

Title: Lecture - Strong Gravity, PHYS 777

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Subject: Strong Gravity

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Linear perturbations of the metric

$$g_{ab} = \hat{g}_{ab} + h_{ab} \quad |h_{ab}| \ll |\hat{g}_{ab}|, \quad \text{Assume } \hat{R}_{ab} = 0$$

$$\delta \Gamma^a_{bc} = \frac{1}{2} \hat{g}^{ad} (\hat{\nabla}_c h_{bd} + \hat{\nabla}_b h_{dc} - \hat{\nabla}_d h_{bc})$$

$$\delta R^a_{bcd} = \hat{\nabla}_c \delta \Gamma^a_{bd} - \hat{\nabla}_d \delta \Gamma^a_{bc}$$

$$\Rightarrow \delta R^a_{bcd} = \frac{1}{2} [\hat{\nabla}_c \hat{\nabla}_b h^a_d + \hat{\nabla}_c \hat{\nabla}_d h^a_b - \hat{\nabla}_c \hat{\nabla}^a h_{bd} - \hat{\nabla}_d \hat{\nabla}_b h^a_c - \hat{\nabla}_d \hat{\nabla}_c h^a_b + \hat{\nabla}_d \hat{\nabla}^a h_{bc}]$$

$$\rightarrow \partial^a \partial^b h_{cd} = \partial^a \partial^b h_{cd} + \partial^a \partial^b h_{cb} + \partial^a \partial^b h_{bc} - \partial^a \partial^b h_{bc} - \partial^a \partial^b h_{cb} + \partial^a \partial^b h_{cd}$$

$$\partial_\mu T_{bd} = \delta G_{bd} = -\frac{1}{2} \square \bar{h}_{bd} + \hat{R}_{adbc} \bar{h}^{ac} - \frac{1}{2} \hat{g}^{cd} \hat{\nabla}_a \hat{\nabla}_c \bar{h}^{ab} + \frac{1}{2} \hat{\nabla}_b \hat{\nabla}_a \bar{h}^a{}_d + \frac{1}{2} \hat{\nabla}_d \hat{\nabla}_a \bar{h}^a{}_b$$

$$\hat{\nabla}_a \bar{h}^{ab} = 0 \quad (\text{Lorenz Gauge})$$

$$\square \bar{h}_{bd} - 2 \hat{R}_{adbc} \bar{h}^{ac} = -16\pi T_{bd} \quad (*)$$

$$x^a \rightarrow x'^a = x^a + \xi^a$$

$$h_{ab} \rightarrow h'_{ab} = h_{ab} - 2 \nabla_{(a} \xi_{b)}$$

$$\hat{\nabla}^a \bar{h}'_{ab} = \hat{\nabla}^a \bar{h}_{ab} - \square \xi_b$$

Original $\xi^a(t, x')$
 In Lorenz gauge $\xi^a(t=0, x'), \partial_t \xi^a(t=0)$

$$\bar{h}^{ac} + \frac{1}{2} \nabla_b \nabla_a \bar{h}^a_d + \frac{1}{2} \nabla_a \nabla_b \bar{h}^a_b$$

$$\bar{h}_{ab} = h_{ab} - \frac{1}{2} h g_{ab}$$

$$\bar{h} = -h$$

e)

$$\hat{\square} \bar{h} = -16\pi T$$

$\bar{h}^{\mu\nu}(t, x')$

use $\int \bar{h}^{\mu\nu}(t=0, x'), \partial_+ \bar{h}^{\mu\nu}(t=0, x')$

Particular cases

(i) Vacuum ($T_{ab} = 0$)

Choose 2 free functions from, $\xi_a, \partial_t \xi_a$ at $t=0$
to have $\bar{h} = 0$, $h_{as} = \bar{h}_{as}$

(1a) Vacuum, flat background $\hat{g}_{ab} = \eta_{ab}$
 $h_{tt} = \partial_t h_{tt} = 0$ at $t=0$

(1b) $\hat{g}_{ab} = \hat{g}_{ab}^{\text{Kerr}}$, $\hat{g}_{ab} = \eta_{ab}$

(ii) Non-vacuum ($T_{ab} \neq 0$)

(11) Non vacuum (1.5)

Propagation of GWs

Use 4 gauge DDF to set $h_{+i} = 0$, $\partial_{\mu} h^{\mu\nu} = 0$

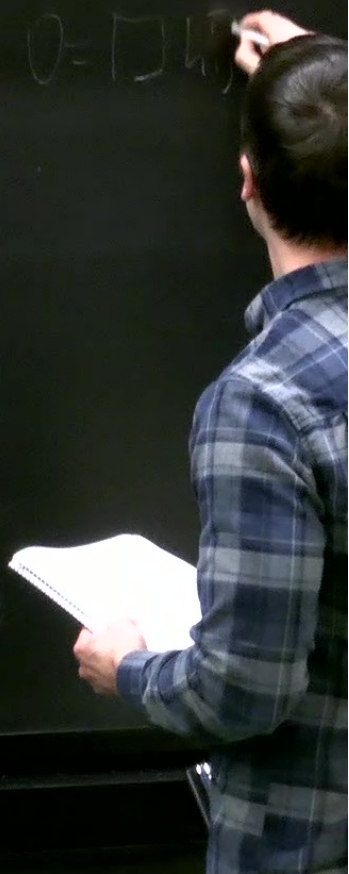
$$0 = \partial_s h_{TT}^{s0} = \partial_0 h_{TT}^{s0}$$

