Title: Future Astrophysical Targets for Intensity Interferometry

Speakers: Norman Murray

Collection/Series: Future Prospects of Intensity Interferometry

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#### Abstract:

Interferometry is the use of wave interference to measure the properties of a source observed by two or more detectors. For example, the Event Horizon Telescope measures the phase and amplitude of 1.3 mm wavelength radiation at telescopes up to ten thousand kilometers apart to reveal event horizon scale images of supermassive black holes. Measuring wave phases in the optical has been demonstrated for baselines no longer than hundreds of meters. Intensity interferometry dispenses with the need to measure phases, allowing much larger baselines, and hence much higher spatial resolution. The technique has been in use for seven decades, but recent advances in detector technology have reinvigorated interest in the method. I will discuss the basics of intensity interferometry, the characteristics of the new detectors, and possible applications of broad astrophysical and cosmological interest. The latter include estimates of the Hubble constant from observations of the disks of active galactic nuclei (AGN), with possible impact on the Hubble tension. The same observations will provide detailed information on the AGN disk and line emission regions; the latter may be crucial for estimating the mass loss rates in AGN winds, which are believed to impact their host galaxies. Other possible applications include spatially resolved measurements of stellar oscillations, which, by analogy with helioseismology, would provide constraints on the run of temperature in stellar interiors, as well as the interior differential rotation.

# Future Astrophysical Targets for Intensity Interferometry

Neal Dalal, Marios Galanis, Charles Gammie, Samuel Gralla,







#### Intensity Interferometry Hanbury Brown & Twiss

- Similar to interferometry in the radio or millimeter band (amplitude interferometry)
- Use large base lines **B** and short wavelengths  $\lambda$  to get high angular resolution
  - θ ~ λ/B
  - $\lambda \sim 5000$  angstroms or  $5 \times 10^{-5}$  cm, B  $\sim 10^4$  km,  $\theta \sim 0.01 \mu$  arcseconds
  - \* AGN disks  $R_s \sim 3x10^8\,km,\,D \sim 100Mpc,\,\theta \sim 20\mu$  arcseconds
  - Stellar disks  $R_{\odot} \sim 7x10^{10}$  cm, D ~ 10 pc,  $\theta \sim 0.5$  milliarcseconds
    - Can get many resolution elements across the stellar disk

# **Intensity Interferometry**

#### Hanbury Brown & Twiss

- The count rate of photons can vary by order unity
- Nearby AGN will have count rates of 10<sup>6</sup> photons per second

This is not shot noise!



1.

### **Intensity Interferometry**

Hanbury Brown & Twiss (1956)

Hanbury Brown and Twiss measured the correlation at different separations d, ranging from 2.5 to 9.2 meters, to find the angular size of Sirius, 0.0063" Note that the correlation of an extended source falls off more rapidly than that of a point source

Base-line  $d = |\mathbf{B}_{\perp}/\lambda| = |(\mathbf{u},\mathbf{v})|$ 



Fig. 2. Comparison between the values of the normalized correlation coefficient  $l^{2}(d)$  observed from Sirius and the theoretical values for a star of angular diameter  $0.0063^{"}$ . The errors shown are the probable errors of the observations

#### **Photon correlations**



Slide credit: Neal Dalal

#### **Photon correlations**



To maximize SNR, we want lots of photons and precise timing.

Slide credit: Neal Dalal

## **SPADs**

#### **Single Photon Avalance Diodes**



#### **SPADs**

#### **Single Photon Avalance Diodes ADVANCED** SCIENCE NEWS **ADVANCED** QUANTUM TECHNOLOGIES www.advancedsciencenews.com www.advquantumtech.com S,₹ b) a) (2)0\_ Upper neutral X<sub>n</sub> X, Multiplication e p-Drift (3) (1) 0 V Xp Lower neutral x\_end $V_{\rm bd}$ V ×\* d) C) 10<sup>5</sup> 104 FWHM = 30 ps Engineered Not engineered 103 104 Diffusion tail (cbs) 10<sup>2</sup> 10<sup>1</sup> Counts 103 **Tunneling-limited** 10<sup>2</sup> 10<sup>0</sup> Ceccarelli+ (2021) SRH-limited 10<sup>1</sup> 10-1 0 20 0.5 1.5 -80 -20 0 1 2 2.5 -60 -40 Time (ns) Temperature (°C)

### **SNSPDs**

#### Superconducting Nanowire Single Photon

#### Detector

Rev. Sci. Instrum. 82, 071101 (2011)



FIG. 8. (Color online) A section of a superconducting nanowire singlephoton detector is shown with a bias current just below the critical current density that would drive the wire normal. (a) An incoming photon creates a small normal region within the nanowire. (b) The superconducting current is expelled from the normal region, increasing the current density in the adjacent areas of the nanowire. (c) That increase in current density is enough to drive those adjacent regions normal, which in turn results in a measurable voltage drop across the detector. Eisaman+ 2024

#### **SNSPDs**

Oripov+ 2023

#### Superconducting Nanowire Single Photon Detector Array



**Fig. 1** | **Overview of the 800** × **500 camera. a**, Imaging at 370 nm, with raw time-delay data from the buses shown as individual dots in red and binned 2D histogram data shown in black and white. **b**, Count rate as a function of bias current for various wavelengths of light as well as dark counts. **c**, False-colour scanning electron micrograph of the lower-right corner of the array,

highlighting the interleaved row and column detectors. Lower-left inset, schematic diagram showing detector-to-bus connectivity. Lower-right inset, close-up showing 1.1- $\mu$ m detector width and effective 5 × 5- $\mu$ m pixel size. Scale bar, 5  $\mu$ m.

