

Title: Lecture - Quantum Gravity, PHYS 644

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Subject: Quantum Gravity

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Recap

Def Lagrangian sym
 $f: g \rightarrow \mathcal{X}(F)$

such that $\mathbb{L}_{p(\xi)} \underline{L} = d\underline{R}(\xi)$
 \uparrow
boundary!

Thm • $\mathbb{L}_{p(\xi)} E_I \approx 0 \Leftrightarrow p(\xi)$ descends to (F_{EC}, ω)

(N1) $\rightarrow \nabla_a \tilde{j}^a \approx 0$ for $\tilde{j}(\xi) = \overset{\circ}{i}_{p(\xi)} \omega - \underline{R}(\xi)$
Noether current

• $\partial \Sigma = \phi$ on F_{EC} $\overset{\circ}{i}_{p(\xi)} \omega \approx -d\underline{Q}_\Sigma(\xi)$ $Q_\Sigma = \int_\Sigma \tilde{j}$

Ex $g =$ Poincaré sym of \mathbb{R}^4

LOCAL (GAUGE)

$G = (\Gamma(E \rightarrow M))$

Ex \mathfrak{g} = Poincaré symms of relativity

LOCAL (GAUGE) SYMMETRIES

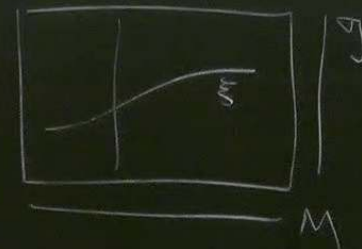
$$\mathcal{G} = \left(\Gamma(\underline{E} \rightarrow M)_{\text{vector bundle}}, [\cdot, \cdot] \right)$$

infinite dimensional Lie algebra

Ex: ① YM $\underline{E} \sim M \times \mathfrak{g}$

$$\xi \in \mathcal{G} \sim C^\infty(M, \mathfrak{g})$$

$$[\xi, \eta]_{\mathcal{G}}(x) = [\xi(x), \eta(x)]_{\mathfrak{g}}$$



(ξ)

underlying!

descends to (F_{EL}, ω)

$$L(\xi) = \int_M \rho(\xi) \omega - R(\xi)$$

Noether current

$$Q_\xi = -dQ_\xi(\xi) \quad Q_\xi = \int_M \underline{J}_\xi$$

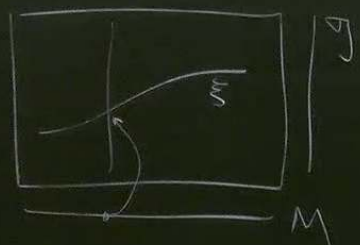
②

of scalar th

SYMMETRIES

$(M), [,]$

and Lie algebra



$$E \sim M \times \mathfrak{g}$$

$$G \sim C^\infty(M, \mathfrak{g})$$

$$[\xi, \eta]_G(x) = [\xi(x), \eta(x)]_{\mathfrak{g}}$$

② $G \approx \text{diff}(M)$
 $E = TM$
 $\xi \in G = \Gamma(TM) = \mathcal{X}'(M)$
 $[\xi, \eta]_G = L_\xi \eta = -L_\eta \xi$

Def (local action)

$$\rho: G \rightarrow \mathcal{X}'(F)$$

$$\xi \mapsto \rho(\xi) = \int_M \left(\frac{\delta \varphi^I}{\delta \varphi^I} \right) \frac{\delta}{\delta \varphi^I(x)}$$

depend only on ξ, φ and a finite no of deriv. at x .

$$\partial \Sigma = \phi$$

on Γ_{EC}

$$\int_{p(\xi)} \omega \approx -dQ_{\Sigma}(\xi)$$

$$Q_{\Sigma} = \int_{\Sigma} \underline{J}$$

$$[\xi, \eta]_{\Gamma}(\infty) =$$

Typically:

$$\delta_{\xi} \varphi^I = A_{\alpha}^I \xi^{\alpha} + B_{\alpha}^I \nabla_{\alpha} \xi^{\alpha} \equiv D_{\alpha}^I \xi^{\alpha}$$

\uparrow $(\varphi, \partial \varphi, \dots)$

Take a $\xi(\infty)$ of compact support on M .



$$\int_M (\delta_{\xi} \varphi^I) E_I \equiv \int_M (D_{\alpha}^I \xi^{\alpha}) E_I \stackrel{\text{lib.p.}}{=} \int_M \xi^{\alpha} (\overset{+}{D}_{\alpha}^I E_I)$$

$$\overset{+}{D} = (A - \mathcal{R}_{\alpha} B^{\alpha}) - B^{\alpha} \nabla_{\alpha}$$

Γ $(\psi(x), \eta(x))$

independent only
on ξ, φ and a finite nb of deriv. at x .

