

Title: Lecture - Gravitational Physics, PHYS 636

Speakers: Ruth Gregory

Collection/Series: Gravitational Physics (Elective), PHYS 636, January 6 - February 5, 2025

Subject: Cosmology, Strong Gravity

Date: February 03, 2025 - 9:00 AM

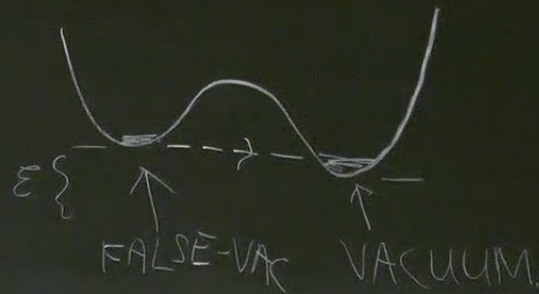
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LECTURE 13 Gravitational Instantons

Coleman "Fate of the false vacuum"

PRD 15 2929 (77)

Coleman de Luccia PRD 21 3305 (80)

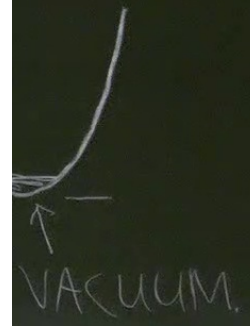


ational Instantons

false vacuum"

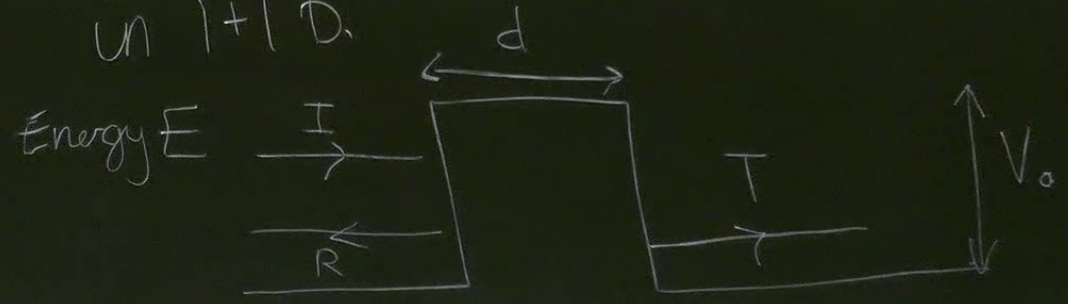
2929 (77)

21 3305 (80)



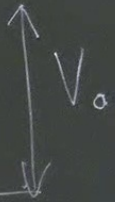
Tunnelling is perhaps the most startling phenomenon in quantum physics.

- Recall Schrödinger tunneling in 1+1 D.



ps the
phenomenon

r tunneling

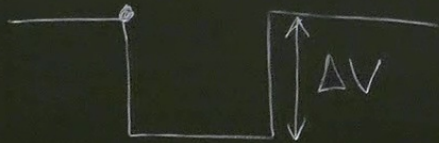


$$\text{Find } |T|^2 = \frac{1}{1 + \frac{V_0^2 \sinh^2 \Omega d}{4E(V_0 - E)}} \sim e^{-2\Omega d}$$

$$\text{where } \Omega^2 = \frac{2m(V_0 - E)}{\hbar^2} \sim 2\Delta V$$

$$\Omega d = \int_0^d \sqrt{2\Delta V} dx$$

Consider a different problem - classical
particle



$$\frac{1}{2} \dot{x}^2 = \Delta V \quad \text{in transit}$$

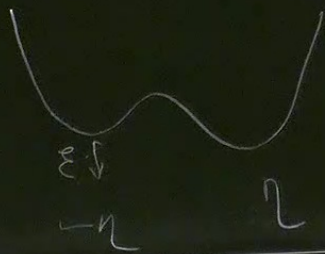
$$\begin{aligned} \text{So } \int \sqrt{2\Delta V} dx &= \int \dot{x} \sqrt{2\Delta V} dt \\ &= \int \left(\Delta V + \frac{1}{2} \dot{x}^2 \right) dt \end{aligned} \quad \begin{array}{l} \text{Euclidean} \\ \text{action} \end{array}$$

