

Title: Lecture - QFT II, PHYS 603

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Collection/Series: Quantum Field Theory II (Core), PHYS 603, November 12 - December 11, 2024

Subject: Condensed Matter, Particle Physics, Quantum Fields and Strings

Date: November 28, 2024 - 9:00 AM

URL: <https://pirsa.org/24110016>

$G = SU(2)$ Gauge Group

Gauge Field $A \rightarrow \{A_\mu^a(x)\}$

Covariant Derivative $D_\mu A_\nu = \partial_\mu A_\nu - i[A_\mu, A_\nu]$

Action $S_{YM}[A] = -\frac{1}{2g^2} \int dx \text{Tr}[F_{\mu\nu} F^{\mu\nu}]$

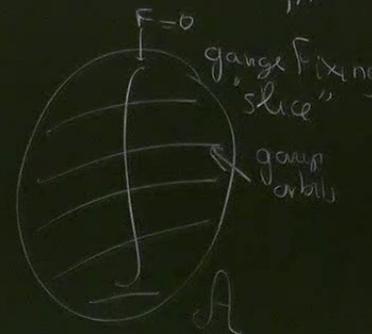
Field Strength $F_{\mu\nu} = [D_\mu, D_\nu] =$ Lorentz Gauge

Gauge Fixing $F[A] = \{F^a_\mu(x) = \partial^\mu A^a_\mu(x)\}$

$$\int D[A] e^{i S_{YM}[A]} \delta[F[A]] \cdot ?$$

Dirac

$$F[A] = 0 \quad \prod_{a,\mu} \delta[F^a_\mu(x)]$$



How to write 1 as an integral?

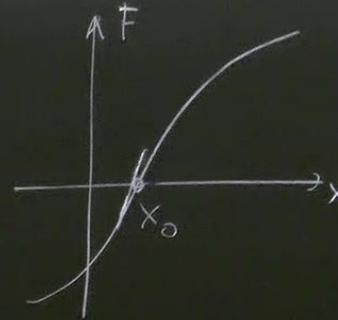
① $\int_{-\infty}^{+\infty} dx \delta(x-x_0) = 1$

$F(x) \mathbb{R} \rightarrow \mathbb{R}$

$F'(x) > 0$

only one x_0 , $F(x_0) = 0$

← derivative



if $F' < 0$

$1 = \int dx \delta(F(x)) F'(x_0) = \int dx \delta(F(x)) |F'(x)| = 1$

② $F(x) \mathbb{R}^N \rightarrow \mathbb{R}^N$

N variables

a unique point in \mathbb{R}^N

$x = (x_1, \dots, x_N)$

$F(x) = (F^1(x), \dots, F^N(x))$

$F'(x) = \left\{ \frac{\partial F^i}{\partial x_j}(x) \right\} > 0, F(x_0) = 0$

$N \times N$ matrix

$$F: \mathbb{R}^N \rightarrow \mathbb{R}^N$$

Lee Group theory

$$\int d^N x \delta(x - x_0) = \int d^N x \delta(F(x)) |\det(F'(x))| = 1$$

\uparrow
N-dim "Dirac Funct"

$$F'(x) = (F^i_j(x)) \quad (F^i_j)' = \frac{\partial F^i}{\partial x^j}$$

$N \times N \text{ matrix}$ partial derivative

$$(3) F: \text{Group } G \rightarrow \mathbb{R}^N$$

$\dim N \quad \text{SU}(2) \rightarrow \mathbb{R}^3$

what is $F'(g)$?

Lee Group theory

G has an invariant measure

Haar measure $\mathbb{1} \rightarrow \mathbb{1} + i\delta\alpha + \dots = \mathbb{1}(\mathbb{1} + S_3)$



$$\delta\alpha = \alpha_a t^a$$

$$d\mu(g) = \prod_a da$$

$$\text{tr}(t_a t_b) = \frac{1}{2} \delta_{ab}$$

$$g \xrightarrow{g'} g'g \quad \text{or} \quad g \cdot g'$$

L R

$$F: G \rightarrow \mathbb{R}^N \quad N = \dim \text{ of the group}$$

$$g \rightarrow F(g) = (F_a(g))_{a=1, \dots, N} \quad \delta \alpha = \delta \alpha_b t^b$$

b) $g \rightarrow g(1 + i \delta \alpha)$ inf. deformation of g

$$F_a(g(1 + i \delta \alpha)) - F_a(g) = \delta F_a = \frac{\partial F_a(g(1 + i \alpha))}{\partial \alpha_b} \delta \alpha_b +$$

