

Title: Lecture - Gravitational Physics, PHYS 636

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Collection/Series: Gravitational Physics (Elective), PHYS 636, January 6 - February 5, 2025

Subject: Cosmology, Strong Gravity

Date: January 09, 2025 - 9:00 AM

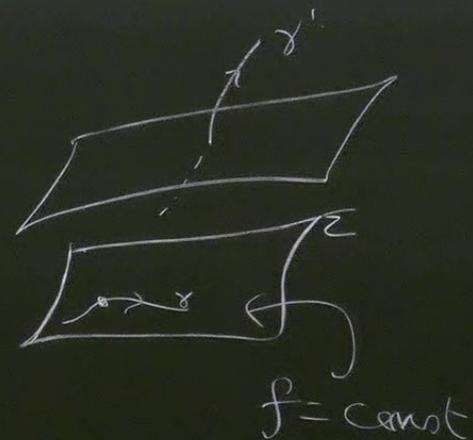
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L3 Forms & Duals

Consider $f \in C^\infty(M)$

If $\gamma \subset \Sigma$ then $\int_\gamma f = 0$

as f is const on Σ , hence γ



Define a covector, df

$$\langle df | I \rangle = If \quad \forall I \in T_p(M)$$

any f

$$d: C^\infty(M) \rightarrow T^*(M)$$

In components $df = \frac{\partial f}{\partial x^m} dx^m$

(Think "grad")

A p-form is a totally antisymmetric rank p covariant tensor. Construct similarly to tensors.

$$\underline{A} \wedge \underline{B} = \underline{A} \otimes \underline{B} - \underline{B} \otimes \underline{A} \leftarrow \text{2-form.}$$

↑
covectors / 1-forms

The wedge product is anti-symmetric.

In components: $(\underline{A}^{(p)} \wedge \underline{B}^{(q)})_{a_1 \dots a_{p+q}} = \frac{(p+q)!}{p! q!} A_{[a_1 \dots a_p} B_{a_{p+1} \dots a_{p+q}]}$

where (eg) $\underline{A}^{(p)} = \frac{1}{p!} A_{a_1 \dots a_p} dx^{a_1} \wedge \dots \wedge dx^{a_p}$

$$\underline{A}^{(p)} \wedge \underline{B}^{(q)} = (-1)^{pq} \underline{B}^{(q)} \wedge \underline{A}^{(p)}$$

ⁿ
covectors / 1-forms

There are only a finite number of forms, $p \leq n$, we denote the bundle of p -forms as $\Lambda^p(M)$.

The n -form is unique up to a factor

2-form.

where (eg) $\underline{A}^{(p)} = \frac{1}{p!} A_{a_1 \dots a_p} dx^{a_1} \wedge \dots \wedge dx^{a_p}$

$$\underline{A}^{(p)} \wedge \underline{B}^{(q)} = (-1)^{pq} \underline{B} \wedge \underline{A}$$

of

$$\epsilon_{a_1 \dots a_n} = \pm 1 \quad (\text{sign of permutation})$$

$$\epsilon_{123} = \epsilon_{231} = \epsilon_{312} = +1$$

$$\epsilon_{132} = \epsilon_{213} = \epsilon_{321} = -1$$

This ϵ is actually a tensor density

$$\underline{\epsilon} = \frac{1}{n!} \epsilon_{\mu_1 \dots \mu_n} dx^{\mu_1} \wedge \dots \wedge dx^{\mu_n}$$

$$= \frac{1}{n!} \underbrace{\det\left(\frac{\partial x}{\partial x'}\right)} \epsilon_{\mu_1 \dots \mu_n} dx^{\mu_1} \wedge \dots \wedge dx^{\mu_n}$$

In 4D have 4 lin. indep 1-forms,
6 2-forms, 4 3-forms & 1 4-form

? Can we relate 1 and 3-forms?

In 3D 3 1-forms
3 2-forms. $\left. \begin{array}{l} \rightarrow \\ \leftarrow \end{array} \right\} \text{relate via} \\ \times \text{ product.}$

Recall in 3D.

$$(\underline{a} \times \underline{b})_i = \epsilon_{ijk} a_j b_k$$

With a metric, note that

$$\det g \rightarrow \det g' = \det \left(g_{mn} \frac{\partial x^m}{\partial x'^n} \right)$$

$$\det g \cdot \det \left(\frac{\partial x}{\partial x'} \right)^2$$

Recall in 3D.

$$(\underline{a} \times \underline{b})_i = \epsilon_{ijk} a_j b_k$$

With a metric, note that

$$\det g \rightarrow \det g' = \det \left(g_{\mu\nu} \frac{\partial x^\mu}{\partial x'^\alpha} \frac{\partial x^\nu}{\partial x'^\beta} \right)$$

$$\det g \cdot \det \left(\frac{\partial x}{\partial x'} \right)^2$$

Hence

$$\sqrt{|g'|} \epsilon_{\mu_1 \dots \mu_n} = \sqrt{|g|} \epsilon_{\mu_1 \dots \mu_n} \frac{\partial x^{\mu_1}}{\partial x'^{\mu_1}} \dots \frac{\partial x^{\mu_n}}{\partial x'^{\mu_n}}$$

Recall in 3D.

