Title: Workshop

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**Collection/Series:** Emmy Noether Workshop: Quantum Space Time

**Subject:** Quantum Gravity

**Date:** March 11, 2025 - 1:30 PM

**URL:** https://pirsa.org/25030062

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# Quantum Gravity in the era of Gravitational-Wave astronomy

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# Question: QG phenomenology in relation to GWs

- model-dependent approach
- model-independent approach

Cosmological models motivated by QG proposals

- Cosmological GW observables:

   amplitude of the GW background
   propagation speed of GWs
   luminosity distance of GW source

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## Question: QG phenomenology in relation to GWs

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- Cosmological GW observables:

   amplitude of the GW background
   propagation speed of GWs
   luminosity distance of GW source

Many QG models : no observable signal

But some predict a blue-tilted GWB or a modified luminosity distance (detectable by future GW interferometers)

Difficult, but still possible, to test QG with GW observations

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## Observations of CBCs by LVK: confirmation of all GR predictions

The low curvature, low energies and large distances characterising production and propagation of GWs difficult to probe Planck-size effects, unless cumulatively amplified by some cosmological mechanism

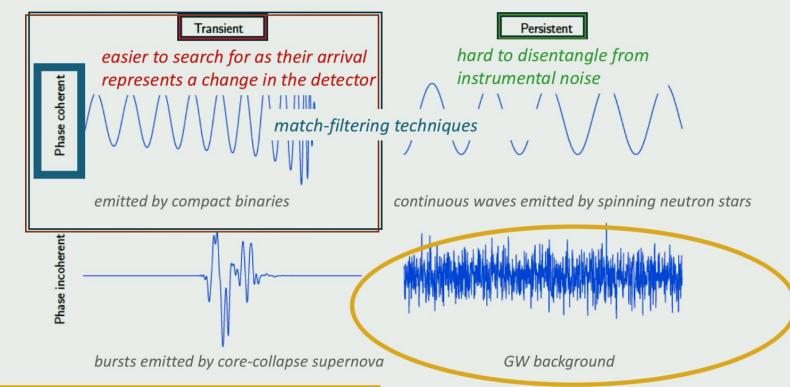
QG proposals where GW observables have been calculated:

- Stelle gravity (issues about unitarity)
- String theory
- Asymptotic safety
- Causal dynamical triangulations
- Non-local gravity
- Horava-Lifshitz gravity (issues about violation of Lorentz invariance)
- Canonical quantum cosmology
- Pre-big-gang cosmology
- String cosmology
- New ekpyrotic scenario
- Brandenberger-Ho non-commutative inflation
- Non-commutative κ-Minkowski
- Padmanabhan's non-local field theory
- Multi-fractional spacetimes

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### Classification of GW signal morphologies

LVK Collaboration, O1, O2, O3



$$\Omega_{\mathrm{gw}}(f) = rac{1}{
ho_c} rac{d
ho_{\mathrm{gw}}}{d\ln f} \cdot rac{
ho_{\mathrm{GW}} \sim \dot{h}^2}{\Omega_{\mathrm{GW}}(f) = \Omega_{\mathrm{ref}} \left(rac{f}{f_{\mathrm{ref}}}
ight)^{lpha}}$$

random nondeterministic phase evolution from a large number of distant sources

### GWB of early universe origin

primordial tensor spectrum of tensor perturbations of the metric  $\mathscr{P}_{\mathsf{t}}(f)$ 

$$ds^{2} = -dt^{2} + a^{2}(t) \delta_{ij} dx^{i} dx^{j} = a^{2}(\tau) \left( -d\tau^{2} + \delta_{ij} dx^{i} dx^{j} \right)$$

with proper wavelength  $\lambda$  :  $f = k/(2\pi) = a/\lambda$ 

 $\mathscr{P}_{\mathsf{t}}(f)$  small compared to  $\mathscr{P}_{\mathsf{s}}(f)$ 

$$r := \mathscr{P}_{\mathrm{t}}/\mathscr{P}_{\mathrm{s}} \ < 0.068$$
 at 95% CL

Planck Collaboration (2020)

$$\mathscr{P}_{\mathsf{t}}(f) = \mathscr{P}_{\mathsf{t}}(f_0) \left(\frac{f}{f_0}\right) \int_{0}^{n_{\mathsf{t}}(f_0) + \frac{1}{2}\alpha_{\mathsf{t}}(f_0) \ln \frac{f}{f_0}} n_{\mathsf{t}} := \mathrm{d}\ln \mathscr{P}_{\mathsf{t}}/\mathrm{d}\ln f$$

