

**Title:** Lecture - QFT II, PHYS 603

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**Collection/Series:** Quantum Field Theory II (Core), PHYS 603, November 12 - December 11, 2024

**Subject:** Condensed Matter, Particle Physics, Quantum Fields and Strings


**Date:** December 03, 2024 - 9:00 AM

**URL:** <https://pirsa.org/24120002>

$\phi^4$  theory massless  $d=4$  1 loop  
 $|k| < \Lambda$  cut-off in momenta

renormalization scale (momentum)  $\mu$

renormalized coupling  $g_R$

$g_R :=$   or S-matrix  $2 \rightarrow 2$  at <sup>reference</sup> momenta  $|p| \sim \mu$

1PI-function  $\rightarrow \Gamma_R$  effective action  $\rightarrow \langle \phi \phi \rangle_R$

op  $\int D[\phi] \exp(-S_R[\phi])$

$B \leftarrow Q$  (8.1) Notes  
 $C \leftarrow X$  (8.2)

$S_R[\phi]$  renormalized action

$$S_R[\phi] = \int d^4x \frac{A}{2} (\partial_\mu \phi)^2 + \frac{B}{2} \phi^2 + \frac{C}{4!} \phi^4$$

$A = 1$

$$C = g_R + g_R^2 \frac{3}{2} \frac{1}{(4\pi)^2} \log\left(\frac{\Lambda^2}{\mu^2}\right)$$

$V$  finite  $B = 0 - g_R \frac{1}{2} \frac{1}{(4\pi)^2} \Lambda^2$   
 $\Lambda \rightarrow \infty, g_R$  finite

$$S_R[\phi] \quad \mathcal{B} \leftarrow \mathcal{Q} \quad (8.1) \text{ Notes}$$

$$\mathcal{G} \leftarrow \mathcal{X} \quad (8.2)$$

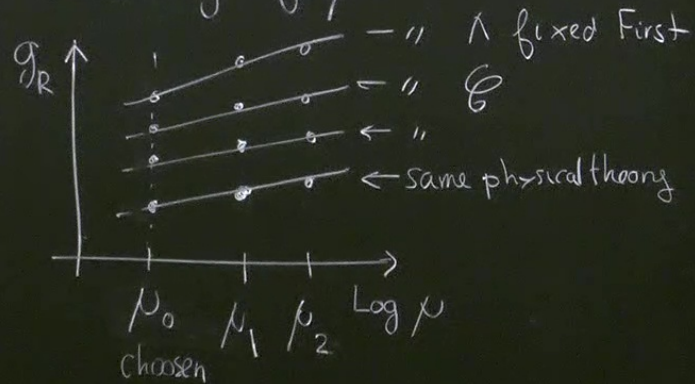
red action

$$(\partial_\mu \phi)^2 + \frac{\mathcal{B}}{2} \phi^2 + \frac{\mathcal{G}}{4!} \phi^4$$

$$\mathcal{G} = g_R + g_R^2 \frac{3}{2} \frac{1}{(4\pi)^2} \log\left(\frac{\Lambda^2}{\mu^2}\right)$$

$$\frac{1}{(4\pi)^2} \Lambda^2$$

Meaning of  $\mu$ ?



$g_R(\mu)$  depends on  $\mu$   $\Lambda, \mathcal{G}$  fixed

$\Lambda \rightarrow \infty$   $g_R(\mu_0)$  Fixed

Coupling constants  $g_R(\mu)$  are "running"  
with the energy/momentum scale

$$\mu_0 \quad \text{vs} \quad \mu_1$$

$$g_R = g_0 \quad \quad g_1$$

$$g_0 + g_0^2 \left( \text{Log} \frac{\Lambda^2}{\mu^2} \right) = g_1 + g_1^2 * \text{Log} \frac{\Lambda^2}{\mu^2}$$

$$g_1 - g_1^2 \frac{3}{2} \frac{1}{(4\pi)^2} \log \left( \frac{\mu_1}{\mu_0} \right)^2 + \dots = g_0$$