In general, $P \propto e^{-S_E}$ where
 S_E is the action of a Euclidean
trajectory from initial state to
emergence (+ back again - the bounce)



Euclidean
action

Translating to field theory,
are looking for Euclidean
solns that interpolate (in
some sense) between false &
true vacua.



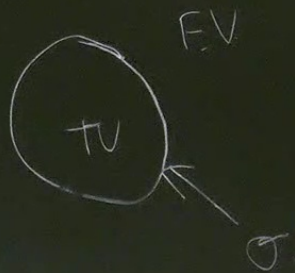
$$\mathcal{L}_\phi = \frac{1}{2}(\partial\phi)^2 - V(\phi)$$

$$V = \frac{\lambda}{2}(\phi^2 - \eta^2)^2 - \frac{\epsilon}{2\eta}(\phi - \eta)$$

$$= \int (\Delta V + \frac{1}{2} \dot{\phi}^2) dt$$

action

If we have a F.V. configuration
by quantum uncertainty a bubble of
T.V. could fluctuate into existence



Since scalar (in model)
has to transit between
"vacua" the boundary
will have energy - domain
wall

in figuration
a bubble of
existence
lar (in model)
ansit between
the boundary
energy - domain
Wall

TOO SMALL - contracts
TOO BIG - larger "cost"
In Euclidean space

$$\text{Energy in wall: } 2\pi^2 R^3 \delta$$

$$\text{Difference in vacuum energy: } \frac{\pi^2 \epsilon R^4}{2}$$

$$\delta E = 2\pi^2 R^3 \delta$$
$$\frac{d}{dR} \delta E = 6\pi^2 R^2 \delta$$

ost"

$$2\pi^2 R^3 \sigma$$

$$\frac{\pi^2}{2} \epsilon R^4$$

$$\delta E = 2\pi^2 R^3 \sigma - \frac{\pi^2}{2} \epsilon R^4$$

$$\frac{d}{dR} \delta E = 6\pi^2 R^2 \sigma - 2\pi^2 R^3 \epsilon = 0 \text{ at } R_0 = \frac{3\sigma}{\epsilon}$$

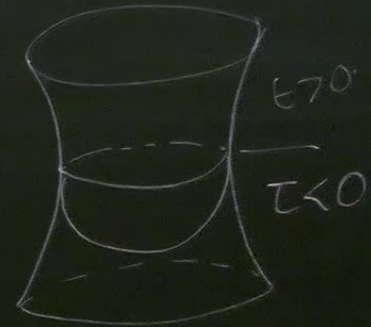
Scalar model : $\phi'' + 3\frac{\phi'}{\rho} = 2\lambda\phi(\phi^2 - \eta^2)$
 $\phi = \eta \tanh[\sqrt{\lambda}\eta(\rho - R_0)]$
 $\sigma = \frac{4}{3} \sqrt{\lambda} \eta^3$

Action : $B = \frac{\pi^2 R_0^3}{2} (4\sigma - \epsilon R_0)$
 $= \frac{27\pi^2}{2} \frac{\sigma^4}{\epsilon^3}$

What happens in real time?

$$r^2 + t^2 = R_0^2$$

$$\Rightarrow r^2 - t^2 = R_0^2$$



Bubble expands
in real time

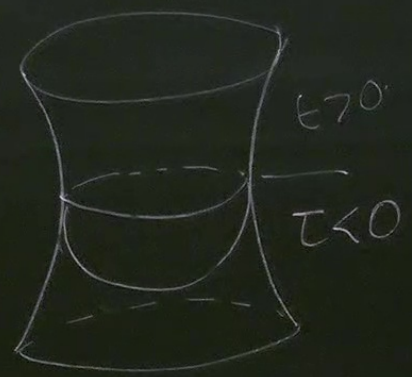
$$(4\sigma - \epsilon R_0)$$

$$\frac{\sigma^4}{\epsilon^3}$$

real time?

$$R_0^2$$

$$R_0^2$$



Bubble expands
in real time.

