

Title: Probes of cosmic inflation: from the CMB to quantum systems

Speakers: Emilie Hertig

Collection/Series: Cosmology and Gravitation

Subject: Cosmology

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Abstract:

Polarization anisotropies of the cosmic microwave background (CMB) encode a wealth of information on fundamental physics. In the coming decade, a new generation of instruments starting with the Simons Observatory (SO) will either detect or tightly constrain the amplitude of B-mode patterns produced by inflationary gravitational waves. The first part of my talk will focus on techniques developed to mitigate secondary B-modes induced by Galactic foregrounds and weak gravitational lensing, in order to extract the primordial signal with optimal precision. I will present resulting performance forecasts for SO, as well as initial efforts to apply these methods to the new data currently being collected.

At the other end of the scale, complementary approaches based on numerical simulations and cold-atom analogue experiments are emerging as a way of probing early-Universe quantum dynamics in real time. The second part of my talk will introduce ongoing work on lattice simulations of false vacuum decay, aiming to understand their range of validity by investigating renormalization effects. Finally, I will outline future avenues for combining cosmological and quantum probes of inflation, exploiting the deep connection between the smallest and largest scales to gain a new perspective on the early Universe.

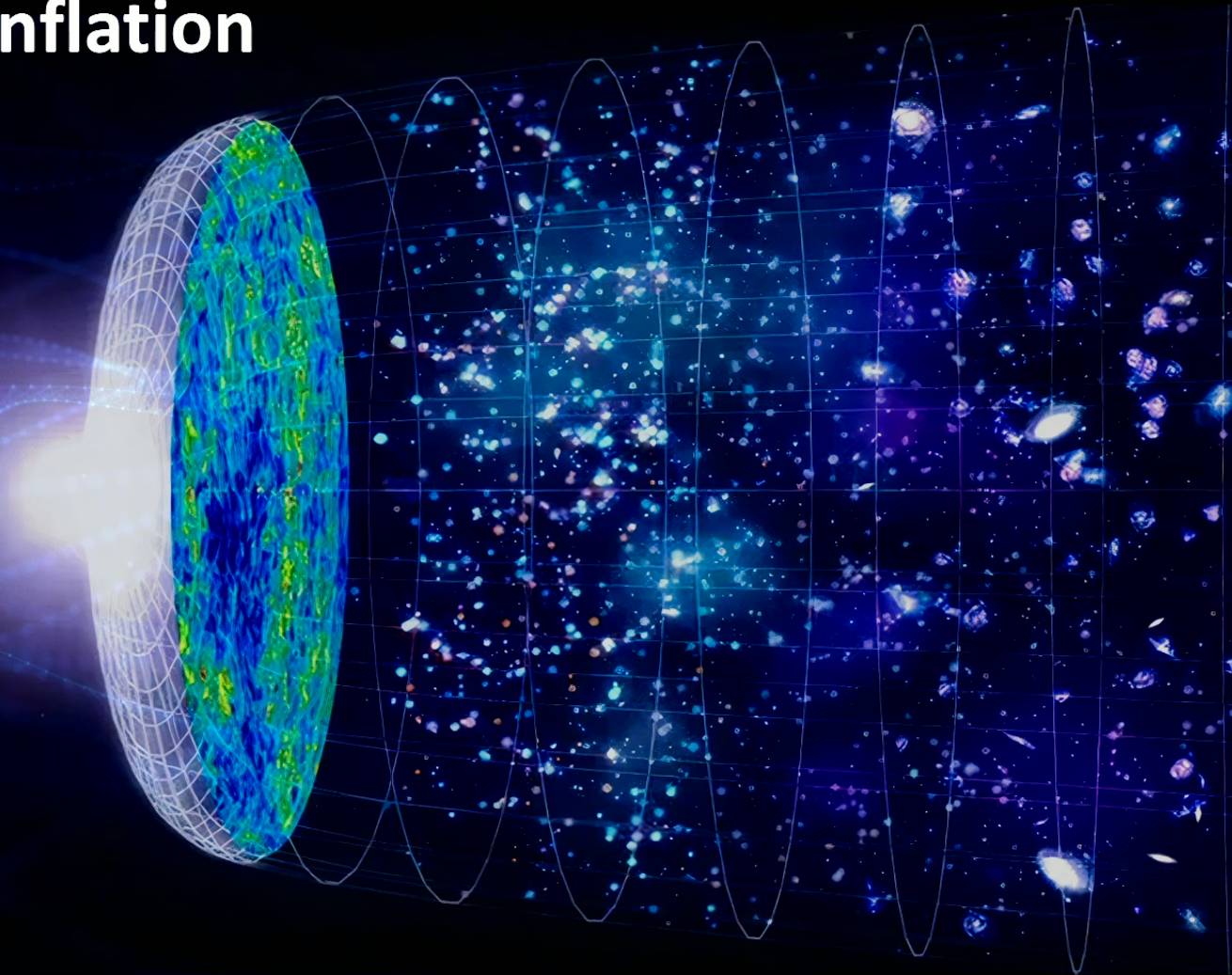
Probes of Cosmic Inflation

*From the CMB to
quantum systems*

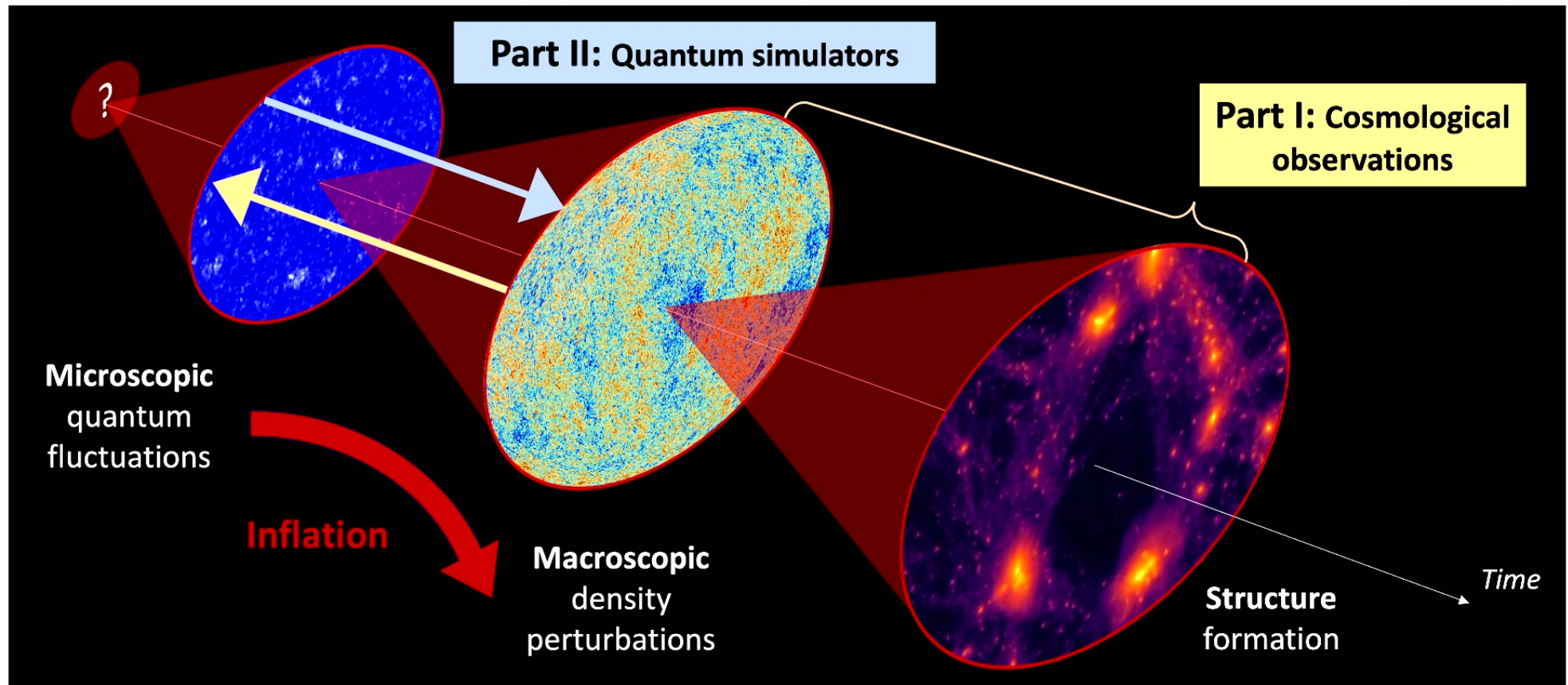
Emilie Hertig

With A. Challinor, B. Sherwin,
H. Peiris *et al.*

Perimeter Institute, 12.12.2024

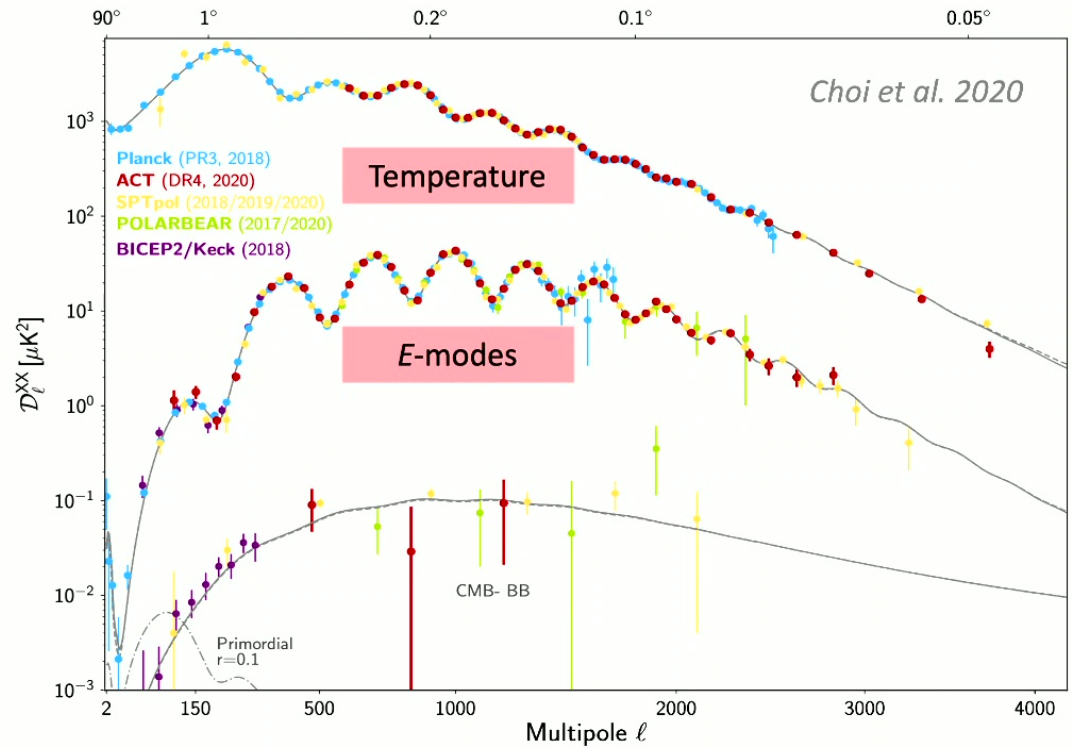
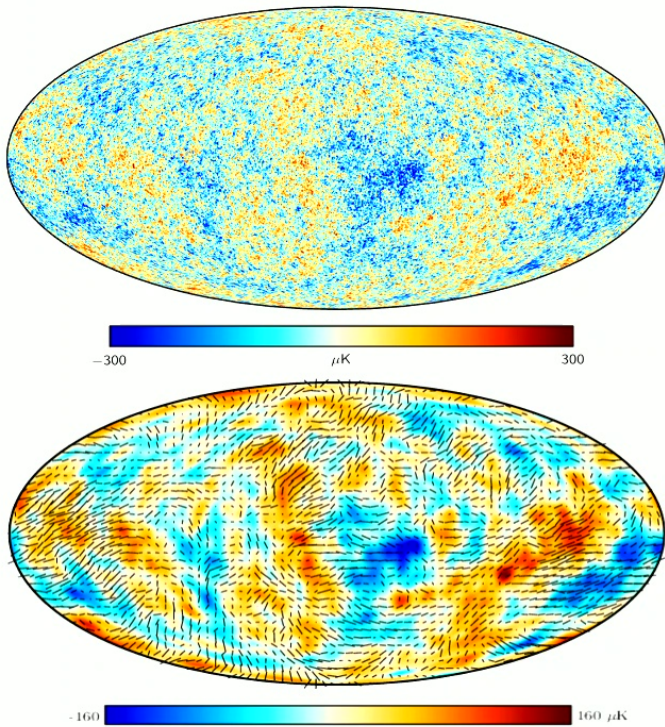