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Hence

$$\sqrt{|g|} \epsilon_{\mu_1 \dots \mu_n} = \sqrt{|g'|} \epsilon_{\mu'_1 \dots \mu'_n} \frac{\partial x^{\mu'_1}}{\partial x^{\mu_1}} \dots \frac{\partial x^{\mu'_n}}{\partial x^{\mu_n}}$$

is covariant

Write as $\epsilon_{\mu_1 \dots \mu_n}$

We can then define a map

$$* \Lambda^p \rightarrow \Lambda^{n-p}$$

$$\underline{A} \mapsto * \underline{A}$$

$$\text{via } (*A)_{a_1 \dots a_{n-p}} = \frac{1}{p!} \sum_{b_1 \dots b_p} \epsilon_{a_1 \dots a_{n-p} b_1 \dots b_p} A_{b_1 \dots b_p}$$

eg \mathbb{R}^3
 $dx^1 \leftrightarrow dx^2 \wedge dx^3$

This is the Hodge dual

Can now extend d
to forms

$$d: \Lambda^p \rightarrow \Lambda^{p+1}$$

br-bp

d reduces to df on fns

pseudo-Leibnizian

$$d(\underline{A} \wedge \underline{B}) = d\underline{A} \wedge \underline{B} + (-1)^p \underline{A} \wedge d\underline{B}$$

eg \mathbb{R}^3
 $dx^1 \leftrightarrow dx^2 \wedge dx^3$

This is the Hodge dual

Can now extend d
to forms

$$d: \Lambda^p \rightarrow \Lambda^{p+1}$$

b1-bp

d reduces to df on fns

pseudo-Leibnizian $d(\underline{A}^{(p)} \wedge \underline{B}^{(q)}) = d\underline{A}^{(p)} \wedge \underline{B}^{(q)}$

$$(-)^p \underline{A} \wedge d\underline{B}$$

$$\wedge d^2 = 0$$

in components

$$(dA)_{a_1 \dots a_{p+1}} = \frac{(p+1)!}{p!} \partial_{[a_1} A_{a_2 \dots a_{p+1}]}$$

$$(\dots) d_1 \dots d_{n-p} = \frac{1}{p!} \epsilon_{a_1 \dots a_p} \dots$$

e.g. Electromagnetism

A_μ 1-form gauge field

$$\underline{A} = A_\mu dx^\mu$$

$$\underline{F} = d\underline{A}$$

F_μ

$$F_{\mu\nu} = \frac{2!}{1!} \partial_{[\mu} A_{\nu]} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}$$

If we add $\underline{d}f$ to \underline{A}

$$\underline{d}(\underline{A} + \underline{d}f) = \underline{d}\underline{A} + \cancel{\underline{d}^2 f}$$
$$= \underline{F}$$

correctors / 1-forms

e.g. $B_{\mu\nu}$ RR 2-form

A_μ couples to point charges

$B_{\mu\nu}$ " " line charges

$$\underline{H} = \underline{d} \underline{B} \quad (\rightarrow) \quad \partial_\mu B_{\nu\lambda} + \partial_\nu B_{\lambda\mu} + \partial_\lambda B_{\mu\nu}$$

$$\underline{A} \wedge \underline{B} = (-1)^{pq} \underline{B} \wedge \underline{A}$$

The Hodge dual gives rise to a dual derivation.

$$\delta : \Lambda^p \rightarrow \Lambda^{p-1} \quad \delta = *d*$$

$$A_{a_1 \dots a_p} \rightarrow (-1)^p \nabla^{a_1} A_{a_1 \dots a_p}$$

Free Maxwell: $dF = 0$ (identity) $\delta F = 0$

$$\underline{H} = d\underline{B} \Leftrightarrow \partial_\mu B_{\nu\lambda} + \partial_\nu B_{\lambda\mu} + \partial_\lambda B_{\mu\nu}$$

Free Maxwell: $dF=0$ (iden

Useful identity

$$\langle d\underline{\omega} | \underline{u}, \underline{v} \rangle = \underline{u}(\langle \underline{\omega} | \underline{v} \rangle) - \underline{v}(\langle \underline{\omega} | \underline{u} \rangle) - \langle \underline{\omega} | [\underline{u}, \underline{v}] \rangle$$