

**Title:** Matter-Antimatter Asymmetry: With the SM and Beyond

**Speakers:** Seyda Ipek

**Collection/Series:** Particle Physics

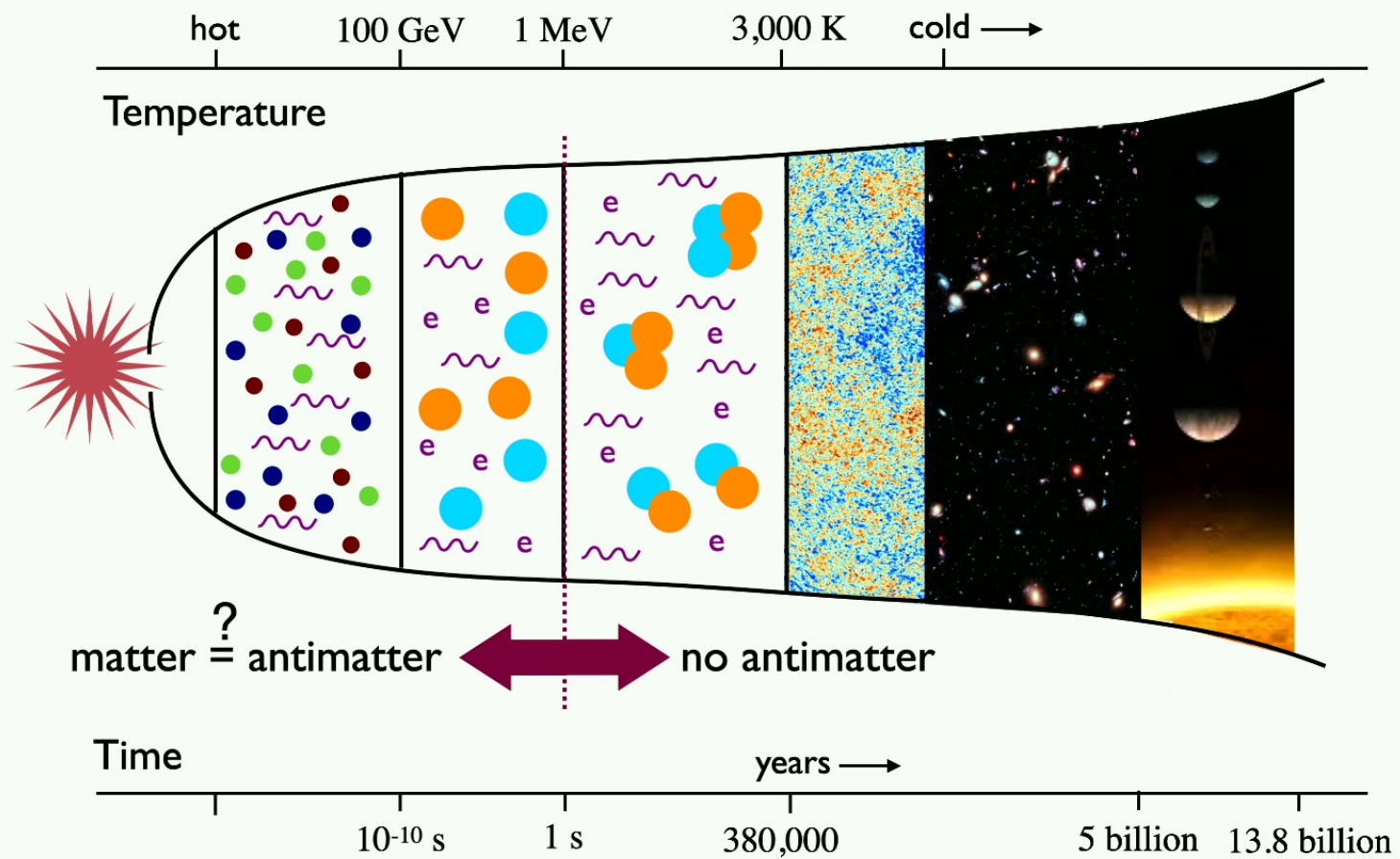
**Subject:** Particle Physics

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**URL:** <https://pirsa.org/25010079>

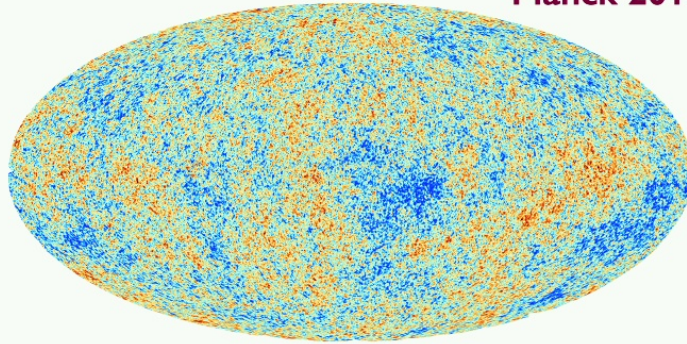
**Abstract:**

There is more matter than antimatter in the universe. This asymmetry requires three conditions: 1- Baryon (or lepton) number violation, 2- C and CP violation and 3- out-of-thermal equilibrium conditions in the early universe, before BBN. Although the SM does not have any out-of-equilibrium process, it does provide baryon number violation and CP violation. Still, it is commonly accepted that the SM CP violation is not enough for producing the observed baryon asymmetry. I will present one new physics model in which the SM CP violation, that is measured in the B meson system, is in fact enough to generate the baryon asymmetry.



# Cosmic Microwave Background

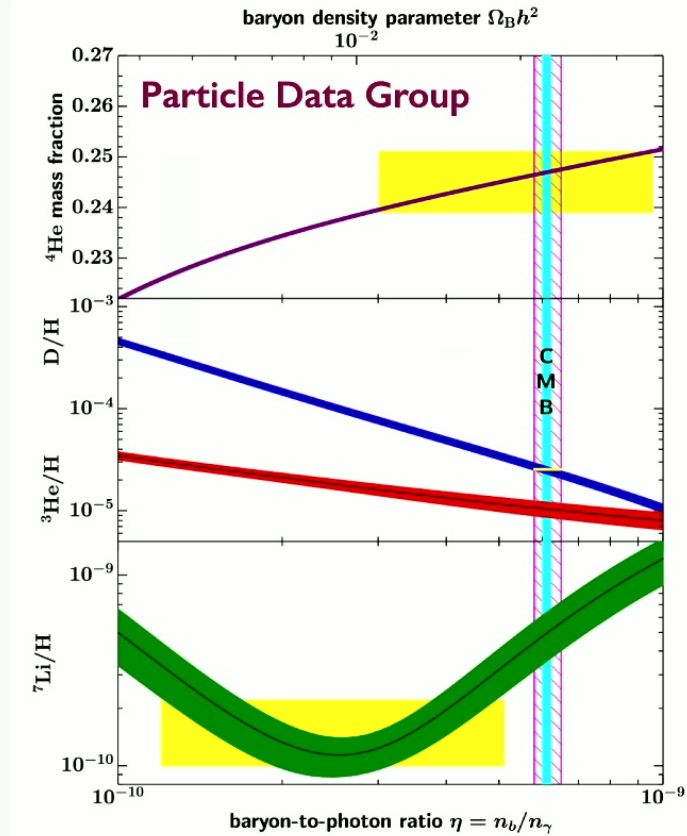
Planck 2015



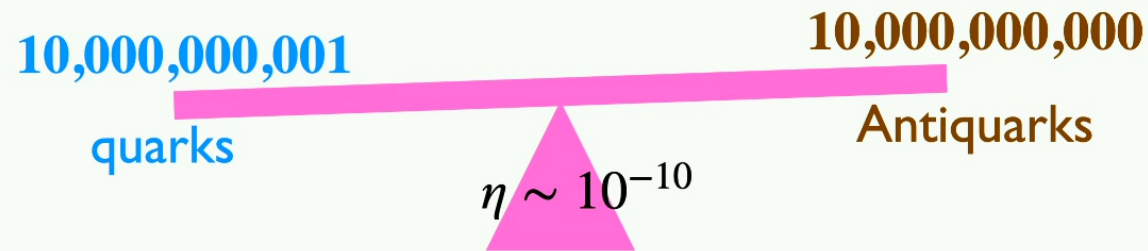
Baryon-to-photon ratio:

$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma} \simeq 6 \times 10^{-10}$$

# Primordial light element abundances



How do we make sure there are more quarks than antiquarks in the early Universe?