$$\alpha_{\mathsf{t}} := \mathrm{d}n_{\mathsf{t}}/\mathrm{d}\ln f$$

$$f_0 = 7.7 \times 10^{-17} \,\mathrm{H}$$

GR + inflation:

negative

$$n_{\mathsf{t}} := \mathrm{d} \ln \mathscr{P}_{\mathsf{t}} / \mathrm{d} \ln f$$

$$\alpha_{\mathsf{t}} := \mathrm{d} n_{\mathsf{t}} / \mathrm{d} \ln f$$

$$f_0 = 7.7 \times 10^{-17} \text{Hz} \longrightarrow k_0 = 0.05 \,\text{Mpc}^{-1}$$

red-tilted spectrum

#### GWB of early universe origin

$$\Omega_{\text{GW}}(f) := \frac{1}{\rho_{\text{crit}}} \frac{\mathrm{d}\rho_{\text{GW}}}{\mathrm{d}\ln f} = \frac{\pi^2 f^2}{3H_0^2} \mathscr{P}_{\mathsf{t}}(f) \mathscr{T}^2(f)$$

$$\rho_{\text{GW}} := [M_{\text{Pl}}^2/(8a^2)] \langle (\partial_{\tau} h_{ij})^2 + (\nabla h_{ij})^2 \rangle$$

transfer function (depends on expansion history of the universe)

GW amplitude of the mode which enters the horizon during RDE:

$$\Omega_{\text{GW}}(f) = \frac{\mathscr{A}}{h^2} r \left(\frac{f}{f_0}\right)^{n_{\text{t}} + \frac{\alpha_{\text{t}}}{2} \ln \frac{f}{f_0}}, \qquad \mathscr{A} \approx 1.4 \times 10^{-15}.$$

A blue-tilted spectrum can produce a stochastic GWB with increasing amplitude that can reach the sensitivity thresholds of GW interferometers at high frequencies

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### Flux-compactification models in string cosmology

nt < 0

red-tilted spectrum

Baumann, McAllister (2015)

#### Note:

- old ekpyrotic scenario: strongly blue-tilted tensor index  $n_t = 2$
- recent single-field version: perturbations are generated before the ekpyrotic phase, and the scalar spectrum is red-tilted (previous versions are ruled out because of a blue-tilted scalar spectrum)
- in all the realisations of the ekpyrotic model: **tensor-to-scalar ratio is exceptionally small** and the resulting **stochastic GWB is well below the detection threshold of any present or future interferometer**

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### Flux-compactification models in string cosmology

red-tilted spectrum

Baumann, McAllister (2015)

Wheeler-De Witt canonical quantum cosmology

pivot scale of the experiment

Kiefer, Krammer (2012) Kramenshchik, et al (2018), 2015)

Bini, et al (2013)

Brizuela, et al (106)

 $\propto H^2$ spectrum at HC: k=aH known >0 numerical

<<1, strong suppression term; increased at late times by (ko/k) LISA or DECIGO:  $k_0/k \sim 10^{-15} - 10^{-13}$ 

constant

Quantum correction is unobservable

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-: blue tilt

+: red tilt

### **Loop Quantum Cosmology**

Dressed-metric approach

quantum fluctuations are described as propagating on a quantum background spacetime which can be described by a dressed metric

Agullo, Morris (2015) Li, Singh, Wang (2020)



red-tilted spectrum

 Effective-constraints approach (quantum corrections from inverse-volume operators, from holonomies, or from both)

Red-tilted tensor spectrum

Bojowald, calcagni, Tsujikawa (2011) Zhu, et al (2016) Inverse-volume corrections react to the discreteness scale of QG

- They depend only on triad variables and not also on connections
- They are not suppressed at the inflationary density scale

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### **Loop Quantum Cosmology**

- Dressed-metric approach
- Effective-constraints approach (quantum corrections from inverse-volume operators, from holonomies, or from both)
- Hybrid-quantisation approach

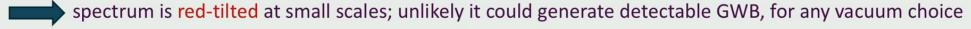
combine different types of quantum representations for the different degrees of freedom

nt(k) depends on the vacuum on which to perturb the background solution

- for some choices of the vacuum: Nt < 0 red-tilted spectrum
- for others the spectrum oscillates rapidly and for some frequencies, there is a blue tilt

