

Title: Adventures in Flat Holography

Speakers: Sabrina Pasterski

Collection/Series: Colloquium

Subject: Other

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Abstract:

Celestial Holography encompasses a decade-long endeavor to understand a flat space realization of the holographic principle starting from symmetries in the infrared. But where does it fit within other attempts at constructing a flat hologram? This colloquium delves into some fun tensions in the literature and hopes for resolving them.

ADVENTURES IN FLAT HOLOGRAPHY

PERIMETER INSTITUTE COLLOQUIUM

SABRINA PASTERSKI

11-27-24



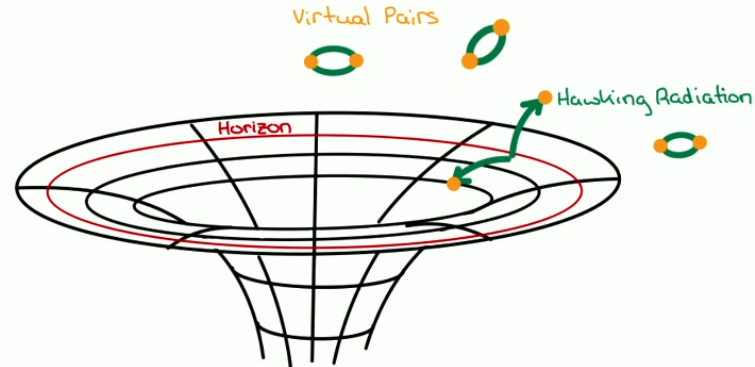
Outline

- I. Context for Flat Holography
- II. The Celestial Holography Program
- III. Fun with Swing Surfaces

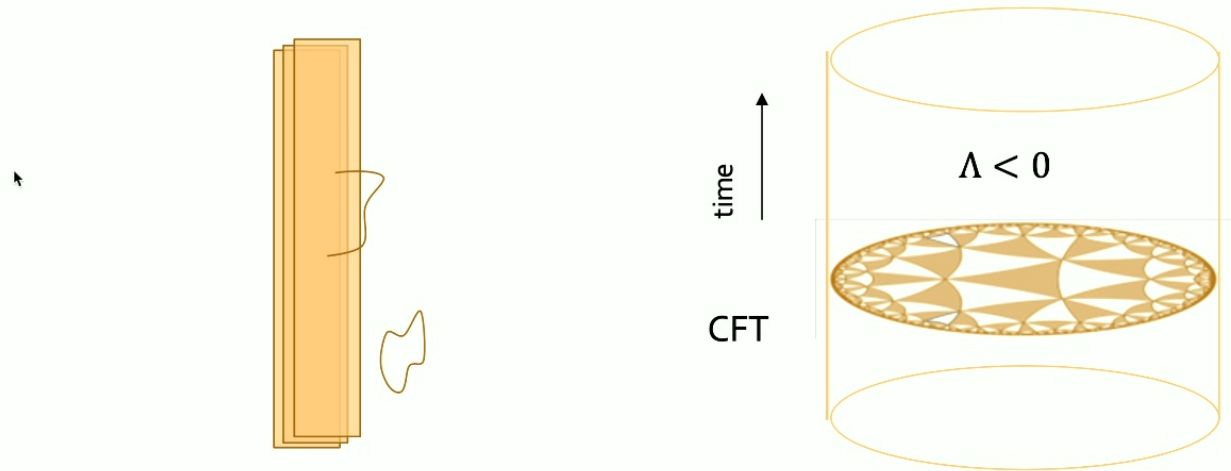
Celestial Holography encompasses a decade-long endeavor to understand a flat space realization of the holographic principle starting from symmetries in the infrared. But where does it fit within **other attempts** at constructing a flat hologram? This colloquium delves into some fun **tensions in the literature** and hopes for resolving them.

The **Holographic Principle** proposes that a quantum theory of gravity has an **equivalent** lower dimensional description **without gravity**.

$$S_{BH} = \frac{c^3 \text{Area}_{\text{Horizon}}}{4G_N \hbar}$$



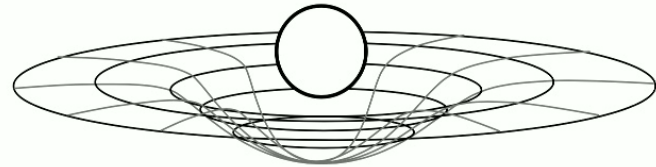
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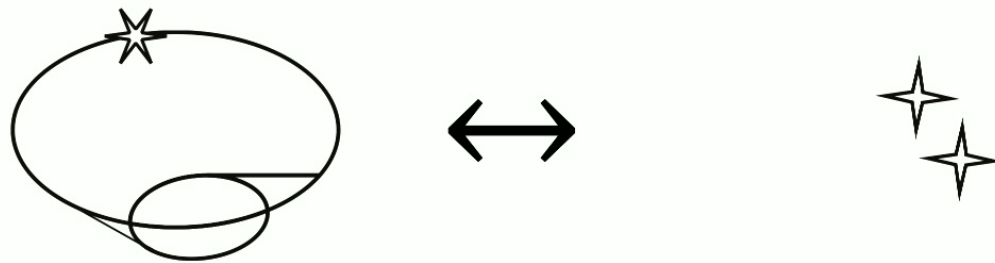
The **Holographic Principle** proposes that a quantum theory of gravity has an **equivalent** lower dimensional description **without gravity**.

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle = H|\psi\rangle$$

+



Asymptotically flat spacetimes, approximate a broad range of physics: from collider experiments to gravitational wave astronomy.

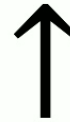


Can we find a dual **hologram?**

Approaches

Bottom Up

- Matching symmetries
ASG generators on both sides
- Consistency conditions
unitary bulk S-matrix
- Reasonable expectations
should include bulk graviton...