Physics need to be a little bit different between matter and antimatter!



JETP Lett. 6 (1967) 4

### Sakharov conditions

Andrei Sakharov  
1921-1989

1. Baryon (matter) number cannot be a conserved quantity
2. Charge and Charge-Parity (CP) symmetries must be violated
3. Out-of-equilibrium processes

?

?

?

?

Can the Standard Model of particle physics explain the baryon asymmetry of the Universe?

?

?

?

Does the Standard Model satisfy the Sakharov Conditions?

?

?

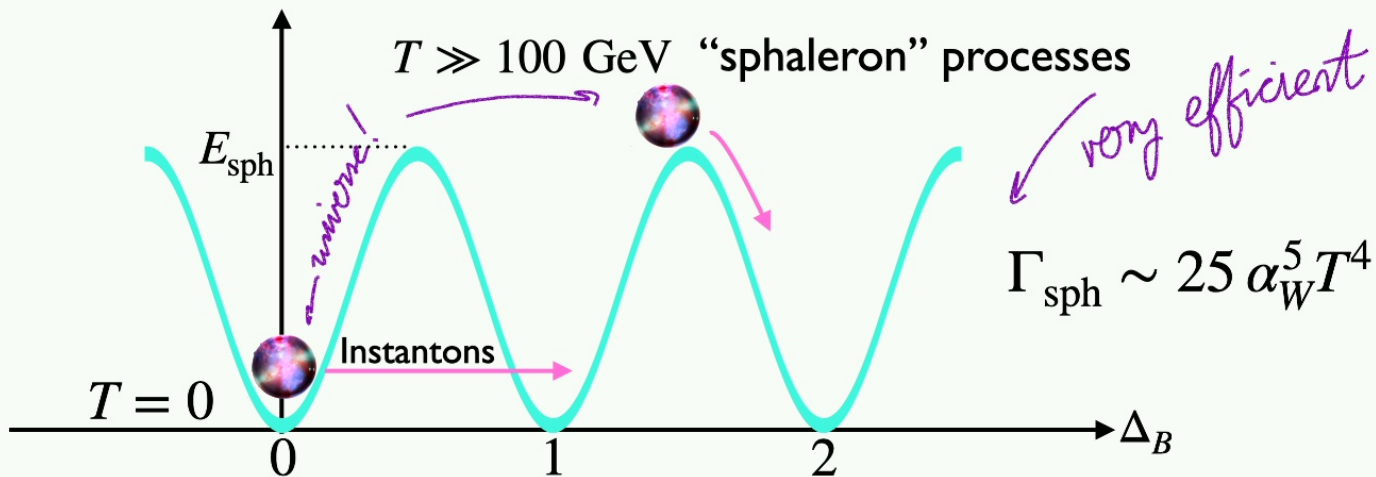
?

?

# Baryon number is violated in weak interactions 😊

only left-handed particles interact via the weak nuclear force

$$\partial^\mu j_\mu^B = 3 \partial^\mu j_\mu^{L_i} = 3 \frac{g^2}{32\pi^2} W^{\mu\nu,a} \tilde{W}_{\mu\nu}^a \quad \rightarrow \quad \Delta_B = \int d^4x \partial^\mu j_\mu^B = 3 \frac{g^2}{32\pi^2} \int d^4x W^{\mu\nu,a} \tilde{W}_{\mu\nu}^a$$

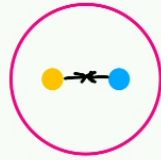


Quantum tunneling is hard!

$$\Gamma \sim e^{-4\pi/\alpha_W} \sim e^{-160}$$

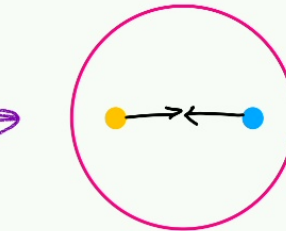
$$E_{\text{sph}} \sim \frac{M_W}{\alpha_W} \sim 10 \text{ TeV}$$

# Equilibrium?



Rate of  
(weak) interactions

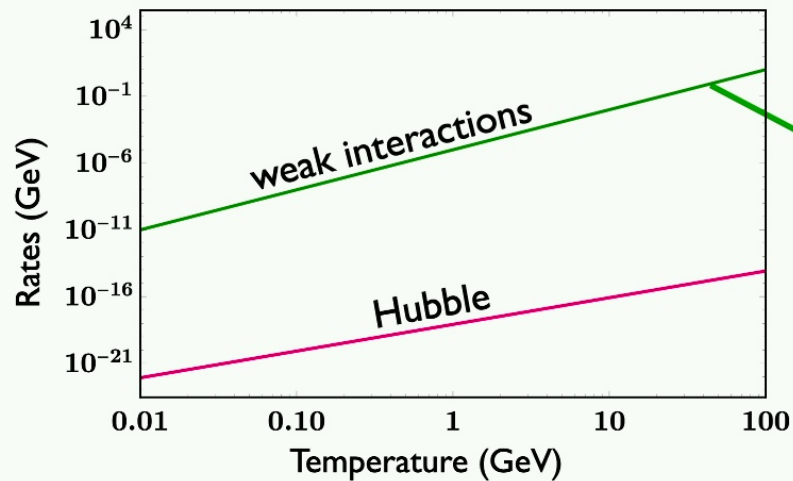
vs



Expansion rate  
of the universe

$$\Gamma_{\text{weak}} \sim G_F^2 \times T^3 \sim \frac{T^3}{10^{10} \text{ GeV}^2}$$

$$H \sim \frac{T^2}{M_{\text{Planck}}} \sim \frac{T^2}{10^{19} \text{ GeV}}$$



Too fast!



SM Universe  
always  
equilibrates!



# CP Violation in the Standard Model



## CP is violated in **weak interactions**

$$K_L \rightarrow 2\pi \quad \text{AND} \quad K_L \rightarrow 3\pi$$

A historical review: Cronin, *Eur. Phys. J. H* 36 (2012) pp.487-508

Entirely because there is a complex phase in the CKM matrix

Great! BUT not enough for the baryon asymmetry



handwavey:

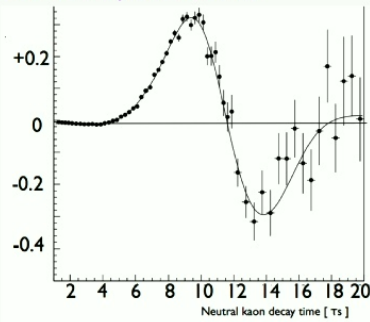
$$\eta \sim J \prod_i \left( \frac{m_i}{M_W} \right)^2$$

more detailed calculations:

$$\eta_{\text{SM CP}} \sim 10^{-20}$$

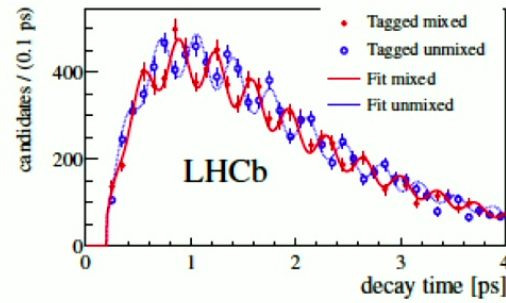
Gavela, Hernandez, Orloff, Pene, CERN 93/7081