GWB decreases in k when k is sufficiently large

de Blas, Olmedo (2016) Castello Gomar, et al (2017)



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Non-local QG (embedding of Starobinsky gravity into a fundamental theory)

primordial tensor spectrum as a function of the comoving wave-number

$$\mathscr{P}_{\mathsf{t}}(k) \sim \left[1 - \frac{3}{2\ln(k_{\mathsf{e}}/k)}\right] e^{-\tilde{\mathsf{H}}_{2}\left[\frac{2H^{2}(k)}{M_{*}^{2}}\right]} \int_{t}^{t}$$

at high
frequencies
(high k) the
non-local
term tends to
unity

non-local Starobinsky gravity is out of reach of any present or future GW interferometer: at high frequencies it reduces to standard Starobinsky inflation (a red-tilted spectrum)

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### **Pre-Big-Bang scenario**

A dynamical dilaton field that satisfies the cosmological equations following from the string effective action

Gasperini, Veneziano (1993)

Isocurvature spectrum of scalar perturbations: inconsistency with CMB data for scalar perturbations

Durrer, Gasperini, Sakellariadou, Veneziano (1998)

Way out via the axion/curvaton mechanism: a viable spectrum of scalar perturbations

r < 0.01 at the pivot scale k0=0.05/Mpc

Blue-tilted GWB that can reach the sensitivity of present and future interferometers

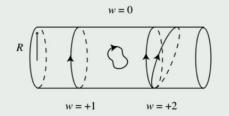
Gasperini (2016)

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### String gas cosmology

T-duality:

compact space



excitation modes of a thermal ensemble of strings

momentum modes  $E_n = n\frac{1}{R}$ 

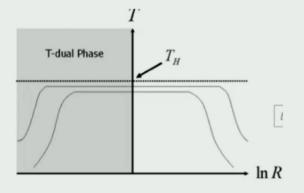
$$E_n = n\frac{1}{R}$$

winding modes  $E_m = mR$ 

$$E_m = mR$$

Brandenberger, Vafa (1989)

temperature cannot rise indefinitely and reaches a maximal: Hagedorn temp.



for an adiabatic process, winding modes dominate the thermal bath in a small space

Explanation of ST dimensionality

Sakellariadou (1996)

### String gas cosmology

inflation is replaced by a quasi-static era where thermal fluctuations generate an almost scale-invariant primordial scalar and tensor spectrum

$$\mathscr{P}_{\mathsf{t}}(k) \simeq \frac{1}{4(M_{\mathrm{Pl}}l_{\mathrm{st}})^4} \hat{T}(k) \left[ 1 - \hat{T}(k) \right] \ln^2 \left[ \frac{1 - \hat{T}(k)}{l_{\mathrm{st}}^2 k^2} \right]$$

during Hagedorn phase 
$$T pprox {
m const} \lesssim T_{
m H}$$
  $n_{
m t} \simeq (1-n_{
m s}) - rac{4}{\ln[(1-\hat{T})(l_{
m s} k)^{-2}]} > 0$  during RDE  $T \sim 1/a$ ,  $\alpha_{
m t} \simeq -lpha_{
m s}$ .

 $\hat{T}(k) := T(k)/T_{\rm H}$ 

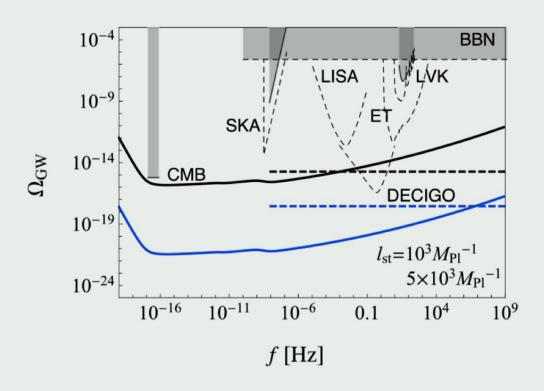
 $\alpha_{\rm f} \simeq -\alpha_{\rm s}$ .

