

Title: Lecture - Quantum Gravity, PHYS 644

Speakers: Aldo Riello

Collection/Series: Quantum Gravity (Elective), PHYS 644, February 24 - March 28, 2025

Subject: Quantum Gravity

Date: March 20, 2025 - 9:00 AM

URL: <https://pirsa.org/25030031>

From COVARIANT to CANONICAL ph. sp.

(F, Ω_Σ)
off shell, presympl.

↓ on-shell

(F_{EL}, ω)
on shell, presympl. if gauge

↓ mod out gauge

" $(F_{EL}, \omega) / \text{Gauge}$ "
on shell, symplectic

presympl
red
→

(P_{can}, W_{can})
off shell, sympl.

↓ on shell

$(e, \iota_e^* W_{can})$
on shell, presympl.

↓ mod out gauge

" $(e, \iota_e^* W_{can}) / \text{Gauge}$ "
on shell, symplectic

usually
cotangent
bundle

impose
"constraints"

E & M

$F \sim$

$F(A)$

$G \sim$

$\rho(\xi)$

$M = \mathbb{R}$

to CANONICAL ph. sp

presympl
red \rightarrow

$$(P_{can}, W_{can})$$

off shell, sympl.

\downarrow on shell

$$(e, \int_C^* W_{can})$$

on shell, presympl.

\downarrow mod out gauge

$$\cong \left(e, \int_C^* W_{can} \right) / \text{Gauge}$$

on shell, symplectic

usually
cotangent
bundle

impose
"constraints"

gauge

gauge

E & M

$$\mathcal{F} \sim \Omega^1(M) \ni A$$

$$F(A) = dA \quad \text{field strength (gauge inv.)}$$

$$\mathcal{G} \sim C^\infty(M) \text{ Abelian}$$

$$\rho(\xi) A = d\xi$$

$$M = \mathbb{R}^{1,3}, \quad \Sigma = \{t=0\}$$

$$\begin{aligned}\Omega_\Sigma &= \int \mathbb{d}F^{0i}(A) \wedge \mathbb{d}A_i \\ &= \int \mathbb{d}(\dot{A}_i - \partial_i A_0) \wedge \mathbb{d}A_i\end{aligned}$$

$$X \in \ker(\Omega_\Sigma) \quad X = \int X_i \frac{\delta}{\delta A_i} + X_0 \frac{\delta}{\delta A_0}$$

$$0 \stackrel{!}{=} \int_\Sigma (\underbrace{\dot{X}^i - \partial_i X_0}_{=0}) \mathbb{d}A_i - \underbrace{X_i}_{=0} \mathbb{d}\dot{A}_i - \underbrace{\nabla^i X_i}_{=0} \mathbb{d}A_0$$

S^2)

$$\rightarrow X_i|_\Sigma = 0$$

X_0 is unconstrained

as long as $\dot{X}^i|_\Sigma = \partial_i X_0$

$$\leftrightarrow [A_\mu] \sim (A_i, E^i) = (A_i|_\Sigma, \dot{A}_i - \partial_i A_0|_\Sigma)$$

$F^{0i}|_\Sigma$
" "

$$\rightarrow \Omega_{\text{can}} = \int_{\Sigma} dE^i \wedge dA_i$$

$$\mathcal{P}_{\text{can}} \sim T^* \Omega(\Sigma) \ni (E^i, A_i)$$

$$\text{GAUGE: } \delta_{\xi} A_{\mu} = \partial_{\mu} \xi + [A_{\mu}, \xi]$$

$$[\delta_{\xi} A_{\mu}] \stackrel{?}{=} \delta_{\xi} [A_{\mu}]$$

$$\text{it works f.c. } \begin{cases} \delta_{\xi} A_i = \partial_i \xi + [A_i, \xi] \\ \delta_{\xi} E^i = 0 + [E^i, \xi] \end{cases}$$

\Rightarrow GAUGE tr. descent to $(\mathcal{P}_{\text{can}}, \omega_{\text{can}})$
 $\rho_{\text{can}}(\xi)(A, E) = (d\xi, 0)$

EOM: are there residual eom that constrain initial data?

$$n_{\nu} \nabla_{\mu} F^{\mu\nu} = 0 \sim \nabla_i F^{i0} = 0$$

↓ on Σ

$$\text{Gauss constr. } \boxed{\nabla_i E^i = 0}$$

indep. of A_0, A_i, \dot{A}_i V
 except through E^i

GENERAL RE

More conceptually:

on \mathcal{F} : $\int_{\rho(\xi)} \Omega_{\xi} = - \int_{\Sigma} Q_{\xi}(\xi) \quad (\text{YM})$
 \uparrow
 $Q_{\xi}(\xi) \approx 0 \quad (\partial \Sigma = \emptyset)$

on \mathcal{P}_{can} : $\int_{\rho(\xi)} \omega_{\text{can}} = - \int_{\Sigma} E^i \partial_i \xi \stackrel{\text{ibp.}}{=} + \int_{\Sigma} (\nabla_i E^i) \xi \approx 0$
 \uparrow Gauss constraint

$(\mathcal{P}_{\text{can}}, \mathcal{P}_{\text{can}}, Q \equiv H)$
 give us the setup for
 sympl. reduction
 $\mathcal{P}_{\text{can}} \rightarrow \mathcal{P}_{\text{reduced}} = e/gay$

\Rightarrow Constraints are the Hamilton generators of gauge transf. on $(\mathcal{P}_{\text{can}}, \omega_{\text{can}})$

$$\{H(\xi), H(\eta)\} = H([\xi, \eta])$$

\uparrow
 $\equiv Q_{\xi}(\xi)$

constraint algebra represents the gauge algebra
 (off shell.)