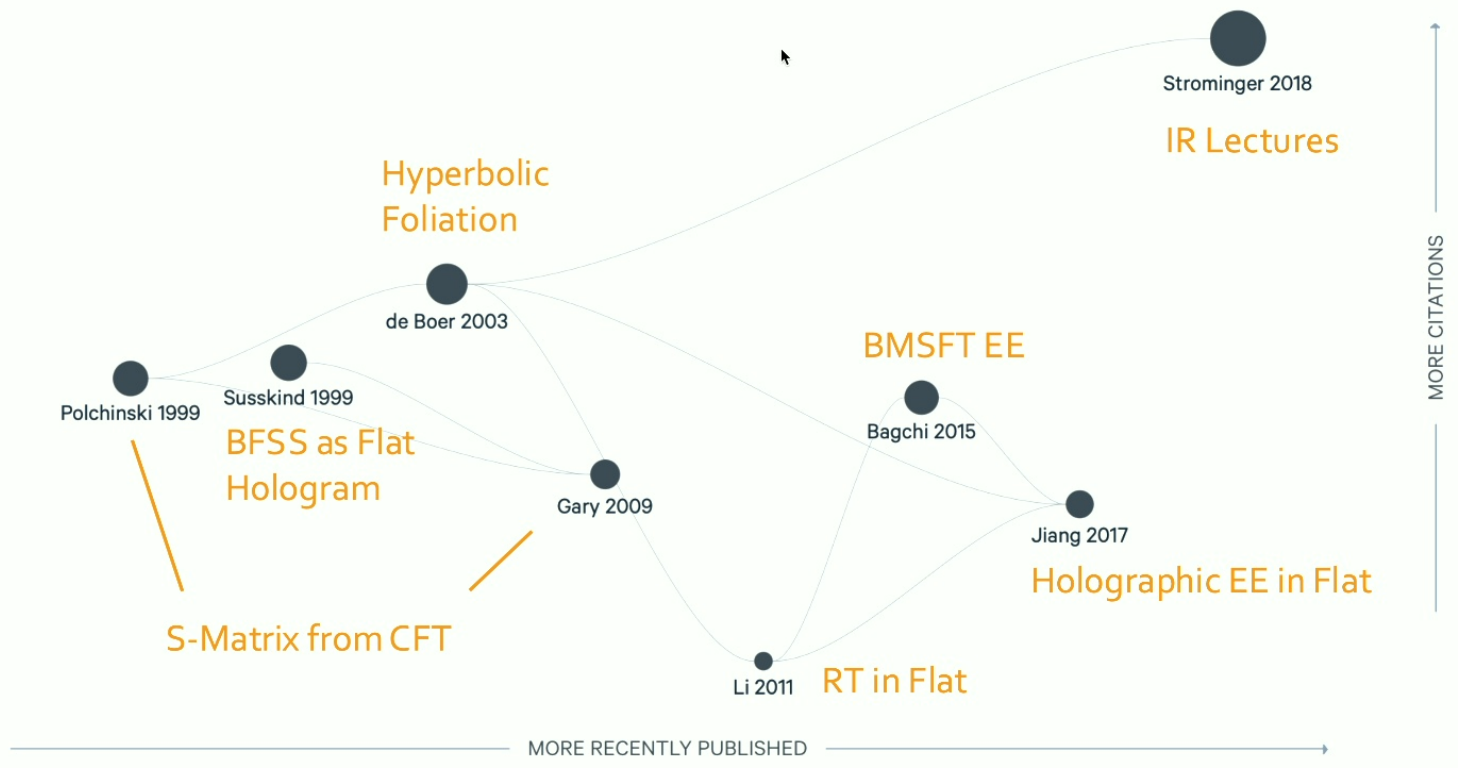


Top Down

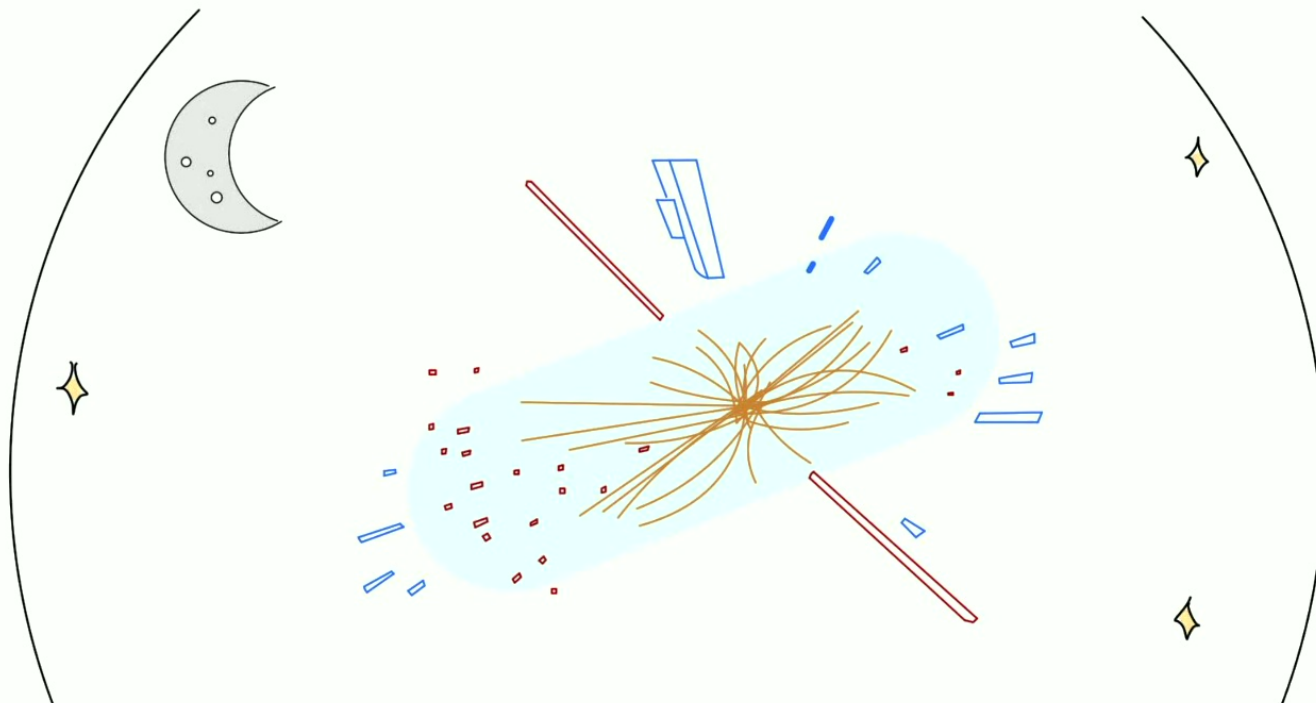
- Limit of UV complete theory
stringy constructions of AdS/CFT



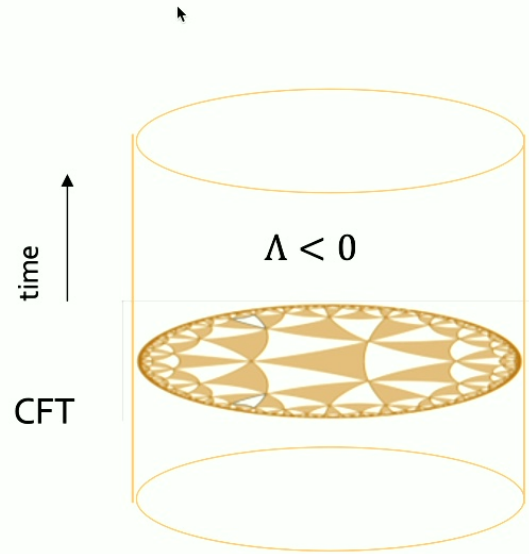
Gold Standard



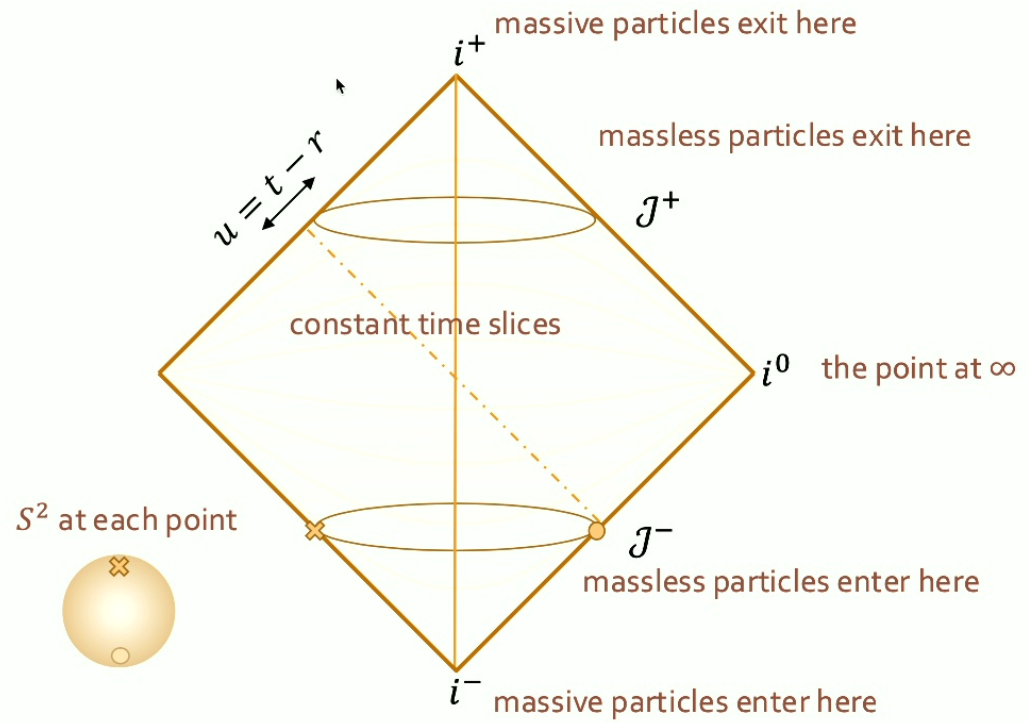
Celestial Holography proposes that scattering in asymptotically flat spacetimes is dual to a CFT living on the celestial sphere.



