

Title: Lecture - Quantum Gravity, PHYS 644

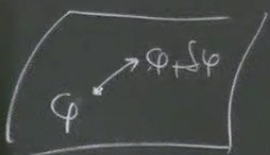
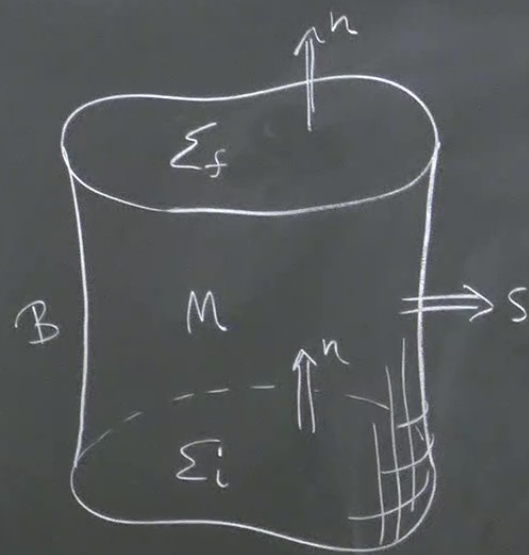
Speakers: Aldo Riello

Collection/Series: Quantum Gravity (Elective), PHYS 644, February 24 - March 28, 2025

Subject: Quantum Gravity

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Spacetime

mfd

M

coords

$x^a, a=1 \dots n$

Cartan
calculus

$(d, i., L.)$

$$df = \sum_a \frac{\partial f}{\partial x^a} dx^a$$

$$X = \sum_a X^a \frac{\partial}{\partial x^a}$$

$$L_X = i_X d + d i_X$$

Field space

$$\mathcal{F} = \Gamma(F \rightarrow M) \sim C^\infty(M, V)$$

$$\varphi^I(x), \quad \left\{ \begin{array}{l} I=1, \dots, m \\ x \in M \end{array} \right. \leftarrow \infty!$$

$(d, i., L.)$

$$dS = \int_M \sum_I \frac{\delta S}{\delta \varphi^I(x)} [d\varphi^I(x)]$$

↑ diff operator!
($d^2\varphi = \varphi d^2x$)

$$X = \int_M \sum_I \left(\int_X \varphi^I(x) \right) \frac{\delta S}{\delta \varphi^I(x)}$$

$$L_X = i_X d + d i_X$$

Mixed forms

$$\Omega^{\bullet}(M \times F) = \sum_{k=0}^{\dim(M)} \sum_{l=0}^{\infty} \Omega^{k,l}(M \times F) \quad \uparrow \subset \Omega^k(M) \otimes \Omega^l(F)$$

w/ differential $d := d + d'$

$$0 = d^2 = dd' + d'd$$

$$0 = i_X d' + d' i_X$$

etc.

Eg. $\Omega^1(M \times F)$

$$\ast \quad \alpha_1(x, \varphi) = \nabla_a \varphi(x) dx^a = -d\varphi \wedge d\varphi(x) \xrightarrow{(+)} \int_{\text{line}} \alpha_1 = \int_{\text{line}} d\varphi \wedge d\varphi \in \Omega^1(F)$$

$$\xrightarrow{(-)} = d\varphi \wedge d\varphi(x)$$

$$\ast \quad \alpha_2(x, \varphi) = \nabla_a d\varphi \wedge dx^a \xrightarrow{(+)} \int_{\text{line}} \alpha_2 = \int_{\text{line}} dd\varphi = d \int_{\text{line}} d\varphi$$

$$\xrightarrow{(-)} = dd\varphi = -dd\varphi$$

Takens thm

$$\Omega^{\text{top}, 1}(M \times F) = \underbrace{\Omega_{\text{src}}^{\text{top}, 1}(M \times F)}_{\text{"source"}} \oplus \underbrace{d\Omega^{\text{top}, 1}(M \times F)}_{\text{"boundary"}}$$

↳ no derivatives on $d\varphi^{\mathbb{I}}$

E.g. d_1 src if $\dim(M)=1$
 d_2 isn't.

[Anderson's homotopy operators]

$$\begin{aligned} \text{Rmk } J \in \Omega^{\text{top}, k} \\ *J &= (-1)^{\text{top}} \tilde{J}_a dx^a \\ *dJ &= (-1)^{k+1} (\nabla_a \tilde{J}^a) \end{aligned}$$