Inflation: a link between cosmic scales

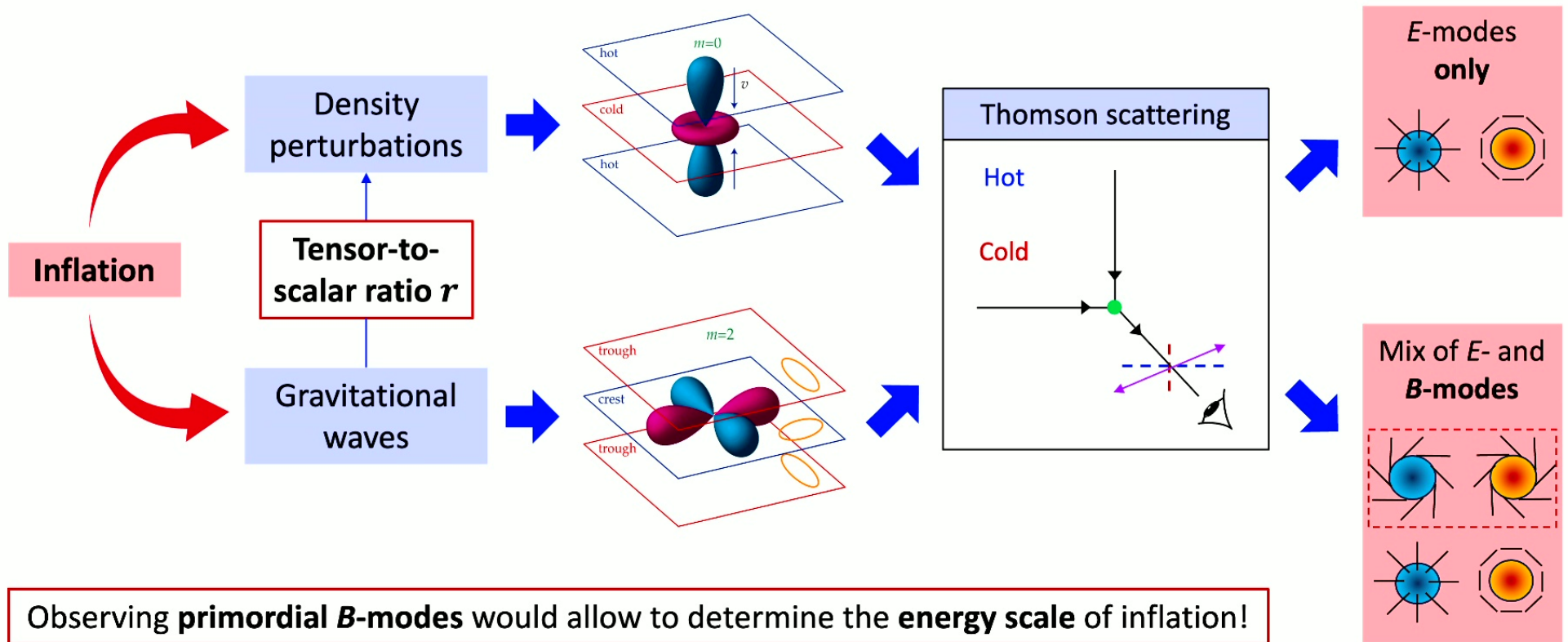


Part I – CMB constraints

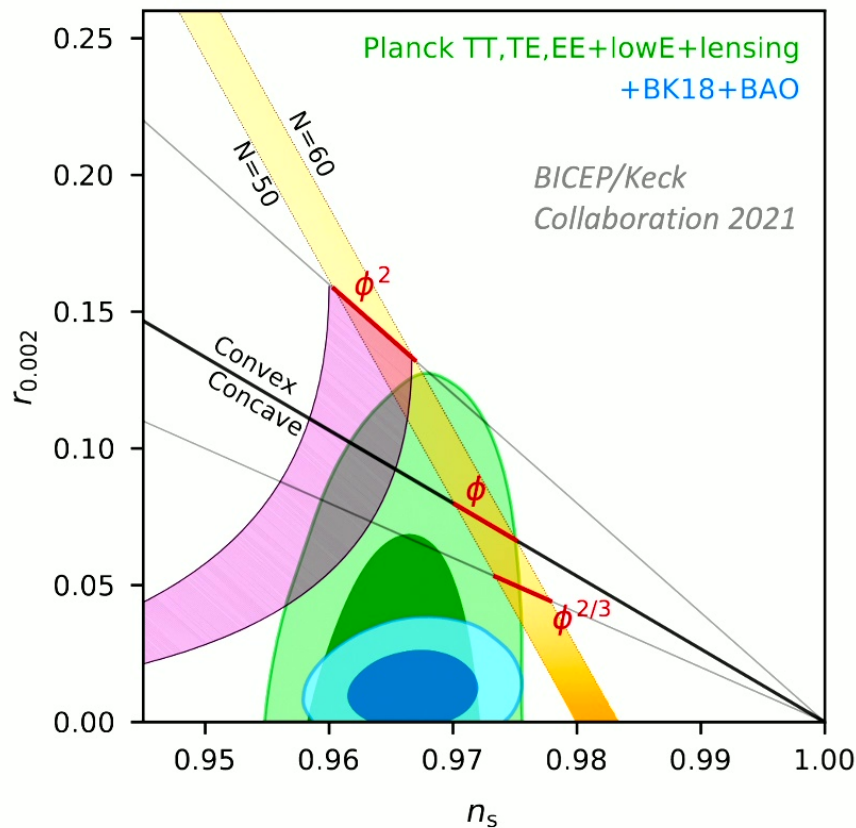
Inflationary prediction: nearly scale-invariant spectrum of primordial density fluctuations



From inflation to CMB polarization



Recent constraints on primordial GW



Current upper bound (95% confidence):

$$r < 0.036$$

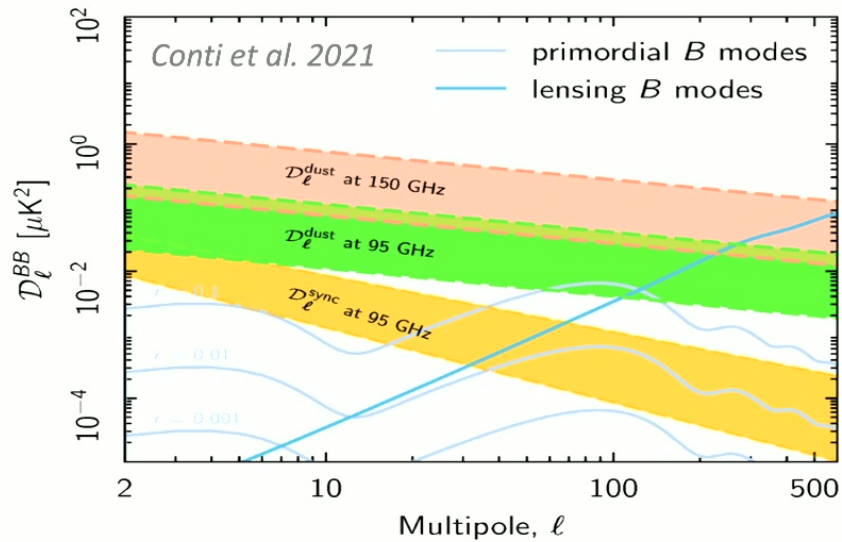
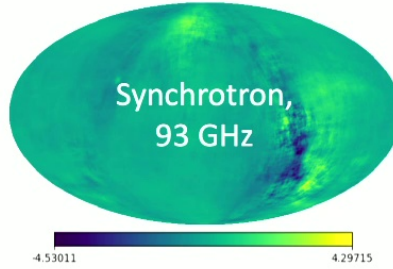
Corresponding statistical uncertainty:

$$\sigma(r) = 0.009$$

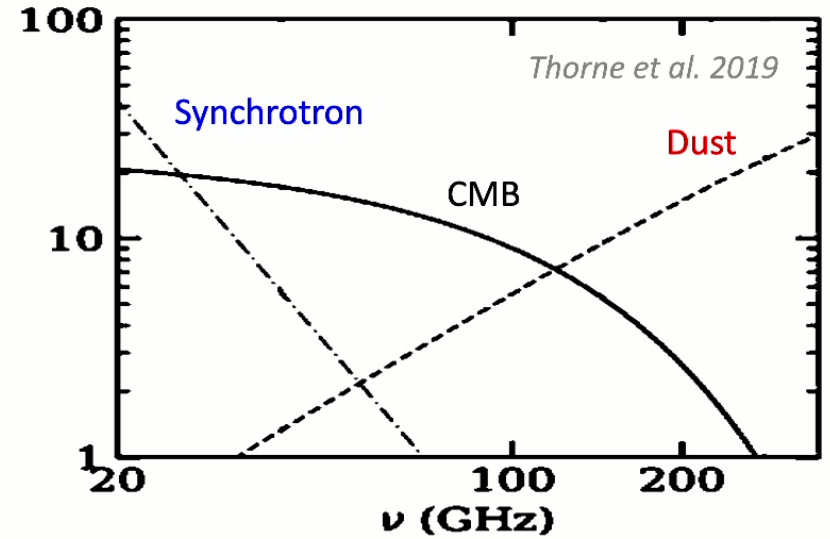
- Some of the **simplest** inflation models now **excluded**, e.g. natural inflation (purple) and monomial potentials (orange)

Challenge 1: Galactic foregrounds

- Dust thermal emission
- Synchrotron radiation from cosmic ray electrons



Rescaled spectral energy distributions (SEDs)

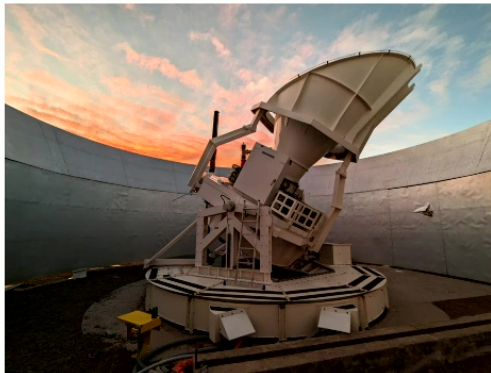
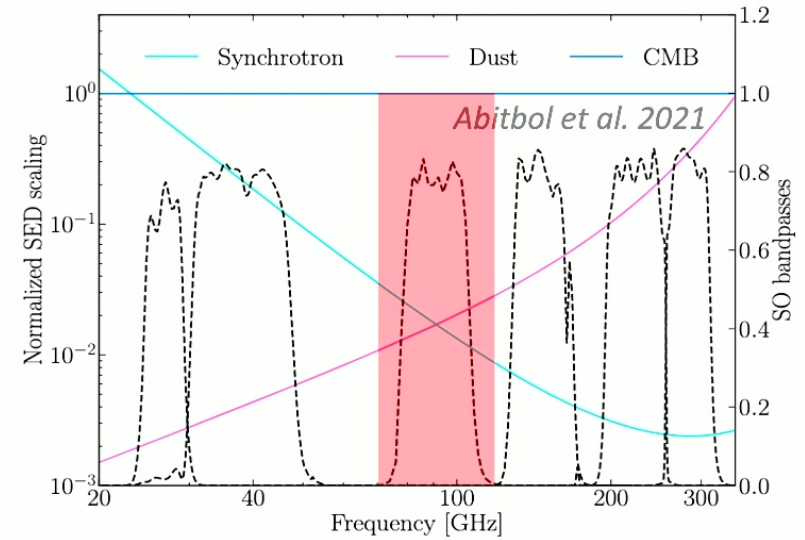
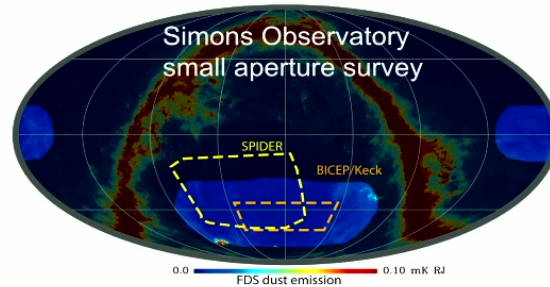


Foregrounds can be distinguished from the CMB by multifrequency observations

The Simons Observatory (SO): SATs

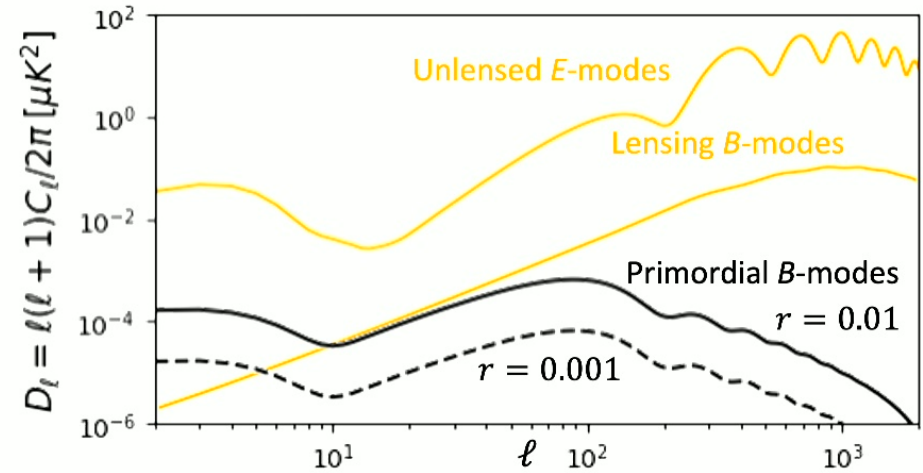
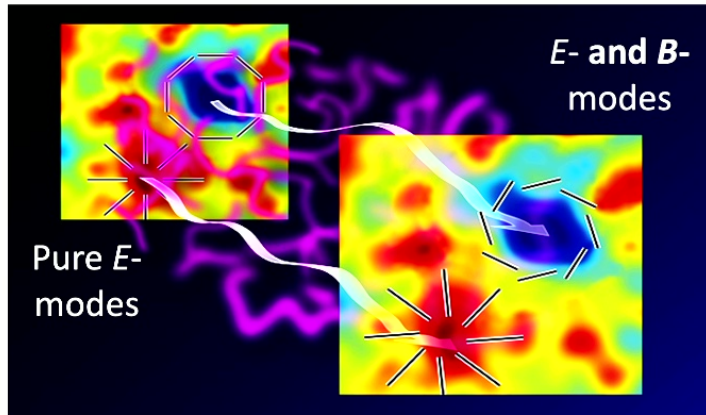
Small Aperture Telescopes

- Targeting faint, **large-scale B-modes** over approx. **10%** of the sky
- **Operational!**



- **6 frequency bands** for component separation
- Most sensitive one at **93 GHz**, where foreground intensity is the lowest

Challenge 2: weak lensing



$$B_l^{\text{lens}} = \sum_{l', L} f(l, l', L) E_{l'}^* \kappa_L^*$$

We need to build a model (**template**) of the lensing **B-modes** from measurements of E and κ

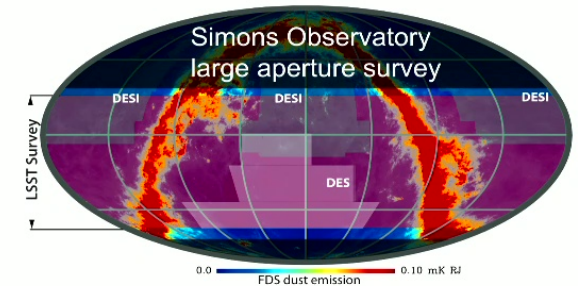
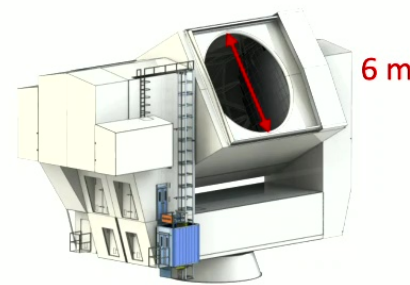
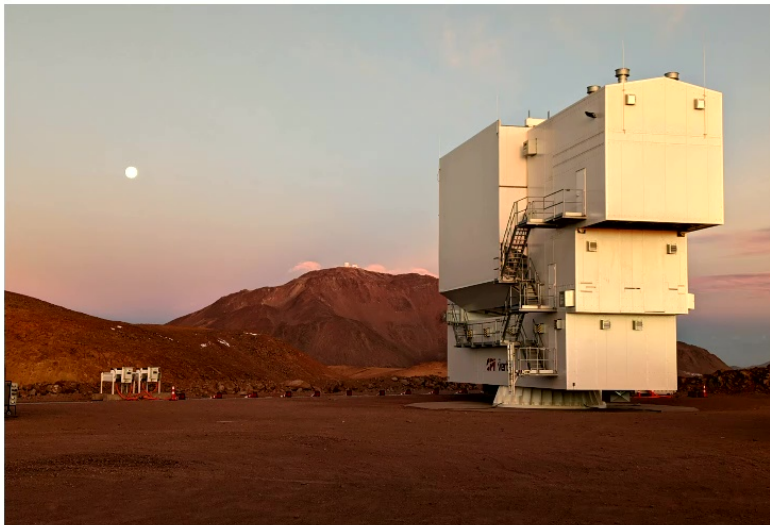
Variance on r from Fisher information (Gaussian likelihood):

$$\sigma^{-2}(r) \Big|_{r=0} = \sum_l \frac{2l+1}{2} \left(\frac{\partial_r C_l}{C_l^{\text{lens}} + N_l} \right)^2$$

The Simons Observatory (SO): LAT

Large Aperture Telescope

- Still under construction, will **start operations in 2025**

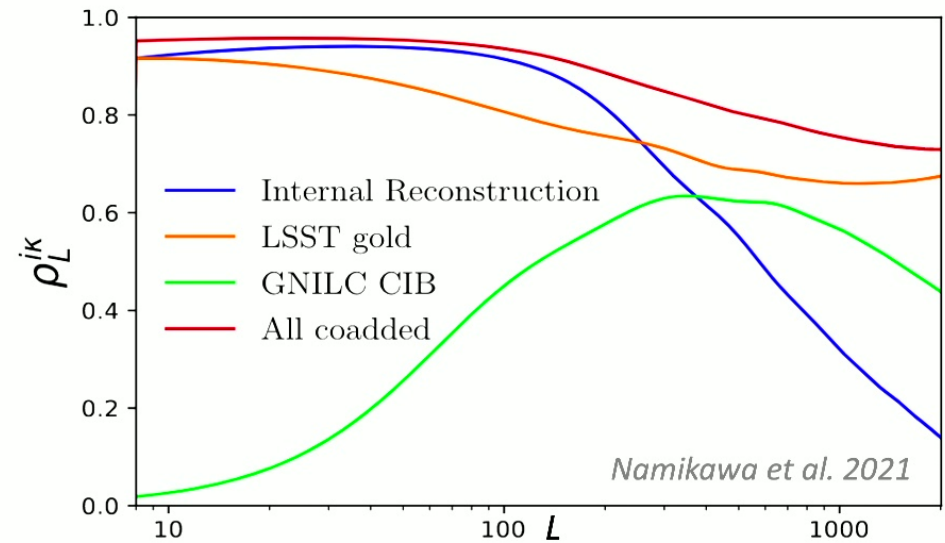


- Will provide **high-resolution** temperature and **E-mode** measurements needed for **lensing template**
- Will observe a larger area (approx. **40% of the sky**), **overlapping** with several galaxy surveys

