Title: Symmetry and Causality Constraints on Fermi Liquids

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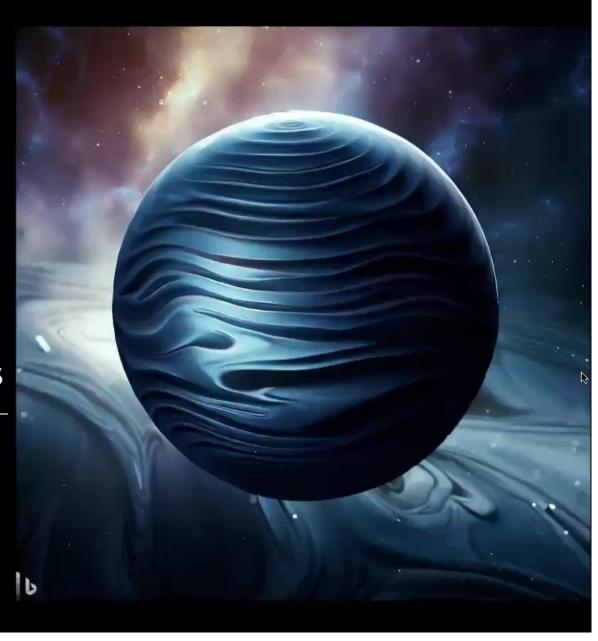
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Symmetry and causality constraints on Fermi liquids

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in collaboration with



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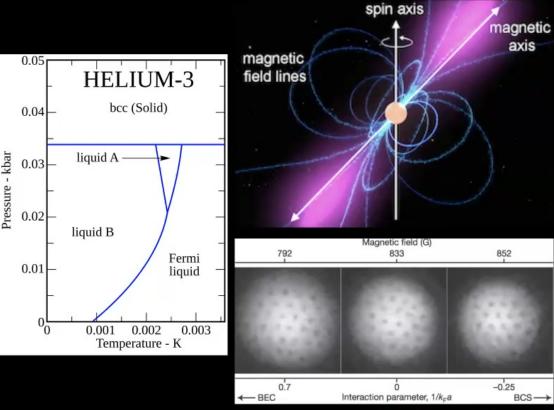
Luca Delacretaz

Chowdhury, Delacretaz, UM, arXiv:2501.02073

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Space-time symmetries in Fermi liquid phases

- He-3 (Galilean)
- Neutron stars, dense QCD (Lorentz)
- Unitary Fermi gas, nonrelativistic anyon gas (Schroedinger)
- Interacting fermionic CFTs at finite density? (conformal)



Zwierlein et al, Nature 05

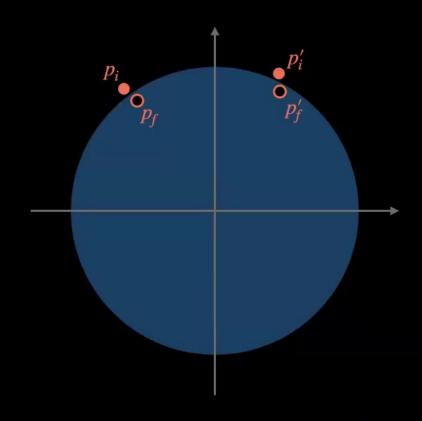
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History: Boosts constrain forward-scattering

• Effective mass relation:

$$m_* \equiv \frac{p_F}{v_F} = \begin{cases} (1 + F_1)m & \text{Galilei} \\ (1 + F_1)\mu & \text{Lorentz} \end{cases}$$

• F_1 measures the strength of current-current interactions



Landau, Baym+Chin

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Constraints from thermodynamics

Boost constraints from momentum susceptibility

$$\chi_{\pi\pi} = \begin{cases} m\rho & \text{Galilei} \\ \varepsilon + p = \mu\rho & \text{Lorentz} \end{cases}$$

Scale constraints from charge susceptibility

$$\rho \propto \mu^{d/z}, \qquad \chi_{\rho\rho} = \frac{\partial \rho}{\partial \mu}, \qquad \Longrightarrow \qquad 1 + F_0 = z \frac{\mu}{p_F v_F}$$

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Are thermodynamic constraints sufficient?

A

- Need for a field theoretic approach (will use the coadjoint orbit action)
 - How do symmetries act on the space of states?
 - What constraints are obtained from requiring the action to be invariant under such transformations?