Applic
 $\Omega^{\text{top}, 1}$
 $\Omega^{\text{top}, 1}$

Application

$\Omega^{top,0}(M \times F) \ni \underline{L}(\varphi, x)$ assume depends only on $\varphi(x), \partial\varphi(x)$

$$\begin{aligned} \Omega^{top,1}(M \times F) \ni \underline{dL} &= \frac{\partial L}{\partial \varphi^I} \underline{d\varphi^I} + \frac{\partial L}{\partial (\partial \varphi^I)} \partial \underline{d\varphi^I} \\ &= \underbrace{\left(\frac{\partial L}{\partial \varphi^I} - \partial \frac{\partial L}{\partial (\partial \varphi^I)} \right)}_{\text{src}} \underline{d\varphi^I} + \underbrace{\partial \left(\frac{\partial L}{\partial (\partial \varphi^I)} \right)}_{\text{bdry}} \underline{d\varphi^I} \\ &\quad \begin{array}{l} \text{Euler-Lagrange} \\ \pi_{\text{src}}(\underline{dL}) \\ \pi_{\text{bdry}}(\underline{dL}) \end{array} \\ &= \underline{E}_I(\varphi) \underline{d\varphi^I} - d \underline{Q} \end{aligned}$$

$$p^I + \underbrace{\rho_e \left(\frac{\partial L}{\partial (\rho_e \psi^I)} \right)}_{\text{body}}$$

$$\vec{a} = \pm \frac{\partial L}{\partial (\rho_e \psi^I)} \psi^I$$

lepraye

$\pi_{\text{body}}(dL)$

(dL)

COVARIANT PHASE SPACE

acha

current
 $J \in \Omega^{top-1}(M) \otimes W$

$$*J = \pm (g_{ab} \tilde{J}^a dx^b)$$

$$top \leftarrow dJ = \pm (\nabla_a \tilde{J}^a) \epsilon$$

$$\int_{\partial M} J = \int_M dJ = \int_M \nabla_a \tilde{J}^a \sqrt{g} = \int_{\partial M} n_a \tilde{J}^a \sqrt{h}$$

$\uparrow \int \sqrt{g} dx^1 \dots dx^r$

action \rightarrow $S = \int_M \underline{L}$ \uparrow Lagrangian $\in \Omega^{top,0}$

$$d\underline{L} = E_I d\varphi^I - d\underline{\omega}$$

\uparrow e.o.m. \uparrow pre-symplectic current

the shell $\mathcal{F}_{EL} := \left\{ \varphi \in \mathcal{F} : E_I(\varphi) = 0 \right\} \xrightarrow{\mathcal{r}_{EL}} \mathcal{F}$

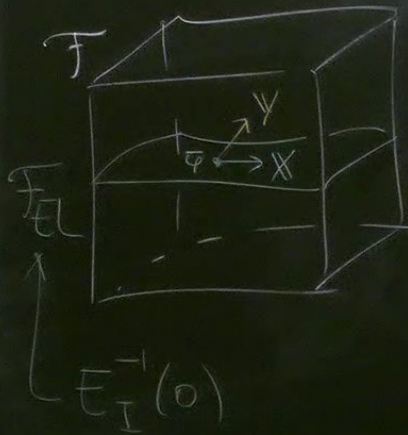
Notation: going on shell means pulling back to \mathcal{F}_{EL}

α, β are mixed forms $\alpha \approx \beta$ iff $\mathcal{r}_{EL}^*(\alpha - \beta) = 0$

What does it mean to pull back $d\varphi^I$?

→ It can only be contracted w/ $X \in T_{\bar{\varphi}} \mathcal{F}_{EL}$, $\bar{\varphi} \in \mathcal{F}_{EL}$

iff $\mathbb{L}_X E_I|_{\bar{\varphi}} = 0$



Y X
X ✓

$$\delta_X E_I|_{\bar{\varphi}} = \frac{\partial E_I}{\partial \varphi^J} (\delta_X \varphi^J) + \frac{\partial E_I}{\partial (\partial_a \varphi^J)} \partial_a (\delta_X \varphi^J) + \dots |_{\bar{\varphi}}$$

→ $\delta_X \varphi^J$ is a solution of the linearized eom of $\bar{\varphi}$
 aka "on shell perturbation"

$$= E_I(\bar{\varphi} + \delta_X \varphi) - E_I(\bar{\varphi}) + \mathcal{O}(\delta_X \varphi^2)$$

to pull back $d\varphi^I$?

contracted w/ $X \in T_{\bar{\varphi}} \mathcal{F}_{EL}$, $\bar{\varphi} \in \mathcal{F}_{EL}$

iff $L_X E_I|_{\bar{\varphi}} = 0$

X X
X ✓

$$\delta_X E_I|_{\bar{\varphi}} = \frac{\partial E_I}{\partial \varphi^J} (\delta_X \varphi^J) + \frac{\partial E_I}{\partial (\partial_a \varphi^J)} \partial_a (\delta_X \varphi^J) + \frac{\partial E_I}{\partial (\partial_a \partial_b \varphi^J)} \partial_a \partial_b (\delta_X \varphi^J) + \dots$$