Multitracer lensing reconstruction

- **Quadratic estimators** using data from the LAT
 - Off-diagonal correlations between pairs of lensed fields (TT , TE , EE and EB)
 - Accurate on large scales, noisy at high multipoles
- Complement with cosmic infrared background (**CIB**) and **galaxy survey data**

The **optimal convergence estimator** is a **weighted sum** of internal and external tracers

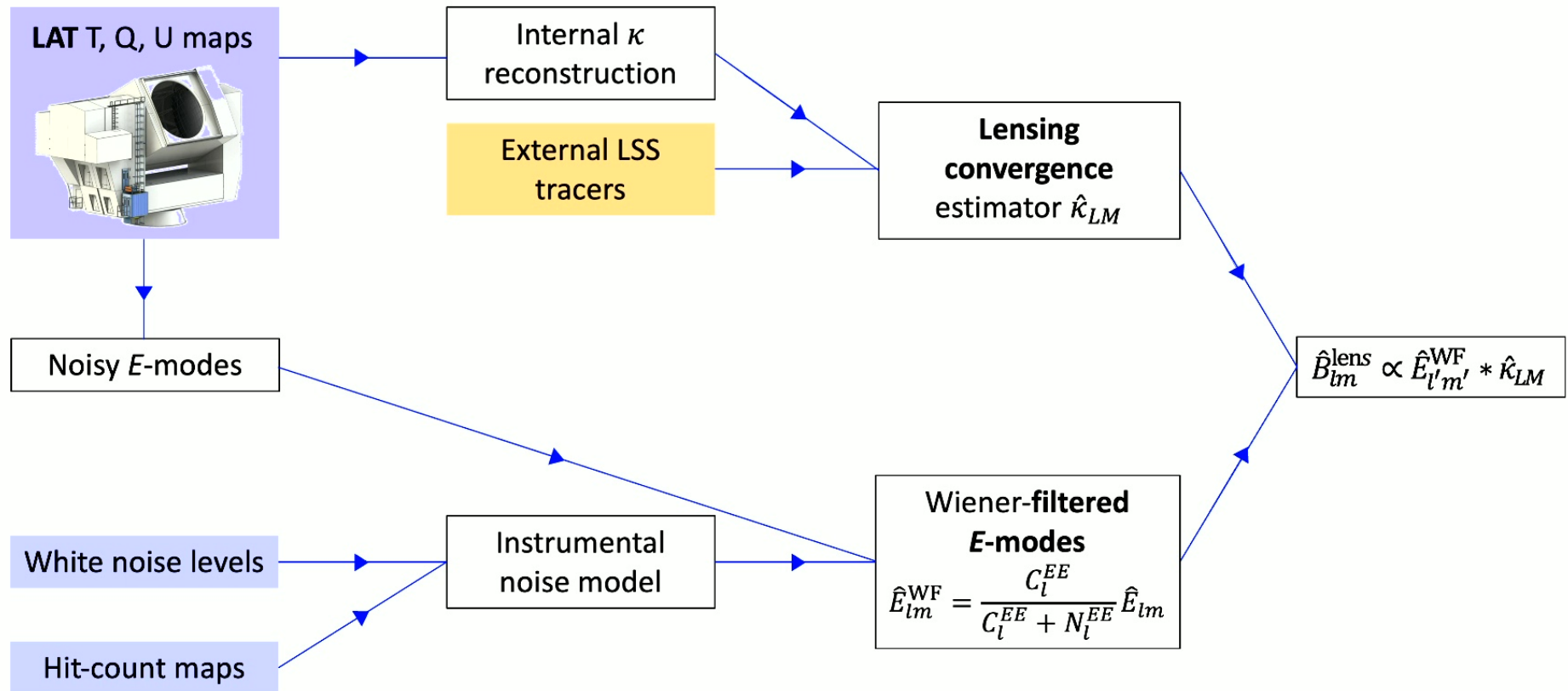


$$\hat{\kappa}_{LM}^{\text{comb}} = \sum_i c_i \hat{\kappa}_{LM}^i$$

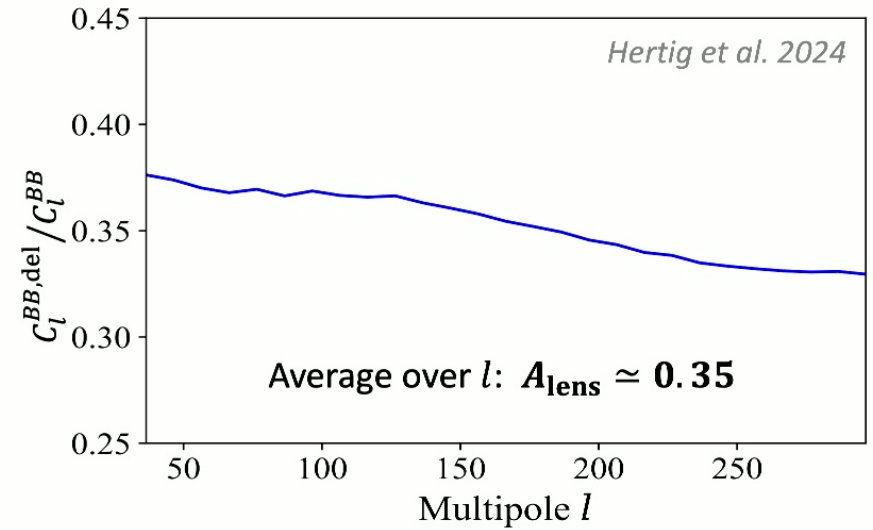
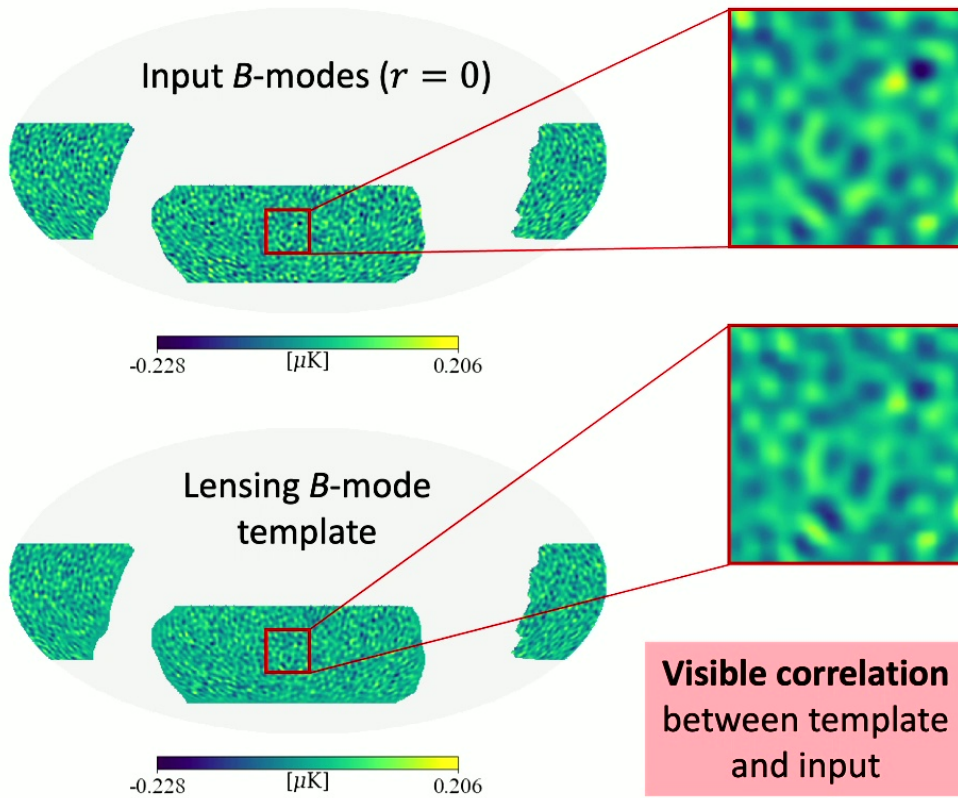
$$c_i = \sum_j (\rho^{-1})_{ij} \rho_{j\kappa} \sqrt{C_l^{\kappa\kappa} / C_l^{ij}}$$

Maximizing $\rho^{\hat{\kappa}_{\text{comb}}\kappa}$

Template construction pipeline

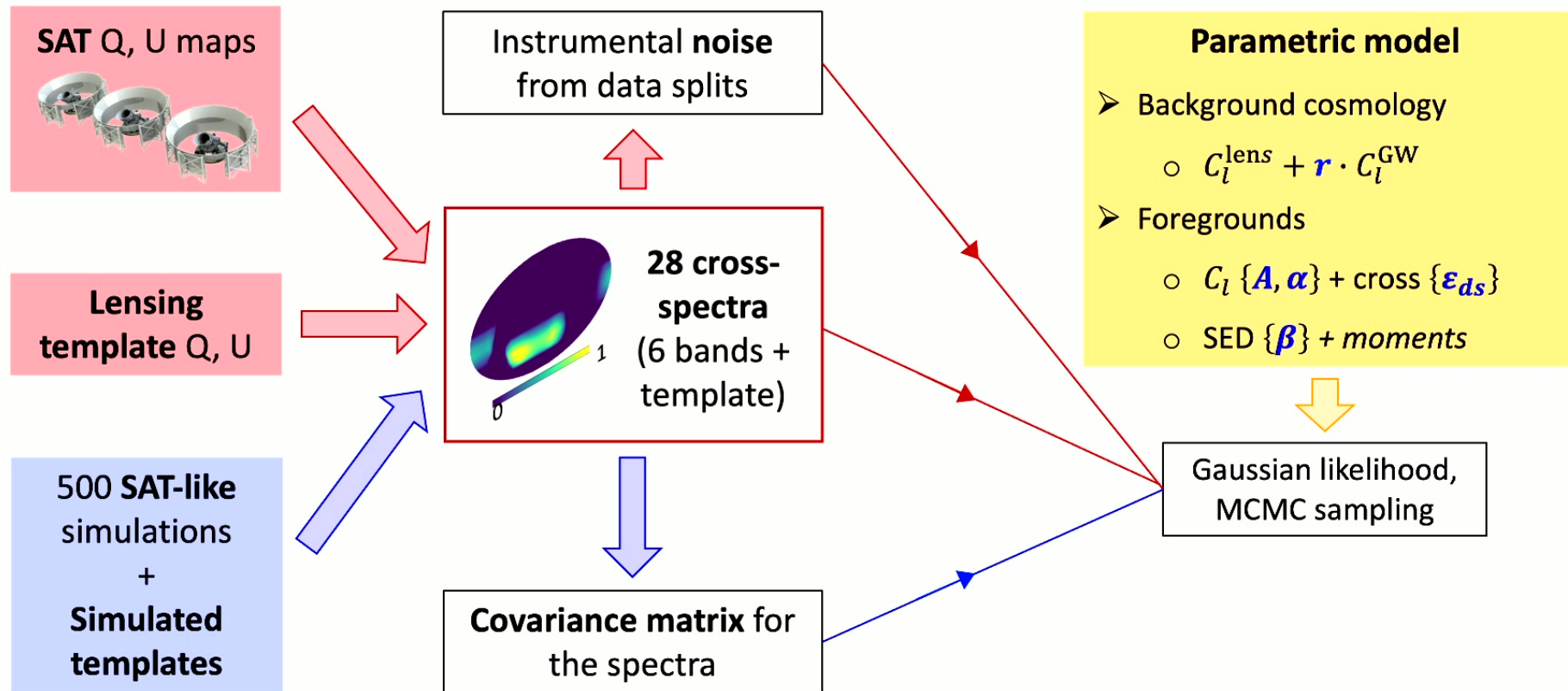


Forecasted delensing efficiency



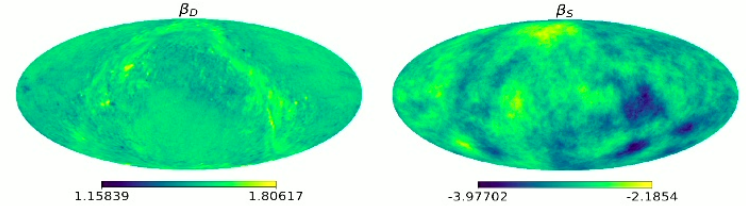
65% delensing efficiency achieved with LAT-like simulations at goal noise level

Component separation pipeline



Foreground model

➤ Power spectra modelled as **power laws** $C_l^{d/s} = A_{d/s} \left(\frac{l}{l_0}\right)^{\alpha_{d/s}-2}$



➤ Spectral energy distributions (SEDs):

- Dust: $\bar{S}_\nu^d \propto \nu^{\beta_d} B_\nu(T_d)$
- Synchrotron: $\bar{S}_\nu^s \propto \nu^{\beta_s}$

Foreground SEDs vary in space due to inhomogeneous dust temperature / grain types and fluctuations in the cosmic ray energy distribution.

➤ Leading order term (sky average):

$$C_{l,0}^{\nu\nu'} = \bar{S}_\nu^d \bar{S}_{\nu'}^d C_l^d + \bar{S}_\nu^s \bar{S}_{\nu'}^s C_l^s + (\bar{S}_\nu^d \bar{S}_{\nu'}^s + \bar{S}_\nu^s \bar{S}_{\nu'}^d) \epsilon_{ds} \sqrt{C_l^d C_l^s}$$