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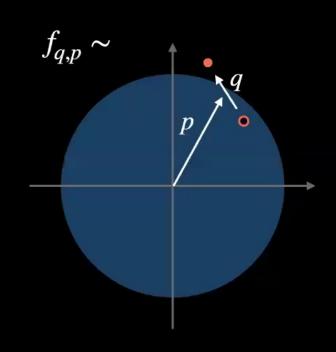
Beyond forward scattering?

$$H = \int_{xp} \epsilon_{p} f_{x,p}$$

$$+ \int_{xpp'} F_{pp'}^{(2,0)} \delta f_{x,p} \delta f_{x,p'} + F_{pp'}^{(2,1)} \left(\nabla_{x} \delta f_{x,p} \right) \delta f_{x,p'} + \dots$$

$$+ \int_{xpp'p''} F_{pp'p''}^{(3,0)} \delta f_{x,p} \delta f_{x,p'} \delta f_{x,p'} \delta f_{x,p''} + \dots$$

$$+ \dots$$



Constraints on higher Landau paramaters?

Delacretaz, Du, UM, Son, PRR 23

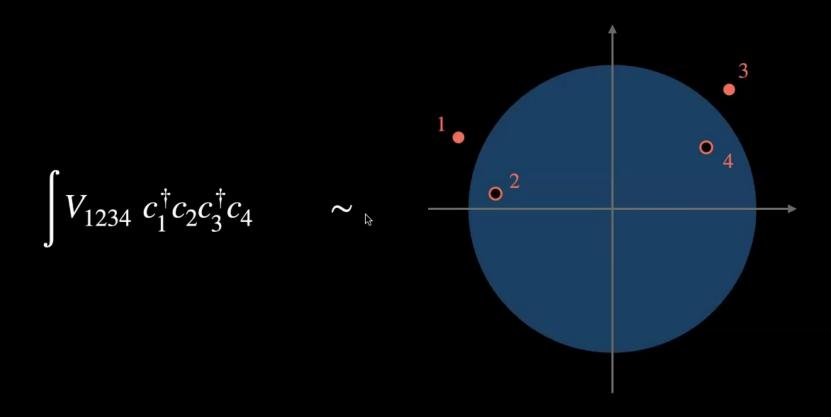
UM, arXiv: 2307.02536

Outline

- Higher Landau parameters and irrelevant interactions
- Parameterizing the space of states using canonical transformations
- Symmetries as canonical transformations
- Constraints for boosts and dilatations
- (Micro)causality bounds on Landau parameters and consequences for conformal Fermi liquids

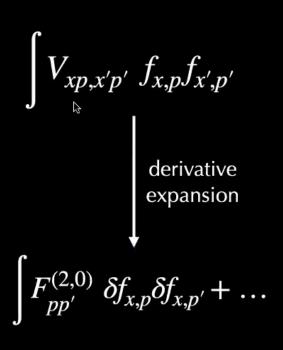
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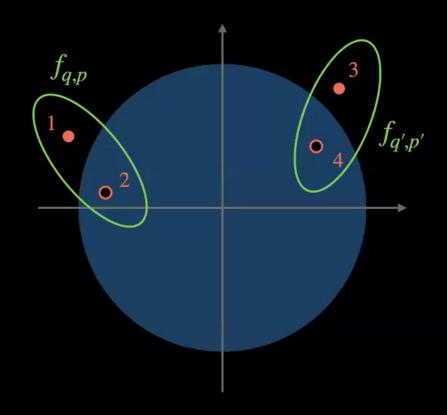
Interpreting the higher Landau functions



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Interpreting the higher Landau functions



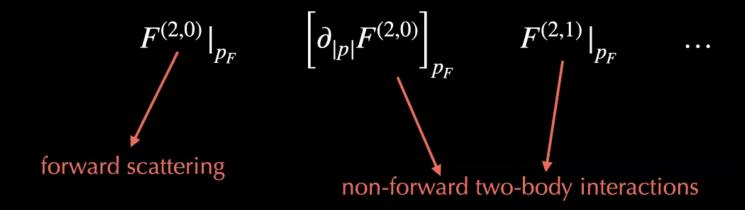


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Interpreting the higher Landau functions

$$\delta f_{x,p} \sim \delta_{p_F}(p) + \delta_{p_F}'(p) + \dots$$

localizes on the Fermi surface under an expansion in fluctuations



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Higher Landau parameters (2D)

$$F^{(2,0)}|_{p_F} \sim \sum_l F_l e^{il(\theta-\theta')}$$

Landau parameters

$$\left[\partial_{|p|}F^{(2,0)}\right]_{p_F} \sim \sum_l \tilde{F}_l e^{il(\theta-\theta')}$$

higher Landau parameters

$$_{\text{R}}F^{(3,0)}|_{p_F} \sim \sum_{l,l'} G_{l,l'} e^{il(\theta-\theta')+il'(\theta-\theta'')}$$

