

**Title:** Lecture - Beautiful Papers

**Speakers:** Pedro Vieira

**Collection/Series:** Beautiful Papers - October 7, 2024 - January 31, 2025

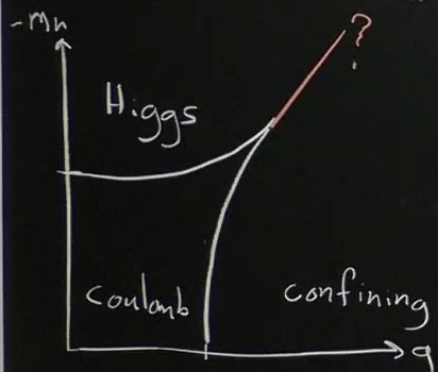
**Subject:** Other

**Date:** November 04, 2024 - 9:15 AM

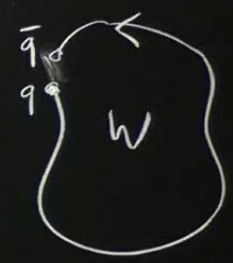
**URL:** <https://pirsa.org/24110032>

# DISCUSSION GUIDE

$$\mathcal{L} = \text{tr}_F \left[ -\frac{1}{2g^2} F_{\mu\nu}^A F_{\mu\nu}^A - m_h |h|^2 - \lambda |h|^4 \right]$$



$$\begin{aligned} & (\mathcal{D}_\mu h)^2 \\ & \cup \\ & A_\mu A_\mu \\ & h = v + \tilde{h} \end{aligned}$$



$$\langle W \rangle \sim e^{-\mu P} \quad \langle W \rangle \sim e^{-\mu P}$$

$$\mathcal{L}_F = \bar{\psi} \not{\partial} \psi + m(\bar{\psi} \psi + \psi \bar{\psi})$$

$$\mathbb{Z}[e^{i\theta} m] = \mathbb{Z}[m] e^{\frac{i\theta}{8\pi^2} \int_M R \wedge R}$$

$$U(1)^A: \psi \rightarrow e^{i\theta/2} \psi$$

T: enforces  $m \in \mathbb{R}$

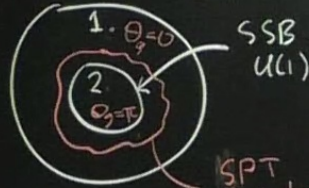
$$m \rightarrow e^{i\theta} m$$

$$\mathbb{Z}_2^A: m \rightarrow -m$$

SU(2)

$$e^{\frac{i\theta}{4\pi^2}} \int_M R \wedge R$$

Neutron Star



1. confining

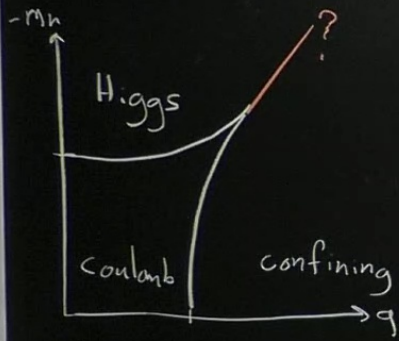
2. asymptotically free quarks

$$\mathcal{L} = \mathcal{E} \det(h_i^a)$$

$$\Rightarrow U(1)_B \text{ NGB}$$

• UW resources

$$\mathcal{L}_h = \text{tr}_F \left[ -\frac{1}{2g^2} F_{\mu\nu}^A F_{\mu\nu}^{A'} - m_h |h|^2 - \lambda |h|^4 \right]$$