Energy gravitates.

The positive energy vac
will be a de Sitter
universe $\Lambda = 8\pi G \epsilon \sqrt{}$

EoS

radius l .



$$l^2 = \frac{3}{8\pi G \epsilon}$$

True

gravitates.

the energy vac
de Sitter

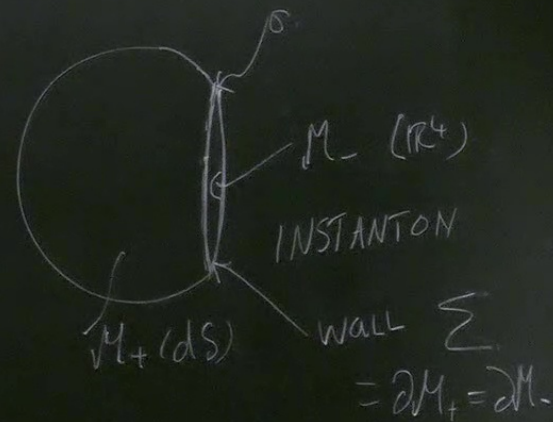
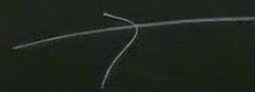
$$\Lambda = 8\pi G \epsilon \cdot v$$

$$l^2 = \frac{3}{8\pi G \epsilon}$$

True vacuum is Minkowski/Euclidean \mathbb{R}^4



ds_+



gravitates.

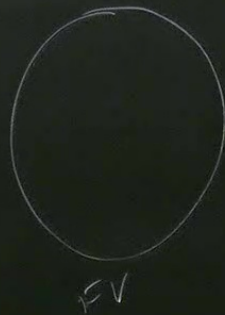
vacuum energy

de Sitter

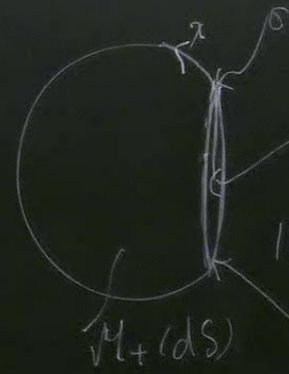
$$\Lambda = 8\pi G \epsilon_v$$

$$l^2 = \frac{3}{8\pi G \epsilon_v}$$

True vacuum is Minkowski/Euclidean \mathbb{R}^4



EV



$M_-(\mathbb{R}^4)$

INSTANTON

WALL Σ

$= \partial M_+ = \partial M_-$

(ds)

$$ds_+^2 = l^2 (d\chi^2 + \sin^2 \chi d\Omega_{III}^2)$$

(\mathbb{R}^4)

$$ds_-^2 = dp^2 + p^2 d\Omega_{III}^2$$

Match at $p_0 = R_0 = l \sin \chi_0$ (Σ)

$$n_+ = \frac{1}{l} \frac{\partial}{\partial \chi} \quad n_- = \frac{\partial}{\partial p}$$

$$K_{+\alpha\beta} = \frac{1}{l} \Gamma_{\alpha\beta}^{\chi} = -\frac{1}{2l} g^{\chi\chi} g_{\alpha\beta,\chi} = -\frac{\cot \chi_0}{l} g_{\alpha\beta}$$

$$K_{-\alpha\beta} = \Gamma_{\alpha\beta}^p = -\frac{1}{p_0} g_{\alpha\beta}$$

$$\Delta K_{\alpha\beta} - \Delta K h_{\alpha\beta}$$

$$(\Sigma) \quad \Delta K_{\alpha\beta} - \Delta K h_{\alpha\beta} = -\frac{2h_{\alpha\beta}}{l \sin \chi_0} (1 - \cos \chi_0) = -8\pi G \sigma h_{\alpha\beta}.$$

$$\Rightarrow \frac{(1 - \cos \chi_0)}{l \sin \chi_0} = \frac{(1 - \sqrt{1 - R^2/l^2})}{R} = 4\pi G \sigma.$$