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### **String gas cosmology**

In the most optimistic case of a tensor-to-scalar ratio saturating the CMB bound, the model reaches DECIGO sensitivity

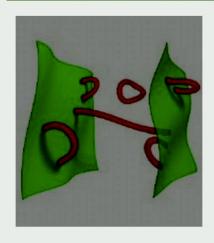
0.1Hz-10Hz (future Japanese GW space mission)



Calcagni, Kuroyanagi (2021)

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### New ekpyrotic scenario



2 flat 3dim branes (boundary of a 5dim ST)

They interact with attractive  $V(\varphi)$  along a compact  $5^{th}$  dim parametrised by the radion  $\varphi$ 

The branes get closed

gravitational energy in the bulk is converted to brane KE

the branes collide

part of the brane KE is converted into matter/radiation

An observer on one of the branes experiences the brane collision as a big bang after a period of contraction: **ekpyrosis** 

During the slow contraction phase: inhomogeneities

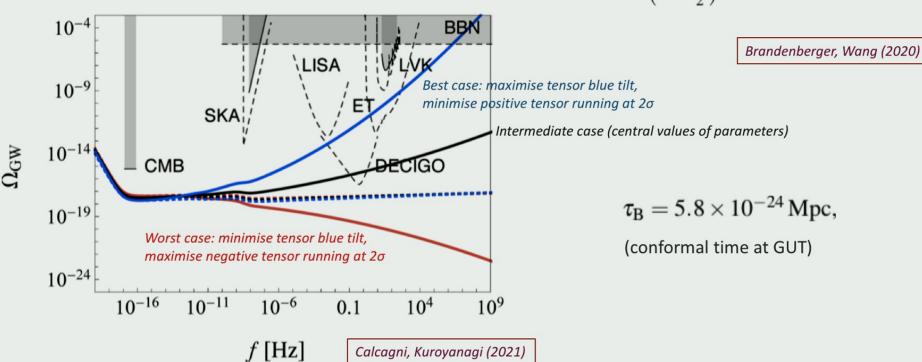
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### New ekpyrotic scenario

$$n_{\rm t} = 1 - n_{\rm s} > 0$$

$$\alpha_{\rm t} = -\alpha_{\rm s}$$

$$r(k_0) \simeq \frac{2^{n_{\rm s}} 25}{\Gamma^2 \left(1 - \frac{n_{\rm s}}{2}\right)} (k_0 \tau_{\rm B})^{2(1 - n_{\rm s})} (1 - n_{\rm s})^2$$



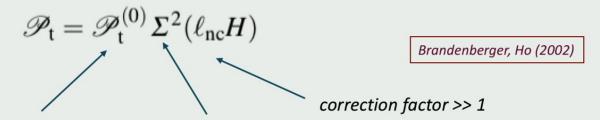
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### **Brandenberger-Ho noncommutative inflation**

Time and space coordinates do not commute



inflaton obeys modified dynamics



*amplitude* in the commutative limit

$$\mathcal{P}_{\mathrm{t}}^{(0)} = \mathcal{P}_{\mathrm{t}}(\Sigma = 1)$$

function encoding the *noncommutative* effects

In natural inflation, one can achieve a blue-tilted tensor spectrum while the scalar spectrum can stay red-tilted

$$n_{\rm t} \simeq \frac{r}{4} \simeq 4\varepsilon, \qquad \qquad \alpha_{\rm t} \simeq 8\varepsilon(\varepsilon - \eta)$$

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### **Multifractional spacetimes**

Different scaling laws in copies of the same object with different sizes

Calcagni (2017)

Dimensional flow: typical for ST arising in QG

't Hoft (1993); carlip (2017)

$$\mathscr{P}_{t}(k) = r(k_0)\,\mathscr{P}_{s}(k_0)\,\exp\left\{n_{t}(k_0)\ln\frac{p(k)}{p(k_0)} + \frac{\alpha_{t}(k_0)}{2}\left[\ln\frac{p(k)}{p(k_0)}\right]^2\right\}$$

$$p(k) \simeq k \left[ 1 + \frac{1}{|\alpha|} \left( \frac{k}{k_*} \right)^{1-\alpha} \right]^{-1}$$

related to the Hausdorff dimension of space

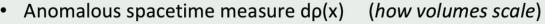
fundamental comoving scale

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Dimensional flow: feature common to all QG theories

change of ST dimension with probed scale

Quantisation of ST geometry



Kinetic operator  $\mathcal{K}$  (modified dispersion relations)