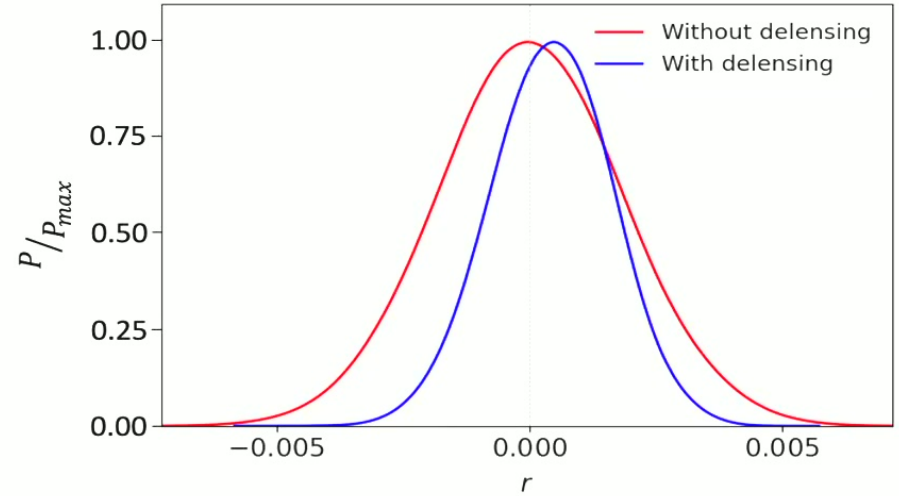
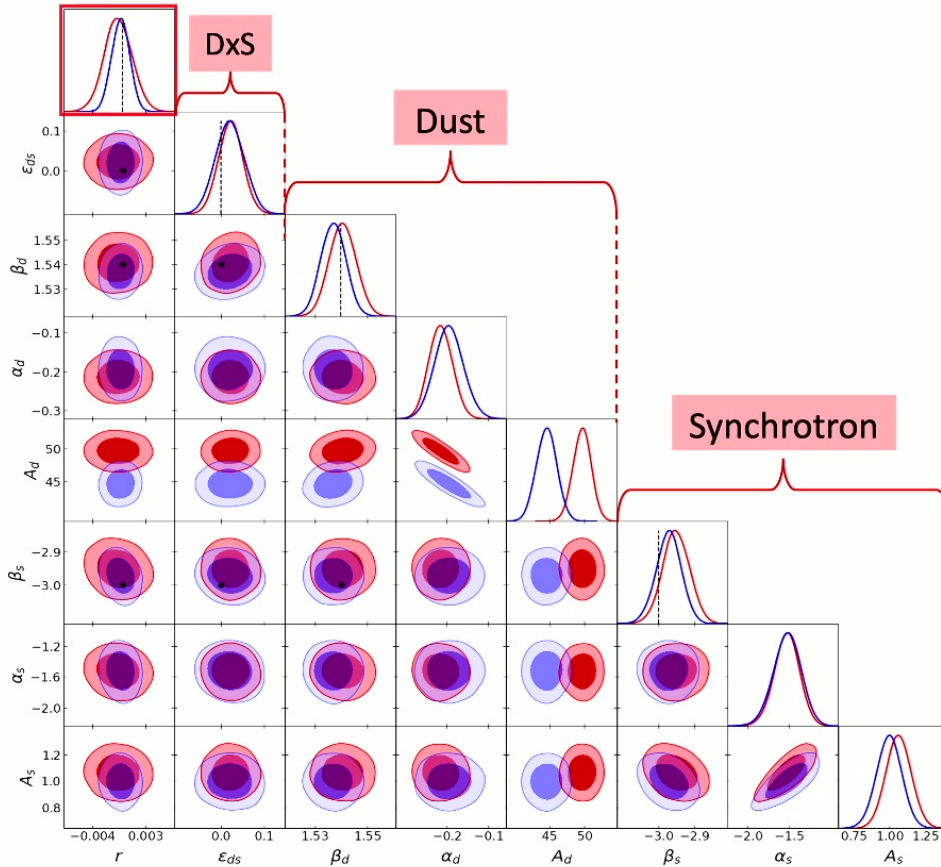
Second order expansion of SEDs:

$$S_\nu(\beta(\hat{n})) = S_\nu(\bar{\beta}) + \partial_\beta S_\nu \cdot \delta\beta(\hat{n}) + \frac{1}{2} \partial_\beta^2 S_\nu \cdot (\delta\beta(\hat{n}))^2$$

« Moment expansion » of foreground power spectra:

extra terms containing $C_l^{\beta\beta} = B_{d/s}(l/l_0)^{\gamma_{d/s}}$ *Chluba et al. 2017*

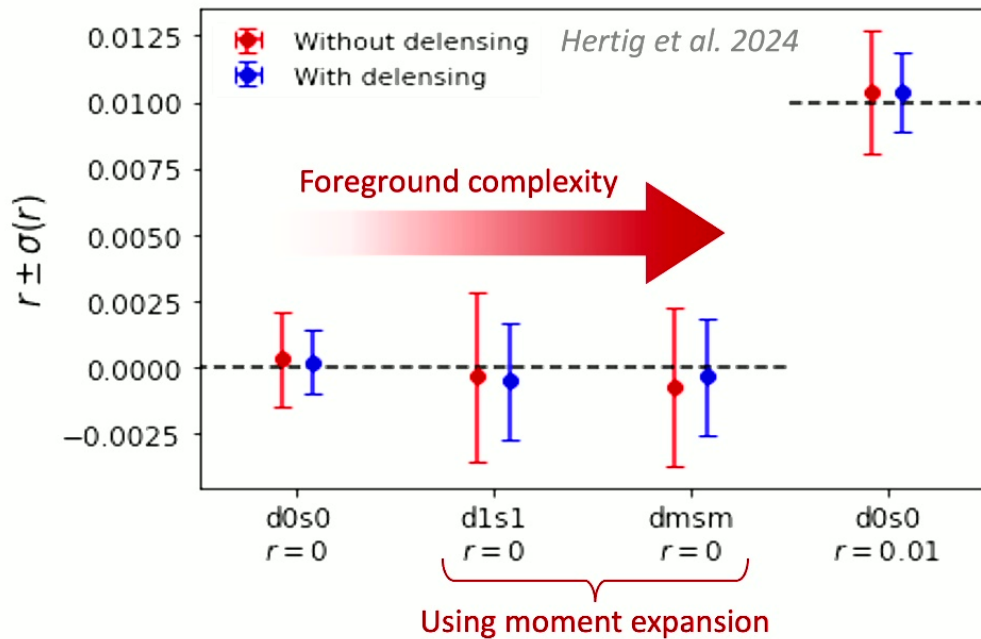
Parameter inference example



	$\sigma(r)$
Without delensing	$1.9 \cdot 10^{-3}$
With delensing	$1.2 \cdot 10^{-3}$
$A_{\text{lens}} = 0.35$ input	$1.2 \cdot 10^{-3}$

Performance forecasts for SO Nominal

Mean of best-fit r and MCMC standard deviations for 100 realizations



Main takeaways

- **No significant bias** within statistical uncertainties
- Error bars are reduced by **~30% to 40%** depending on foreground complexity
- **Nonzero r** successfully detected

SO's target precision $\sigma(r)|_{r=0} = 0.003$ is achieved for all models after delensing

Towards application to real data

Goal: build a **lensing template from external tracers** for early SO analysis

ACT DR6 lensing reconstruction

$$W^{\kappa}(z) = \frac{3}{2H(z)} \Omega_m H_0^2 (1+z) \chi(z) \left(\frac{\chi_* - \chi(z)}{\chi_*} \right)$$

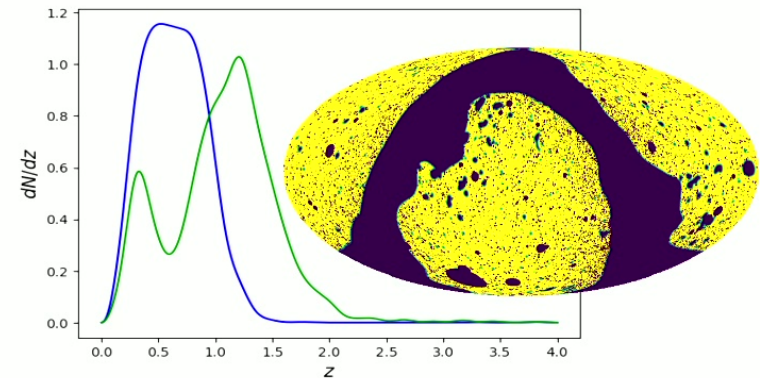
- Minimum-variance combination of TT, TE, EE and EB **quadratic estimators**



unWISE blue and **green** galaxy samples

$$W^g(z) = \frac{b(z) dN/dz}{\int dz' (dN/dz')}$$

- Fit **linear bias** model to $C_l^{\kappa g}$



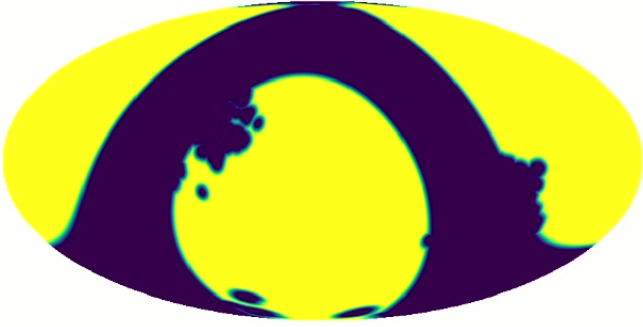
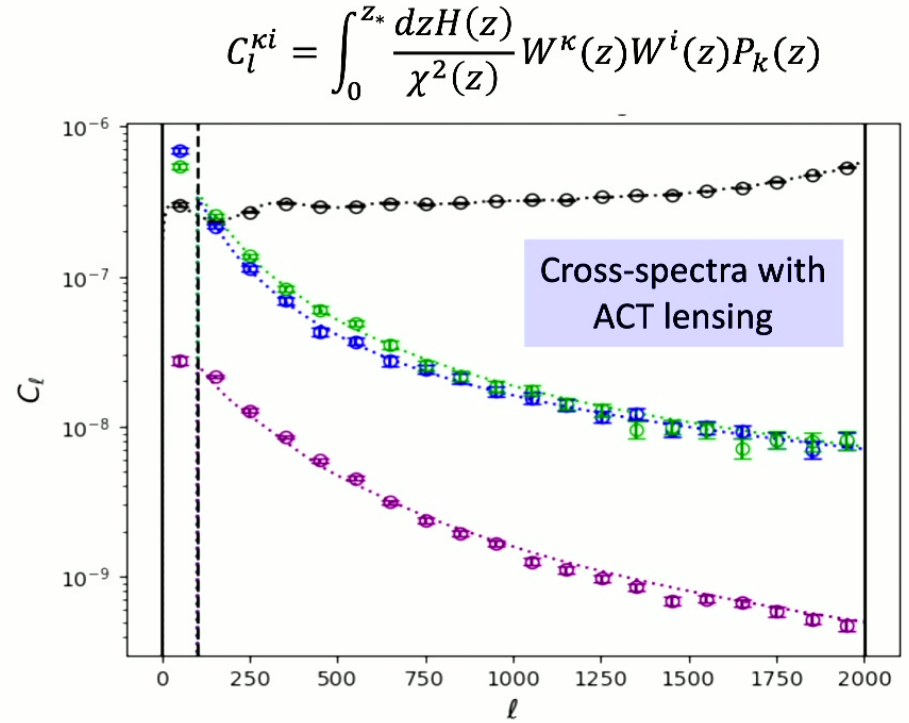
Towards application to real data

Goal: build a lensing template from external tracers for early SO analysis

Planck GNILC CIB

$$W^I(z) = b_c \frac{\chi^2(z)}{H(z)(1+z)^2} e^{-\frac{(z-z_c)^2}{2\sigma_z^2}} f_{\nu(1+z)}$$

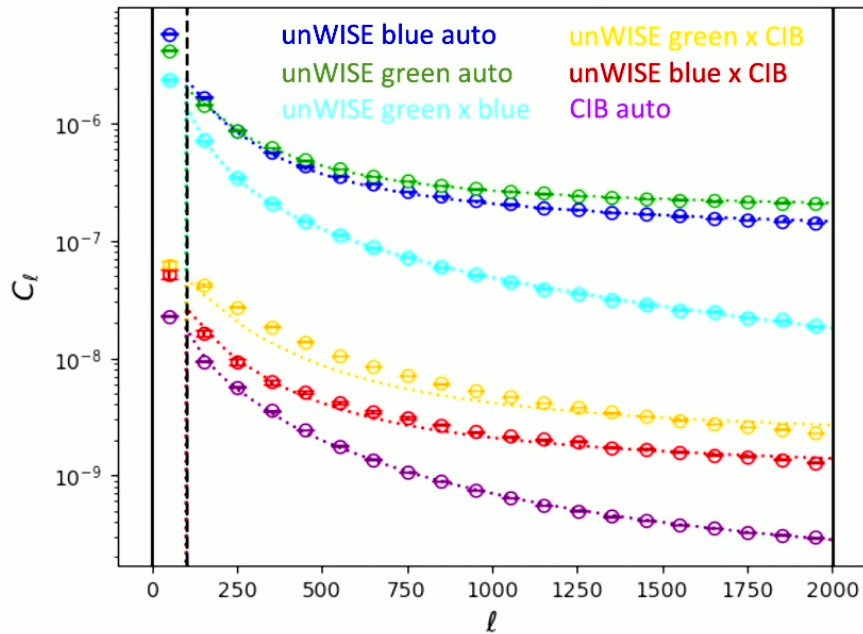
➤ Fit amplitude b_c to C_l^{KI}

Combined lensing estimator

Cross-spectra between external tracers

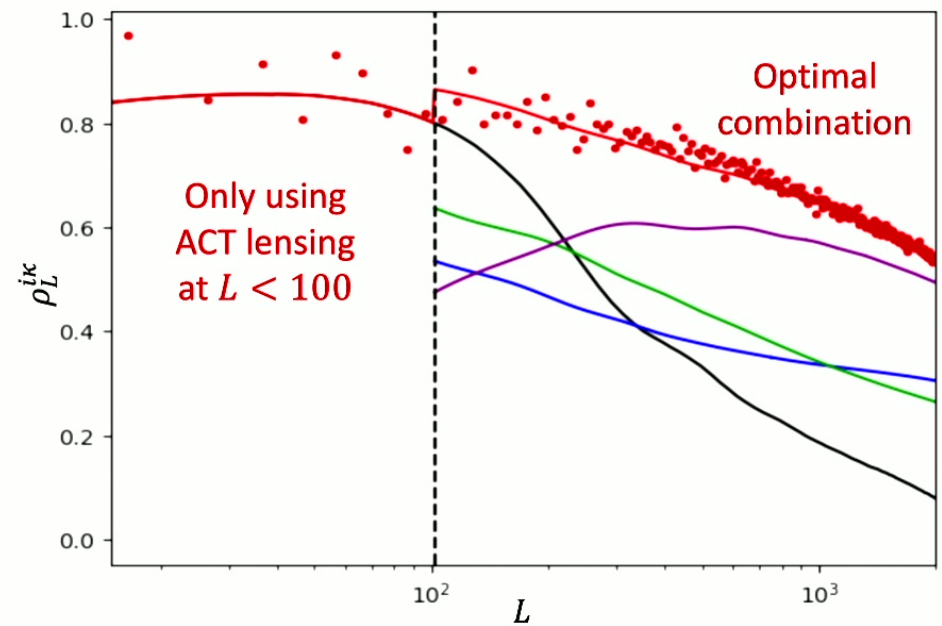
- Include shot noise and dust contamination
- Too simplistic for C_l between galaxy samples



Correlation coefficients with true κ

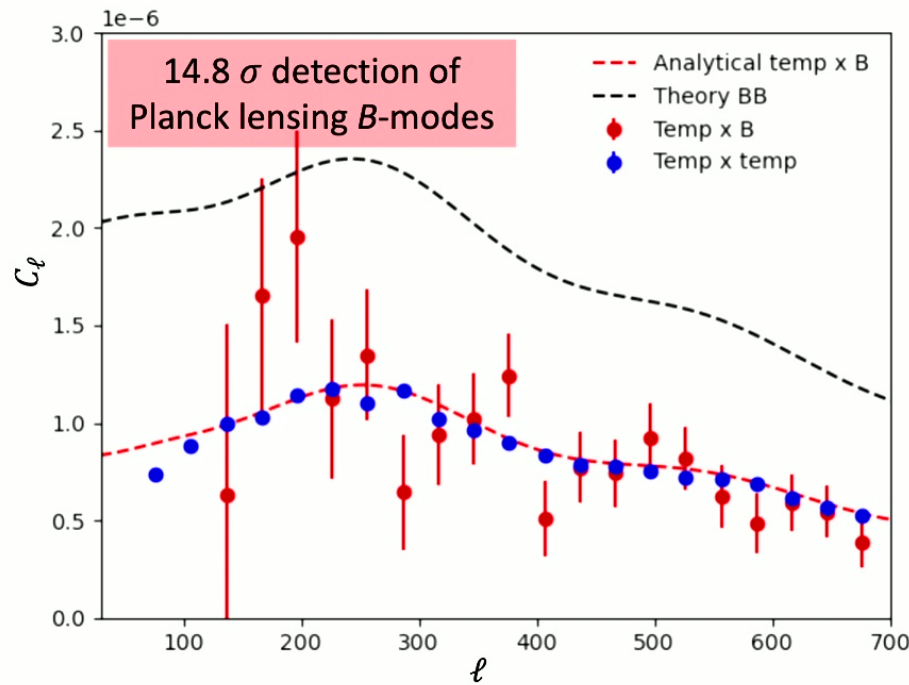
- For combined tracer, estimated using

$$\rho_L^{\kappa \hat{\kappa}_{\text{comb}}} = \rho_L^{\hat{\kappa}_{\text{comb}} \hat{\kappa}_{\text{comb}}}$$