Choose $h^{00} = 0 \Rightarrow h^{00} = 0$ for all time

$$\square h_{ij}^{TT} = 0$$

$$h_{ij}^{TT} = A_{ij} \exp(i k_{\mu} x^{\mu})$$

↑
constant



(ab + 0)

$= 0, \quad \partial_t^{ab} h_{ab} = 0$

all time

$$\begin{aligned}
 0 &= \square h_{ij} \\
 &= \partial_j^{ab} (k_a k_b) h_{ij} \\
 &= -k_a k^a h_{ij}
 \end{aligned}$$

$k_a k^a = 0 \Leftrightarrow k^a$ is null

$0 = \partial_a h_{ij}^{ab} = k_a h_{ij}^{ab} = 0$

Choose $k^a = (\omega, 0, 0, \omega)$ (in z-direction)

$$A_{ij} = \begin{pmatrix} h_+ & h_x & 0 \\ h_x & -h_+ & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

(in z-direction)

In TT gauge

$$U^a = (U^t, \vec{0})$$

$$U^a \nabla_a U^b = U^t \partial_t U^b + U^a \Gamma_{ac}^b U^c$$

$$+ U^t \Gamma_{++}^b U^t = 0$$

geodesics in TT gauge

(A) $x_A^c = (t, 0, 0, 0)$

(B) $x_B^a = (t, L, 0, 0)$

gauge

$$U^a = (U^+, \vec{0})$$

$$U^b = U^+ \Gamma_{++}^b U^+ = U^+ \Gamma_{++}^b U^+$$

$$+ U^+ \Gamma_{++}^b U^+ = 0$$

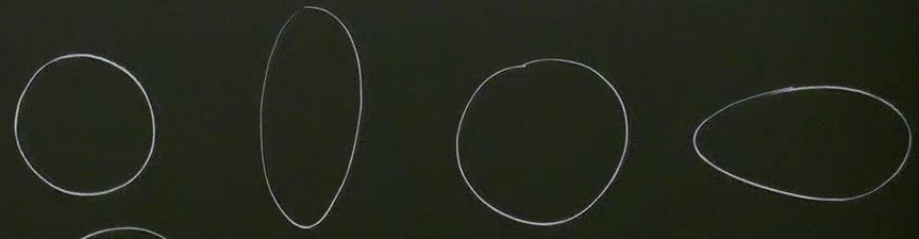
(A) $x_A^a = (t, 0, 0, 0)$

(B) $x_B^a = (t, L, 0, 0)$

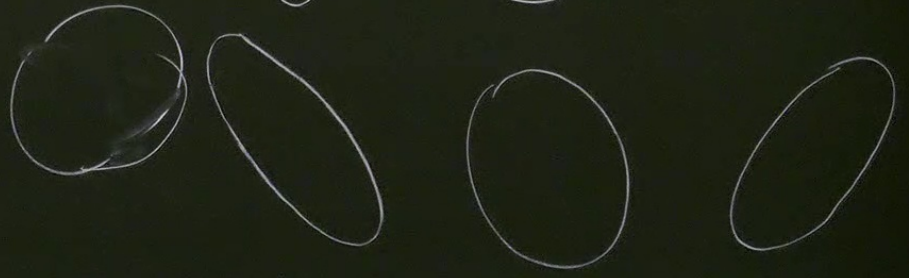
For h_+ wave

$$S_{AB}^2 = [L(1 + h_+ \cos(\omega t))]^2$$

h_+



h_x



constant

Rotate by ϕ about z-axis

$$h_x \pm i h_y \rightarrow (h_x \pm i h_y) e^{\pm 2i\phi}$$

\uparrow spin-2 \leftarrow

$$\square \bar{T}_{ab} = -16\pi T_{ab}$$

Green function

$$\square G(x^a - \bar{x}^a) = \delta^4(x^a - \bar{x}^a)$$

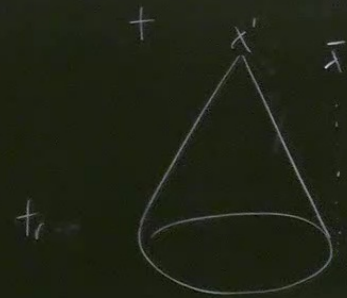
$$G(x^a - \bar{x}^a) = \frac{-1}{4\pi|x^a - \bar{x}^a|} \delta[|x^a - \bar{x}^a| - (x^0 - \bar{x}^0)] \Theta(x^0 - \bar{x}^0)$$