Perturbed action for a small perturbation  $h_{\mu\nu}$  over the background  $g_{\mu\nu}^{(0)}$  up to a source term

$$S = \frac{1}{2\ell_*^{2\Gamma}} \int \! \mathrm{d}\rho \, \sqrt{|g^{(0)}|} \left[ h_{\mu\nu} \mathcal{K} h^{\mu\nu} + O(h_{\mu\nu}^2) \right]$$

characteristic length scale around which QG effects become relevant (if there is only one fundamental scale, then it is near to the Planch scale)

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### Dimensional flow: feature common to all QG theories

change of ST dimension with probed scale

Geometric observable  $d_{\rm H}(\ell) := {\rm d} \ln \rho(\ell)/{\rm d} \ln \ell$ 

Quantisation of ST geometry

Anomalous spacetime measure dp(x) (how volumes scale)

Kinetic operator  $\mathcal{K}$  (modified dispersion relations)

Related to Hausdorff dim in Fourier space and the spectral dim (observable)

> In any plateau of dimensional flow, where all dim are approximately constant:

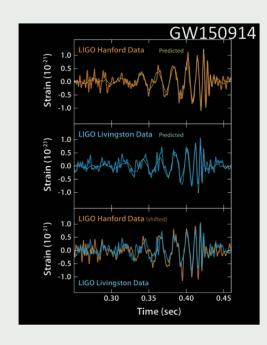
$$\mathscr{K}(-k^2) = -\ell_*^{2-2\beta} k^2 + k^{2\beta}$$

The value of Hausdorff dimensions, spectral dimension and of the parameter Γ for different QG theories is known

$$\beta = \frac{d_{
m H}^k}{d_{
m S}}$$
  $\Gamma \simeq \frac{d_{
m H}}{2} - \frac{d_{
m H}^k}{d_{
m S}} \approx {
m const}$ 

Calcagni, Kuroyanagi, Marsat, Sakellariadou, Tamanini, Tasinato (2019)

$$c_{\mathrm{GW}} := \frac{\mathrm{d}\omega}{\mathrm{d}|\mathbf{k}|}$$



Peaked at frequency: 
$$f = \omega/(2\pi) = 100\,\mathrm{Hz}$$

$$m < 1.2 \times 10^{-22} \,\mathrm{eV}.$$



$$\omega \approx 630 \,\mathrm{Hz} \approx 4.1 \times 10^{-13} \,\mathrm{eV}$$

At that frequency: 
$$\Delta c := c_{\text{GW}} - 1$$

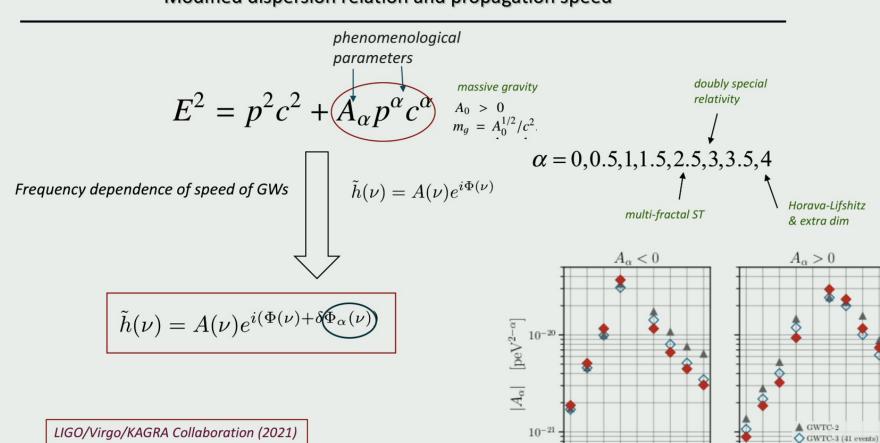
$$|\Delta c| < 4.2 \times 10^{-20}$$

$$m < 1.2 \times 10^{-22} \,\mathrm{eV}$$

LIGO/Virgo Collaboration (2016)

$$\omega^2 = |\mathbf{k}|^2 + m^2$$

$$arDelta c \simeq -m^2/(2\omega^2)$$
 if the mass is small



 $\alpha$ 

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In several QG proposals. the propagation speed of GWs may vary as a function of the energy scale

LVC: BNS GW170817 
$$\longrightarrow$$
  $-3 \times 10^{-15} \le c_T - 1 \le 7 \times 10^{-16} \text{ (in } c = 1 \text{ units)}$ 