What's special
about flat?



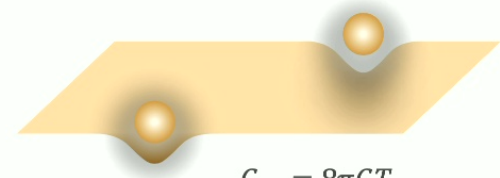
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What's special about flat?



$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$



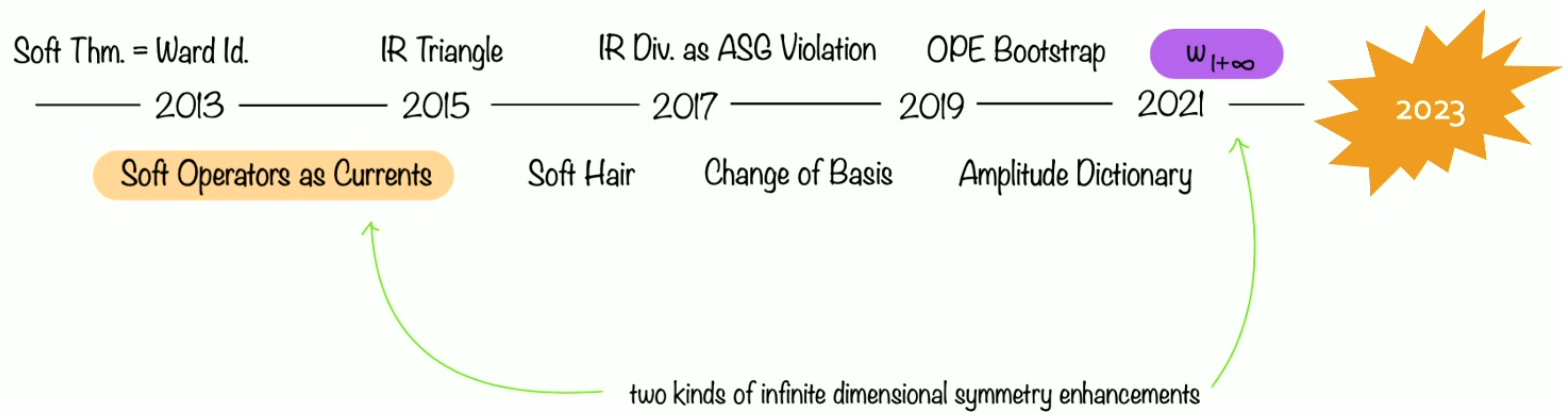
$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

How to identify the Asymptotic Symmetry Group:

1. Pick a gauge.
2. Identify physical falloffs.
3. Ask which diffeomorphisms preserve those falloffs.

BMS is an **infinite dimensional** enhancement of Poincare!

Celestial Holography proposes that scattering in asymptotically flat spacetimes is dual to a CFT living on the celestial sphere.