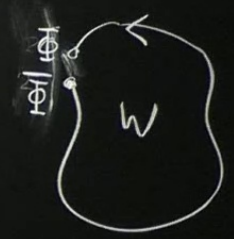
$$\langle W \rangle \sim e^{-\mu P} \quad \langle W \rangle \sim e^{-\mu P}$$

$$\begin{aligned} & (D_\mu h)^2 \\ & \psi \\ & A h A h \\ & h = v + \tilde{h} \end{aligned}$$

$$G = U(1)$$

$$\phi \rightarrow e^{i\alpha} \phi$$

$$\Phi \rightarrow e^{i2\alpha} \Phi$$



$$W = e^{i\int A} \rightarrow W e^{i\alpha}$$

$$\mathcal{L}_f = \bar{\psi} \not{\partial} \psi$$

$$U(1)^k: \psi \rightarrow e^{i\alpha} \psi$$

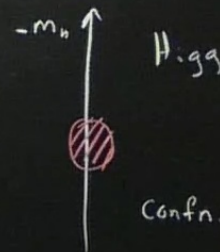
$$\mathbb{Z}_2^A: m \rightarrow -m$$

$$\mathcal{L}_f = \bar{\Psi} \not{\partial} \Psi + m_f (\bar{\Psi} \Psi + \Psi \bar{\Psi})$$

$$U(1)_f: \Psi \rightarrow e^{i\theta/2} \Psi \quad T: \text{enforces } m \in \mathbb{R}$$

$$\mathbb{Z}_2^A: \left. \begin{array}{l} m \rightarrow e^{i\theta} m \\ m \rightarrow -m \end{array} \right\} \Rightarrow \theta = 0, \pi$$

$$\mathcal{Z}[e^{i\theta} m] = \mathcal{Z}[m] e^{\frac{i\theta}{8\pi^2} \int_M R \wedge R}$$



$$SU(2)_c \times SU(2)_f$$

$$\Psi_a^i: 2 \oplus 2 = 3 \oplus 1$$

$$\langle h_a^i \rangle = v \delta_a^i$$

Vacuum Structure of  
Charge -k 2D QCD

Based on:

[1912.10064] Armoni, Sugimoto.

[2009.07567] Komargodski et al



## Part I : The Model

$$S_{QED} = \int d^4x \left\{ \frac{-1}{4e^2} F^2 + i\bar{\psi} \gamma^\mu (\partial_\mu + i\kappa A_\mu) \psi \right.$$

$$\kappa \in \mathbb{Z} > 0$$

$$\text{Reps: } \gamma^0 = \sigma^1, \gamma^1 = i\sigma^2 \quad \text{and} \quad \psi = \begin{pmatrix} \psi_R \\ \psi_L \end{pmatrix}$$

with  $\psi_{L/R}$  handed spinors.

$$\text{U(1) Gauge Trans: } A \rightarrow A + d\lambda, \quad \psi \rightarrow e^{-i\kappa\lambda} \psi \\ \lambda \sim \lambda + 2\pi$$

$$\text{flux Quantization: } \frac{1}{2\pi} \int F \in \mathbb{Z}.$$

Classical Symmetries:



Reps:  $\gamma^0 = \sigma^1$ ,  $\gamma^1 = i\sigma^2$  and  $\psi = (\psi_L)$

with  $\psi_{L/R}$  handed spinors.

U(1) Gauge Trans:  $A \rightarrow A + d\lambda$ ,  $\psi \rightarrow e^{-ik\lambda} \psi$   
 $\lambda \sim \lambda + 2\pi$

Flux Quantization:  $\frac{1}{2\pi} \int F \in \mathbb{Z}$ .

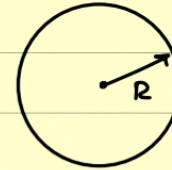
Classical Symmetries:  $\mathbb{Z}_2$

$\mathbb{Z}_2$  Axial for  $e^{i\alpha} \in \text{U(1)A} / \mathbb{Z}_2$ .