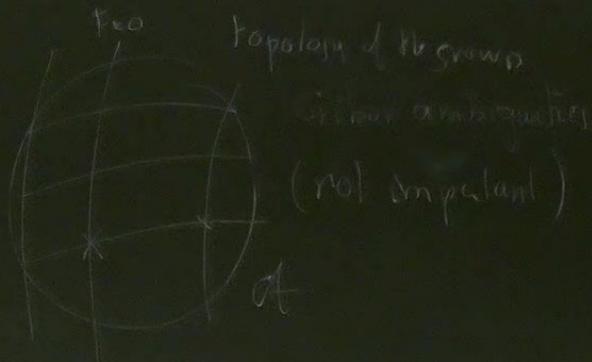
define $F'_{ab}(g) = \frac{\partial F_a(g(1 + i \alpha))}{\partial \alpha_b} \Big|_{\alpha=0}$

$N \times N$ matrix

$F: G \rightarrow \mathbb{R}^N$ a unique g_0 s.t. $F(g) = 0$

$$\int_G d\mu_{\text{Haar}}(g) \delta_{\text{Haar}}(g, g_0) = \int_G d\mu(g) \delta[F(g)] \left| \det [F'(g)] \right| = 1$$

δ -function on \mathbb{R}^N $N \times N$ matrix
 δ -function on the group



$$\boxed{[F'(g)] = 1}$$

N x N matrix

an assumption

Gauge Fixing

$$G = SU(2) \rightarrow \mathcal{G} = \sum_{a=1}^3 G_a \tau_a$$

start from A $F[A] \neq 0$
 gauge field configuration

gauge transf. $A \rightarrow A_g = g A g^{-1} + \partial g \cdot g^{-1}$

a unique g_0 such that A_{g_0} satisfies
 my gauge fixing condition $F[A] = 0$

$$1 = \int_{\mathcal{G}} D[g] \delta[F[A_g]] |\det(F'[A_g])|$$

\mathcal{G} \uparrow all gauge transformations

$N \times N$ matrix

$$= S_0(a) \Rightarrow g = \prod_{2 \times 2} G_v$$

$\neq 0$

$$+ \partial g \cdot g^{-1}$$

g_0 satisfies

$$F[A] = 0$$

$$\left| \det(F'[A_0]) \right|$$

$$Z = \int D[A] \exp(i S_{YM}[A])$$

$$= \int_{\mathcal{A}} D[A] \int_{\mathcal{G}} D[g] \delta[F[A_g]] \left| \det[F'[A_g]] \right| \exp(-i S_{YM}[A])$$

$$= \int_{\mathcal{G}} [g] \int_{\mathcal{A}} D[A] \quad "$$

$$D[A] = D[A_g], \quad S_{YM}[A] = S[A_g]$$

gauge invariance

$$g = SU(2) \Rightarrow g = \prod_{z \in \mathbb{H}} G_z$$

$J \neq 0$

$g^{-1} + \partial g \cdot g^{-1}$ most general gauge transf

A_{g_0} satisfies

$$F[A] = 0$$

$$\left| \det(F'[A_{g_0}]) \right|$$

only

$$Z = \int D[A] \exp(i S_{YM}[A])$$

(3)

$$= \int_D D[A] \int_G D[g] \delta[F[A_g]] \left| \det[F'[A_g]] \right| \exp(-i S_{YM}[A])$$

$$= \int_G D[g] \int_A D[A] \quad "$$

$$D[A] = D[A_g], \quad S_{YM}[A] = S[A_g]$$

gauge invariance
Hdar meas of G

Lattice Theory (K. Willich)

②

④ g fixed, A_g is a dummy variable $A_g \rightarrow A$ Faddeev-Popov Determinant

$$Z = \int D[g] \times \left[D[A] S[F[A]] \left| \det[F'[A]] \right| \exp(i S_{\text{YM}}[A]) \right]$$

$$= \cancel{\text{Vol}(\mathcal{G})} \times \square$$

Gauge Fixed correct Functional integral

sect. 6.5 of notes

$$F[A] \rightarrow F[A](x) = \partial^\mu A_\mu^a(x) \quad A_\mu = A_\mu^a t_a, \delta\alpha = \delta\alpha^a t_a$$

infinit. gauge transf.