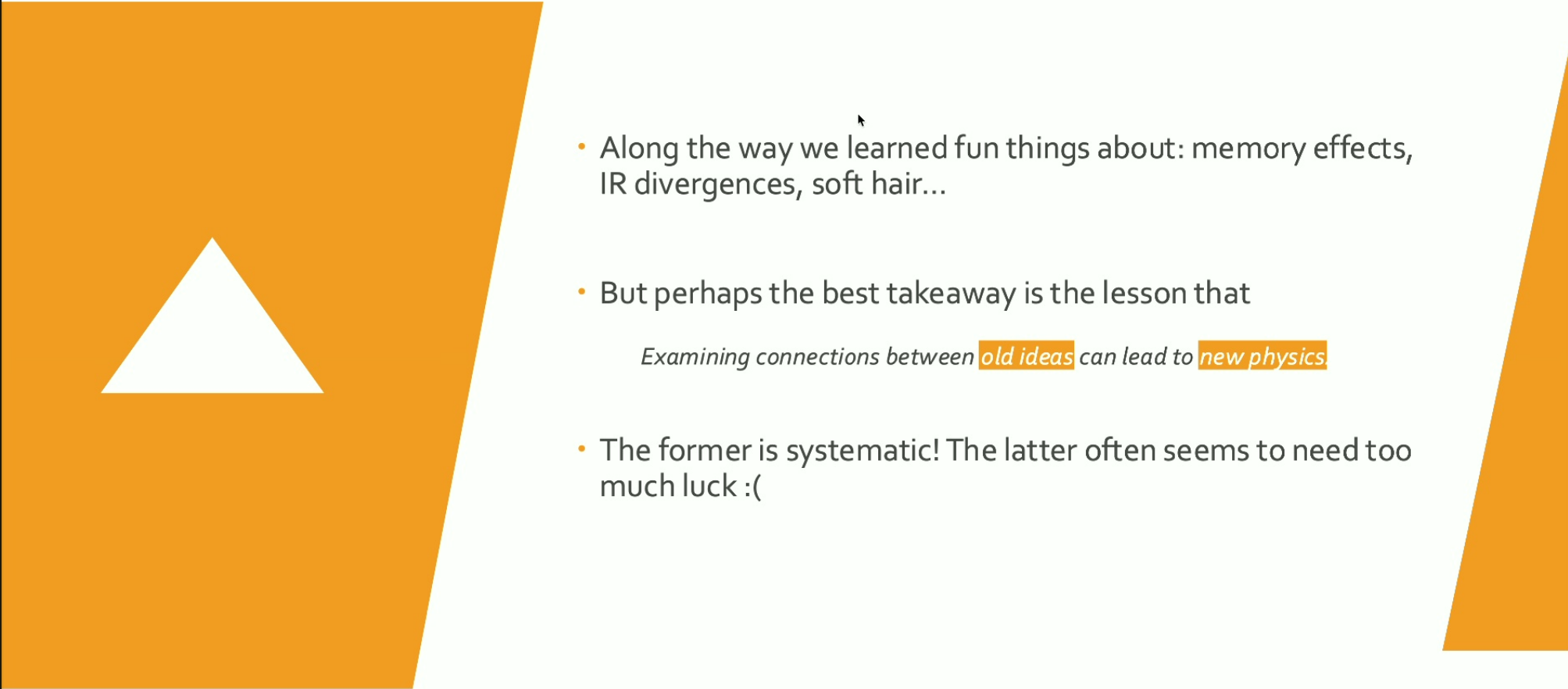


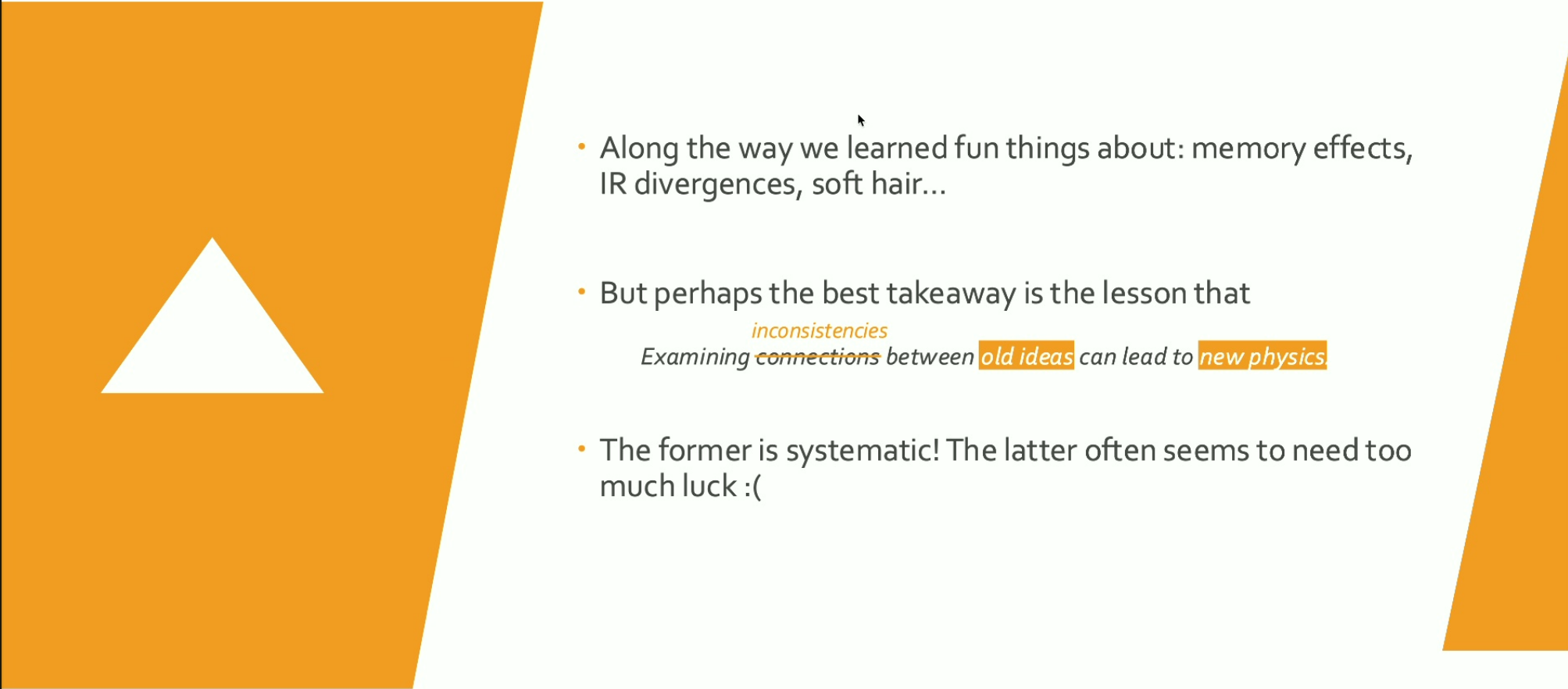
The program started as a **bottom up** approach, asking how the BMS group constrained scattering.

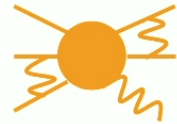
For a holographic dual living on null infinity the constraints point to a **2D CFT** structure.

Treating the dual as a 2D CFT led to recognizing collinear limits as encoding chiral algebras from twisted holography, providing a **top down** example.

→ **Simons Collaboration!**

- 
- Along the way we learned fun things about: memory effects, IR divergences, soft hair...
 - But perhaps the best takeaway is the lesson that
*Examining connections between **old ideas** can lead to **new physics***
 - The former is systematic! The latter often seems to need too much luck :(

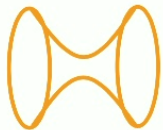
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Amplitudes

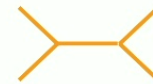
The S-matrix is a boundary observable.

Many subfields of hep-th nominally care about flat holography, even if its less explicit...



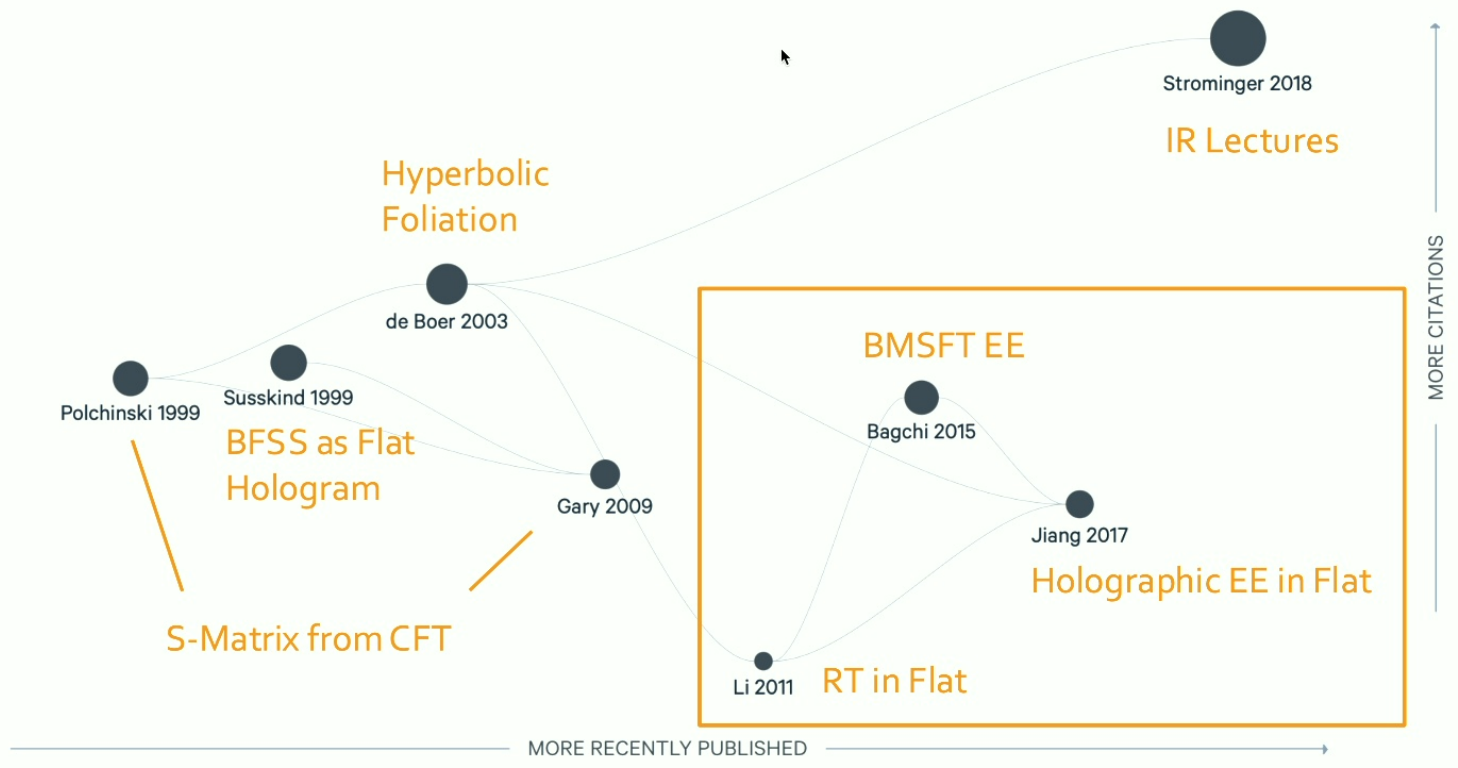
ItFromQubit

BH entropy has no Λ !



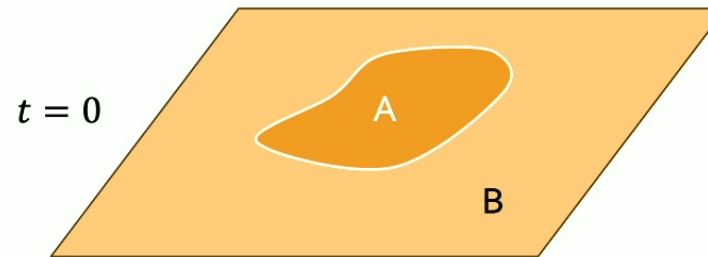
Bootstrap

S-matrix insights from CFTs.
What CFTs have bulk duals?