$\psi_L \rightarrow e^{-i\alpha} \psi_L$ ,  $\psi_R \rightarrow e^{i\alpha} \psi_R$

Anomaly free symmetry:  $\mathbb{Z}_k^{\text{Axial}}$

2) 1-form symmetry :  $\mathbb{Z}_k^{1\text{-form}}$



$$A \rightarrow A + \frac{1}{k} dz, \quad \psi \rightarrow e^{-i\theta} \psi$$

$$\text{with: } z(x^0, x^1 + 2\pi R) = z(x^0, x^1) + 2\pi l$$

$$\text{and: } l = 1, 2, \dots, k.$$



NOT a gauge transformation  $\forall l \neq k$ .

$$\text{e.g: } z(x) = \frac{l}{R} x^1 \Rightarrow A_1 \rightarrow A_1 + \frac{l}{kR}$$

But  $A_1 \rightarrow A_1 + \frac{1}{R}$  is a gauge transformation.

Action on Wilson Loops

$$\mathbb{Z}_k^{1\text{-form}} : \omega \rightarrow e^{i2\pi l/k} \omega$$

Working in  $A_0 = 0$  gauge:

$$S_{\text{ce}} = \int d^3x \left\{ \frac{1}{2e^2} (\partial_0 A_i)^2 + \frac{1}{8\pi} (\partial_\mu \varphi)^2 + \frac{k}{2\pi} \varphi A_i \right\}$$

Canonical Momenta:

$$\pi_{A_i} = \frac{1}{e^2} \partial_0 A_i + \frac{k}{2\pi} \varphi \quad ; \quad \partial_i \pi_{A_i} = 0$$

$$\pi_{\varphi} = \frac{1}{4\pi} \partial_0 \varphi$$

Hamiltonian:

$$H = \int dx^i \left\{ \frac{e^2}{2} \left( \pi_{A_i} - \frac{k}{2\pi} \varphi \right)^2 + 2\pi \pi_{\varphi}^2 + \frac{1}{8\pi} (\partial_i \varphi)^2 \right\}$$

Algebra:  $[A_i, \pi_{A_i}] \propto \delta$   
 $[\varphi, \pi_{\varphi}] \propto \delta$

Energies of  $|\theta\rangle$

$$E(\theta) = \langle \theta | H | \theta \rangle \propto \tilde{f}(\theta) M_0 + (\dots)$$

Indep of  $(\theta, M_0)$

where  $M_0$  is a mass perturbation:

$$M_0 \bar{\psi} \psi.$$

①  $M_0 = 0 \Rightarrow |\theta + 2\pi i\rangle$  all degenerate

②  $M_0 \neq 0$  but small, lifted degeneracy

For  $\theta \neq \pi$ , single lowest energy state

## Part II: Confinement of Wilson Loops

Focusing on  $M_0 = 0$ :

①  $\hat{U}$  order parameter:

$$\langle \theta | \hat{U} | \theta \rangle = e^{i\theta/k} \neq 0 \quad \text{but.} \quad \text{⚡}$$

mixed anomaly  $\Rightarrow \hat{V} \hat{U} \hat{V}^{-1} = e^{-2\pi i/k} \hat{U}$ .

Thus,  $\langle \theta | \hat{U} | \theta \rangle$  allows us to identify the vacuum.

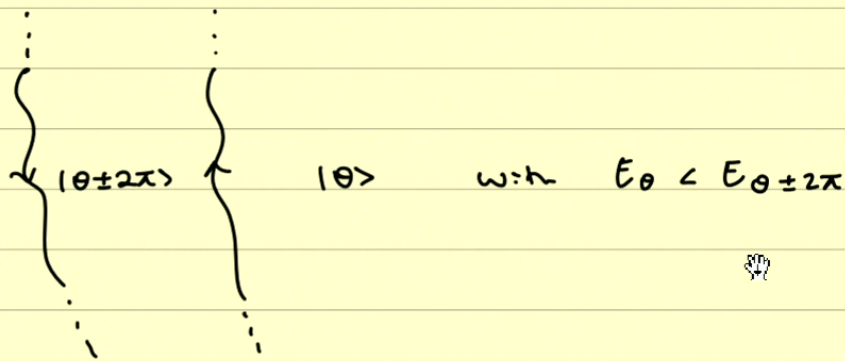
② Action on Wilson Loops

$$\hat{U} \mathcal{W} \hat{U}^{-1} = e^{i2\pi/k} \mathcal{W}.$$

$\Rightarrow \mathcal{W}$  op shifts  $\theta \rightarrow \theta \pm 2\pi$  depending on orientation.