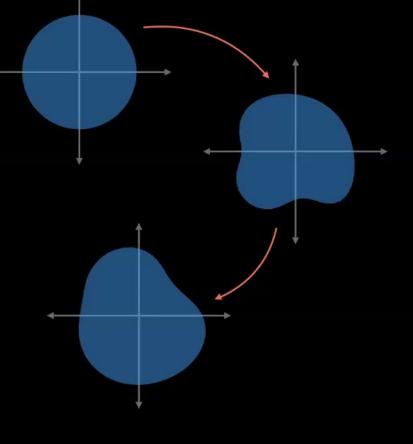
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The space of states

 Particle-hole coherent states are generated by canonical transformations

$$|\phi_{xp}\rangle = e^{i\int\phi_{xp}[\psi^{\dagger}\psi]_{xp}}|FS\rangle$$

The space of states forms a coset G/H
 where G is the group of canonical
 transformations and H consists of
 transformations that leave the ground state
 invariant



Delacretaz, Du, UM, Son, PRR 23

UM, arXiv: 2307.02536

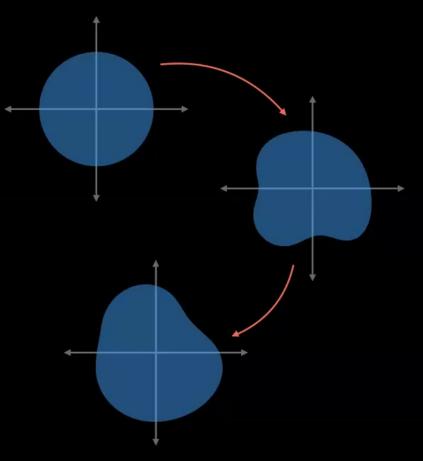
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The space of states

 All states are related by canonical transformations, whose action on the distribution function is given by

$$f_{xp} \xrightarrow{F_{xp}} f + \{F,f\} + \frac{1}{2} \{F, \{F,f\}\} + \dots$$

 Space-time symmetries may be embedded into canonical transformations to deduce their action on the distribution function



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Symmetries as canonical transformations

Symmetries that don't act on time:

Translations	$\{p_i, \cdot\}$
Rotations	$\{\epsilon_{i_1i_{d-2}jk}x_jp_k,\cdot\}$
Galilean boosts	$\{mx_i-tp_i,\cdot\}$

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Symmetries as canonical transformations

Symmetries that don't mix time and space:

Dilatations

 $\{-x_ip_i,\cdot\}+zt\partial_t$

Non-relativistic special conformal transformations

$$\left\{\frac{mx^2}{2} - tx_i p_i, \cdot\right\} + t^2 \partial_t$$

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Deriving the constraints

- Use transformation of f_{xp} to impose invariance of bosonized action $S[f_{xp}]$
- Expand f_{xp} in fluctuations $\phi_{x\theta}$ or $\delta p_F(x,\theta)$ (any any other parametrization)
- Set variation to zero order by order to obtain linear and non-linear constraints
- Check: Galilean effective mass constraint and scale invariance constraint can be derived in this way at leading order in fluctuations

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Non-linear (quadratic) constraints

Galilei:

$$mv_F \tilde{F}_1 + mp_F \epsilon_F'' - p_F = 0$$

$$3G_{1,l} + v_F(\tilde{F}_l + \tilde{F}_{l+1}) + [lF_l - (l+1)F_{l+1}] = 0$$

Dilatations:

$$v_F \tilde{F}_0 + p_F \epsilon_F'' - (z - 1)v_F = 0$$

$$3G_{l,0} + v_F(\tilde{F}_l + \tilde{F}_{-l}) + (z - 2)v_F F_l = 0$$

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An alternate derivation

- The constraints can also be derived from Ward identities by identifying conserved currents $j^{\mu}[f_{xp}]$ and $T^{\mu\nu}[f_{xp}]$
- Galilean invariance: $T^{0i} = mj^i$
- Scale (+conformal) invariance: $zT_0^0 + T_i^i = 0$
- Imposing only scale invariance at the level of the action (for both z=1,2) gives the same result as the Ward identity for scale+conformal invariance
 - Scale invariance \Longrightarrow conformal invariance even for z = 2?