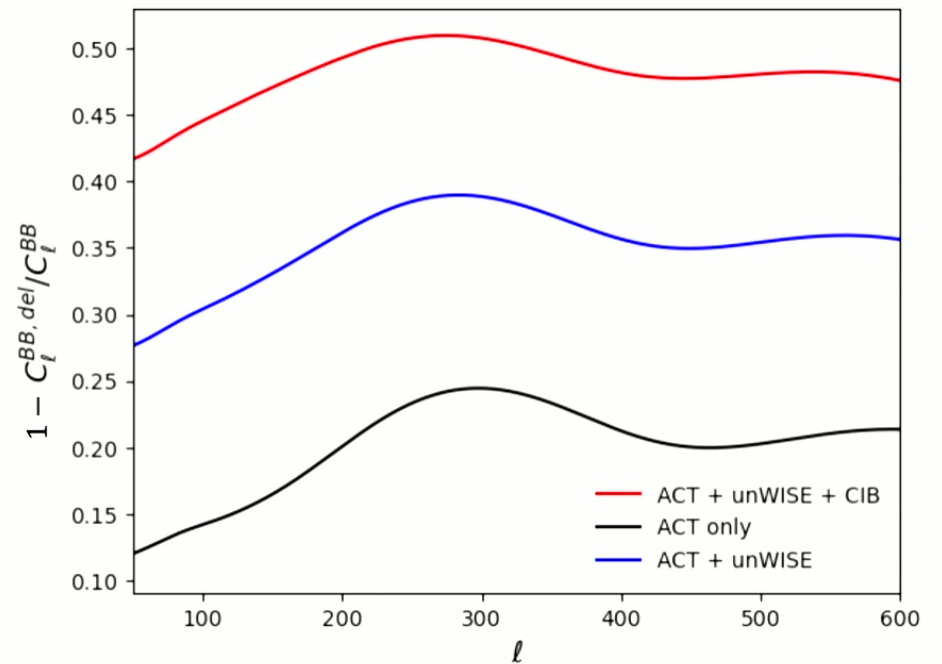


Template characterization

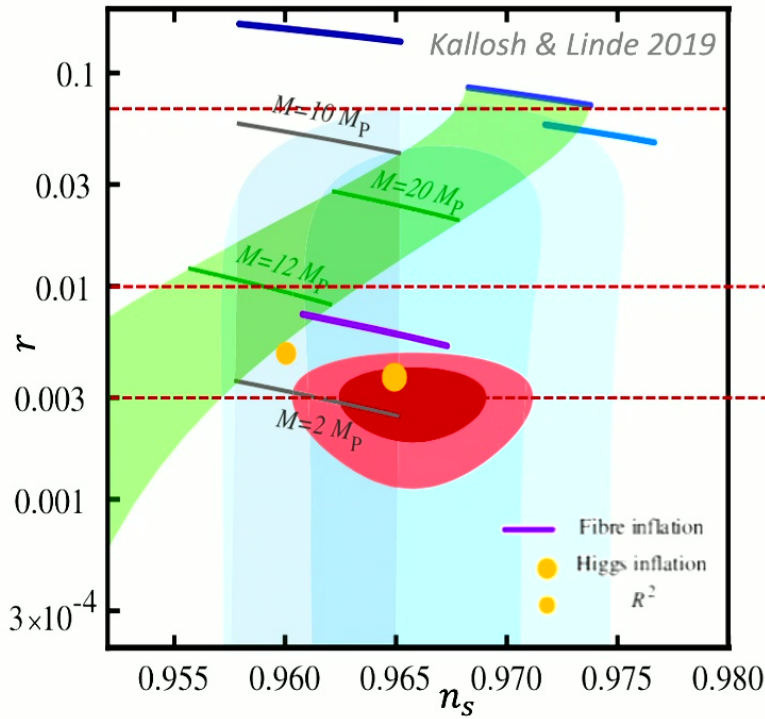
Template auto-spectrum and cross-spectrum with Planck SMICA B-modes



Delensing efficiency: fractional decrease in lensing B-mode power



Future prospects

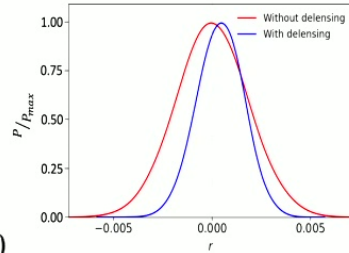


Delensing will become **increasingly important** as noise levels continue diminishing!

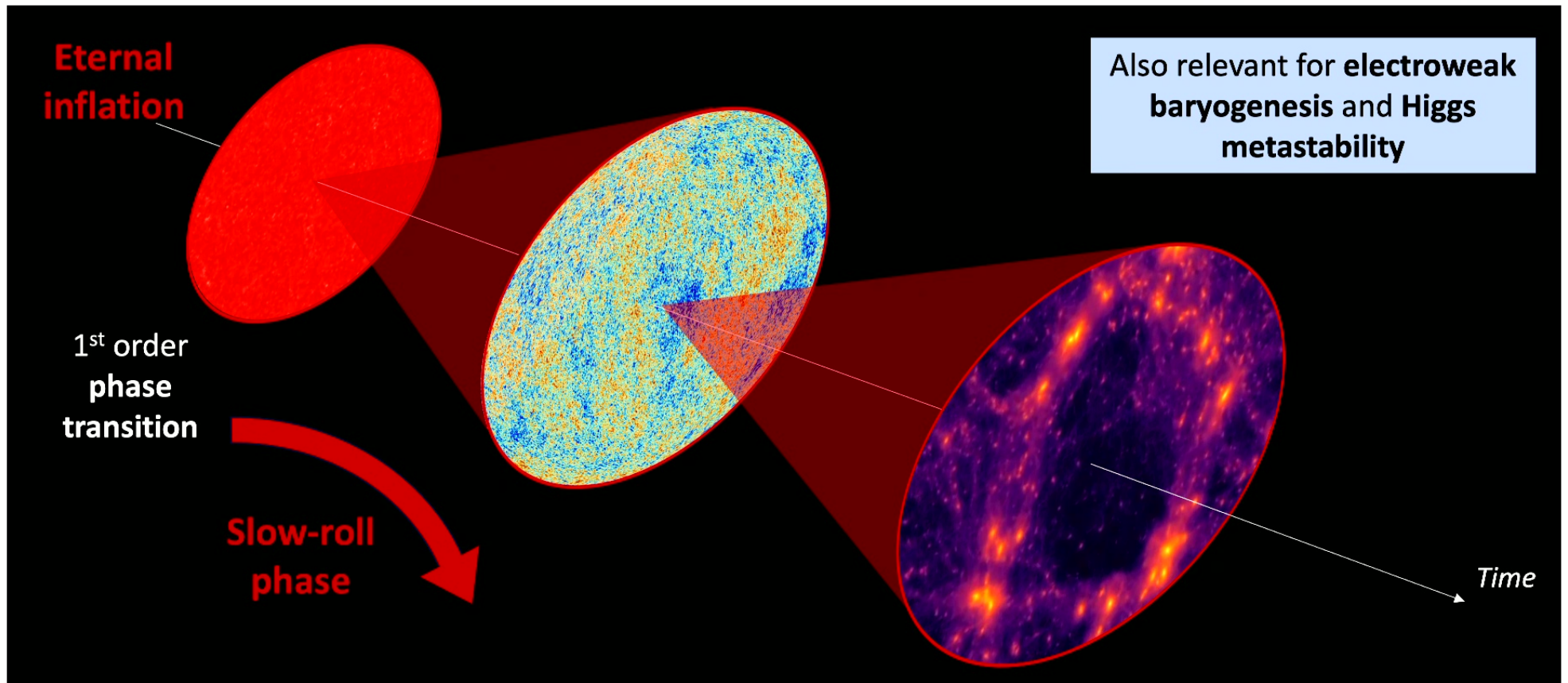
Our forecasts
 $r < 0.01$

Additional SO SATs (2026)

LiteBIRD, CMB-S4, ...
 $r < 0.003$



Part II – Simulating false vacuum decay



Instanton formalism

Relativistic scalar field in $D + 1$ dimensions: $\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi)$

➤ Path integral in **Euclidean time** $\tau = it$

$$\langle \phi_{TV} | e^{-\frac{\hat{H}\tau}{\hbar}} | \phi_{FV} \rangle = \int \mathcal{D}[\phi] e^{-\frac{S_E}{\hbar}}$$

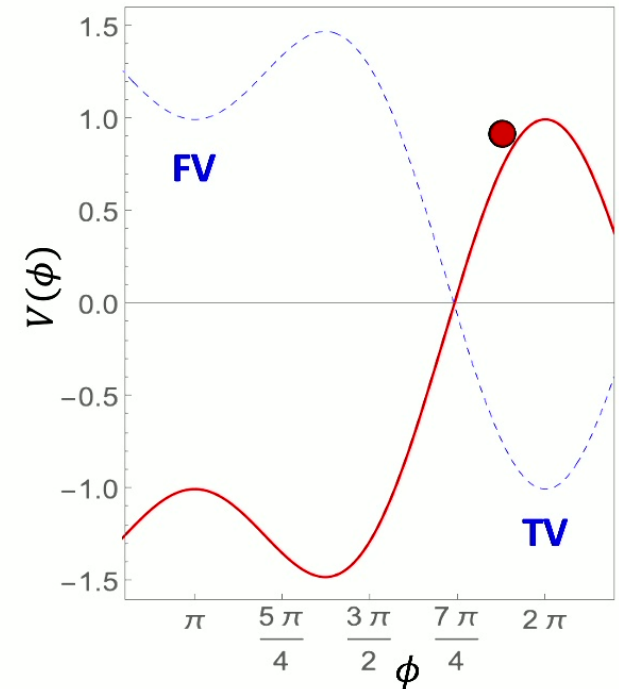
$$S_E = \int d\tau d^D \vec{x} \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial \tau} \right)^2 + \frac{1}{2} (\vec{\nabla} \phi)^2 + V(\phi) \right]$$

Minimizing S_E :

$$(\partial_\tau^2 + \nabla^2) \phi = V'(\phi)$$

➤ Assuming **$O(D + 1)$ symmetry**

$$\frac{d^2 \phi}{d\rho^2} + \frac{D}{\rho} \frac{d\phi}{d\rho} = V'(\phi) \quad \left\{ \begin{array}{l} \rho = \sqrt{\tau^2 + |\vec{x}|^2} \\ d\phi/d\rho \Big|_{\rho=0} = 0 \\ \phi(\rho = \infty) = \phi_{FV} \end{array} \right.$$



Classical particle path in
inverted potential

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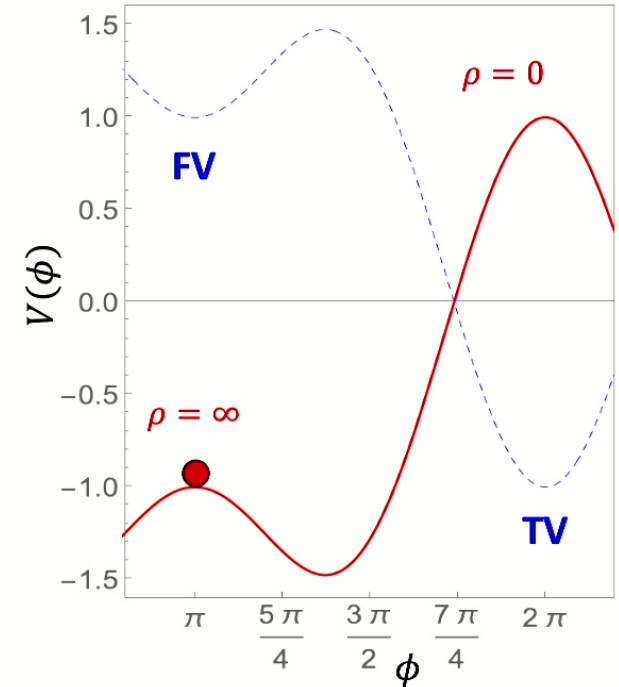
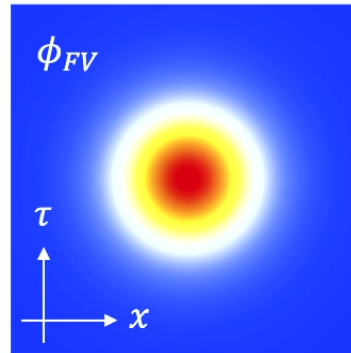
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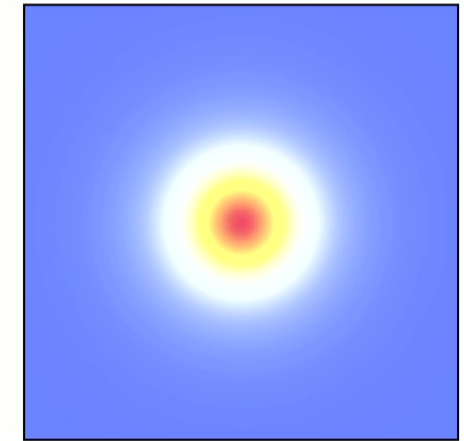
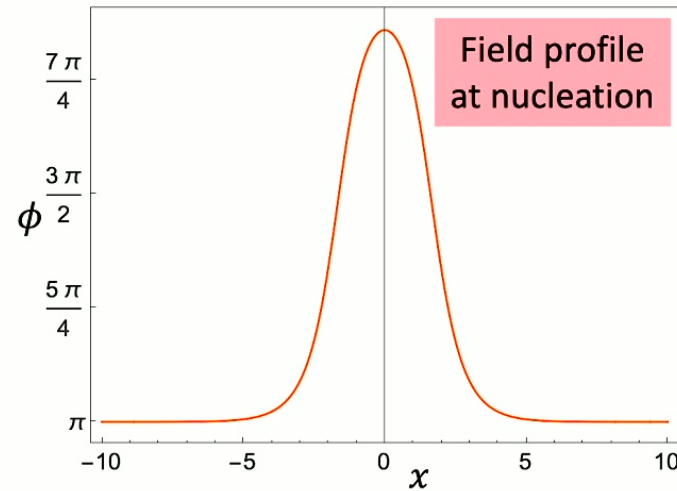
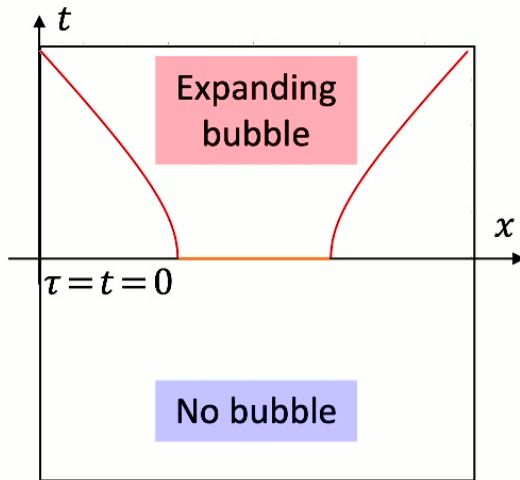
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Classical particle path in inverted potential

