 2$: marginal.

op.
only one γ_L, γ_R : mass. $\Delta_m = 1. < 2 \Leftrightarrow Z_i Z_{i+1} - X_i$

\Rightarrow Renormalization group (RG) fixed point.

ing dim.

$$b^{2-\Delta_0} \lambda$$

$\Delta_0 < 2$: relevant

$\Delta_0 > 2$: irrelevant

$\Delta_0 = 2$: marginal.

When $m > 0$, energy gap $\Delta(m) > 0$.

Under scaling: $m \rightarrow b^{2-\Delta_m} m = b m$

$$\Delta \rightarrow b \Delta$$

$$t \rightarrow b t, \quad \varepsilon \rightarrow b^{-1} \varepsilon$$

$$\Rightarrow \Delta \propto m^{\# = 1}$$

Lesson: in $(1+1)d$,

Ising mode " = " Majorana fermion.
critical pt Massless.

Q

Ferromagnet $\xrightarrow{\text{Phase transition}}$ Paramagnet

Q: ① How could a spin model ("bosonic", $[O_i, O_j] = 0$) be equivalent to a fermion model?
 ② What is Z_i operator in fermion model? How to compute Δ_Z ?

Back to JW transform: $\gamma_{2i-1} = \left(\prod_{|s|<i} X_j \right) Z_i$

$\gamma_{2i} = \left(\prod_{|s|<i} X_j \right) Y_i$ Fermions are non-local

p.b.c if $X = -1$
 Anti-p.b.c if $X = +1$

P.b.c $-J Z_{2i} Z_{2i+1} = -J \left(\prod_{j=1}^L X_j \right) (-i \gamma_{2i} \gamma_{2i+1}) = -J (-X) i \gamma_{2i} \gamma_{2i+1}$, can absorb $(-X)$ by defining $\gamma_{2i+1} = \gamma_i(X)$

B.C. of fermion is a dynamical d.o.f. \rightarrow Signs of a gauge theory (Gauge group = Z_2)