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Lorentz invariance

- Lorentz boosts mix time and space can't be implemented as canonical transformations
- Can still impose Lorentz invariance as a Ward identity $T^{0i} = T^{i0}$
- Linear constraint agrees with before + new quadratic constraint

$$\mu v_F \tilde{F}_1 + \mu p_F e_F'' + p_F v_F^2 - p_F = 0$$

$$3\pi^{2}\mu G_{1,l} + \pi^{2}\mu v_{F}(\tilde{F}_{l} + \tilde{F}_{l+1})$$

$$+\pi^{2}v_{F}(p_{F}v_{F} + (l+1)\mu)F_{l+1}$$

$$+\pi^{2}v_{F}(p_{F}v_{F} + l\mu)F_{l}$$

$$+p_{F}v_{F}^{2}F_{l}F_{l+1} = 0$$

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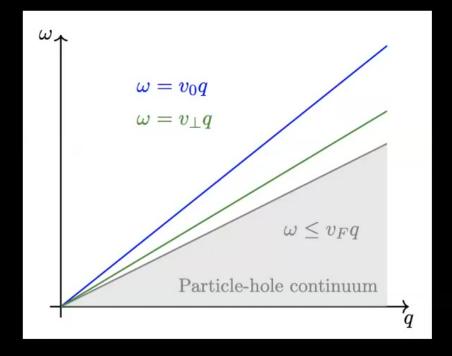
Infinite towers of constraints

- ullet The variation of the action under a symmetry transformation mixes orders in δp_F or ϕ
- ullet Constraints relate higher Landau parameters like $G_{l,l'}$ and \tilde{F}_l to lower ones like F_l mixing happens at all orders
- We obtain an infinite tower of constraints between various higher Landau parameters
 - In particular, if e.g. $F_1 \neq 0$ in a boost invariant system, an infinite set of irrelevant interactions are forced to be nonzero. These will contribute to leading order in nonlinear response, e.g. $F^{(3,0)}|_{p_F}$ and $\left[\partial_{|p|}F^{(2,0)}\right]_{p_F}$ to $\langle\rho\rho\rho\rangle$

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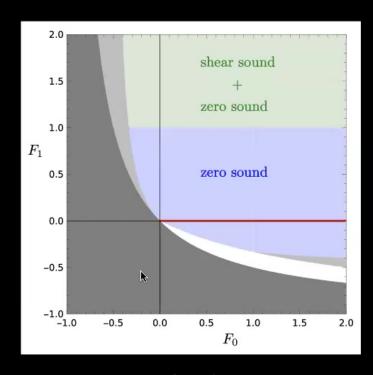
Collective modes and causality

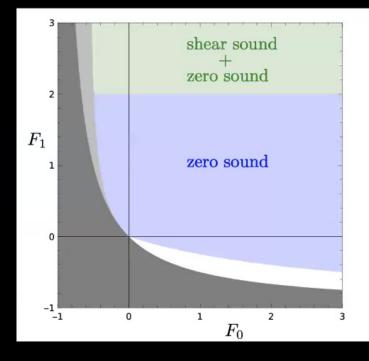
- ullet Zero and shear sounds in Fermi liquids, corresponding to the presence of F_0 and F_1
- $v_0(F_0, F_1) \le c$, $v_{\perp}(F_1) \le c$, $v_F \le c$ as a result of "microcausality"
- Relativistic conformal invariance fixes v_F in terms of F_0 and F_1 , forbidding certain values of (F_0, F_1)



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Causality constraints





Dark grey: excluded by $v_F \le c$ Light grey: excluded by bounds on sound

d = 2

d = 3

Outlook

- A (mostly) systematic way to impose microscopic space-time symmetries on Fermi liquids
- Infinite towers of symmetry constraints on interactions, organized by (ir)relevance of interactions

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- Marginal interactions necessitate the presence of higher Landau parameters, which contribute to leading order in non-linear response
- Constraining the parameter space for conformal Fermi liquids using causality
 new causality bounds on conformal couplings?

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