Analytic continuation



$t > 0$: classical evolution

Bubble wall at $\rho = R$

\Rightarrow hyperbolic trajectory

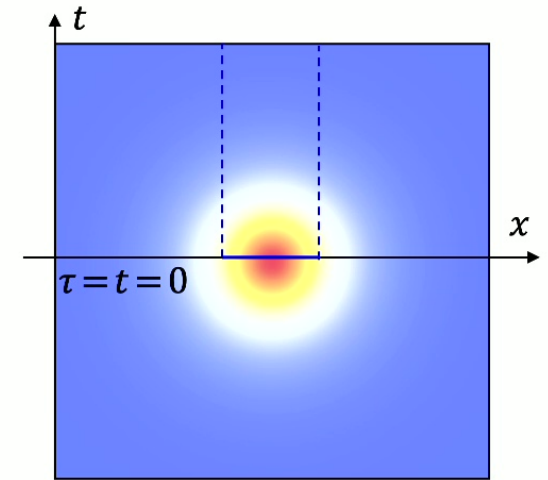
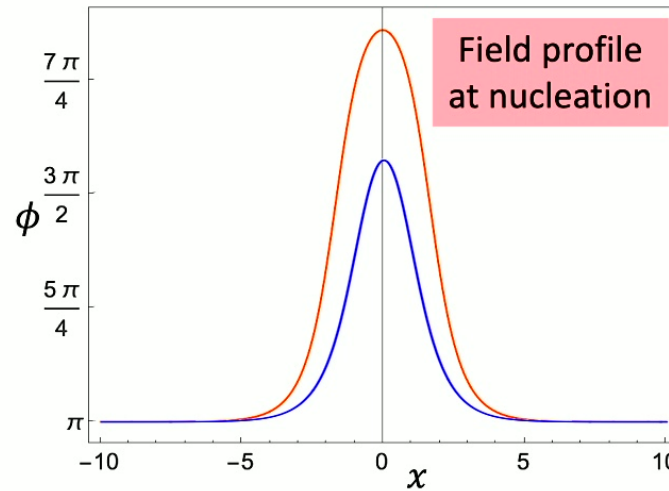
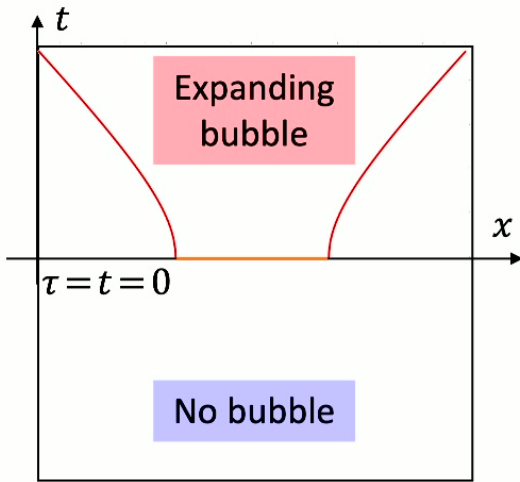
$$-t^2 + |\vec{x}|^2 = R^2$$

Thermal case

\triangleright Independent of τ at $T \gg 0$

$$\frac{d^2\phi}{dx^2} = V'(\phi)$$

Analytic continuation



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Bubble wall at $\rho = R$

\Rightarrow hyperbolic trajectory

$$-t^2 + |\vec{x}|^2 = R^2$$

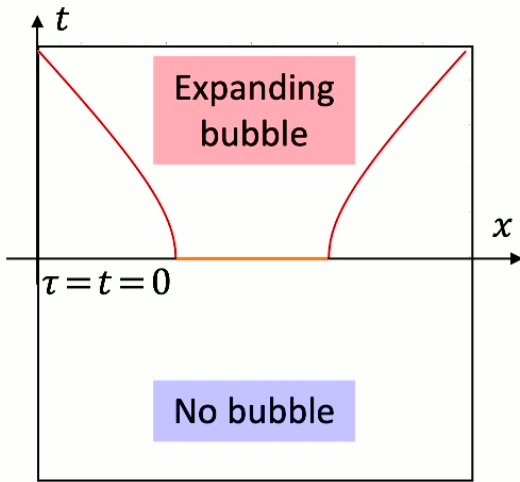
Thermal case

➤ Independent of τ at $T \gg 0$

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➤ Fixed but **unstable** solution

Analytic continuation

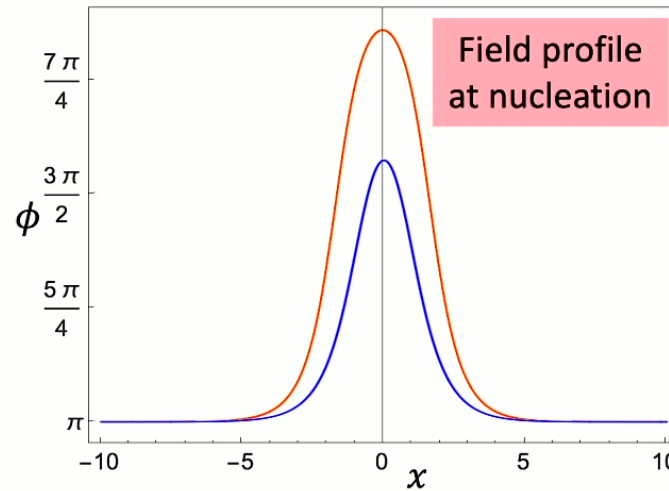


$t > 0$: classical evolution

Bubble wall at $\rho = R$

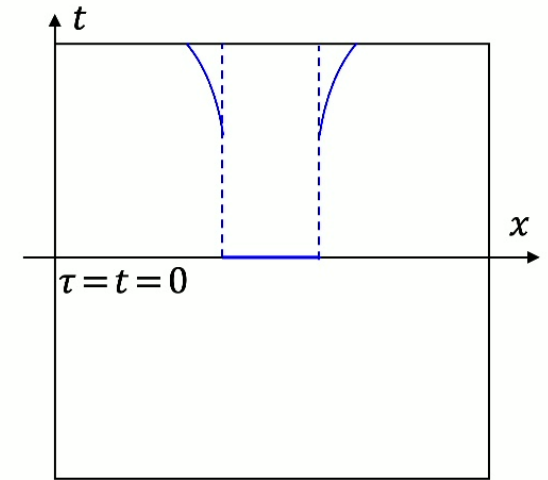
\Rightarrow hyperbolic trajectory

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Missing information

⊗ Real-time nucleation dynamics



Thermal case

➤ Independent of τ at $T \gg 0$

$$\frac{d^2\phi}{dx^2} = V'(\phi)$$

➤ Fixed but **unstable** solution

Semiclassical lattice simulations

Step 1 – Quantum initial conditions

- Sample Gaussian fluctuations of free massive scalar field ($\hbar = c = 1$)

$$\langle \delta\phi_k^* \delta\phi_{k'} \rangle = \frac{1}{2\omega_k} \delta(k - k') \quad \langle \dot{\phi}_k^* \dot{\phi}_{k'} \rangle = \frac{\omega_k}{2} \delta(k - k')$$

$$\text{with } \omega_k^2 = k^2 + m_{FV}^2 = k^2 + V''(\phi)|_{\phi_{FV}} \text{ and } \mathbf{k} \leq \mathbf{k}_{UV}$$

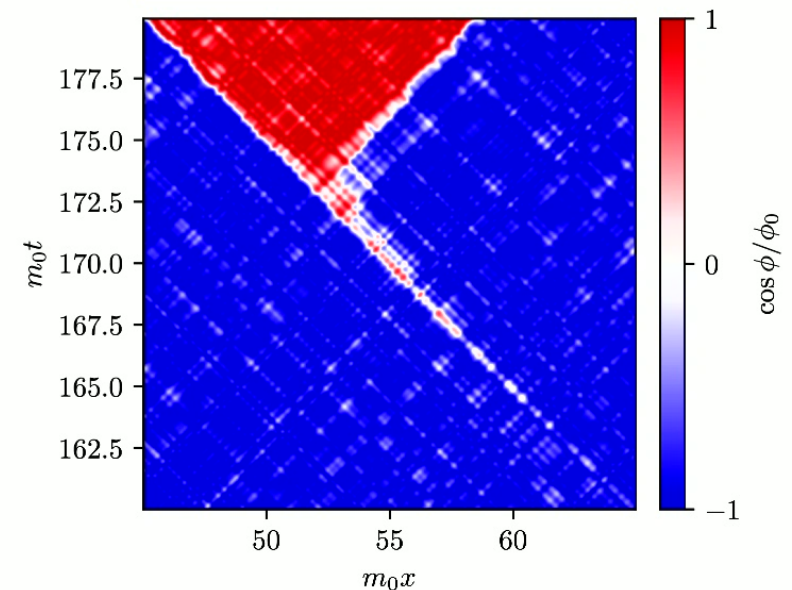
Step 2 – Classical evolution

- Solve real-time Hamiltonian EoM with periodic BCs in space

$$\left. \begin{aligned} \Pi &= \dot{\phi} \\ \dot{\Pi} &= \nabla^2 \phi - V'(\phi) \end{aligned} \right\} \text{Real-time dynamics } \phi(x_i, t_i)$$

See Braden et al. 2019, arXiv:1806.06069

Example in 1 + 1D using A. Jenkins' lattice-fvd code



Semiclassical lattice simulations

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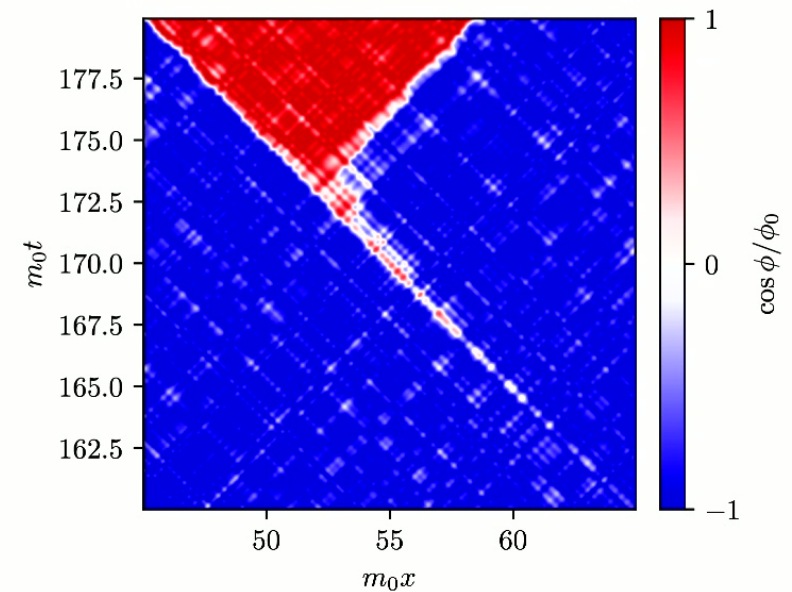
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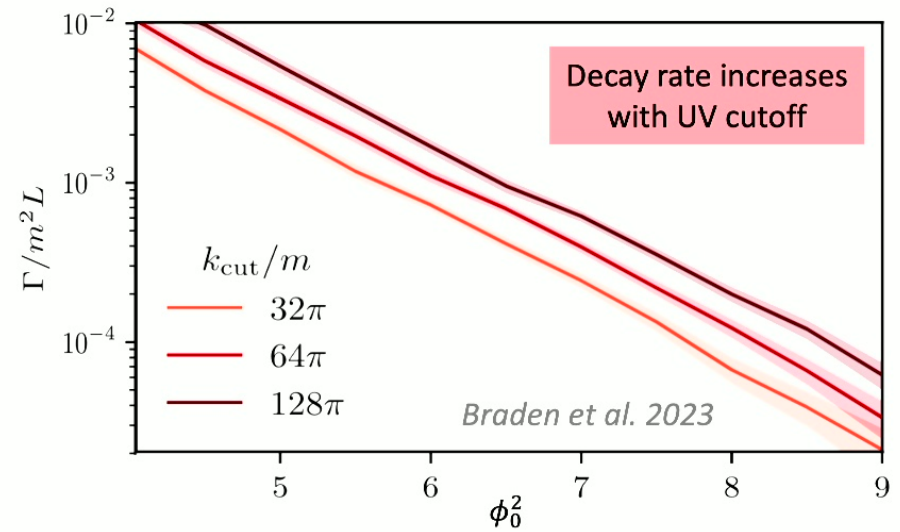
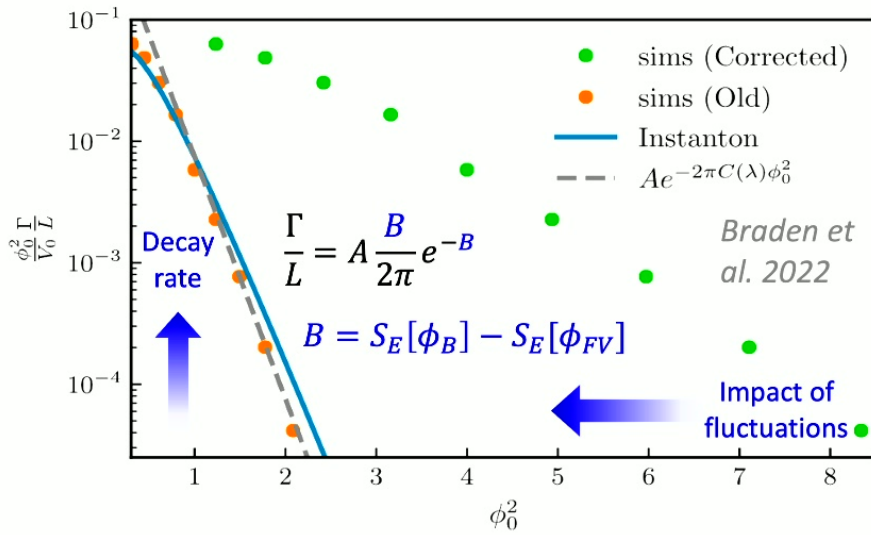
$$\left. \begin{aligned} \Pi &= \dot{\phi} \\ \dot{\Pi} &= \nabla^2 \phi - V'(\phi) \end{aligned} \right\} \begin{array}{l} \checkmark \text{ Real-time dynamics } \phi(x_i, t_i) \\ \checkmark \text{ Truncated Wigner approximation} \end{array}$$

See Braden et al. 2019, arXiv:1806.06069

Example in 1 + 1D using A. Jenkins' lattice-fvd code



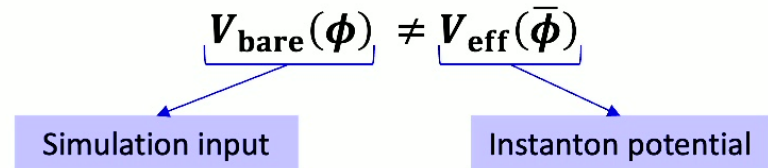
Decay rate comparison



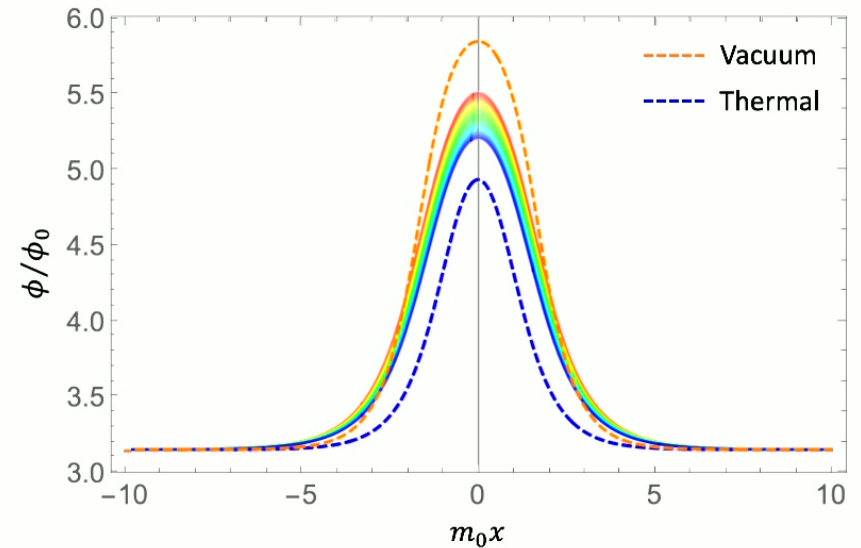
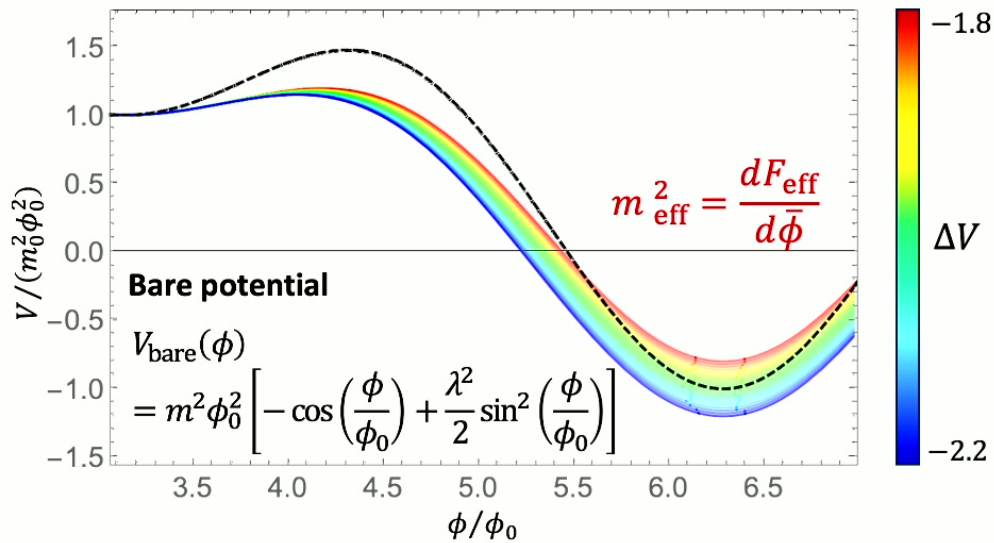
⚡ Discrepancy between lattice results and instanton prediction