CPLear, 1990-96



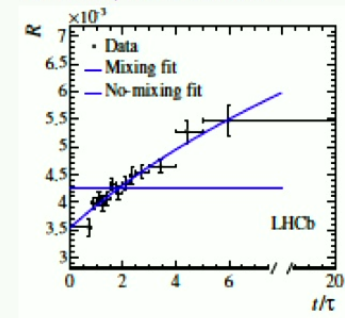
Kaon

LHCb, 1304.4741

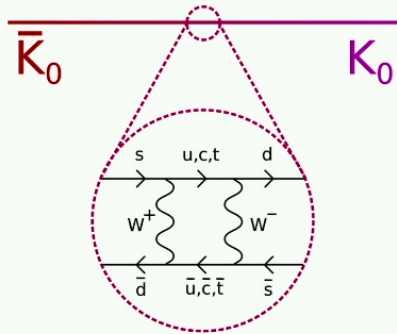


B mesons

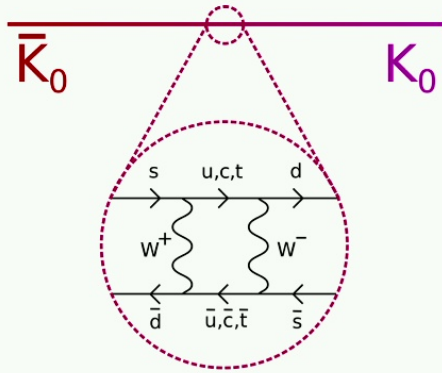
LHCb, 1211.1230



D meson



Are particle—antiparticle oscillations special for CP violation?



Hamiltonian: 
$$\mathbf{H} = \mathbf{M} - \frac{i}{2}\mathbf{\Gamma}$$

$$\mathbf{M} = \begin{pmatrix} M_{11} & M_{12} \\ M_{12} & M_{22} \end{pmatrix}$$

$$\mathbf{\Gamma} = \begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{12} & \Gamma_{22} \end{pmatrix}$$

in the  
 $\{|B^0\rangle, |\bar{B}^0\rangle\}$   
 basis

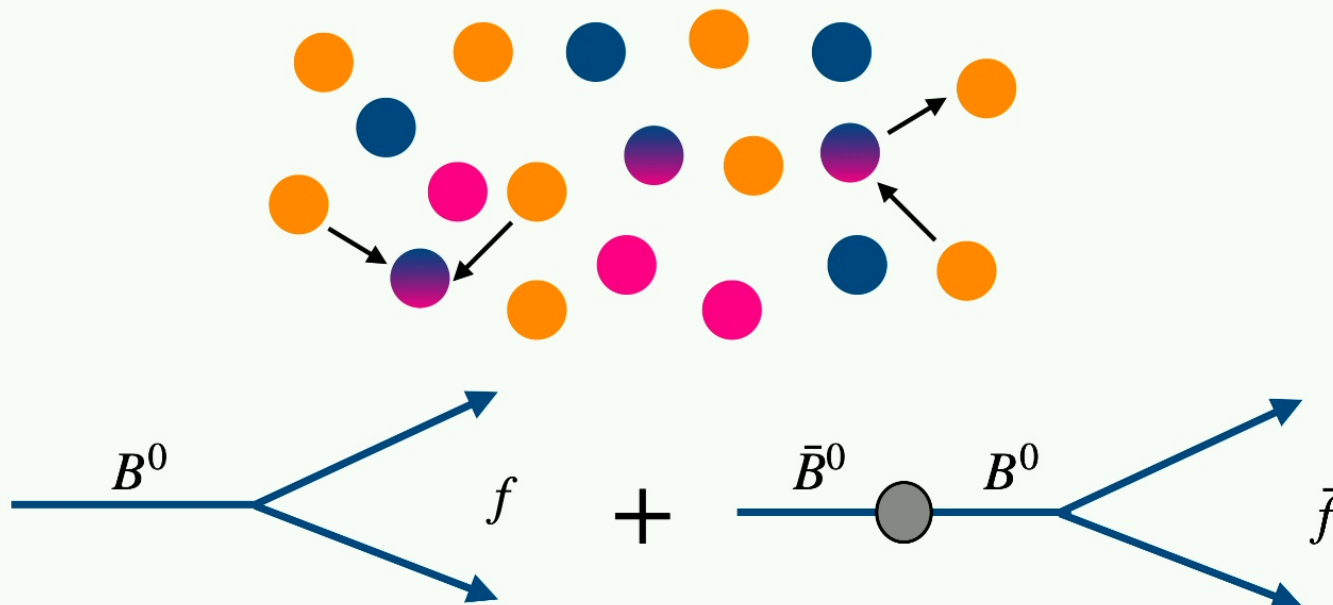
eigenvalues:  $|B_{H,L}\rangle = p|B\rangle \pm q|\bar{B}\rangle$   $\left(\frac{q}{p}\right)^2 = \frac{M_{12}^* - (i/2)\Gamma_{12}^*}{M_{12} - (i/2)\Gamma_{12}}$

mass states  $\neq$  interaction states

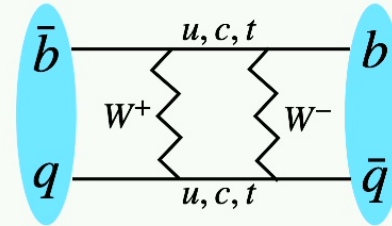
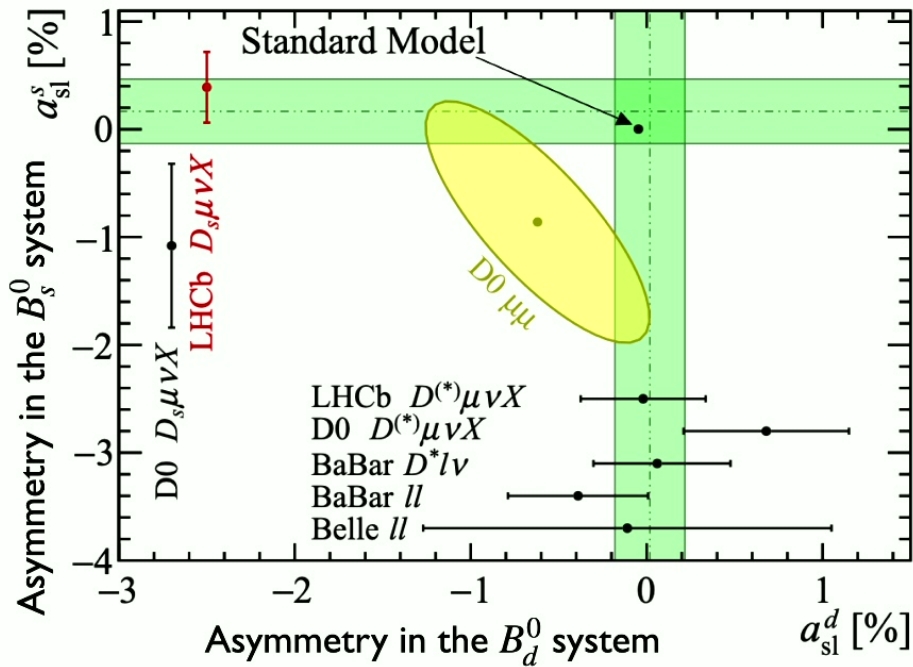
**→ OSCILLATIONS!**

CP violation exists as a phase difference between  $M$  and  $\Gamma$

$$A_{\text{SL}} = \frac{\Gamma(B^0(t) \rightarrow \ell^+ X^-) - \Gamma(\bar{B}^0(t) \rightarrow \ell^- X^+)}{\Gamma(B^0(t) \rightarrow \ell^+ X^-) + \Gamma(\bar{B}^0(t) \rightarrow \ell^- X^+)}$$



Gershon and Gligorov, arXiv:1607.06746



$$q = d, s$$

There is some room for new physics!