$$A_\mu^a \rightarrow A_\mu^{a'} = A_\mu^a + \underbrace{(\partial_\mu \delta\alpha^a - \delta\alpha^b A_\mu^c)}_{\delta A_\mu^a(x)} = A_\mu^a + \delta A_\mu^a(x) \quad A' = A + \delta A$$

$$F^a[A + \delta A](x) = \partial^\mu A_\mu^a(x) + \partial^\mu \delta A_\mu^a(x)$$

$$\delta F^a(x) = \partial^\mu \delta A_\mu^a(x) = \partial^\nu \cdot \underbrace{D_\nu}_{\text{Diff. op. which depends on } A} \cdot \delta\alpha^a$$

$$\Rightarrow [F'(A)]_{xy}^{ab} = \left[\partial^\mu D_\mu \right]_{xy}^{ab}$$

↑
Diff. operator in X

$$F[A] = \left[\partial^\mu D_\nu \right]_{xy}^{ab}$$

↑
Diff operator

$$[A_\mu, \alpha]^a = \epsilon_{bc}^a A_\mu^b \alpha^c$$

↑
structure of $su(2)$

$$\partial^\mu D_\nu \cdot \alpha^a = \partial^\mu \left[\partial_\nu \alpha^a - i [A_\nu, \alpha]^a \right]$$

$$= \partial^\mu \partial_\nu \alpha^a - i \epsilon_{bc}^a \left(\partial^\mu A_\nu^b \cdot \alpha^c + A_\nu^b \partial^\mu \alpha^c \right)$$

$$\left[\partial^\mu D_\nu \right]^{ac} = \delta^{ac} \Delta - i \epsilon_{bc}^a \left(\partial^\mu A_\nu^b \right) \leftarrow \epsilon_{bc}^a A_\nu^b \cdot \partial^\mu$$

diff operator 2nd 1st

$$\det [F[A]] = \det \left[\partial^\mu D_\nu \right]$$

↑
depends on A

← differential operator functions $\alpha: M \rightarrow \text{Lie } su(2)$
which depends on A

all gauge transformations

$D[A] = D[A_g]$, $\int D\psi$ gauge invariance Haar measure g Lattice Theory (K. Wilson)

$$b \cdot \alpha^c + A_{\mu}^b \partial^{\mu} \alpha^c$$

$$\epsilon_{bc}^a A_{\mu}^b \partial^{\mu} \alpha^c$$

↑ 1st

U(1) $\det(F(A)) = \det(\Delta)$ indep of A

Faddeev Popov Ghosts Use Grassman "Calculus"

$$\det[-\partial^{\mu} D_{\mu}] = \int D[\bar{c}, c] \exp(-\bar{c} (-\partial^{\mu} D_{\mu}) c)$$

Algebra OK

↑ ↑
anticommuting fields

c ghosts
 \bar{c} antighost

$$C = \{ C^a(x) \} \quad \begin{array}{l} a=1,2,3 \\ x \in M \end{array}$$

$$\bar{C} = \{ \bar{C}_a(x) \}$$

$$S[C, \bar{C}] = - \int d^4x \bar{C}_a(x) \partial^\mu D_\mu C^a(x)$$

Ghosts
"Achen"

$$= \int d^4x \bar{C}_a(x) (-\partial^\mu \partial_\mu) C^a(x) + i \bar{C}_a \partial^\mu [A_\mu, C]^a(x)$$

$$\boxed{= \int d^4x [\partial^\mu \bar{C}_a(x)] [D_\mu C^a(x)]}$$

↑
contains A

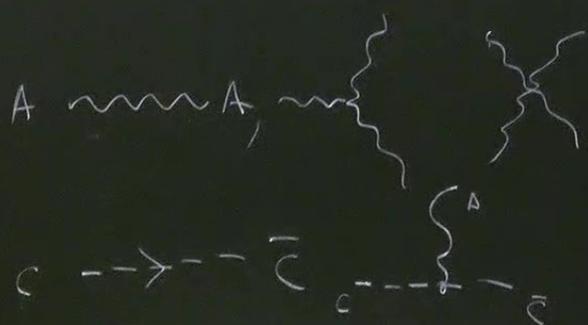
Diff. op. which depends on A

Diff operator in X

$$Z = \int D[A] S[F[A]] \exp(i S_{YM}[A]) \exp(i S_{ghost}[\bar{c}, c, A])$$

↑ local action ↑ local action

QFT Functional Integral



c^a, \bar{c}^a a indices charged
 no μ indices Spin=0
 Fermions
 anticommut

$c^a(x)$

$\frac{gac}{p^2} \approx$ Scalar Field
 interaction

α^+, α^- creation op
 annh.
 $\{\alpha^+, \alpha^-\} = 1$
 Violates Spm Statistics!