💡 Dependence on UV cutoff

Possible cause: renormalization effects



Probing the effective potential



Effective force on $\bar{\phi}$ in symmetric minima

$$F_{\text{eff}}(\bar{\phi}) = \langle V'_{\text{bare}} \rangle = \sum_{n=0} \frac{1}{(2n)!} \frac{d^{2n} V'_{\text{bare}}}{d\phi^{2n}} \Big|_{\bar{\phi}} \langle \delta\phi^{2n} \rangle$$

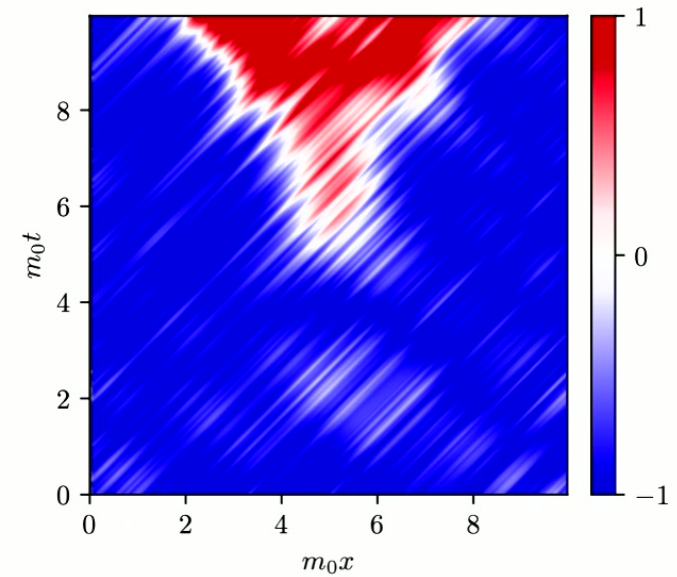
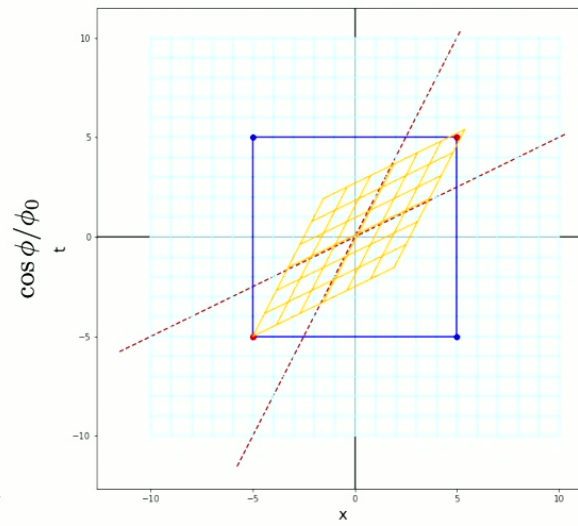
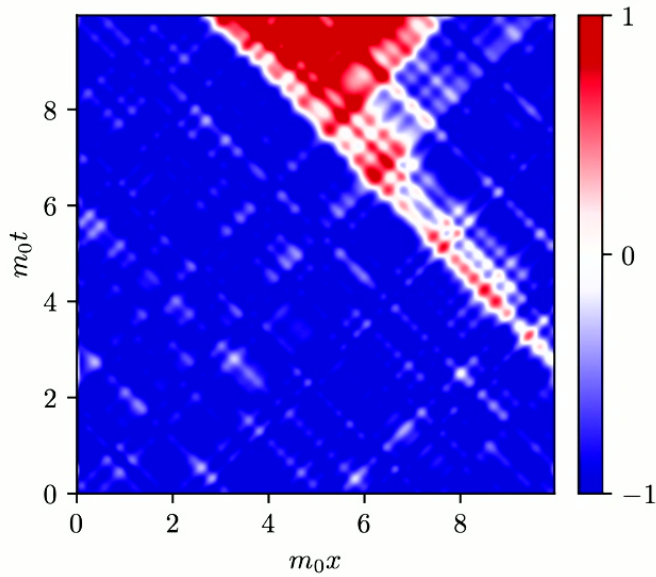
- V_{eff} is **not fully determined** in-between the minima!
- **Bubble profiles** are sensitive to the full shape of the potential

Extracting bubble profiles

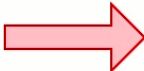
$$\phi(x, t) = \sum_{k_x k_t} \phi_{k_x k_t} e^{i(k_x x + k_t t)}$$

$$(x_B, t_B) \rightarrow (x', t')$$

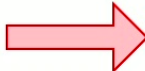
$$\phi_B(x_B, t_B) = \sum_{k_x k_t} \phi_{k_x k_t} e^{i(k_x x' + k_t t')}$$



Simulation frame



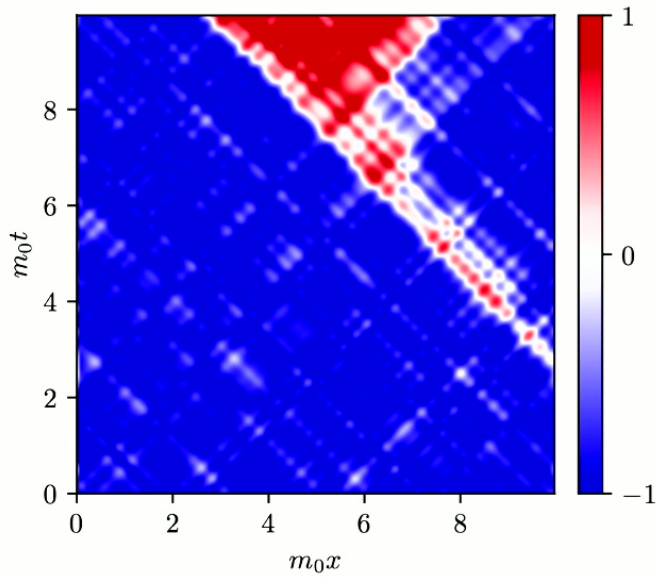
Lorentz transform



Boosted frame

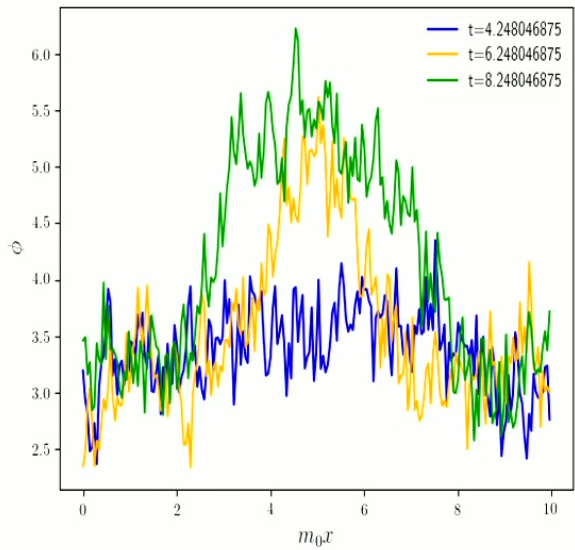
Extracting bubble profiles

$$\phi(x, t) = \sum_{k_x k_t} \phi_{k_x k_t} e^{i(k_x x + k_t t)}$$

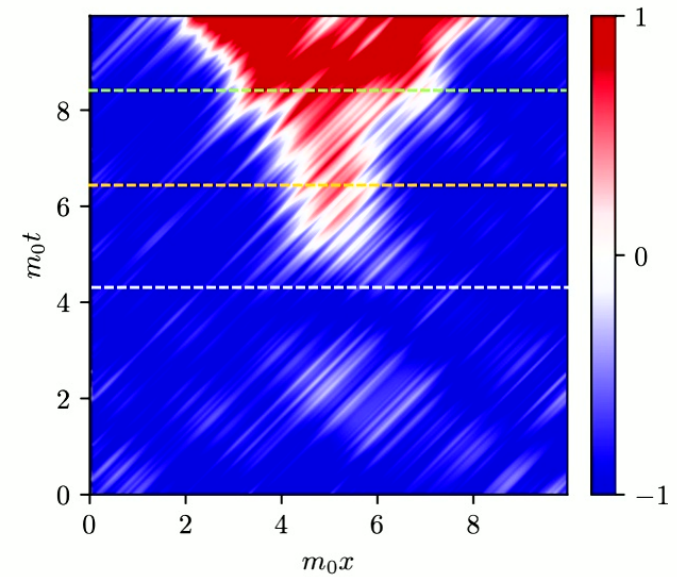


Simulation frame

Field profiles around nucleation



$$\phi_B(x_B, t_B) = \sum_{k_x k_t} \phi_{k_x k_t} e^{i(k_x x' + k_t t')}$$



Boosted frame

Outlook: cold-atom analogues

➤ **BEC** with 2 hyperfine states $\psi_{1,2}(\mathbf{x}) = \sqrt{n_{1,2}(\mathbf{x})} e^{i\phi_{1,2}(\mathbf{x})}$

➤ **Non-relativistic** Hamiltonian:

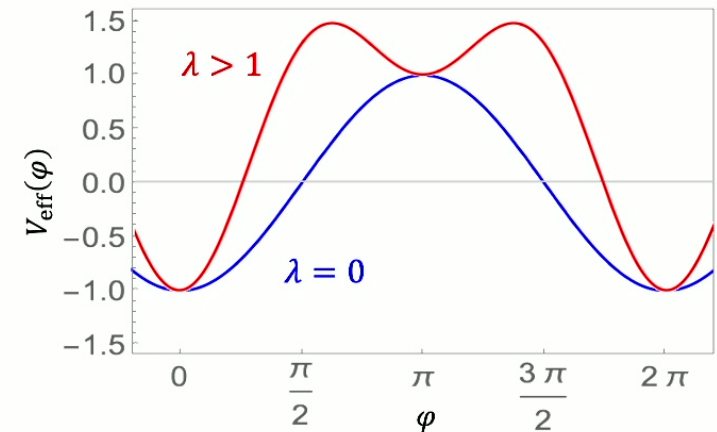
$$\hat{H} = \int d\mathbf{x} \sum_{i=j} \left[-\hat{\psi}_i^\dagger \frac{\hbar^2}{2m} \nabla^2 \hat{\psi}_j + \frac{g}{2} \hat{\psi}_i^\dagger \hat{\psi}_i^\dagger \hat{\psi}_j \hat{\psi}_j \right] - v(t) \sum_{i \neq j} \hat{\psi}_i^\dagger \hat{\psi}_j$$

➤ Low- k EoM for **phase difference** $\hat{\varphi} = \hat{\phi}_1 - \hat{\phi}_2$ is analogue to relativistic scalar field:

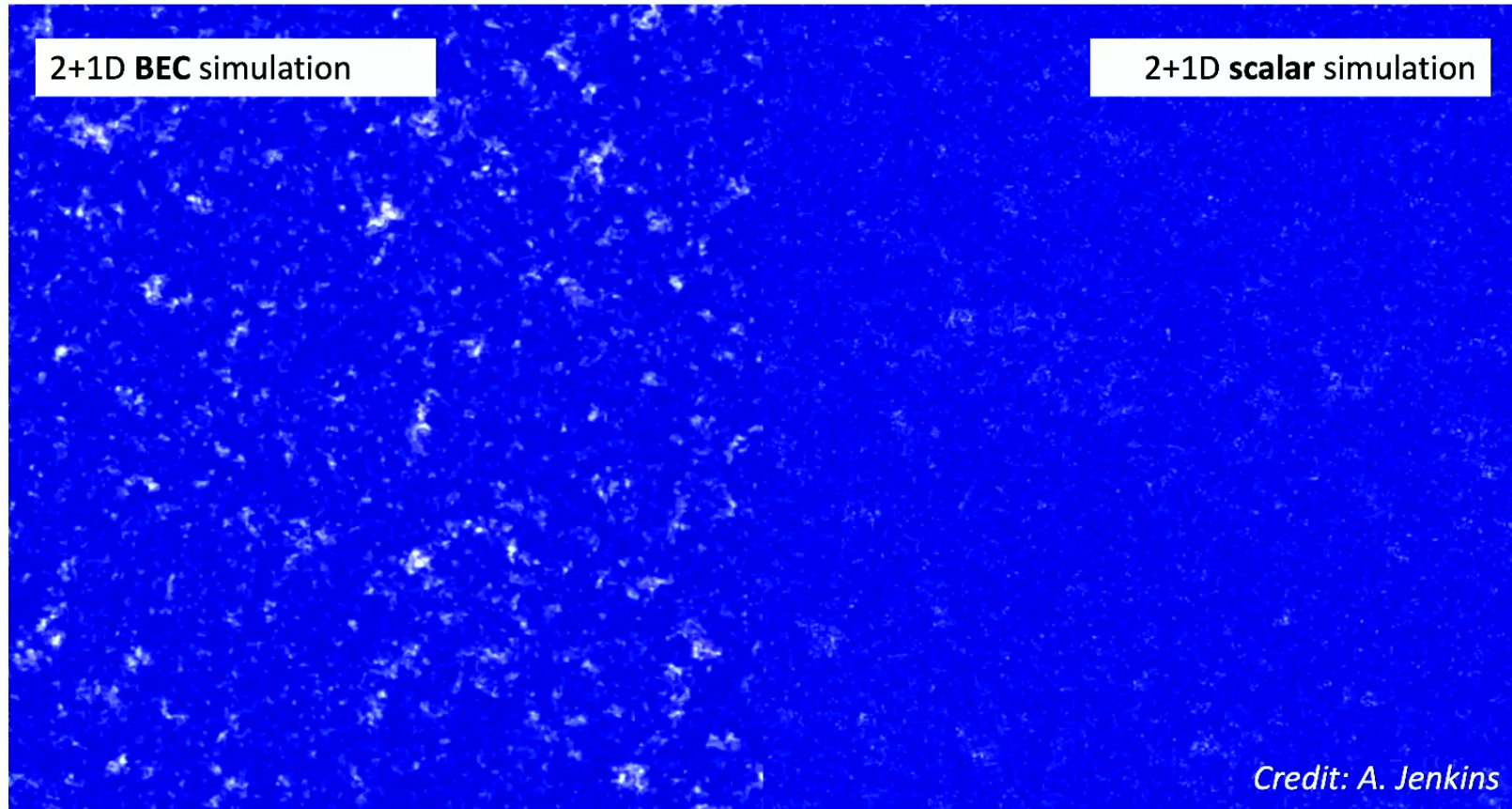
$$(c_s^{-2} \partial_t^2 - \nabla^2) \varphi = -V'(\varphi)$$