$I_{\alpha\beta}$

$$B = I_{\text{INST}} - I_{\text{BACKGROUND}}$$

$$= \frac{1}{16\pi G} \int_{\chi < \chi_0} \sqrt{g} \underbrace{(R - 2\lambda)}_{2\lambda} - \frac{1}{8\pi G} \int \Delta K h$$

$$\Delta K_{\alpha\beta} - \Delta K h_{\alpha\beta} = -\frac{2h_{\alpha\beta}}{l \sin \chi_0} (1 - \cos \chi_0) = -8\pi G \sigma h_{\alpha\beta}$$

$$\Rightarrow \frac{(1 - \cos \chi_0)}{l \sin \chi_0} = \frac{(1 - \sqrt{1 - r^2/l^2})}{R} = 4\pi G \sigma$$

$$B = I_{\text{INS}} - I_{\text{BACKGROUND}}$$

$$= \frac{1}{16\pi G} \int_{\chi < \chi_0} \sqrt{g} (R - 2l) - \frac{1}{8\pi G} \int_{\Sigma} \Delta K h + \int_{\Sigma} \sigma h$$

\downarrow
 $12\pi G \sigma$

$$= \frac{3}{8\pi G l^2} \cdot 2\pi^2 l^4 \int_0^{\chi_0} \sin^3 \chi d\chi - \frac{\sigma}{2} \cdot 2\pi^2 R^3$$

$$= \frac{3\pi l^2}{4G} \left[\frac{2}{3} - \cos \chi_0 + \frac{1}{3} \cos^3 \chi_0 \right]$$

$$- \frac{\pi^2 R^3}{4\pi G} \frac{(1 - \cos \chi_0)}{R}$$

$$= \frac{\pi l^2}{4G} (1 - \cos \chi_0) \left[\frac{2 - \cos \chi_0 - \cos^3 \chi_0}{- \sin^3 \chi_0} \right]$$

$$\Delta K_{\text{ap}} - \Delta K_{\text{h ap}} = -\frac{2h_{\text{ap}}}{l \sin \chi_0} (1 - \cos \chi_0) = -8\pi\epsilon_0 h_{\text{ap}}$$

$$\Rightarrow \frac{(1 - \cos \chi_0)}{l \sin \chi_0} = \frac{(1 - \sqrt{1 - R^2/\epsilon^2})}{R} = 4\pi\epsilon_0$$

$$B = I_{\text{INS}} - I_{\text{BACKGROUND}}$$

$$= \frac{1}{16\pi\epsilon} \int_{\chi < \chi_0} \sqrt{g} (R - 2l) - \frac{1}{8\pi\epsilon} \int_{\Sigma} \Delta K_{\text{IH}} + \int_{\Sigma} \sigma_{\text{IH}}$$

$$= \frac{3}{8\pi\epsilon l^2} \cdot 2\pi l^4 \int_0^{\chi_0} \sin^3 \chi d\chi - \frac{\sigma}{2} \cdot 2\pi^2 R^3$$

$$= \frac{3\pi l^2}{4\epsilon} \left[\frac{2}{3} - \cos \chi_0 + \frac{1}{3} \cos^3 \chi_0 \right]$$

$$- \frac{\pi^2 R^2}{4\pi\epsilon} \frac{(1 - \cos \chi_0)}{R}$$

$$= \frac{\pi l^2}{4\epsilon} (1 - \cos \chi_0) \left[2 - \cos \chi_0 - \cos^3 \chi_0 \right]$$