Better  
variab

Better to consider variations of  $\mu$

$$\mu_1 = \mu_0 + \delta\mu_0$$

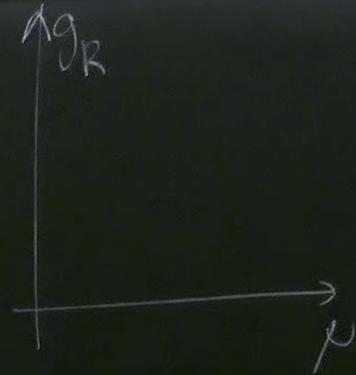
$$g_1 = g_0 + \delta g_0$$

$$\delta g_0 = \frac{\delta\mu}{\mu} \frac{3}{(4\pi)^2} g_0^2 + \dots$$

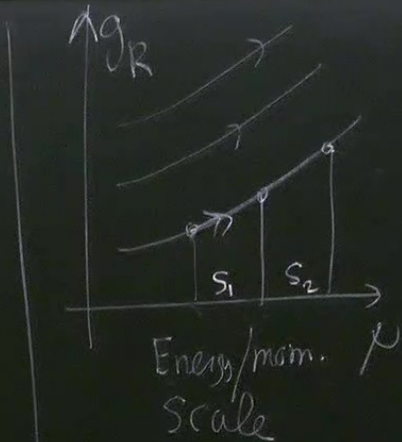
↑ 2 loops

$$g_1 = g(\mu_1; g_0, \mu_0)$$

$$\mu \frac{d}{d\mu} g(\mu) = \frac{3}{(4\pi)^2} g(\mu)^2 + \dots \text{ with i.o. } g(\mu_0) = g_0$$



$\mu_0 + \delta\mu_0$   
 $g_0 + \delta g_0$   
 $+ \dots$   
 $\uparrow 2 \text{ loops}$



Flow equation

$$\mu \frac{d}{d\mu} g(\mu) = \beta_g(g(\mu))$$

Beta-Function

$$\beta(g) = + \frac{3}{(4\pi)^2} g^2$$

$\phi^4$   $d=4$

with i.o.  $g(\mu_0) = g_0$

$g_0, \mu_0 \rightarrow g(\mu)$

$\mu_0 \rightarrow \mu = S \mu_0$  rescaling factor  $S_1, S_2 = S$

Flow equation generates a (semi)group additive in  $\log \mu$

$$\mu_0 \rightarrow \mu_1 = S_1 \mu_0 \rightarrow \mu_2 = S_2 \mu_0$$

$$g_0 \rightarrow g_1 = \mathcal{R}_{S_1}(g_0) \rightarrow g_2 = \mathcal{R}_{S_2}(g_1) = \mathcal{R}_{S_2 S_1}(g_0)$$

Scale transformation in the coupling

group in perturbative QFT  $\leftrightarrow$  Wilson Ren  
relevant couplings irrelevant couplings



1PI-function  $\rightarrow \Gamma_R$  effective action  $\rightarrow \langle \phi \cdot \phi \rangle_R$  UV finite  $\mathcal{D} = 0 - g_R \frac{1}{2(4\pi)^2} \Lambda^2$   
 $\Lambda \rightarrow \infty, g_R$  finite

③ Flow equation generates a (semi)group additive in  $\log \mu$  = Renormalization Group

$$\mu_0 \rightarrow \mu_1 = S_1 \mu_0 \rightarrow \mu_2 = S_2 \mu_0$$

$$g_0 \rightarrow g_1 = \mathcal{R}_{S_1}(g_0) \rightarrow g_2 = \mathcal{R}_{S_2}(g_1) = \mathcal{R}_{S_2 S_1}(g_0)$$

Scale transformation in the coupling

$\Phi$ ED Gell-Mann-Low '54  
 Stueckelberg - Petermann '53  
 Callan - Zymanski '70  
 QFT & Stat. Mech.