$$\Theta(x - \bar{x}^0) = \begin{cases} 1 & \text{when } x^0 - \bar{x}^0 > 0 \\ 0 & \text{otherwise} \end{cases}$$

$$\bar{h}_{ab}(x^a) = -16\pi \int G(x^a - \bar{x}^a) T_{ab}(\bar{x}^a) d^4x$$

$$= 4 \int \frac{T_{ab}(t - |x' - \bar{x}'|, \bar{x}') d^3\bar{x}}{|x' - \bar{x}'|}$$

$$t_r = t - |x' - \bar{x}'|$$



Transform to Fourier Space

$$\begin{aligned}\tilde{h}_{ab} &= \int dt e^{i\omega t} \bar{h}_{ab} \\ &= 4 \int dt_r \int d^3\bar{x} \frac{T_{ab}(t_r, \bar{x}')}{|\bar{x}' - \bar{x}|} e^{i\omega(t_r + |\bar{x}' - \bar{x}|)} \\ &= 4 \int d^3\bar{x} \frac{\tilde{T}_{ab}}{|\bar{x}' - \bar{x}|} e^{i\omega|\bar{x}' - \bar{x}|}\end{aligned}$$

$$\frac{\partial^2 h}{\partial t^2} = \frac{1}{\omega} \partial_t \dot{h}$$

Assume

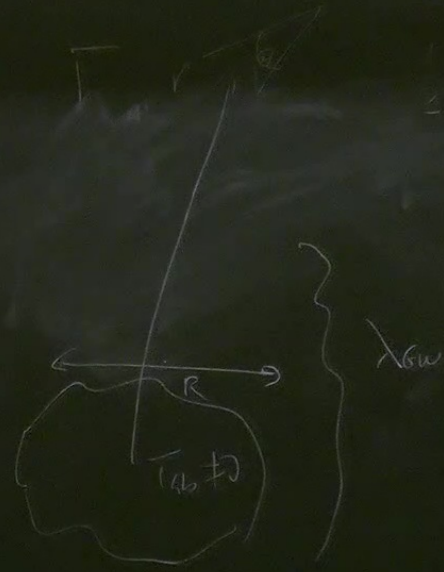
(Lorentz)

$$r \gg \frac{1}{\omega}$$

(wave zone)

$$R_{\text{source}} \ll \lambda_{\text{GW}}$$

$$\lambda_{\text{GW}} = \frac{1}{\omega}$$



$$\frac{1}{r} e^{i\omega r}$$

Transform to Fourier Space

$$\bar{h}^{00} = \frac{1}{\omega} \partial_t \bar{h}$$

$$\bar{h}_{ab} = \int dt e^{i\omega t} \bar{T}_{ab}$$

$$= 4 \int dt_r \int d^3\bar{x} \frac{T_{ab}(t_r, \bar{x}')}{|\bar{x}' - \bar{x}|} e^{i\omega(t_r + |\bar{x}' - \bar{x}|)}$$

$$= 4 \int d^3\bar{x} \frac{\bar{T}_{ab}}{|\bar{x}' - \bar{x}|} e^{i\omega|\bar{x}' - \bar{x}|} \approx \frac{4}{r} e^{i\omega r} \int d^3\bar{x} \bar{T}_{ab}$$

Assume

$$\frac{e^{i\omega|\bar{x}' - \bar{x}|}}{|\bar{x}' - \bar{x}|}$$

constant

$$\int d^3\bar{x} \tilde{T}_{ij} = \int \left[\underbrace{\partial_k (\bar{x}^i \tilde{T}^{kj})}_0 - \partial_k (\tilde{T}^{kj}) \bar{x}^i \right] d^3\bar{x}$$

$$\nabla_a T^{ab} = 0$$

$$i\omega \tilde{T}^{0b} + \partial_i \tilde{T}^{ib} = 0$$

$$= \frac{i\omega}{2} \int (\tilde{T}^{0i} \bar{x}^i + \tilde{T}^{i0} \bar{x}^i) d^3\bar{x}$$

$$= \frac{i\omega}{2} \int [\partial_e (\bar{x}^i \bar{x}^j \tilde{T}^{0e}) - \bar{x}^i \bar{x}^j (\partial_e \tilde{T}^{0e})] d^3\bar{x}$$

$$= \frac{-\omega^2}{2} \int \underbrace{\bar{x}^i \bar{x}^j \tilde{T}^{00}}_{\tilde{T}^{ij}} d^3\bar{x}$$

$$\begin{aligned} \bar{h}_{ij} &= -2\omega^2 \frac{e^{i\omega r}}{r} \bar{I}_{ij} \\ \bar{h}_{ij} &= \frac{2}{r} \ddot{I}_{ij} (+r) \end{aligned}$$