SM predictions:

$$A_{sl}^d|_{SM} = (-4.7 \pm 0.4) \times 10^{-4}$$

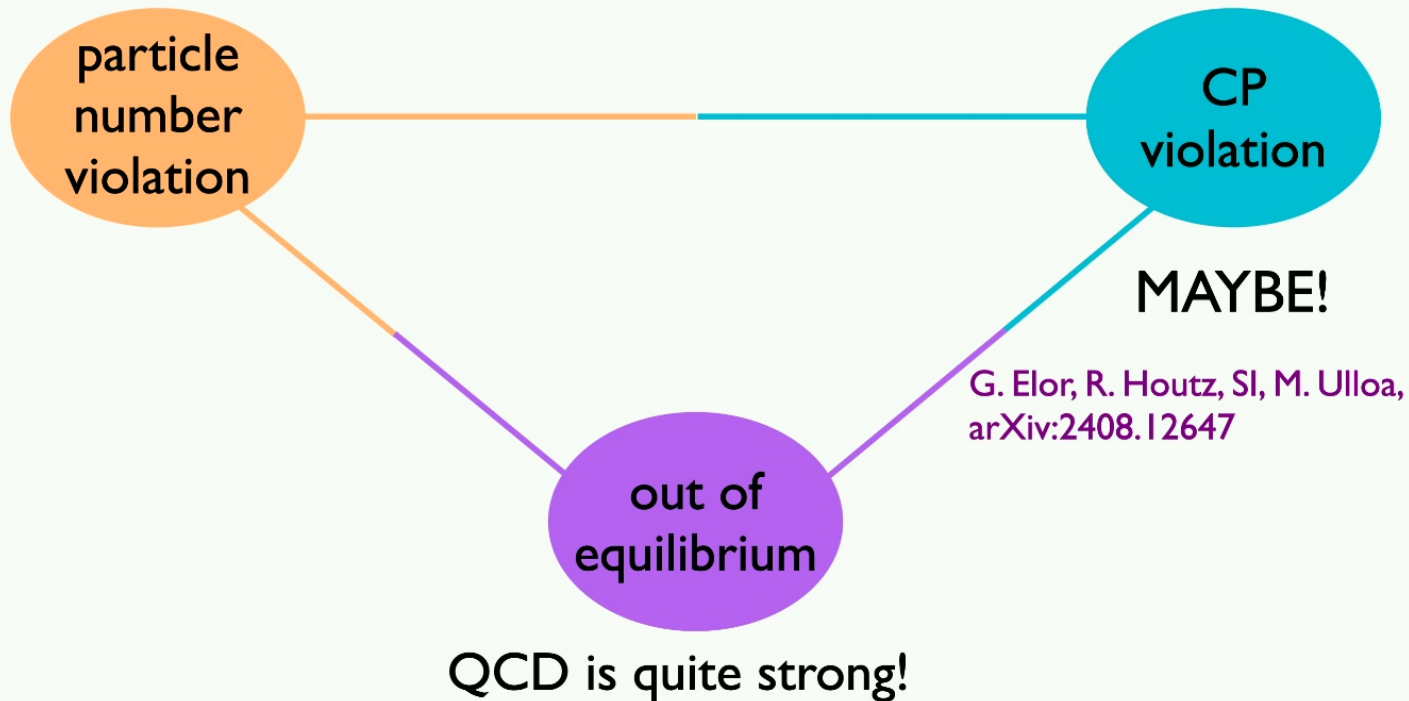
$$A_{sl}^s|_{SM} = (2.1 \pm 0.2) \times 10^{-5}$$

Lenz and Tetlalmatzi-Xolocotzi, arXiv: 1912.07621

Is any of this relevant for the baryon asymmetry?

$B^0$  mesons are formed below  $T \sim \text{GeV}$

no sphalerons and no baryon number violation!



# A Great Model: (B-)Mesogenesis

K. Aitken, D. McKeen, T. Neder, **A. Nelson**, arXiv: 1708.01259

G. Elor, M. Escudero, **A. Nelson**, arXiv: 1810.00880

**A. Nelson**, H. Xiao, arXiv: 1901.08141

G. Alonso-Alvarez, G. Elor, **A. Nelson**, H. Xiao, arXiv: 1907.10612

G. Elor, R. McGehee, arXiv: 2011.06115

G. Alonso-Alvarez, M. Escudero, G. Elor, arXiv: 2101.02706



F. Elahi, G. Elor, R. McGehee, arXiv: 2109.09751

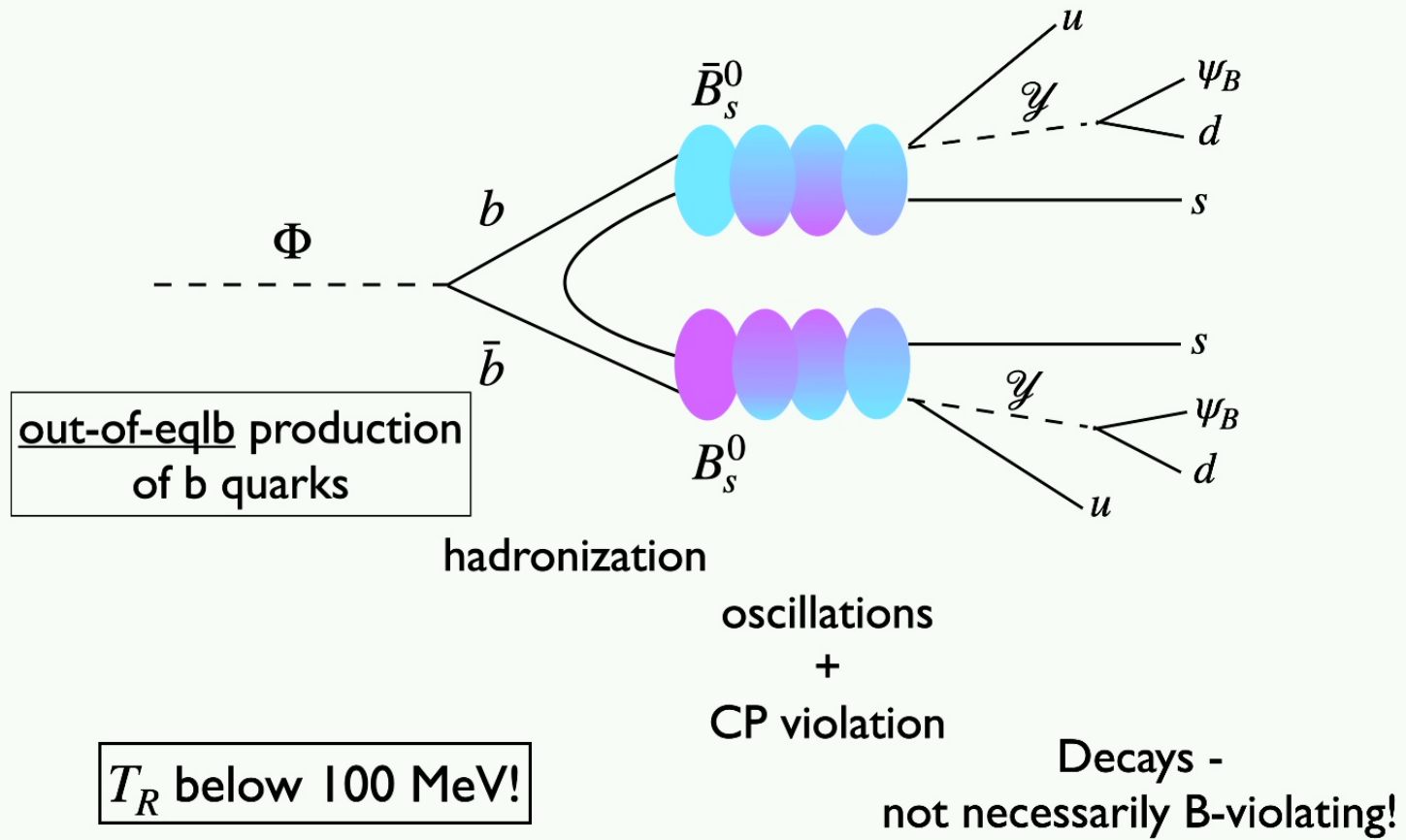
J. Berger, G. Elor, arXiv: 2301.04165

+...