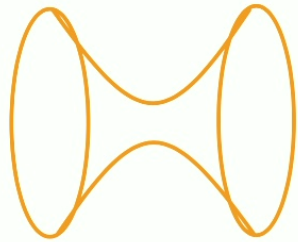


The RT Proposal

The **entanglement entropy** is the von Neumann entropy of the reduced density matrix $\rho_A = \text{Tr}_B(\rho_{AB})$ and measures entanglement between A and B.



$$\mathcal{S}(\rho_A) = -\text{Tr}[\rho_A \log \rho_A]$$

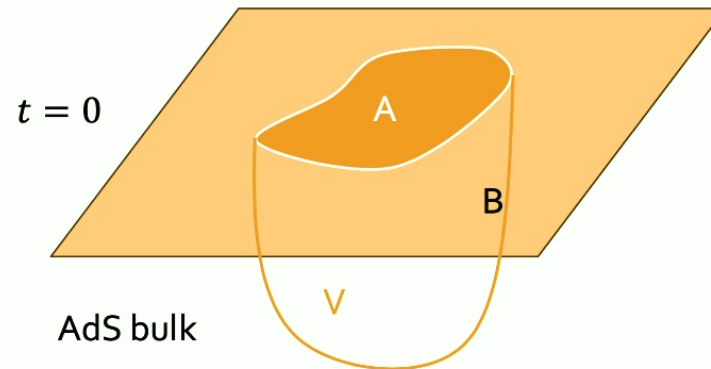


$$\frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

Geometry = Entanglement

The RT Proposal

For holographic CFTs the entanglement entropy is equal to the **minimal area surface** homologous to the boundary region.



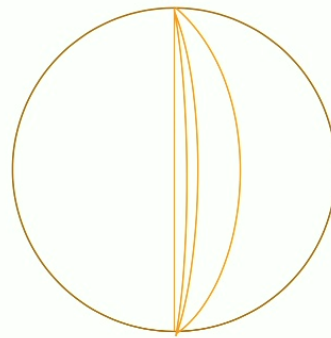
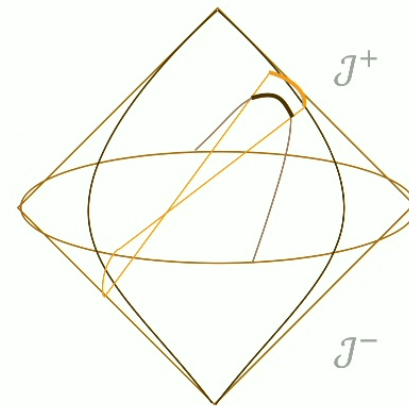
$$S(A) = \text{ext}_{V \sim A} \frac{A_V}{4G_N}$$

[Ryu, Takayanagi '06]

What's scary about flat?

Now a few things become strange when try to take the flat limit:

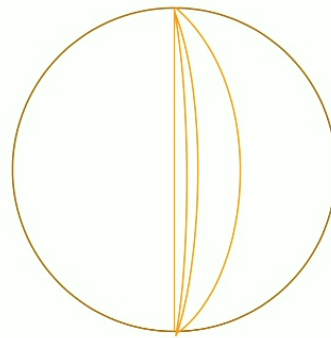
- Minimal surfaces have **restricted ends.**
 - For equal time slices they end on great circles.
 - Only null geodesics reach null infinity.


 i^0


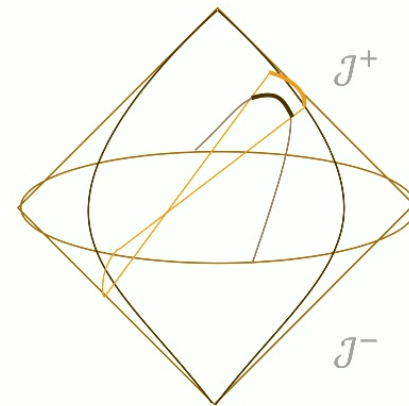
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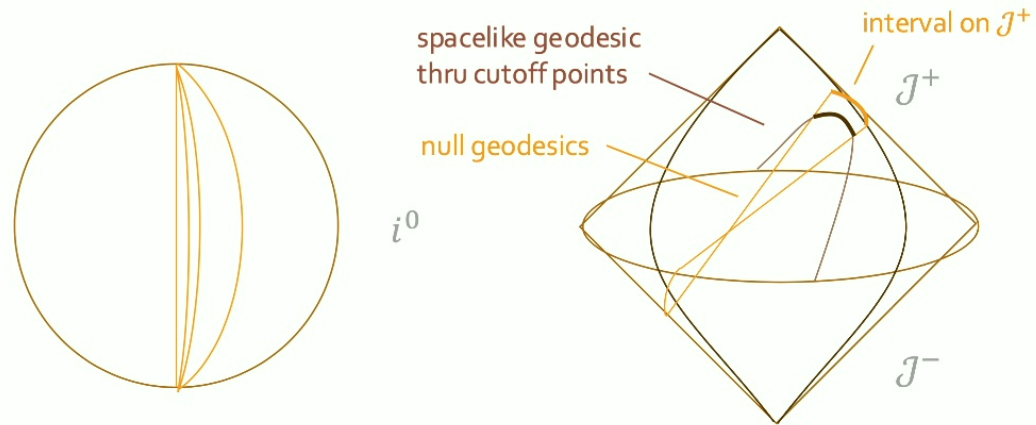
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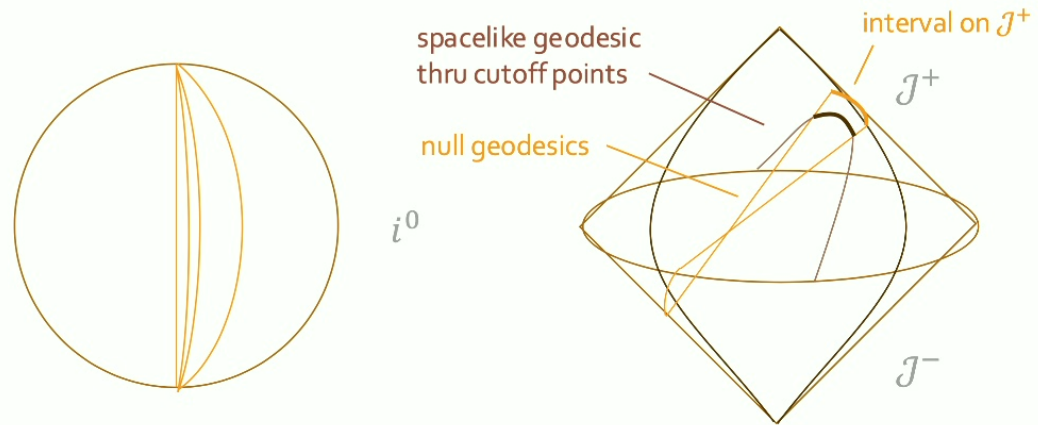
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- Minimal surfaces have **restricted ends**.
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- There is no longer a **natural cutoff scale**.
 - This effects how we regulate the minimal surfaces
 - and leads to problems identifying a finite central charge.

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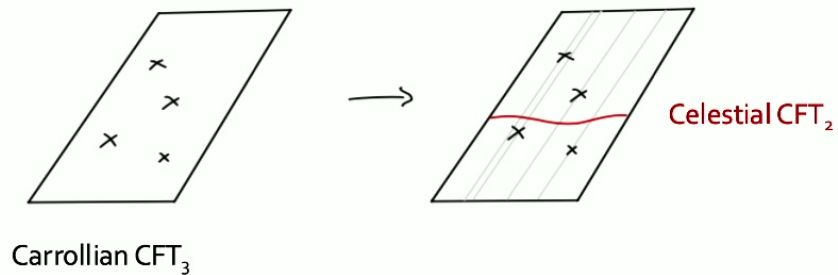
However, there is a holographic EE proposal!

Rant I:

Carrollian vs (=) Celestial

There are two natural ways to think about the flat hologram living at null infinity.