Construct a function for  $c_T(f)$  which satisfies the LIGO-Virgo bounds whilst modifying the millihertz regime significantly

sharp transitions for.  $\, c_T(f)$  in the frequency band between LISA and LIGO frequencies

A general theoretical and numerical toolkit for quantifying the perspective of LISA to measure a frequency-dependent propagation speed of GWs only through its effects on GW waveforms from merging massive BH binaries

Baker, Sakellariadou, et al (2022)

Quadratic action for the linearised transverse-traceless GW modes

$$S_T = \frac{M_{\rm Pl}^2}{8} \int dt \, d^3x \, a^3(t) \, \bar{\alpha} \, \left[ \dot{h}_{ij}^2 - \frac{c_T^2(f)}{a^2(t)} (\vec{\nabla} h_{ij})^2 \right]$$

Sharp transitions for  $\ ^{C}T$  in the frequency band between LISA and LIGO frequencies, to ensure consistency with results from GW170817

Polynomial ansatz

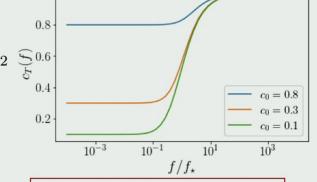
$$c_T(f) = 1 + \sum_n \beta_n \left(\frac{f}{f_*}\right)^n$$

1.0

LIGO bound implies: 
$$|\beta_n| \lesssim 10^{-15-n} (f_*/\mathrm{Hz})^n$$

**EFT-inspired ansatz** 

$$c_T(f) = \left[1 + \frac{f_{\star}^2}{f^2} - \frac{f_{\star}^2}{f^2} \sqrt{1 + 2\left(1 - c_0^2\right) \frac{f^2}{f_{\star}^2}}\right]^{1/2}$$



Baker, Sakellariadou, et al (2022)

### Luminosity distance

$$F=:rac{L}{4\pi(d_L^{
m EM})^2}$$
 power per unit area emitted by the source

flux of light reaching an observer

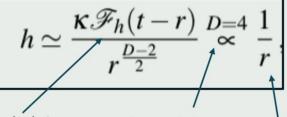
luminosity distance

GR: 
$$d_L^{\text{EM}}(z) = (1+z) \int_{t(z)}^{t_0} \frac{\mathrm{d}t}{a} = (1+z) \int_0^z \frac{\mathrm{d}z}{H_{\text{GR}}} \stackrel{\text{at small } z}{\simeq z/H_0},$$

$$a(z) = (1+z)^{-1}$$

$$H_{\text{GR}}^2 = (\kappa^2/3)\rho^2$$

Amplitude of GWs in the local wave zone of the source



rescaling with scale factor

function in retarded time topological dimensions

Cardoso, Dias, Lemos (2003)

Calcagni, Kuroyanagi, Marsat, Sakellariadou, Tamanini, Tasinato (2019)

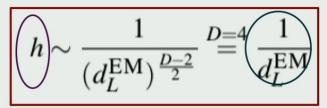
comoving distance from source

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### Luminosity distance

### **Standard sirens**

GR:



In theories beyond GR, this relation can be different

Standard sirens:

strain measured in an interferometer

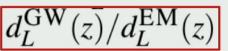
observations in the optical spectrum

arradia sirens.

**GW** luminosity distance

(up to a retarded timedependent function)

$$h \simeq rac{\kappa \mathscr{F}_h(t-r)}{r^{rac{D-2}{2}}} \overset{D=4}{\sim} rac{1}{r} \, .$$



deviates from 1

cosmological observable that can be determined from standard sirens

### Predictions of QG proposals: luminosity distance

Luminosity distance in QG in a model-independent way: dimensional flow

$$rac{d_L^{ ext{GW}}}{d_L^{ ext{EM}}} = 1 \pm |\gamma - 1| \left(rac{d_L^{ ext{EM}}}{\ell_*}
ight)^{\gamma - 1} \qquad \text{and } \qquad$$

Calcagni, Kuroyanagi, Marsat, Sakellariadou, Tamanini, Tasinato (2019)

- Only one fundamental scale:  $m{\gamma} 
  eq m{O}$  Equation for the ratio is exact and  $\ell_* = O(\ell_{ ext{Pl}})$  with  $m{\gamma} = m{\Gamma}_{m{U}m{V}}$
- If If If s a mesoscopic scale much larger than the Planck scale, the equation for the ratio is valid only near the IR:

Observations can place bounds on the two parameters  $\gamma$  and  $\gamma$  in a model-independent way, by constraining the ratio as a function of the redshift of the source

### Predictions of QG proposals: luminosity distance

Luminosity distance in QG in a model-independent way: dimensional flow

$$rac{d_L^{ ext{GW}}}{d_L^{ ext{EM}}} = 1 \pm |\gamma - 1| \left(rac{d_L^{ ext{EM}}}{\ell_*}
ight)^{\gamma - 1} \qquad \text{and } \qquad ext{Constant}$$

Calcagni, Kuroyanagi, Marsat, Sakellariadou, Tamanini, Tasinato (2019)

BNS merger GW170817 and a simulated z = 2 SMBH merging event that could be observed by LISA:

the only theories that can be tested are those for which  $\Gamma_{
m meso} > 1 > \Gamma_{
m UV}$ 

LQC / SF / GFT

Observations can place bounds on the two parameters  $\gamma$  and  $\gamma$  in a model-independent way, by constraining the ratio as a function of the redshift of the source

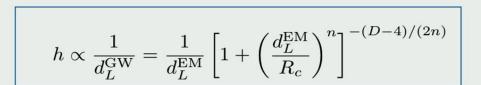
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### **Brane/String theory: Extra dimensions**

Constraints on the number of spacetime dimensions from GWs

Damping of the waveform due to gravitational leakage (beyond R<sub>c</sub>) into extra dim

Deviation depends on the number of dimensions D and would result to a systematic overestimation of the source  $~d_L^{\rm EM}$  inferred from GW data



 $h \propto \frac{1}{d_L^{\rm GW}} = \frac{1}{d_L^{\rm EM}} \left[ 1 + \left( \frac{d_L^{\rm EM}}{R_c} \right)^n \right]^{-(D-4)/(2n)}$ 

**Strain** measured in a GW interferometer

**Luminosity distance** measured for the optical counterpart of the standard siren

- Consistency with GR in D=4 dim
- Some models based on extra dimensions are ruled out

LIGO/Virgo/KAGRA Collaboration(2019)

GRB 170817A and GW170817

GW event 1.7 s before γ-ray observation

BNS merger at 40 Mpc

#### What I have not discussed

#### **Predictions of beyond GR theories:**

- generation: relates the outgoing radiation to the properties of the source, a hard problem even in GR
- propagation: birefringence (I've discussed only dispersion relation and amplitude damping)

GWs propagate from the source to the detector, the amplitudes of left versus right polarised modes are exponentially enhanced or suppressed, a **parity-violating effect** 

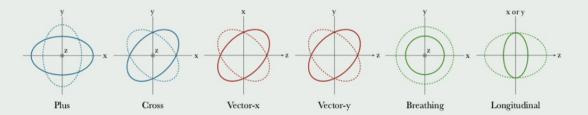
Martinovic, Badger, Sakellariadou, Mandic, PRD 104 (2021) 8, L081101

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### **Predictions of beyond GR theories:**

- generation: relates the outgoing radiation to the properties of the source, a hard problem even in GR
- propagation: birefringence (I've discussed only dispersion relation and amplitude damping)
- polarisation: a metric theory of gravity can allow up to six modes of polarisation: two tensor (helicity ± 2),
   two vector (helicity ± 1), and two scalar (helicity 0) modes



Callister, Sakellariadou et al, PRX 7 (2017) 041058

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- polarisation: a metric theory of gravity can allow up to six modes of polarisation: two tensor (helicity ± 2),
   two vector (helicity ± 1), and two scalar (helicity 0) modes
- remnant properties (exotic ultracompact objects without horizon, gravastars, fuzballs, etc)

compact object with an interior de Sitter condensate phase and an exterior Schwarzschild geometry

a ball-shaped mess of interacting fundamental strings

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#### **Conclusions**

Quite unlikely to have, in the near future, laboratory experiments that test quantum gravity:

$$d_L^{
m EM} \simeq rac{z(1+z)}{H_0} \stackrel{z \ll 1}{\simeq} rac{z}{H_0}$$

Cosmology provides a test of quantum theory of gravity (a quantum theory of spacetime geometry

Difficult, but still possible, to test QG with GW observations

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