$$= \frac{\pi^2}{4\epsilon} (1 - \cos \chi_0)^2 = \frac{\pi l^2}{4\epsilon} (4\pi\epsilon_0 R)^2$$

$$K_{\alpha\beta} = \Gamma_{\alpha\beta}^p = -\frac{1}{\rho_0} g_{\alpha\beta}$$

$$= \frac{1}{16\pi G} \int_{x < x_0} \sqrt{g} (R - 2\Lambda) - 2\Lambda$$

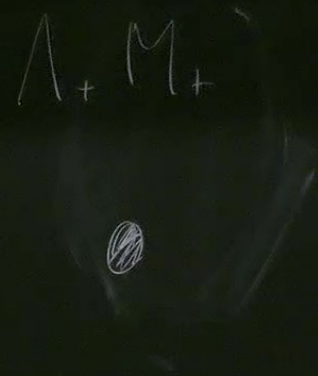
Then find $\sin \chi_0 = \frac{2(4\pi G \sigma l)^2}{1 + (4\pi G \sigma l)^2}$

$$\alpha B = \frac{\pi l^2}{G} \frac{(4\pi G \sigma l)^4}{[1 + (4\pi G \sigma l)^2]^2}$$

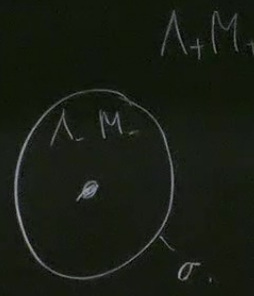
$$\frac{1}{2} \int_{-\pi/2}^{\pi/2} \cos^2 \theta d\theta = \frac{1}{2} \int_{-\pi/2}^{\pi/2} \frac{1 + \cos 2\theta}{2} d\theta = \frac{1}{4} \left[\theta + \frac{\sin 2\theta}{2} \right]_{-\pi/2}^{\pi/2} = \frac{1}{4} \left(\pi + 0 - (-\pi + 0) \right) = \frac{\pi}{2}$$

CDL is the "gold standard" for tunnelling, but is incredibly idealized
 Most phase transitions are catalyzed by impurities

standard" for
edibly idealized
e catalyzed by
black hole.

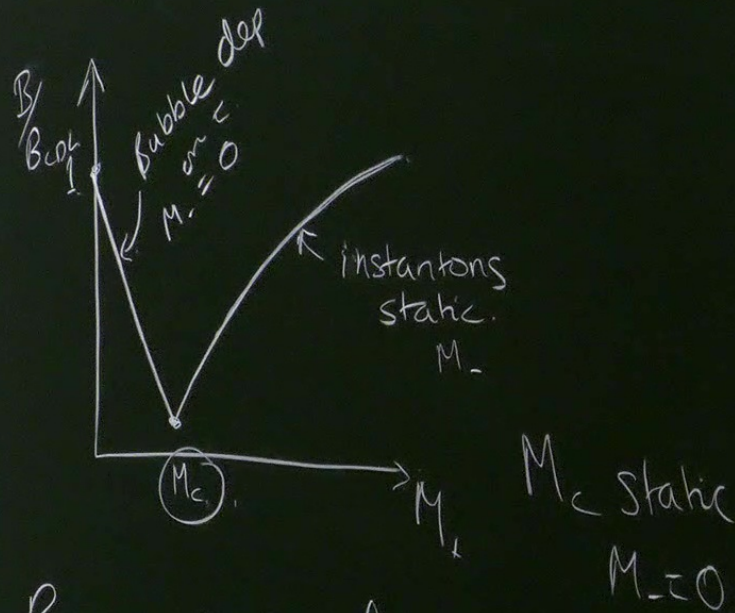


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SDS has
conical singularities





Range of M_- for M_+

Lowest action dominant