$\rightarrow \delta_X \varphi^J$ is a solution of the linearized eom at $\bar{\varphi}$
 aka "on shell perturbation"

$$= E_I(\bar{\varphi} + \delta_X \varphi) - E_I(\bar{\varphi}) + \mathcal{O}(\delta_X \varphi^2)$$

$$E_I(\varphi) =$$

$$\delta_X E_I|_{\bar{\varphi}} =$$

$$L = -\frac{1}{2}(\nabla\varphi)^2 - \frac{1}{2}m^2\varphi^2$$

$$E_I(\varphi) = \square\varphi - m^2\varphi + \lambda\varphi^3$$

$$\delta_{xx} E_I|_{\bar{\varphi}} = \underline{\square\delta\varphi} - \underline{m^2\delta\varphi} + \underline{3\lambda\bar{\varphi}^2\delta\varphi}$$

η_{ab}

↓

$$\frac{E_I}{(\partial_a\varphi^j)}$$

d eom of $\bar{\varphi}$

can shell perturbation,

Thm
 Let $\underline{\Omega} := d\underline{\Theta} = \Omega^{top-1,2}(M \times F)$ presympl current
 then $d\underline{\Omega} \approx 0$ i.e. $\nabla_a \tilde{\Omega}^a \approx 0$ [on shell]

i.e. the presympl current is conserved on-shell

PF: $0 = d^2 \underline{L} = d\underline{E} - \underbrace{d d\underline{\Theta}}_{+d\underline{\Omega}} \approx d\underline{\Omega}$ \square

[on shell]

$$0 \approx \int_M d\underline{\Omega} = \int_{\Sigma_f} \underline{\Omega} - \int_{\Sigma_i} \underline{\Omega}$$

shell

$$\Rightarrow \omega := \int_{\Sigma} \underline{\Omega} \quad \text{is indep. of chosen Cauchy surface } \Sigma \hookrightarrow M$$

\mathbb{A}
 $\Omega^2(\mathcal{F}_{EL})$

no a "canonical" $(\mathcal{F}_{EL}, \omega) \approx \underline{Q}$: symplectic?
 COV. PH. SF.

Side rm k

$$\partial \Sigma \neq \emptyset$$

$$\partial M = \Sigma_f \cup \bar{\Sigma}_i \cup B$$

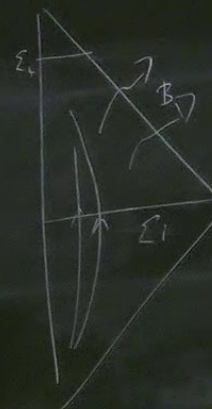
time k

"box"

balance eq

$$\Omega_{\Sigma_i} \approx \Omega_{\Sigma_f} + \underbrace{\int_B \Omega}_{\text{"presymplectic flux"}}$$

"presymplectic flux"



Ex scalar field (Mink)

$$\tilde{\omega} = -(\nabla^\alpha \phi) \partial_\alpha \phi$$

$$\tilde{\omega} = \partial_\alpha \phi \wedge \nabla^\alpha \partial_\beta \phi$$

$$\Omega_\Sigma = - \int_\Sigma \underbrace{n_\alpha \nabla^\alpha \partial_\beta \phi \wedge \partial^\beta \phi}_{\partial_t \partial_\alpha \phi} = - \int_\Sigma \partial \dot{\phi} \wedge \partial \phi$$

\uparrow $\{t=0\}$

initial conds \leftrightarrow sols eom $\sim F_{EL}$

$$\begin{cases} \pi(\vec{x}) = \dot{\phi}(t=0, \vec{x}) \\ \phi(\vec{x}) = \phi(t=0, \vec{x}) \end{cases}$$

$$(F_{EL}, \omega) \simeq \left(\text{cononical ph sp} \right) \int_\Sigma d^3x \wedge \partial \phi$$

no "canonical" $(F_{EL})^\omega \simeq Q$: symplectic?
 COV. PH. SP.

color field (Mink)

$$= -(\nabla^e \varphi) \lrcorner \mathbb{d}\varphi$$

$$= \mathbb{d}\varphi \wedge \nabla^e \mathbb{d}\varphi$$

$$= - \int_{\Sigma} \underbrace{\nabla^e \mathbb{d}\varphi \lrcorner \mathbb{d}\varphi}_{\partial_t \mathbb{d}\varphi} = - \int_{\Sigma} \mathbb{d}\dot{\varphi} \wedge \mathbb{d}\varphi$$

\uparrow initial conds \leftrightarrow sols eom $\sim F_{EL}$

$\omega \simeq$ (canonical ph sp)
 $(\mathbb{d}\pi \lrcorner \mathbb{d}\varphi)$

$$\begin{cases} \gamma(\vec{x}) := \dot{\varphi}(t=0, \vec{x}) \\ \phi(\vec{x}) := \varphi(t=0, \vec{x}) \end{cases}$$

But $(F_{EL})^\omega$ will be degenerate for a gauge theory!