#### **SNSPDs**

<u>Oripov+ 2023</u>

#### Superconducting Nanowire Single Photon Detector Array



**Fig. 2** | **Electrical operation of the detectors and readout bus. a**, Circuit diagram of a bus and one section of 50 detectors with ancillary readout components. SNSPDs are shown in the grey boxes and all other components are placed outside the imaging area. A photon that arrives at time  $t_0$  has its location

determined by a time-of-flight readout process based on the time-of-arrival difference  $t_2 - t_1$ . **b**, Oscilloscope traces from a photon detection showing the arrival of positive (green) and negative (red) pulses at times  $t_1$  and  $t_2$ , respectively.

#### **Possible Astrophysical Targets**

- AGN
- Resolved Asteroseismology
- Photon rings
- Tidal Disruption Events
- Supernovae

#### **Possible Astrophysical Targets**

- AGN
  - Disk angular size
    - Measure H<sub>o</sub>
  - Disk scale height
    - Thin versus thick
  - Map Broad Line Region
    - Determine where outflows emerge
- Resolved Asteroseismology
  - 2D power spectra (velocity versus l)
    - Run of Temperature
    - Rotational splitting  $\Rightarrow$  Internal differential rotation

#### **AGN Broad Line Region**



# **AGN variability**

• AGN luminosity varies over time, for both continuum and lines.



# Geometric measurement of H<sub>0</sub>

#### **AGN Broad Line Region**



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#### **AGN Broad Line Region**



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## Measuring H<sub>0</sub>

# Measuring H<sub>0</sub>

- Time lags between line variability & continuum variability tell us physical size of line-emitting region.
- Interferometry tells us the angular size of the same line-emitting region (same photons)
- Comparing the two tells us the angular diameter distance to the AGN
- Since these are line emitters, we also have redshift
- Distance + redshift =  $H_0$



all these numbers are interesting!







FIG. 3.—Composite quasar spectrum using median combining. Powerlaw fits to the estimated continuum flux are shown. The resolution of the input spectra is  $\approx 1800$ , which gives a wavelength resolution of about 1 Å in the rest frame.





#### Ferrarese & Merritt 2000

**BAL Outflows & Galaxy Evolution** 



FIG. 2.—BH mass vs. the central velocity dispersion  $\sigma_c$  of the host elliptical galaxy or bulge (*filled circles*) or the rms velocity  $v_{\rm rms}$  measured at one-fourth of the effective radius (*open circles*). Crosses represent lower limits in  $v_{\rm rms}$ . The solid and dashed lines are the best linear fits using  $\sigma_c$  (as in Fig. 1b) and  $v_{\rm rms}$ , respectively.





- Wind luminosity = 1/2 dM/dt v<sup>2</sup> = 1/2  $\Omega R^2 \rho v^2$
- Momentum loss rate =  $dM/dt v = \Omega R^2 \rho v$
- Measure v directly
- Estimate  $\rho$
- Need R



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# Asteroseismology

Internal differential rotation

- Solar dynamo is driven by differential rotation
- Stellar magnetic fields produce x-rays and the bulk of the UV flux
- X-rays and UV evaporate protostellar disks
- X-rays and UV can strip planetary atmospheres



**Figure 29** Internal rotation (left) and the corresponding errors (right) derived from the MDI full-disk analysis averaged over all Dynamics Runs. We have erased color from the regions where estimates of rotation are deemed unreliable; contours are retained on the left for ease of labeling.



Figure 30 Internal rotation (left) and the corresponding errors (right) derived from an average over the first six years of the HMI 72-day analysis. We have erased color from the regions where estimates of rotation are deemed unreliable; contours are retained on the left for ease of labeling.

#### Internal differential rotation



Figure 1. The acoustic p-mode spectrum of the Sun, as measured using the first eight months of GOLF data. At 3 mHz the ratio S/N is  $\sim$ 3000.



- 50 100 150 200 250 300 angular degree, *l*
- Resolved stellar oscillation velocity power spectra

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#### **Internal differential rotation**



#### Internal differential rotation

SOLAR *p*-MODE FREQUENCY SPLITTINGS

1093



#### Internal differential rotation

636 GOUGH & TOOMRE



Figure 2 Ray paths in the standard model of the sun represented by the continuous line in Figure 1: (a) for two acoustic waves; the more deeply penetrating wave is  $p_8$  (l = 2) and the shallower wave is  $p_8$  (l = 100); (b) for the gravity wave  $g_{10}$  (l = 5). Note that the number of reflections per revolution is not integral, and indeed is almost never rational, so the ray paths



Figure 3 Lower turning points for p modes of a solar model, determined by the vanishing of  $\kappa$ , plotted against degree l for the three cyclic frequencies  $v = \omega/2\pi = 2$ , 3, 4 mHz. The curves for 2 and 3 mHz terminate at the lowest-order modes, at values of l determined by Equation 3.2 with n = 1.



Sound speed v. R

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**Fig. 3.** The dashed curve is the square of the spherically averaged sound speed in the sun. The solid curve corresponds to a standard theoretical model. The magnitudes of the slopes of the curves are lower immediately beneath the convection zone, where the temperature gradient is too small to drive the instability. The inset shows that the convectively unstable region of relatively high slope extends somewhat more deeply into the sun than it does in the model.