Baryogenesis  
via  
B meson oscillations  
in the early Universe???



# A simplified version of the model



To be a bit more concrete:

		$SU(3)$	$U(1)$	mass (GeV)
$\Phi$	scalar	1	0	11-100
$\mathcal{Y}$	scalar	3	-1/3	will come to this
$\Psi_{\mathcal{B}}$	fermion	1	0	1

Baryon number -1 ←

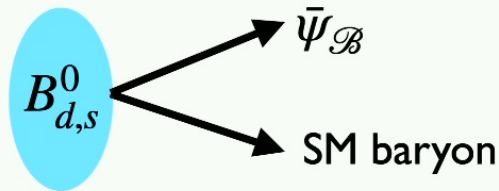
$$\mathcal{L} = - \sum_{i,j} y_{u_i d_j} \mathcal{Y}^* \bar{u}_{iR} d_{jR}^c - \sum_k y_{\psi d_k} \bar{\Psi}_{\mathcal{B}} \mathcal{Y} d_{kR}^c + \text{h.c.}$$

integrating out  $\mathcal{Y}$

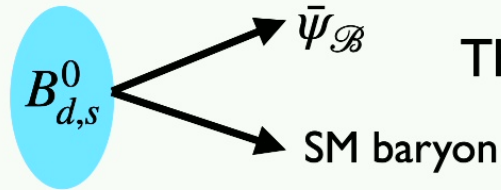
$$\mathcal{O}_{d_k, u_i d_j} = C_{d_k, u_i d_j} \epsilon_{\alpha\beta\gamma} (\bar{\Psi}_{\mathcal{B}} d_k^\alpha) (\bar{d}_j^\beta u_i^\gamma)$$

$$C_{d_k, u_i d_j} \equiv \frac{y_{\psi d_k} y_{u_i d_j}}{M_{\mathcal{Y}}^2}$$

exotic  
B meson decays



Baryon number  
is conserved



The (visible) baryon asymmetry generated:

G. Alonso-Alvarez, M. Escudero, G. Elor, arXiv: 2101.02706

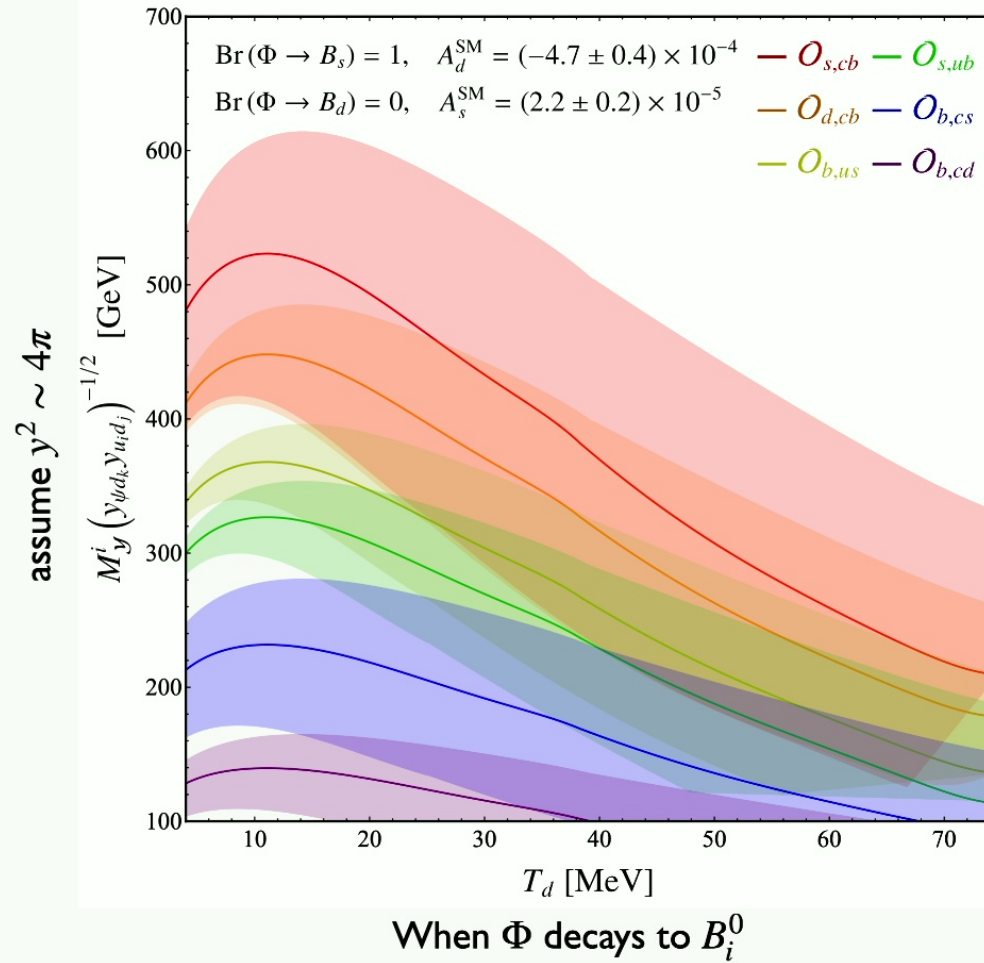
$$Y_{\mathcal{B}} \simeq 5 \times 10^{-5} \sum_{i=d,s} \underbrace{[\text{Br}(B_i^0 \rightarrow \bar{\psi}_{\mathcal{B}} \mathcal{B}_{\text{SM}})]}_{\substack{\text{observable} \\ \sim \frac{1}{M_{\mathcal{Y}}^4}}} \underbrace{A_{sl}^i}_{\substack{\text{observable} \\ \sim 10^{-5} \text{ from the SM}}} \alpha_i(T_d)$$

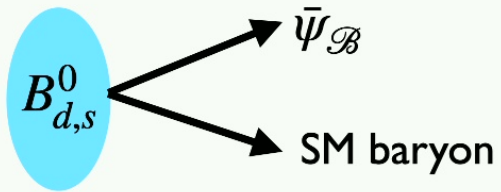
$0 \leq \alpha_I \lesssim 1.5$   
numerical details from solving the Boltzmann equations involved

not much room for exotic branching fractions!