all gauge transformations

$D[A] = D[A_g]$, $S_M[A]$ $S[A]$
gauge invariance
Higgs mass g Lattice Theory (K. Wilson)

$U(1)$ $\det(F(A)) = \det(\Delta)$ indep of A

Faddeev Popov Ghosts

Use Grassman "Calculus"

$$\det[-\partial^\mu D_\mu] = \int D[\bar{c}, c] \exp(-\bar{c} (-\partial^\mu D_\mu) c)$$

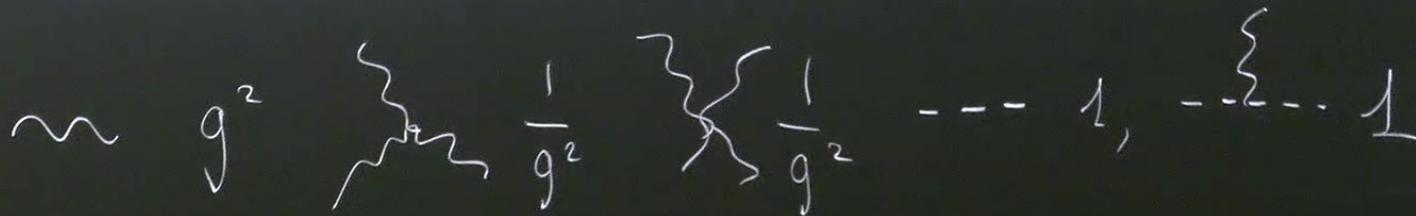
Algebra OK

↑ ↑
anticommuting fields

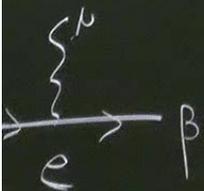
c ghosts
 \bar{c} antighost

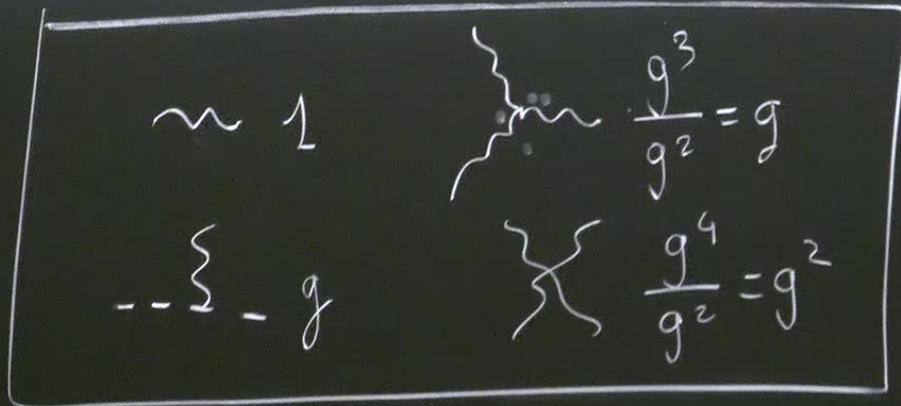
$x^c + A^b_\mu \partial^\mu x^c$
 a b $\vec{\partial}^\mu$
 c μ \uparrow $1st$

$A \rightarrow g A$

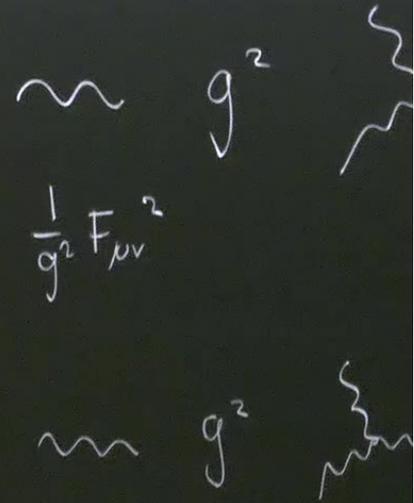


$$\frac{1}{g^2} F_{\mu\nu}^2$$





$$A \rightarrow g A$$



$$\frac{1}{g^2} F_{\mu\nu}^2$$

ϕ spin 0
 ψ spin $\frac{1}{2}$
 A_μ spin 1

coupling $A \rightarrow g A$