group in perturbative QFT  $\leftrightarrow$  Wilson Ren  
 relevant couplings irrelevant couplings

'54  
 rmann '53

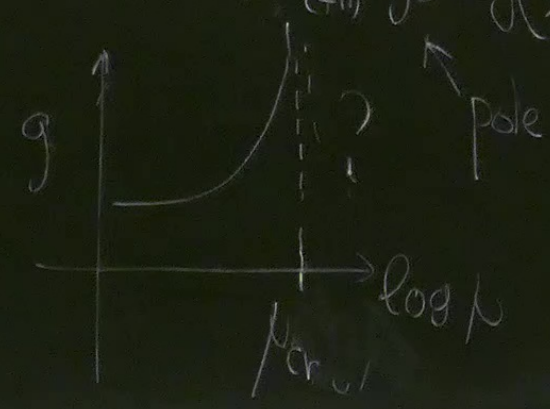
$$\beta(g) = g^2 \frac{3}{(4\pi)^2}$$

$$\frac{dg(\mu)}{d \log \mu} = + \frac{3}{4\pi^3} g(\mu)^2$$

start from  $\mu_0, g_0, g(\mu)$

$$\frac{dg^{-1}}{d \log \mu} = - \frac{3}{4\pi^3}$$

$$g(\mu) = \frac{g_0}{1 - \frac{3}{(4\pi)^3} g_0 \log(\mu/\mu_0)}$$



$$\mu_{\text{cut}} = \mu_0 \exp\left(\frac{(4\pi)^3}{3} \frac{1}{g_0}\right)$$

$$g(\mu) = \infty$$

Landau Pole  $\Rightarrow$  ghost in the spectrum

renormalization Group

Gell-Mann-Low '54

Stueckelberg-Petermann '53

Wittmann '70

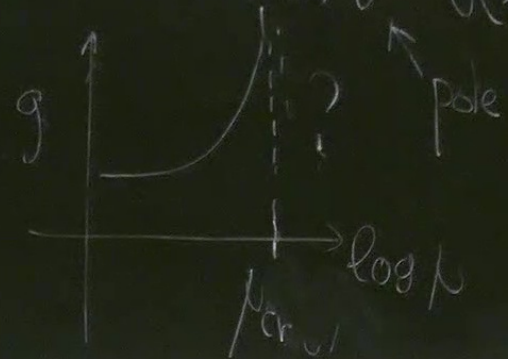
QFT

$$\beta(g) = g^2 \frac{3}{(4\pi)^2}$$

$$\frac{dg(\mu)}{d \log \mu} = + \frac{3}{(4\pi)^2} g(\mu)^2$$

start from  $\mu_0, g_0 \rightarrow g(\mu)$

$$g(\mu) = \frac{g_0}{1 - \frac{3}{(4\pi)^2} g_0 \log(\mu/\mu_0)}$$



$$\frac{dg^{-1}}{d \log \mu} = - \frac{3}{(4\pi)^2}$$

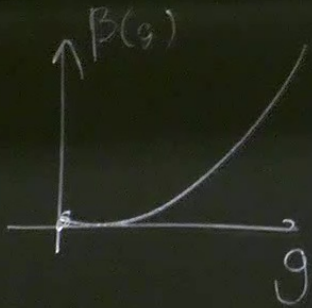
$$\mu_{\text{crit}} = \mu_0 \exp\left(\frac{(4\pi)^2}{3} \frac{1}{g_0}\right)$$

$$g(\mu) = \infty$$

Landau Pole  $\Rightarrow$  ghost in the spectrum

This Problem

$$\beta(g) > 0$$



UV

$g \nearrow$

IR

$g \rightarrow 0$

$$\mu \rightarrow 0 \quad g(\mu) \simeq \frac{(4\pi)^2}{3} \frac{1}{\log(\frac{1}{\mu})}$$