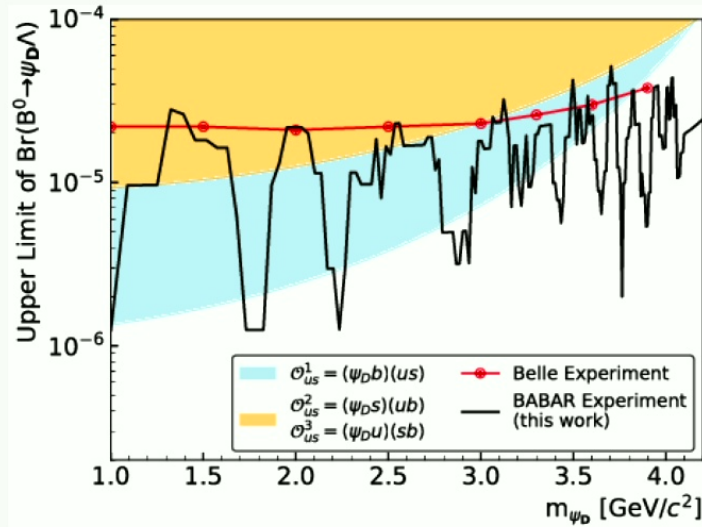
Observed asymmetry:  $Y_{\mathcal{B}}^{\text{meas}} = \frac{n_B - n_{\bar{B}}}{s} \simeq 8 \times 10^{-11}$

## Successful Baryogenesis

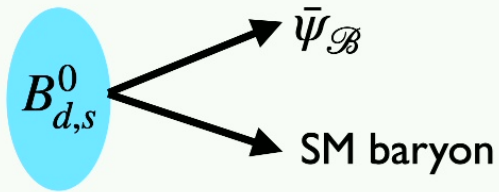




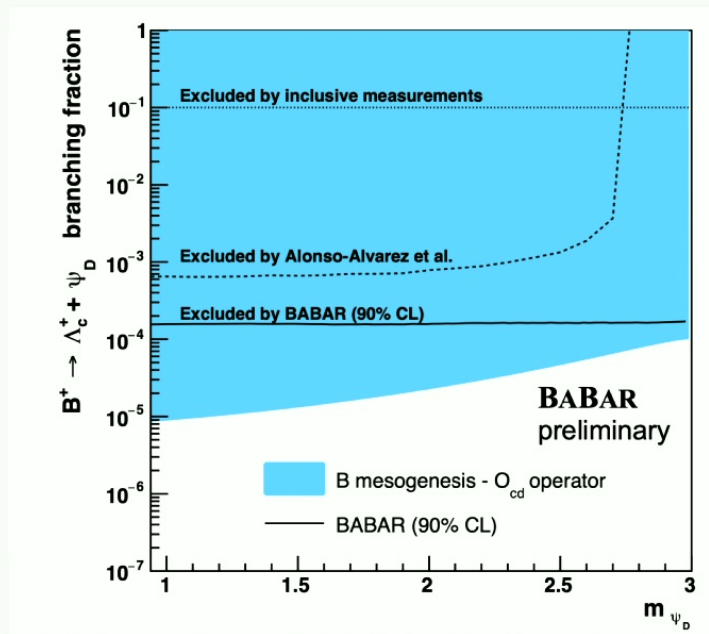
A lot of experimental constraints from exotic B decays + LHC



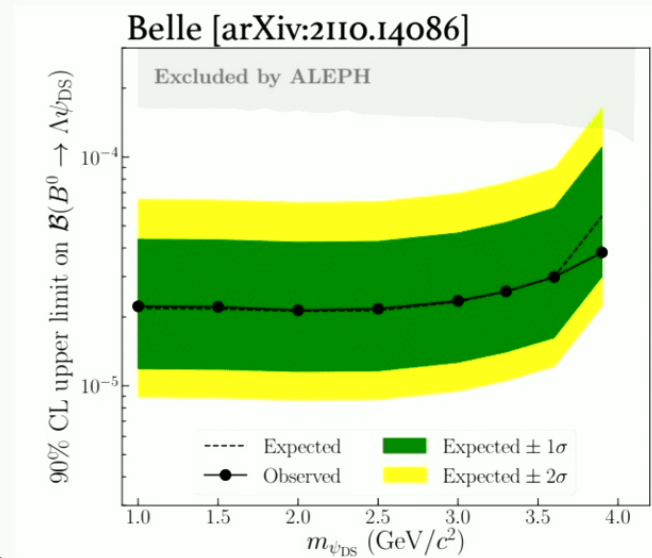
BaBar, *Phys.Rev.D* 107 (2023) 9, 092001, 2302.0028



A lot of experimental constraints from exotic B decays + LHC



collider constraints from squark searches also apply



Operator	$(M_y^f)_{\min}$ [TeV]	Decay	$\Gamma_0$ [GeV <sup>5</sup> ]
$\mathcal{O}_{b,ud}$	$\sim 1.7\sqrt{y_{\psi b} y_{ud}}$	$B_d \rightarrow \bar{\psi}_B n$ $B_s \rightarrow \bar{\psi}_B \Lambda$	$3.5_{\pm 0.4} \cdot 10^{-5}$ n.a.
$\mathcal{O}_{b,us}$	$\sim 1.7\sqrt{y_{\psi b} y_{us}}$	$B_d \rightarrow \bar{\psi}_B \Lambda$ $B_s \rightarrow \bar{\psi}_B \Xi^0$	$1.4_{\pm 0.1} \cdot 10^{-4}$ $3.2_{\pm 0.1} \cdot 10^{-5}$
$\mathcal{O}_{b,cd}$	$\sim 0.9\sqrt{y_{\psi b} y_{cd}}$	$B_d \rightarrow \bar{\psi}_B \Sigma_c^0$ $B_s \rightarrow \bar{\psi}_B \Xi_c^0$	$0.7_{\pm 0.4} \cdot 10^{-6}$ $6.6_{\pm 3.3} \cdot 10^{-7}$
$\mathcal{O}_{b,cs}$	$\sim 0.9\sqrt{y_{\psi b} y_{cs}}$	$B_d \rightarrow \bar{\psi}_B \Xi_c^0$ $B_s \rightarrow \bar{\psi}_B \Omega_c$	$4.7_{\pm 2.0} \cdot 10^{-6}$ $5.0_{\pm 3.0} \cdot 10^{-6}$

Operator	$(M_y^f)_{\min}$ [TeV]	Decay	$\Gamma_0$ [GeV <sup>5</sup> ]
$\mathcal{O}_{d,ub}$	$\sim 3.8\sqrt{y_{\psi d} y_{ub}}$	$B_d \rightarrow \bar{\psi}_B n$ $B_s \rightarrow \bar{\psi}_B \Lambda$	$3.6_{\pm 0.4} \cdot 10^{-5}$ n.a.
$\mathcal{O}_{s,ub}$	$\sim 2.3\sqrt{y_{\psi s} y_{ub}}$	$B_d \rightarrow \bar{\psi}_B \Lambda$ $B_s \rightarrow \bar{\psi}_B \Xi^0$	$1.3_{\pm 0.4} \cdot 10^{-4}$ $2.0_{\pm 0.1} \cdot 10^{-5}$
$\mathcal{O}_{d,cb}$	$\sim 1.1\sqrt{y_{\psi d} y_{cb}}$	$B_d \rightarrow \bar{\psi}_B \Sigma_c^0$ $B_s \rightarrow \bar{\psi}_B \Xi_c^0$	$8.2_{\pm 0.4} \cdot 10^{-5}$ $7.0_{\pm 0.4} \cdot 10^{-5}$
$\mathcal{O}_{s,cb}$	$\sim 1.1\sqrt{y_{\psi s} y_{cb}}$	$B_d \rightarrow \bar{\psi}_B \Xi_c^0$ $B_s \rightarrow \bar{\psi}_B \Omega_c$	$9.7_{\pm 5.0} \cdot 10^{-5}$ $1.3_{\pm 0.6} \cdot 10^{-4}$