Thm (Noether 2)
If Γ is a local Lagrangian sym
then $\dagger D_\alpha^I E_I \equiv 0$.

Pf: ξ of cpt support

$$0 \equiv \int_M d\underline{J}(\xi) \stackrel{\text{Noether 1}}{=} \int_M (\delta_\xi \varphi^I) E_I = \int_M \xi^\alpha (\dagger D_\alpha^I E_I) \quad \forall \xi \Rightarrow \dagger D_\alpha^I E_I \equiv 0 \quad \square$$

↑ arbitrary local function

Physical interpretation: in the presence of local sym there are
LESS com than Dof. (by $\# \alpha$)



$$D = (A - \nabla_0 B^a) - B^a \nabla_a$$

Physical

→ lose deterministic evolution

UNLESS

we somehow "declare" physically indistinguishable
 $\#$ configurations (related by action of G !)
 ↳ REDUNDANCY of GAUGE

But I can choose ξ to vanish at Σ in
 with $\xi|_{\Sigma_f}$ arbitrary!
 → $Q_{\Sigma_f}(\xi) \approx 0$ □

Proposition $\partial\Sigma = \emptyset \Rightarrow Q_\Sigma(\xi) \approx 0 \quad \forall \xi|_{\partial\Sigma}$

Pf $0 \approx \int_M dJ(\xi) = Q_{\Sigma_f}(\xi) - Q_{\Sigma_i}(\xi)$
 (no cpt support, but onshell + NL)



Physical interpretation. in the presence of local sym there are
 LESS eom than Dof. (by #d)

Corollary

From arbitrariness of Σ & ξ
 $Q_\Sigma(\xi) \approx 0 \Rightarrow \underline{J}(\xi) \approx d\underline{j}(\xi)$

And in fact:

$$\underline{J}(\xi) = \xi^d \tilde{C}_d^a + \nabla_a \tilde{j}^{ab}(\xi)$$

s.t. on shell: $\tilde{C}_d^a \approx 0$ and $\tilde{j}^{ab} \approx j^{[ab]}$

$$\underset{\partial\Sigma = \phi}{\sim} Q_\Sigma(\xi) = \int_\Sigma \boxed{n_e \tilde{C}_d^e} \xi^d \approx 0$$

Constraints! Ex: $n_e \tilde{C}_d^e \sim$ Gauss $\nabla_i F^{io}$

Thm local sym, on $(\mathbb{F}_{ELI} \omega)$, $\partial \Sigma = \emptyset$

$$\mathbb{I}_p(\xi) \omega \underset{\substack{\uparrow \\ \text{"old thm"}}}{\approx} -dQ_\xi(\xi) \approx 0$$

\Rightarrow the "Cov. Ph. Sp." $(\mathbb{F}_{ELI} \omega)$
has a degenerate sympl structure
in the "gauge directions"

$$\ker(\omega^\flat) \sim \text{Im}(p)$$

$m, \omega (\mathbb{F}_{ELI} \omega), \partial \mathcal{E} = \phi$

$$- dQ_{\mathcal{E}}(\xi) \approx 0$$

"hm"

ov. Ph. S_p " ($\mathbb{F}_{ELI} \omega$)

degenerate sympl structure

"gauge directions"

$$\ker(\omega^b) \sim \text{Im}(\rho)$$

identifying out gauge fixes

two problems:

- 1) makes $\mathbb{F}_{ELI} \omega$ symplectic
- 2) makes eom deterministic!

Thm local sym, on $(\mathbb{F}_{EL} \omega)$, $\partial \mathcal{E} = \phi$

$$i_{p(z)} \omega \approx -dQ_E(\xi) \approx 0$$

↑
"old thm"

⇒ the "Cov. Ph. Sp" $(\mathbb{F}_{EL} \omega)$
has a degenerate sympl structure
in the "gauge directions"

$$\ker(\omega^b) \sim \text{Im}(p)$$

⇒ Quotienting out gauge fixes

two problems:

- 1) makes \mathbb{F}_{EL}/G symplectic
- 2) makes eom deterministic!