- The **Carrollian** description lives on a codimension 1 null surface
- The **Celestial** description lives on a codimension 2 spacelike surface that is a cross section of this null cone.

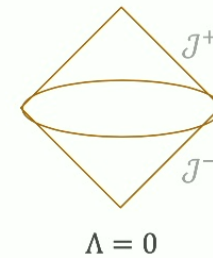
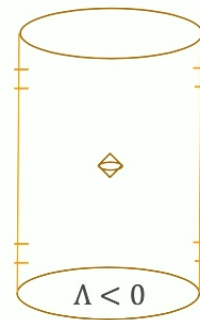


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Rant II: 3D vs 4D

Which bulk dimension is your favorite?

4D

- You have more interesting soft physics.
- It's the real world!

3D

- Everything else is easier.
BMS₃ from $\text{Vir} \times \text{Vir}$, non-zero c , simpler EE computation, simpler RT surfaces ...

Rant II: 3D vs 4D

Starting from the **Virasoro** generators in AdS₃/CFT₂

$$\begin{aligned}[\mathcal{L}_n^\pm, \mathcal{L}_m^\pm] &= (n-m)\mathcal{L}_{n+m}^\pm + \frac{c^\pm}{12}n(n^2-1)\delta_{n+m,0} \\ [\mathcal{L}_n^+, \mathcal{L}_m^-] &= 0\end{aligned}$$

we can form the linear combinations

$$\begin{aligned}\mathcal{L}_n &= \mathcal{L}_n^+ - \mathcal{L}_n^-, & \mathcal{M}_n &= \frac{1}{l}(\mathcal{L}_n^+ + \mathcal{L}_n^-) & \leftarrow \text{here be dragons} \\ c_L &= c^+ - c^-, & c_M &= \frac{1}{l}(c^+ + c^-)\end{aligned}$$

which limit to the **BMS₃** algebra when we take l large.

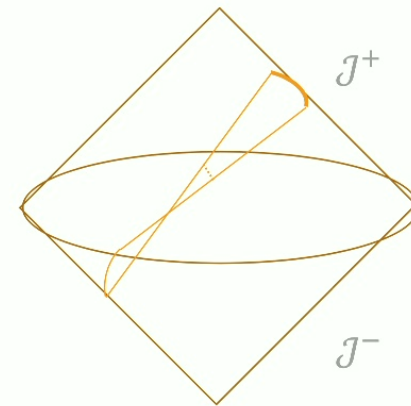
$$\begin{aligned}[\mathcal{L}_n, \mathcal{L}_m] &= (n-m)\mathcal{L}_{n+m} + \frac{c_L}{12}n(n^2-1)\delta_{n+m,0} & \leftarrow \text{rotations} \\ [\mathcal{L}_n, \mathcal{M}_m] &= (n-m)\mathcal{M}_{n+m} + \frac{c_M}{12}n(n^2-1)\delta_{n+m,0} & \leftarrow \text{translations} \\ [\mathcal{M}_n, \mathcal{M}_m] &= 0\end{aligned}$$

RT proposal in flat

- Using some Galilean CFT tricks [Bagchi et al '14] computed the entanglement entropy of an interval in a BMSFT at null infinity.

$$S_{EE}^{BMSFT}(A) = \frac{c_L}{6} \ln \left(2 \sin \frac{l_\phi}{2} \right) + \frac{c_M}{12} l_u \cot \frac{l_\phi}{2}$$

- In [Jiang et al '17] Wei Song and friends then identified a geometric quantity that matches this answer.

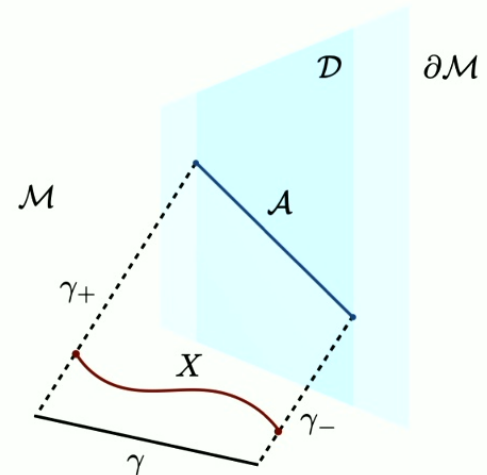


RT proposal in flat

The **Swing Proposal** states that the BMSFT EE can be computed from the minimal length spacelike geodesic between **radial** null geodesics tangent to the approximate **modular flow**.

$$S_{\mathcal{A}} = \min_{X_{\mathcal{A}} \sim \mathcal{A}}^{\text{ext}} \frac{\text{Area}(X_{\mathcal{A}})}{4G},$$
$$X_{\mathcal{A}} = X \cup_{p \in \partial \mathcal{A}} \gamma(p)$$

[Jiang et al '20]

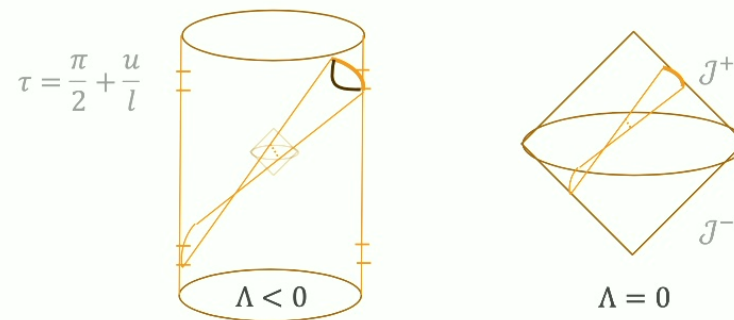


But this sounds very different than RT proposal!

Lifting to AdS

This prescription would seemingly lift to AdS₃.

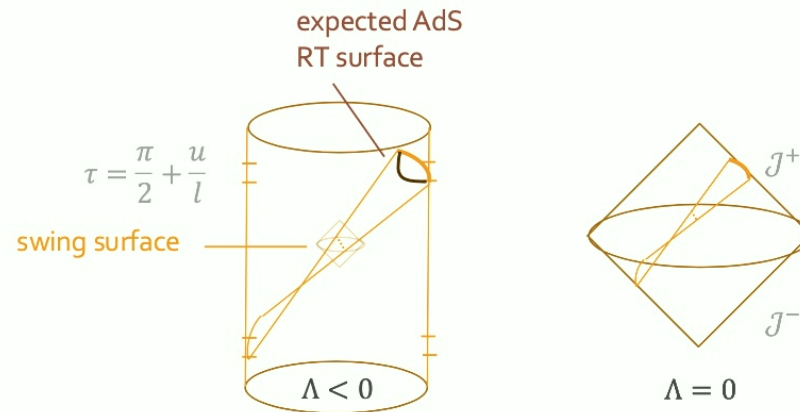
- Null infinity maps to a small window near $\tau = \frac{\pi}{2}$.
- Can send in radial geodesics and construct the swing.