**Constraints:**  
LEP recast +  
designated searches at  
Belle II and BaBar

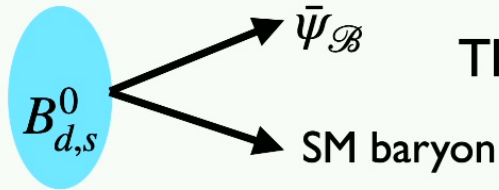
$\psi_B$  counted as  
missing energy

Take  $m_{\psi_B} = 1$  GeV

maximum allowed from proton decay

$$M_y \gtrsim 1 \text{ TeV}$$





The (visible) baryon asymmetry generated:

G. Alonso-Alvarez, M. Escudero, G. Elor, arXiv: 2101.02706

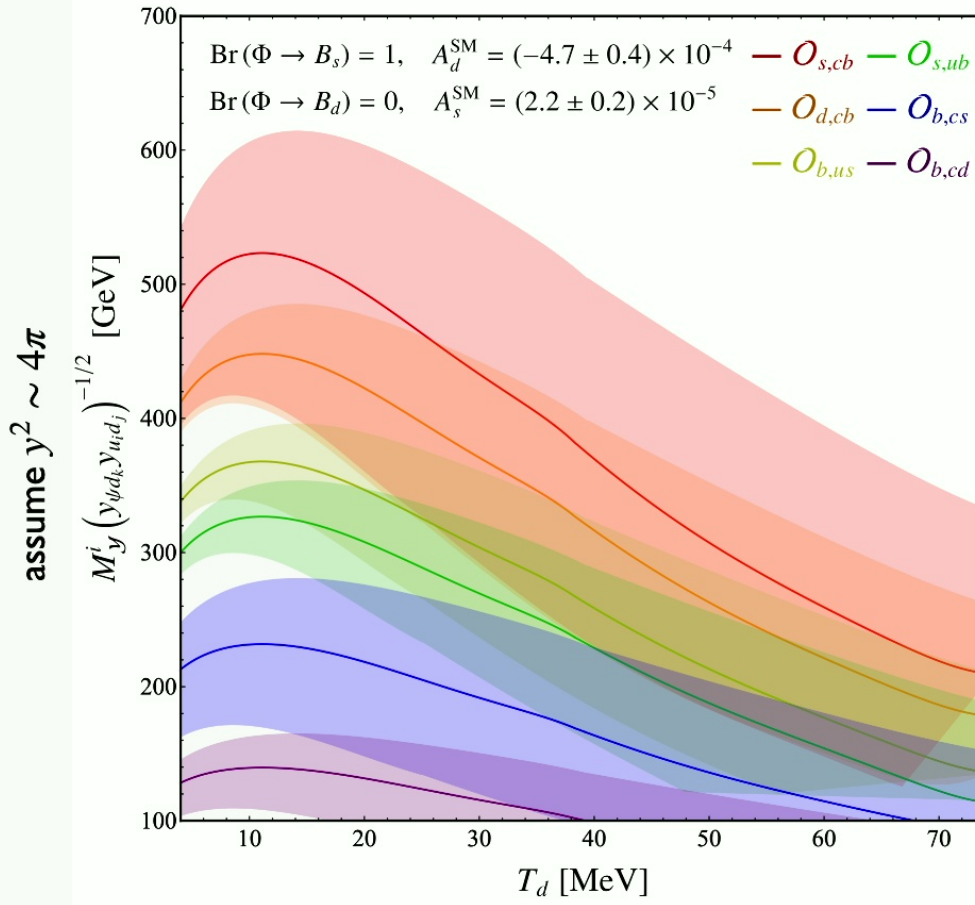
$$Y_{\mathcal{B}} \simeq 5 \times 10^{-5} \sum_{i=d,s} \underbrace{[\text{Br}(B_i^0 \rightarrow \bar{\psi}_{\mathcal{B}} \mathcal{B}_{\text{SM}})]}_{\sim \frac{1}{M_{\mathcal{Y}}^4}} \underbrace{A_{sl}^i}_{\sim 10^{-5} \text{ from the SM}} \alpha_i(T_d)$$

$0 \leq \alpha_I \lesssim 1.5$   
 numerical details from solving the Boltzmann equations involved

not much room for exotic branching fractions!

Observed asymmetry:  $Y_{\mathcal{B}}^{\text{meas}} = \frac{n_B - n_{\bar{B}}}{s} \simeq 8 \times 10^{-11}$

### Successful Baryogenesis



When  $\Phi$  decays to  $B_i^0$

Large BAU wants  
light-ish mediator  $\mathcal{Y}$

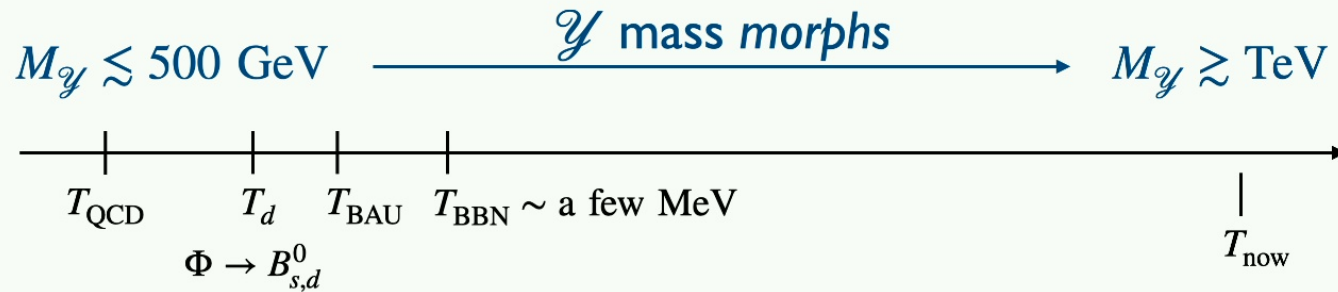
BUT

Today's experiments  
want heavy-ish  $\mathcal{Y}$



What if the  $\mathcal{Y}$  mass  
changed over time?

# The Cosmic Timeline



How does the mass change from 500 GeV to TeV?

at temperatures  $\sim \text{MeV}$ ????

Well... Hmmm... Not easily!

Need some mechanism to change the *morphon* mass

(delayed) phase transition?      some other field triggering a mass/vev change?

An example:

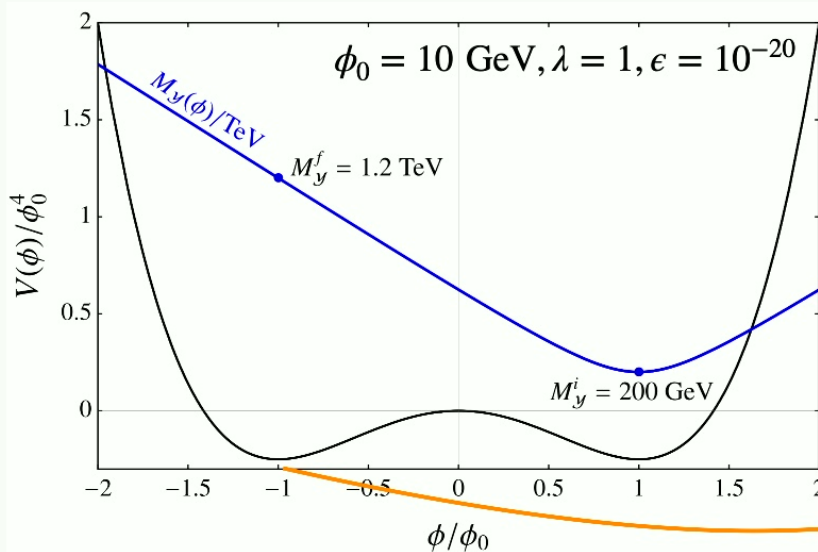
$$V_{\text{scalar}} = m_{\mathcal{Y}}^2 |\mathcal{Y}|^2 + y_{\phi\mathcal{Y}} |\mathcal{Y}|^2 \phi + \frac{1}{2} \lambda_{\phi\mathcal{Y}} |\mathcal{Y}|^2 \phi^2 + \frac{1}{4} \lambda (\phi^2 - \phi_0^2)^2 + \epsilon \phi_0 \phi^3$$