Example: parametrized particle
"time" reparametrization is a local Lga. sym.

$$\underline{p}(\xi) = 0 \text{ off shell } \dot{\xi}$$

General Relativity

$$S = \int_M \sqrt{g} \left(\frac{1}{2} R - \Lambda \right) d^4x = \int_M \mathcal{L}$$

$\mathcal{F} = \{ g_{ab} \text{ metric on } M \}$

Notation:

$$dg^{ab} := g^{aa'} g^{bb'} dg_{a'b'} \left[= -d(g^{ab}) \right]$$

$$dg = g^{ab} dg_{ab} = 2 \frac{d\sqrt{g}}{\sqrt{g}}$$

$$\mathcal{G} = \text{diff}(M) = \mathcal{X}^1(M)$$

$$\mathcal{L}_{p(\xi)} g_{ab} \equiv \delta_{\xi} g_{ab} = \mathcal{L}_{\xi} g_{ab} = 2\nabla_{(a} \xi_{b)}$$

change in field conf. "Active diff" ~ Lie transport

\mathcal{L}

$$\delta g_{ab} = -d(g^{ab})$$

$$= 2 \frac{d\sqrt{g}}{\sqrt{g}}$$

$$\delta g_{ab} = 2 \nabla_{(a} \xi_{b)}$$

Active diffeos ~ Lie transport

$$[p(\xi), p(\eta)] = -p([\xi, \eta]) \quad \uparrow \triangle$$

PROP $[p(\xi)] \mathcal{L} = L_{\xi} \mathcal{L}$

BACKGROUND INDEPENDENCE

Counter ex: scalar field on (M, g_{ab})

$$\mathcal{F} = C^{\infty}(M) \ni \phi$$

$$\mathcal{L} = \frac{1}{2} g^{ab} \nabla_a \phi \nabla_b \phi \sqrt{g} dx^d$$

$$L_{\phi} \mathcal{F} \equiv 0$$

$$L_{\xi} \mathcal{F} = L_{\xi} \phi = \xi^a \nabla_a \phi$$

$$\left[L_{p(\xi)} g_{ab} \equiv d_{\xi} g_{ab} \right] = \left[L_{\xi} g_{ab} - \mathcal{L}_{\xi} g_{ab} \right]$$

change in field conf. "active diffeos" ~ Lie transport

→ Cor $L_{p(\xi)} \underline{L} = L_{\xi} \underline{L} = d \left(\frac{1}{\xi} \underline{L} \right)$
 $\underbrace{\hspace{10em}}_{R(\xi)}$

↳ diffeos on Lagr sym

$$\tilde{\Theta}^e \sim \nabla d g$$

$$\tilde{R}^e = \xi^e L$$

$$\Rightarrow \tilde{J}^e(\xi) = i_{p(\xi)} \tilde{\Theta}^e - \tilde{R}^e(\xi)$$

$$= \tilde{C}_b^e \xi^b + \nabla_b \tilde{J}^{ob}(\xi)$$

↑ $G^e_b + \Lambda \delta^e_b \approx 0$

KOMAR 2-CURRENT

$$\tilde{J}^{ob} = -\frac{1}{2} \left(\nabla^a \xi^b - \nabla^b \xi^a \right)$$

$$\text{or } \tilde{J} = -\frac{1}{2} * d \xi^b$$

$$B^a) - B^a \nabla_a$$

Physical interpretation: in the presence of local sym there are LESS com than Dof. (by #d)

Thm

$$i_{p(\xi)} \underline{\Omega} = - \underline{d} \underline{J}(\xi) + i_{\xi} \underline{E}^{ab} dg_{ab} + d i_{\xi} \underline{\omega}$$

Pf. only need begrund indep!

$$\begin{aligned} \text{bind} \rightarrow \parallel \\ L_{p(\xi)} \underline{\omega} &= i_{p(\xi)} \underline{\Omega} + \underline{d} i_{p(\xi)} \underline{\omega} \\ L_{\xi} \underline{\omega} &= d i_{\xi} \underline{\omega} + i_{\xi} \underline{d} \underline{\omega} \end{aligned}$$

$R(\xi)$
bind

$$\text{but } i_{\xi} \underline{d} \underline{\omega} = i_{\xi} (-\underline{d} L + E) = + \underline{d} (i_{\xi} L) + i_{\xi} E$$