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Lifting to AdS

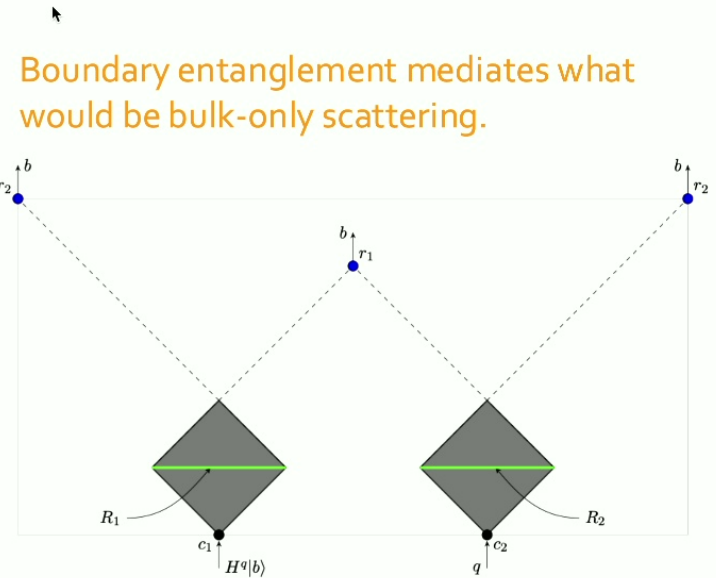
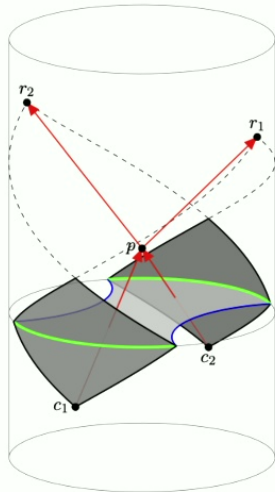
There seem to be various red flags:

- The proposal does not match what we expect in AdS!
- The BMSFT EE is zero for an equal u cut!
- The BMSFT is likely not even positive?
- The radial geodesics should not be so special?!



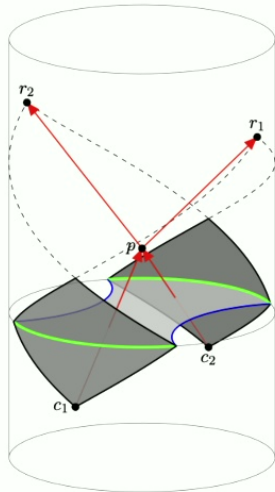
What gives?

Flat space physics naturally lifts to scattering questions in AdS. Thankfully, Alex and friends have given us some relevant information theoretic tools.

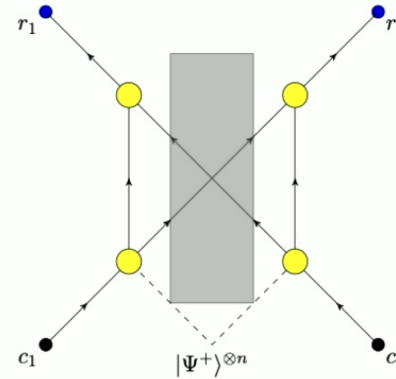


[May '19]

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Boundary entanglement mediates what would be bulk-only scattering.

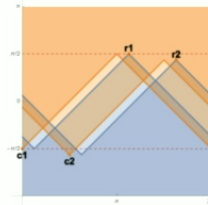


[May '19]

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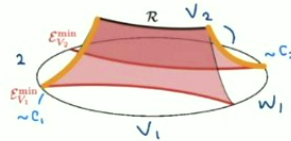
In global AdS₃ the BMSFT EE is the mutual information of a near-forward scattering configuration.

- We can identify the swing as the ridge surface in [May '19]

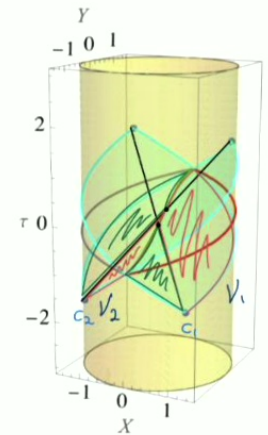


$$|S_{EE}^{BMSFT}| = \frac{1}{2} \max\{I(V_1 : V_2), I(W_1 : W_2)\}$$

- In Alex's proof this would generally be a lower bound.

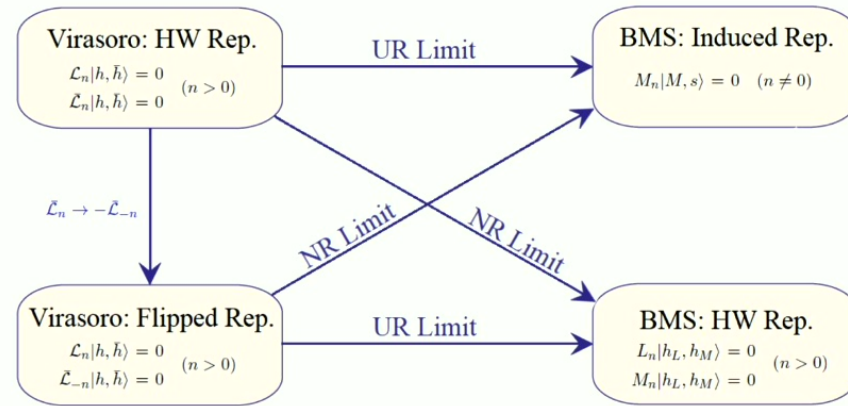


$$I(V_1 : V_2) \geq \frac{\text{area}(\mathcal{R})}{2G_N} + O(1)$$



2

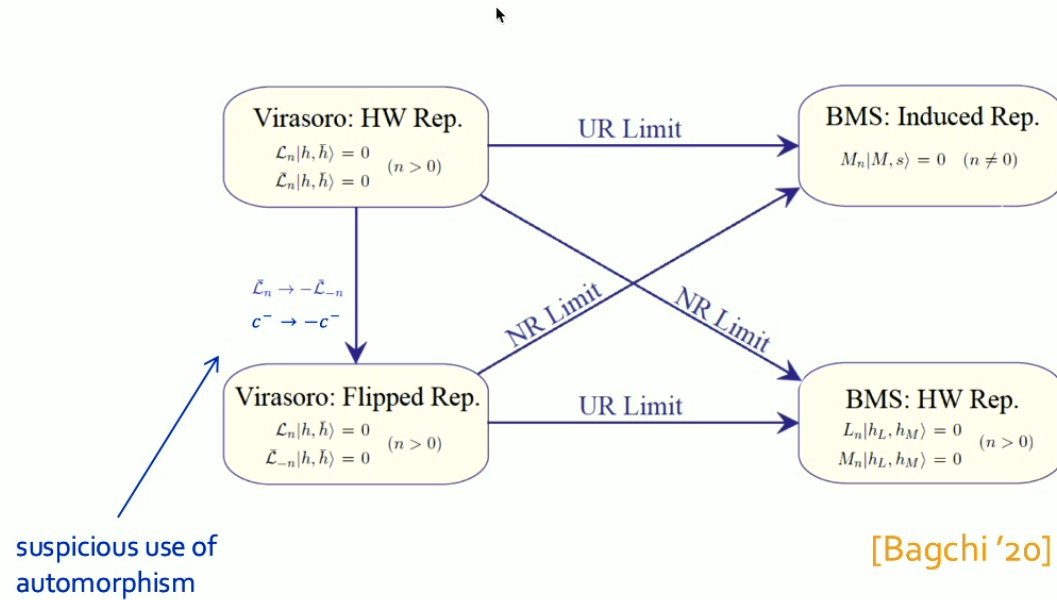
The BMSFT entanglement entropy computation uses some trickery to handle the lowest weight conditions.



[Bagchi '20]

2

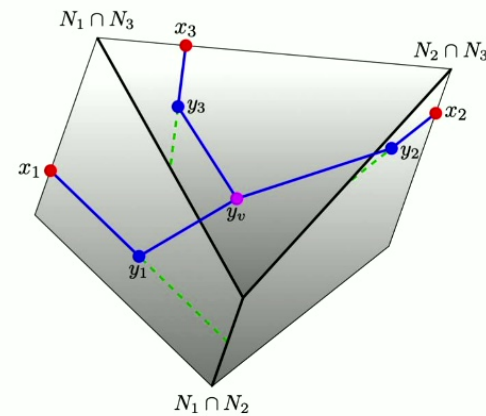
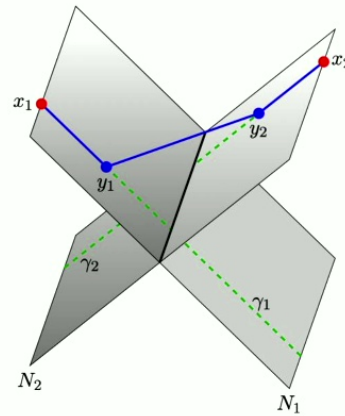
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Lifting to AdS

While we might be happy to point out mistakes some things work!