➤ For **fast enough** modulation $v(t) = v_0 + \lambda \hbar \omega \cos(\omega t)$: $\varphi \equiv \varphi_{\text{slow}} + \varphi_\omega$ and

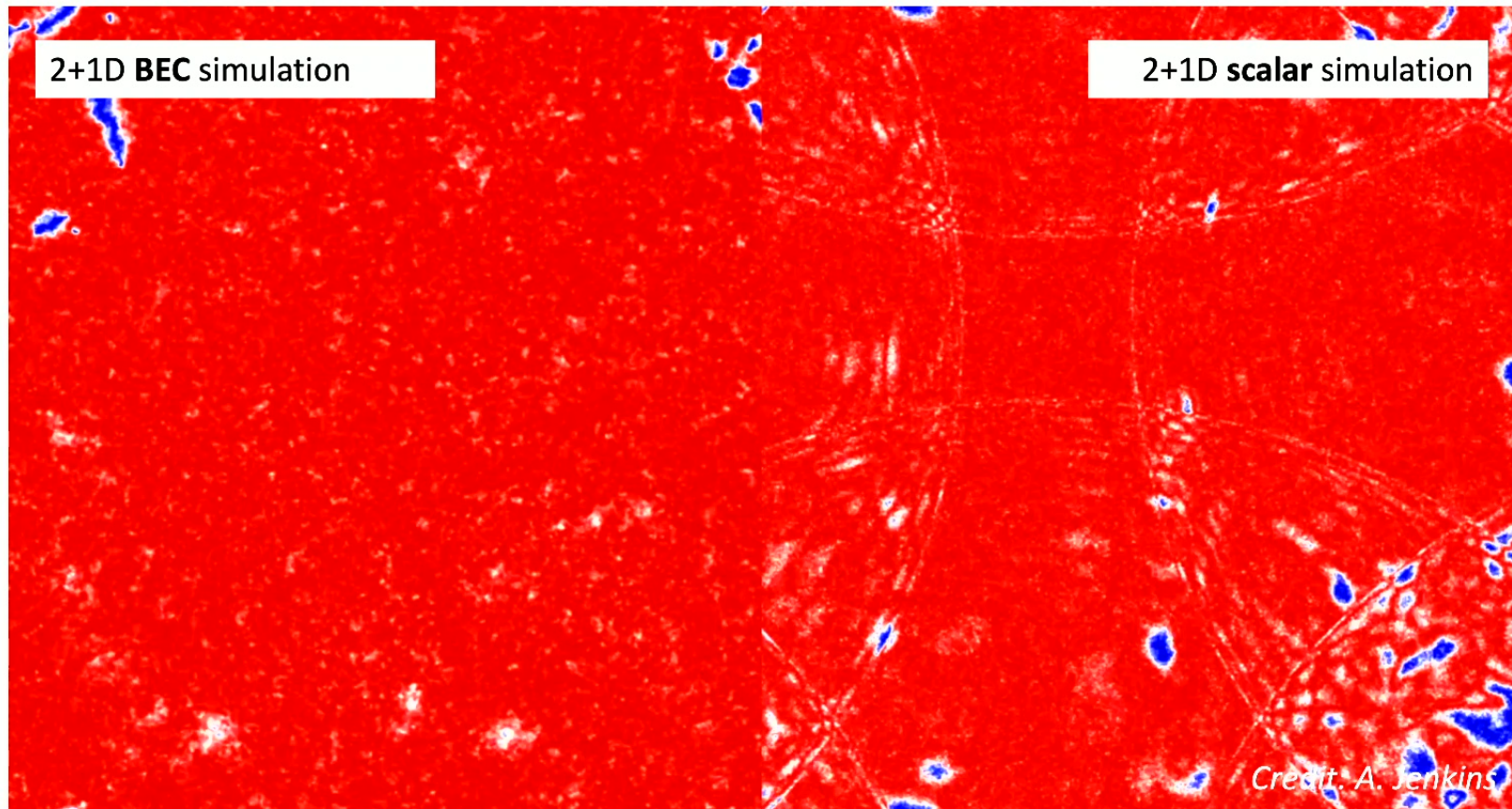
$$V_{\text{eff}}(\varphi_{\text{slow}}) = V_0 \left[-\cos(\varphi_{\text{slow}}) + \frac{\lambda^2}{2} \sin^2(\varphi_{\text{slow}}) \right]$$



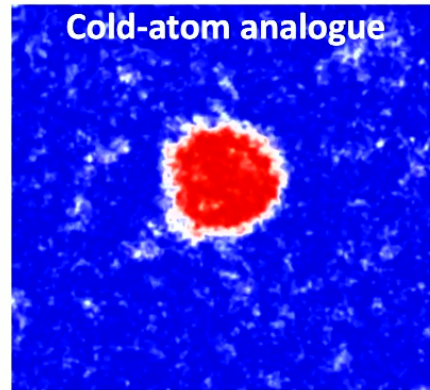
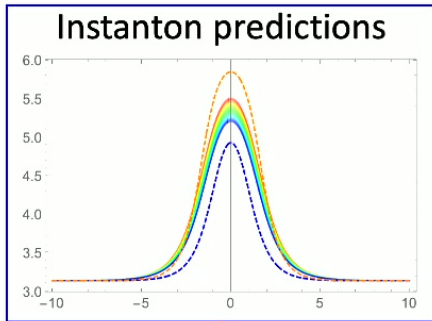
The Big Bang in the lab



The Big Bang in the lab



Next steps: quantum simulators



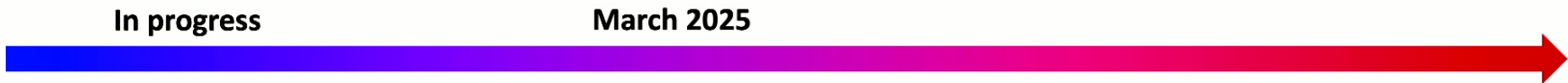
Same decay channels?

Theoretical work on renormalization with stacked bubble profiles

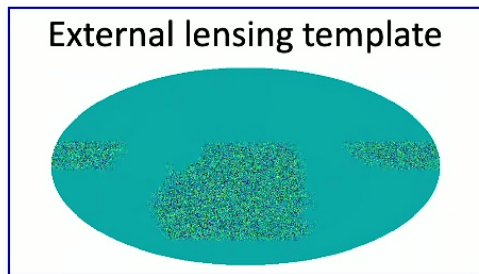
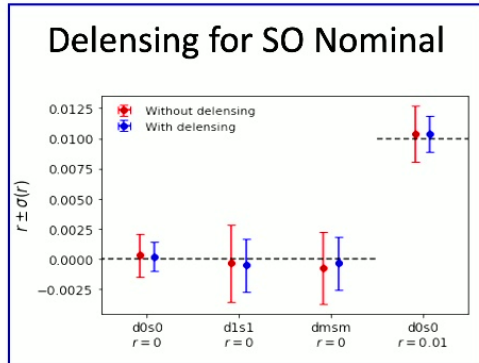
Extend findings to 2 spatial dimensions

Interpretation

- Range of validity of lattice approach
- « Calibration » in a controlled setting



Next steps: CMB



Current status

Forecasts with SO:UK and SO:Japan

Illustration of the SO:UK and SO:Japan satellite configurations, showing multiple satellite units in a circular arrangement.

Modelling of systematics, e.g. LAT foregrounds

Adapt code to real data

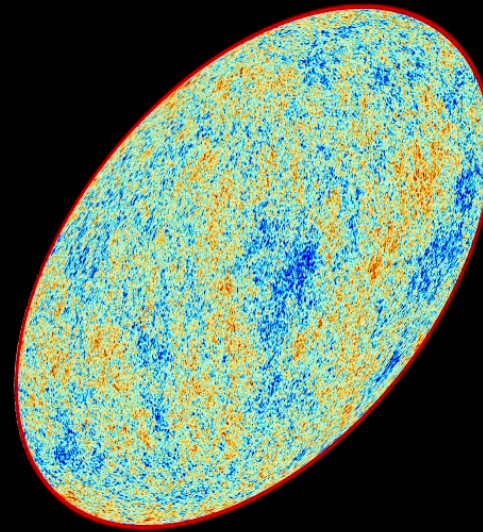
First SO r report with delensing

2025

Decreasing noise and $\sigma(r)$



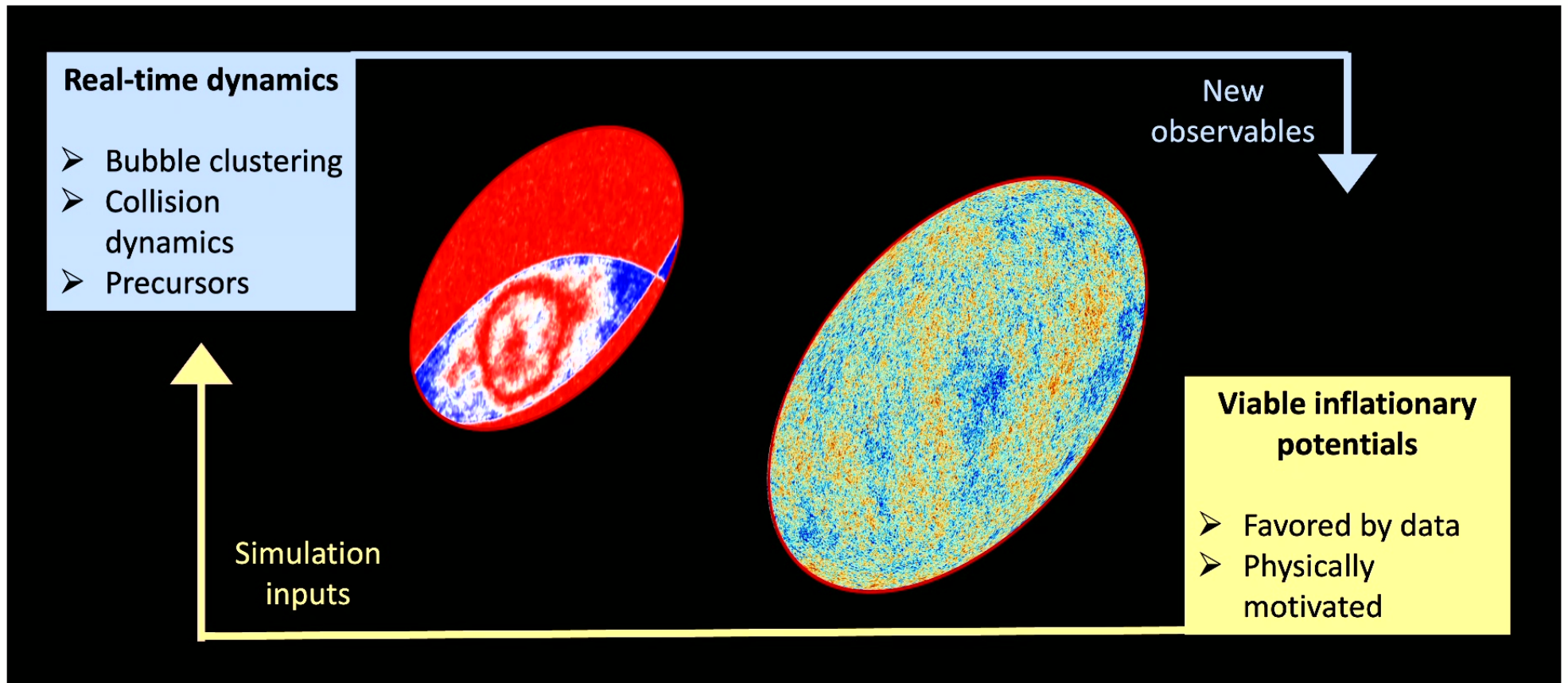
A new perspective on the early Universe



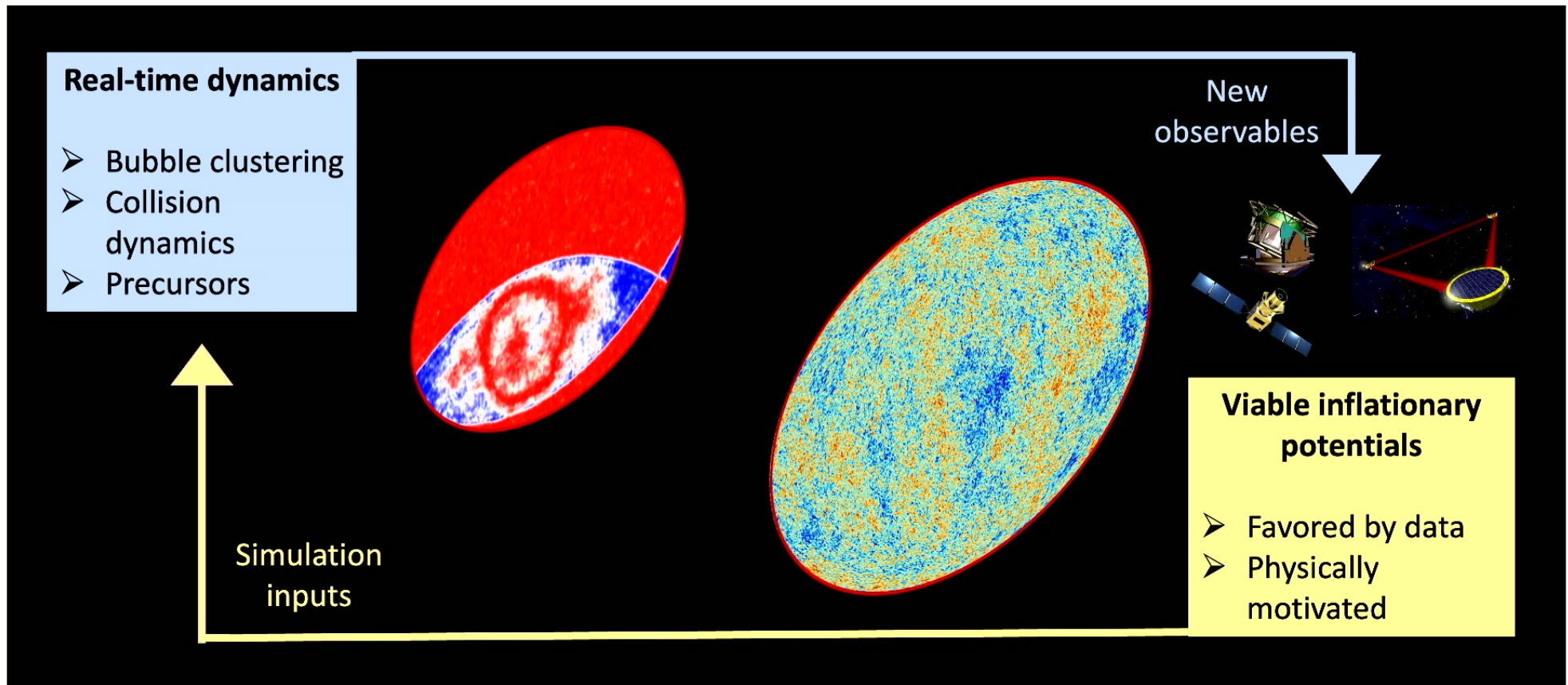
Viable inflationary potentials

- Favored by data
- Physically motivated

A new perspective on the early Universe



A new perspective on the early Universe



The Big Bang in the lab