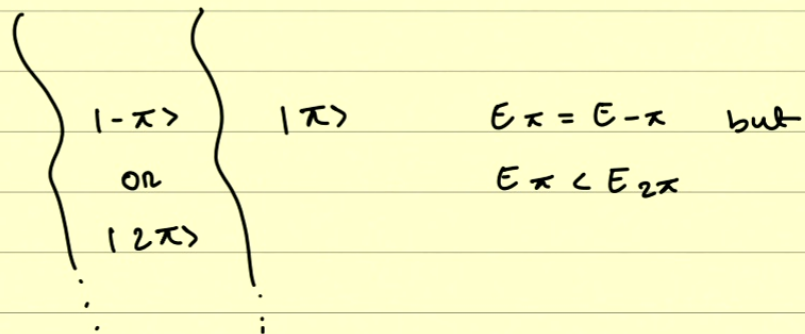
## Confinement vs. Deconfinement

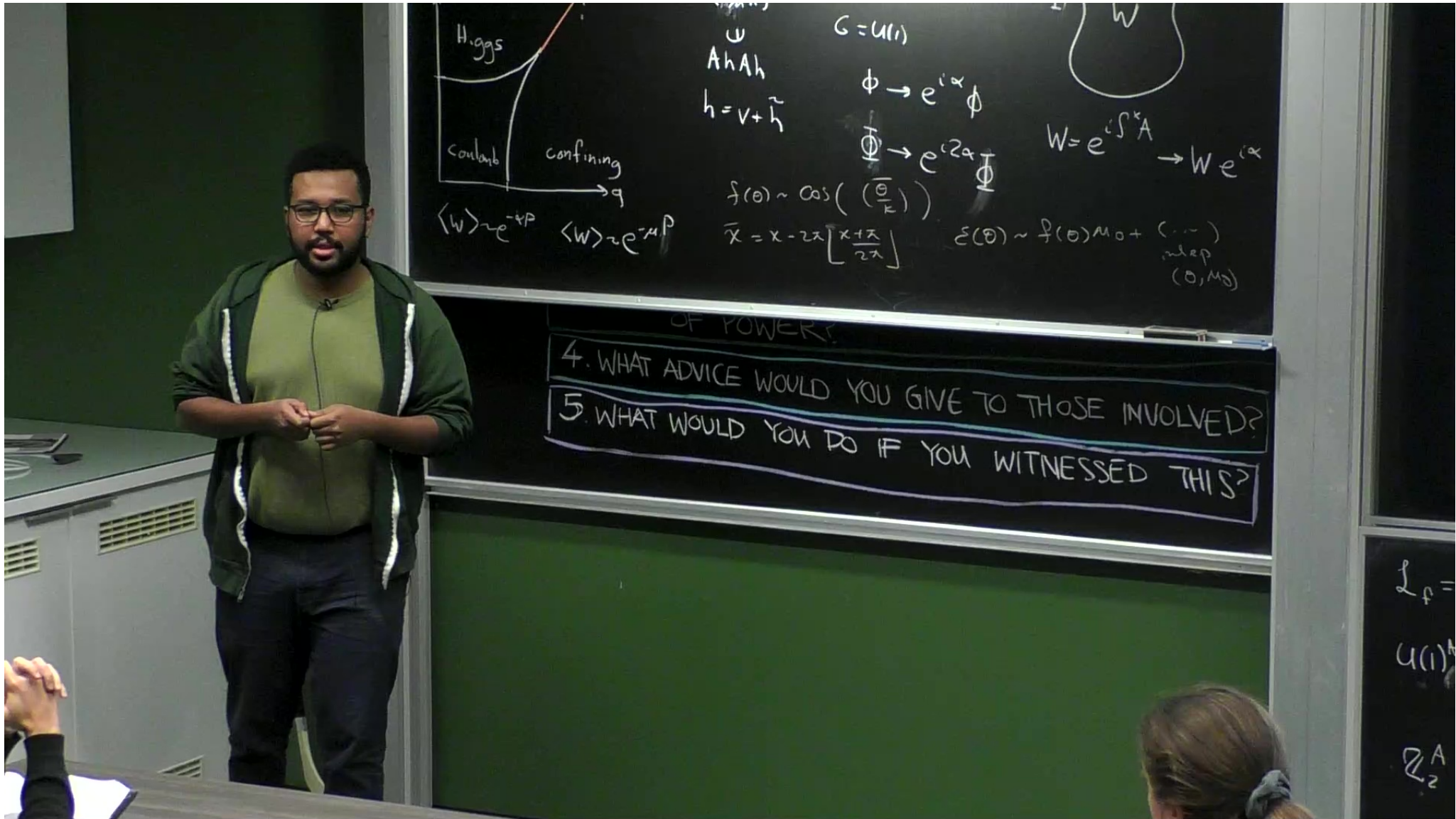
$\theta = 0$



$\Rightarrow \langle W \rangle \sim e^{-\text{Area}} \Rightarrow$  Confined.

$\theta = \pi$ , Two degenerate vacua:  $|\pm\pi\rangle$





$\langle W \rangle = \langle \epsilon \rangle$   $[2\pi]$   $(0, M_0)$

1) Melanese 93'

2) Ambjorn 15'



$$\Delta V^\alpha = -\Gamma_{\mu\beta}^\alpha V^\beta \Delta x^\mu$$

$$V'^\alpha = U^\alpha_\beta V^\beta$$

$$U^\alpha_\beta = \frac{\partial x^\alpha}{\partial x'^\beta}$$

$$U \rightarrow \Lambda U \Lambda^{-1} \Rightarrow \Lambda^\alpha_\beta U^\beta \Lambda^\gamma_\delta \approx \Lambda^\alpha_\gamma U^\delta$$

$$W = -U + \epsilon U$$



$$\int \mathcal{D}g e^{\frac{i}{\hbar} S[g]}$$