$\phi^4$  weakly coupled IR

(too) strongly coupled UV

Consistency of  $\phi^4$  theory in HEP

But this occurs at energies

$$E_{\text{Landau}} \simeq \exp\left(\frac{0}{g}\right) = \infty \text{ large in perturbation theory}$$

non perturbative

mploke "slavery"  
energy:  $\uparrow$   
QCD dichotomy

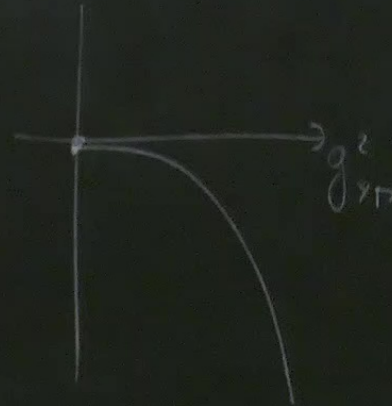
Same in QED  $\leftarrow$  up to  $\pm 10^{-15}$

not consistent in the UV

$$\exp(137) \gg 1$$

QCD - YM

$$\beta(g_{YM}^2) = -C(g_{YM}^2)^2$$



relevant couplings

irrelevant couplings

Why does it work?

Dimensional analysis

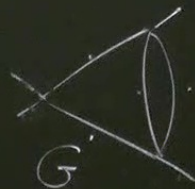
$$\int d^d k \quad \frac{1}{k^2}$$

↑  
internal loop

$$\int \frac{d^d h}{B} \quad \frac{1}{h^2}$$

L

"Superficial" degree of divergence of a Feynman Graph (Integral)



$$N=4$$

$$L=4$$

$$V=3 \quad B=2$$

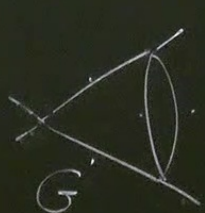
$$\frac{1}{p^2}$$

$$\times g$$

- N external "legs" B
- L internal lines E
- V internal vertices
- $N + 4V = 2L$  for  $\phi^4$

irrelevant couplings

"Superficial" degree divergence of a Feynman Graph (Integral)  $\omega(G)$



$$N=4$$

$$L=4$$

$$V=3 \quad B=2$$

$$\frac{1}{p^2}$$

$N$  external "legs"

$L$  internal lines

$V$  internal vertices

$$N + 4V = 2L \text{ for } \phi^4$$

$B$  internal loops

$$B = 1 + L - V \quad \text{Euler}$$

↑ connected graph

$$B = 1 + L/2 - N/4$$

$$\omega(G)$$

internal loops

$$1 + L - V \quad \text{Euler}$$

↑ connected graph

$$= 1 + L/2 - N/4$$

$$\omega(G) = d B - 2 L$$

$$\omega(G) < 0 \quad \text{superficially convergent } O(1)$$

$$\omega(G) = 0 \quad \text{marginally divergent } (\text{Log } \Lambda)$$

$$\omega(G) > 0 \quad \text{superficially divergent} \rightarrow \Lambda^{\omega(G)}$$

$$\omega(G) = (d - 4) B + (4 - N) \quad \text{external lines}$$

dimension of space



$d=4$     $N=0, N=2$     $\omega=2$     $\phi^2$     $\Lambda^2 + \log \Lambda^* M^2$  CT +  $(\partial\phi)^2 \cdot \log \Lambda^*$  ET  
 $N=4$     $\omega=0$     $\phi^4$     $\log \Lambda^*$  CT   Renormalizable  
 $d>4$    Any  $N>4$ , B large enough  
↓  
divergences    $\phi^6, (\partial^2\phi)^2, \phi^8, \dots$  no closure   Non-renormalizable

$[M^2] = 2$   
 $[g] = d - 4 \frac{d-2}{2} = \boxed{4-d = [g]}$  dim of coupling  
 dimensions in  $[k]$  or  $[E]$   
 $S[\phi]$   $[S] = 0$     $[\phi] = \frac{d-2}{2}$  engineering dim  
 $d^d x \partial_\mu \partial_\nu$