- The swing surface is still equal to the $|\text{EE}'|$.
- [Hijano '17] can reproduce BMSFT correlators from null geodesic networks.

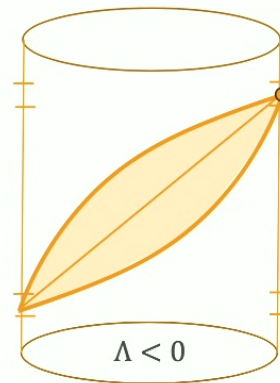


What?!

3

We can spoof BMSFT correlators with highly spinning primaries near the bulk point configuration.

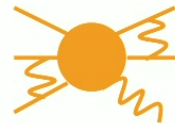
- Light ray operators on the boundary are a natural new candidate for BMSFT primaries.
- Light ray operators through the bulk are linear superpositions.



$$\langle \phi(u_1, \phi_1) \phi(u_2, \phi_2) \rangle \propto \left(\frac{\sin\left(\frac{\phi_1 - \phi_2}{2} + \frac{u_1 - u_2}{2L}\right)}{\sin\left(\frac{\phi_1 - \phi_2}{2} - \frac{u_1 - u_2}{2L}\right)} \right)^J$$

$$\lim_{L \rightarrow \infty} \left(\frac{\sin\left(\frac{\phi_1 - \phi_2}{2} + \frac{u_1 - u_2}{2L}\right)}{\sin\left(\frac{\phi_1 - \phi_2}{2} - \frac{u_1 - u_2}{2L}\right)} \right)^{-L\xi} = e^{-\xi(u_1 - u_2) \cot\left(\frac{\phi_1 - \phi_2}{2}\right)}$$

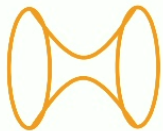
[wip w/ Jackie & Rob '25]



Amplitudes

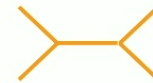
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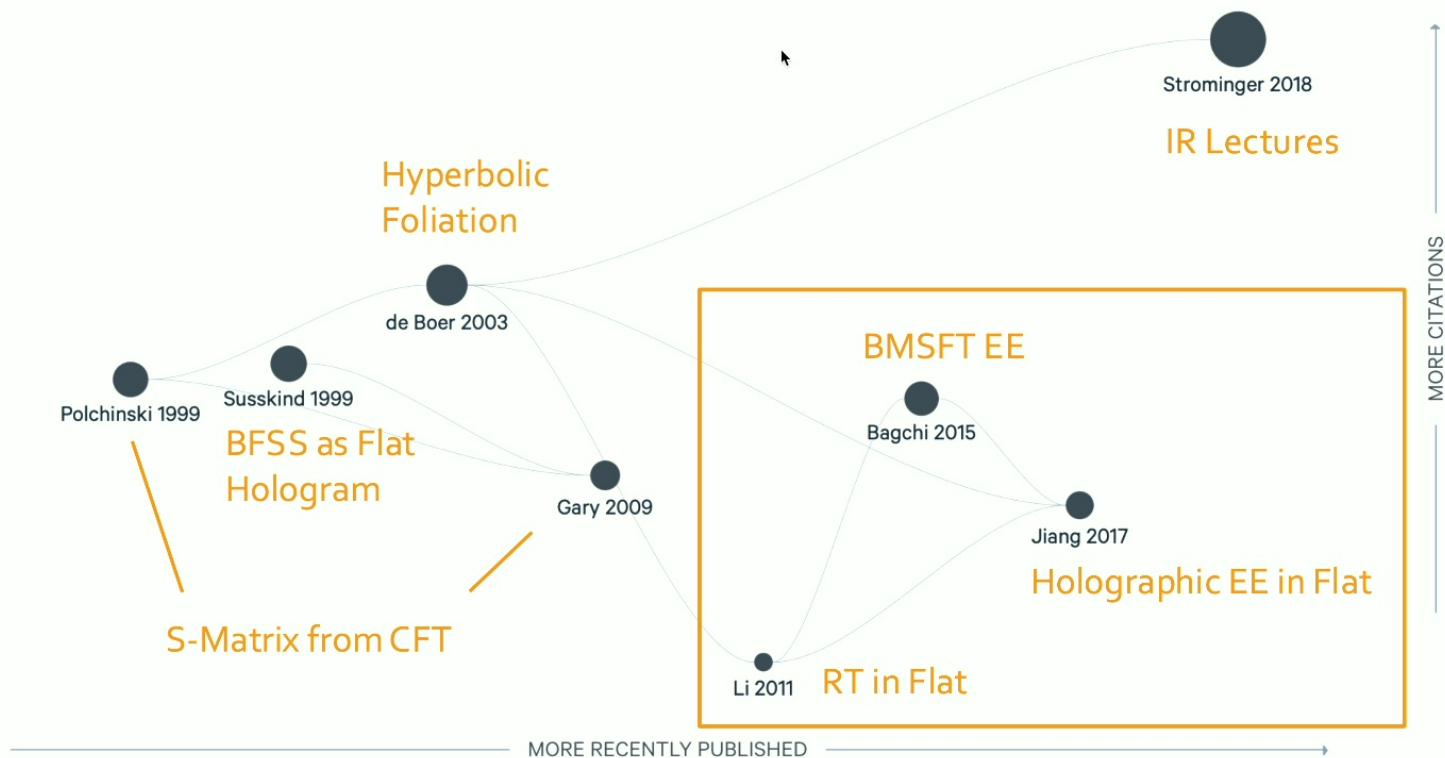
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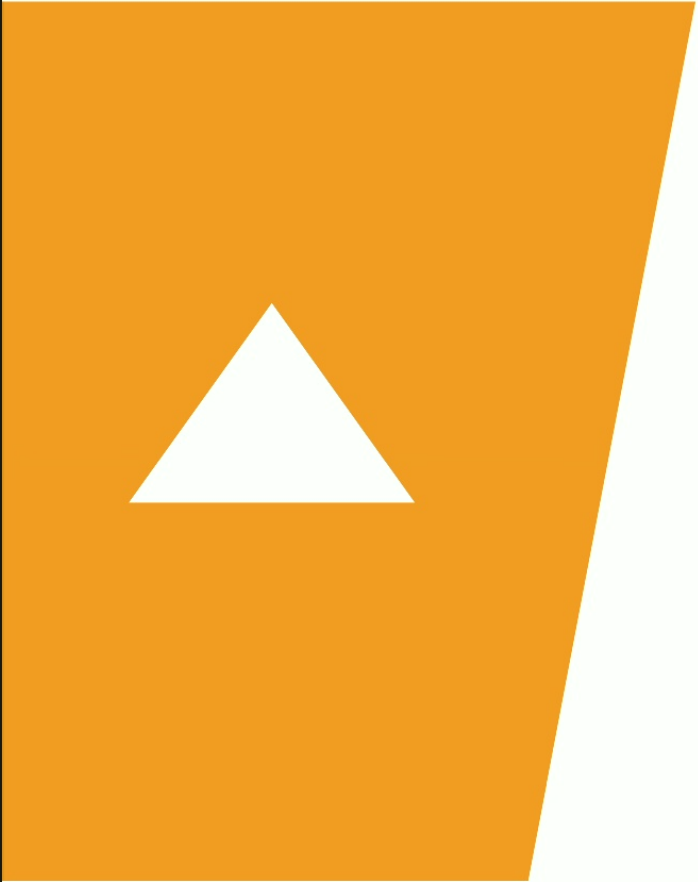
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+ is fun!

Thank You!