& closes on-shell. $H \equiv Q_{\xi} \approx 0$

GENERAL RELATIVITY

$$\mathcal{F} = \text{Riem}_{1,3}(M) \ni g_{ab}$$

$$\mathcal{g} = \text{Diff}(M) \simeq (\mathcal{X}(M), [\cdot, \cdot])$$

$$\mathcal{L} = \frac{1}{2} \sqrt{g} (\text{Ricci}(g) - 2\Lambda)$$

$$\tilde{E}_{ab} = \#(G_{ab} + \Lambda g_{ab}) \approx 0$$

$$\frac{1}{2} g^{-1}(\partial g \dots)$$



$$\tilde{\Theta}^a = -\frac{1}{2} (\nabla_b dg^{ba} - \nabla^a dg) = -\frac{1}{2} (g^{bc} d\Gamma_{bc}^a - g^{eb} d\Gamma_{bc}^c)$$

GENERAL RELATIVITY

$$\mathcal{F} = \text{Riem}_{1,3}(M) \ni g_{ab}$$

$$\mathcal{g} = \text{Diff}(M) \simeq (\mathcal{X}(M), [\cdot, \cdot])$$

$$\mathcal{L} = \frac{1}{2} \sqrt{g} (\text{Ricci}(g) - 2\Lambda)$$

$$\tilde{E}_{ab} = \#(G_{ab} + \Lambda g_{ab}) \approx 0$$

$$\frac{1}{2} g^{-1}(\partial g \dots)$$



$$\tilde{\Theta}^a = -\frac{1}{2} (\nabla_b dg^{ba} - \nabla^a dg) = -\frac{1}{2} (g^{bc} d\Gamma_{bc}^a - g^{ab} d\Gamma_{bc}^c)$$

$$\ominus_{\Sigma} = -\frac{1}{2} \int_{\Sigma} \sqrt{h} \left(n_a \underbrace{d\Gamma_{bc}^a}_{\text{York}} g^{bc} - \underline{\underline{n^a d\Gamma_{ac}^c}} \right)$$

$$d(\nabla_e n^e) = \nabla_e dn^e + \underline{\underline{d\Gamma_{ab}^e}} n^b$$

$$g^{ab} d(\nabla_e n_b) = \nabla^e dn_a - g^{ab} \underline{\underline{d\Gamma_{ab}^c}} n_c$$

$$\beta^e = dn^e + g^{ab} dn_b$$

$$n_a \beta^e = n_a dn^e + dn_a n^e = d(n^2) \equiv 0$$

$$\ominus_{\Sigma} = -\frac{1}{2} \int_{\Sigma} \sqrt{h} \left(\nabla_e n_b \underbrace{dg^{ab}}_{\text{York}} + 2 \underbrace{d(\nabla n)}_{\text{York}} - \nabla_e \beta^e \right)$$

drops out in Ω_{Σ} (York)

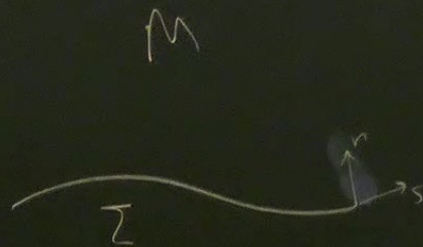
corner term $\rightarrow 0$

$$\nabla_{(e} n_b) = \frac{1}{2} L_n g_{ab}$$

$$\Theta_{\Sigma} \sim \int_{\Sigma} \underbrace{\dot{g} \delta g}_{\downarrow \Omega} + \underbrace{d(\quad) + d(\quad)}_{\substack{\downarrow \\ \text{drop out} \\ (\partial \Sigma = \emptyset)}}$$

$$\rightarrow = d(n^2) \equiv 0$$

$\underbrace{\partial \beta^a}_{\text{corner term}} \rightarrow 0$



$$\underline{d} \Gamma_{cc}^c$$

$$\nabla_e n_b = \frac{1}{2} L_n g_{ab}$$

$$\Theta_{cc} \sim \int_{\Sigma} \underbrace{g \delta g}_{\Omega} + \underbrace{d(\dots) + d(\dots)}_{\text{drop out } (\partial \Sigma = \emptyset)}$$

$$n_e dn^e + dn_e n^e = d(n^2) \stackrel{-1}{=} 0$$

$$\underbrace{d(\nabla n)}_{\text{York}} - \underbrace{\nabla_e \beta^e}_{\text{corner term} \rightarrow 0}$$

$$\int_{\partial \Sigma} \delta_e \beta^e \sim s_e dn^e + s^e dn_e$$