$$S[g] = \int dt \left[ L^{(2)} \dot{h} \dot{h} + \sqrt{G} L^{(3)} \dot{h}^3 + G L^{(4)} \dot{h}^4 \right]$$

$$g = \eta + \sqrt{G} h$$



$$\langle W^{(2)} \rangle = \int \int_{(C)} \langle \Gamma \Gamma \rangle = 0$$

$$W = +i U(A) = 0 \Rightarrow U(C) = \mathbb{1}$$

$\in \text{SU}(4)$

Albay, Malbocena  
07

$$g \begin{pmatrix} C_{01} & 2\theta_1 \\ -2\theta_1 & C_{01} \\ & & 0 \end{pmatrix}$$

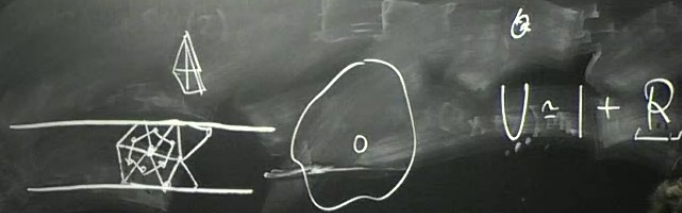
• UW res

5 [g]

$$S[g] = \int L^{(2)} h h + \alpha L^{(3)} h^3 + \beta L^{(4)} h^4$$

$$\langle W^{(2)} \rangle = \int_{(c)} \langle \Gamma \Gamma \rangle = 0$$

$$W = \tau U = 0 \Rightarrow U(c) = 1$$



$$g \begin{pmatrix} c_{01} & s_{01} & 0 \\ -s_{01} & c_{01} & s_{02} \\ 0 & -s_{02} & c_{02} \end{pmatrix} g^{-1} + \tau R R R R$$

$$+ \tau U = (c_{01} + c_{02})$$

• UW resources