field-dependent *morphon* mass

$$M_{\mathcal{Y}}^2(\phi) = m_{\mathcal{Y}}^2 + y_{\phi\mathcal{Y}} \phi + \frac{1}{2} \lambda_{\phi\mathcal{Y}} \phi^2$$

$$M_{\mathcal{Y}}^i \simeq 100 \text{ GeV}$$

$$M_{\mathcal{Y}}^f \simeq 1000 \text{ GeV}$$

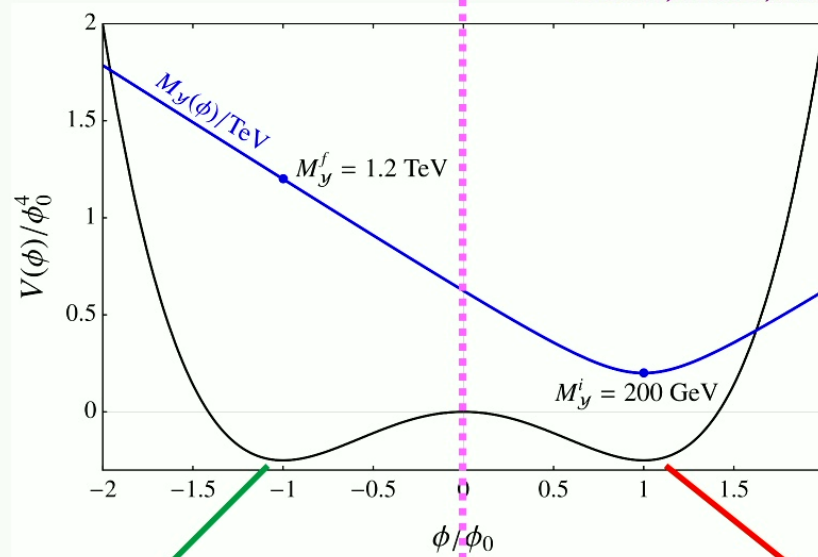


Not easy but can be done!

at the expense of  
some tuning

very slightly lower minimum

Gelmini, Pascoli, Vitagliano, Zhou, arXiv:2009.01903  
 Gelmini, Glesier, Kolb, PRD 39 1558 1989  
 Preskill, Privedi, Wilczek, Wise, Nucl. Phys. B 363 207



Not enough  
BAU  
OK for  $T = 0$

slightly more  
favored!

Domain  
Walls

enough BAU  
too light  
for  $T = 0$

$$\epsilon \lesssim 0.2 \lambda$$

$$\epsilon < \frac{2\sqrt{2}}{3} \sqrt{\frac{8\pi^2 g_*}{90} \frac{T^2}{M_{\text{Pl}}} \frac{\sqrt{\lambda}}{\phi_0}} \Bigg|_{T=T_c}$$

$$\epsilon > \frac{2\sqrt{2}}{3} \sqrt{\frac{8\pi^2 g_*}{90} \frac{T^2}{M_{\text{Pl}}} \frac{\sqrt{\lambda}}{\phi_0}} \Bigg|_{T=T_{\text{BBN}}}$$

$$\epsilon > \left(\frac{4}{3}\right)^3 \frac{4\pi\lambda\phi_0^2}{M_{\text{Pl}}^2}$$

DWs percolate

DWs grow to horizon size

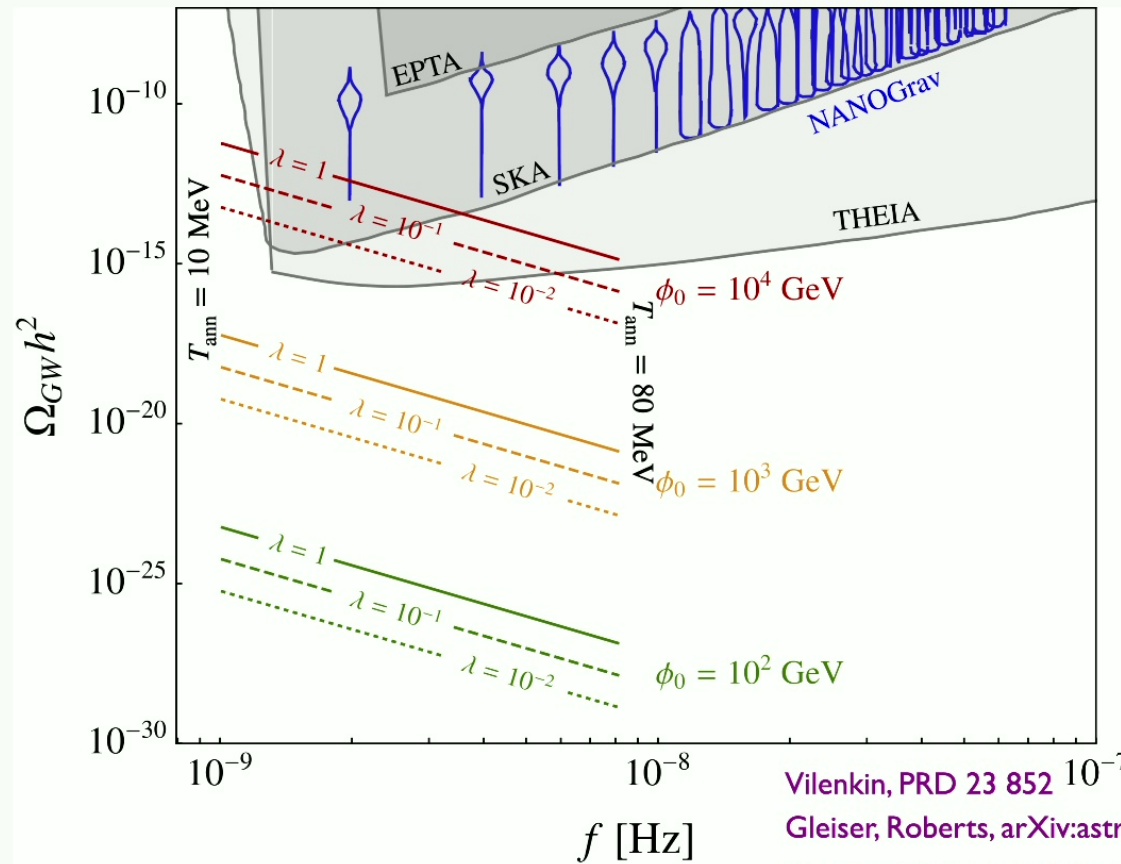
$p_V < p_S$  at  $R \sim 1/2H(T)$   
 vacuum =  $\epsilon\phi_0^4$       due to surface tension =  $\sigma/R$

DWs annihilate at  $T=10$  MeV

$p_V > p_S$  at  $T = 10$  MeV

DWs annihilate before they trigger inflation

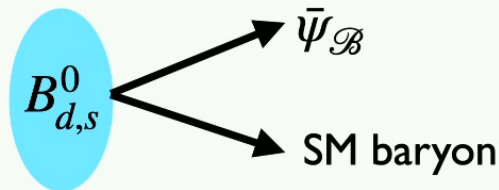
# Gravitational wave signals from domain wall annihilations





## The moral of the story?

SM CP violation *can be* enough for baryogenesis



We considered one way to do this:  
Mesogenesis with a Morphing Mediator

a scalar that changes mass from 100 GeV  $\rightarrow$  TeV  
(at a temp  $\sim$  10 MeV)

Different minima separated  
domain walls

maybe other ways of morphing?