

Title: Big Bang Nucleosynthesis: New Physics and New Tools

Speakers: Cara Giovanetti

Collection/Series: Particle Physics

Subject: Particle Physics

Date: October 01, 2024 - 1:00 PM

URL: <https://pirsa.org/24100075>

Abstract:

Big Bang Nucleosynthesis (BBN) is a powerful tool for probing both new physics and LCDM, and complements analyses utilizing the Cosmic Microwave Background (CMB) and results from particle experiment. I will provide two examples of BBN probes of BSM models. I will then discuss new kinds of analyses that can be performed with the recently-released fast and differentiable BBN code LINX. In particular, LINX can be used to perform full BBN+CMB joint analyses at a level of sophistication that has never been achieved before, even in LCDM analyses.

Big Bang Nucleosynthesis: New Physics and New Tools

Cara Giovanetti (NYU)

October 1st, 2024

Perimeter Institute Particle Physics Seminar

Based on work with Mariangela Lisanti, Hongwan Liu, Siddharth Mishra-Sharma, Joshua T. Ruderman, Martin Schmaltz, and Neal Weiner

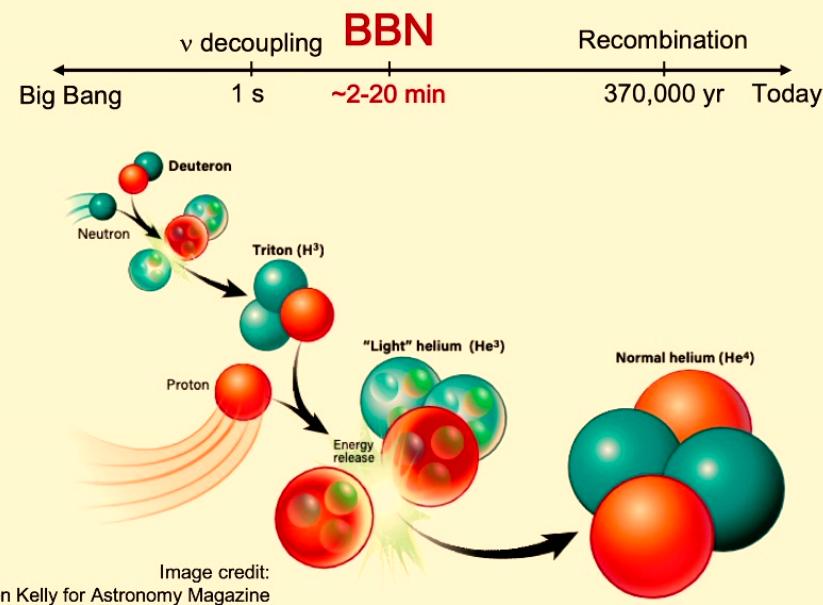
Other work

- Orbital Dynamics of the Solar Basin C.G., R. Lasenby, K. Van Tilburg, 2408.16041
- Neutrino Spectral Distortions in BBN
- Gravitational Wave Constraints on PBHs
- A Fast and Differentiable Recombination Code

Outline

- Why BBN?
 - BBN is a powerful probe of both Λ CDM and new physics
- BBN informs new physics
 - Portal Models
 - Electrophilic dark matter
 - Neutrinophilic dark matter
 - BBN still has a lot to say about new physics
- New tools for BBN
 - We can now perform sophisticated joint analyses for the first time

Why BBN?

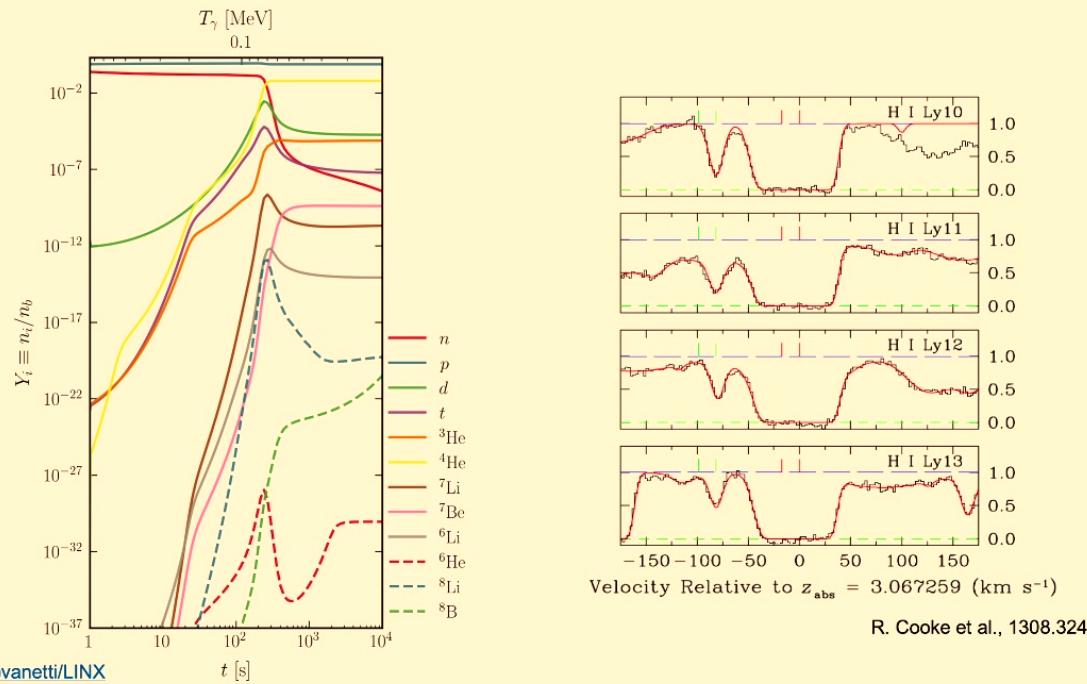


High temperatures
($T_{SM} \sim \text{MeV-keV}$)

High densities
($a \sim 10^{-9}-10^{-11}$)

Long times

Prediction meets measurement



<https://github.com/cgiovannetti/LINX>

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Why BBN?

High
temperatures

Precise
measurements

**BBN is a sensitive
probe of Λ CDM
and new physics**

High densities

Precise
predictions

Why BBN?



BBN is sensitive to the **expansion rate**, the **photon** and **neutrino temperatures**, and the **baryon-to-photon ratio**.

Helium-4 and Deuterium

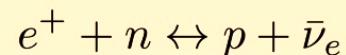
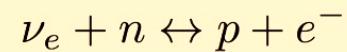
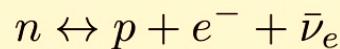


- Large binding energy (28 MeV)
- Most neutrons end up in ${}^4\text{He}$
- Small binding energy (2 MeV)
- Easily broken up

BBN is sensitive to the expansion rate, the photon and neutrino temperatures, and the baryon-to-photon ratio.



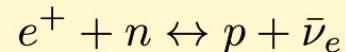
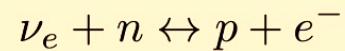
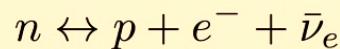
Neutrino temperature and expansion rate determine freeze-out of proton-neutron interconversion.



BBN is sensitive to the **expansion rate**, the **photon** and **neutrino temperatures**, and the baryon-to-photon ratio.



Neutrino temperature and expansion rate determine freeze-out of proton-neutron interconversion.



$$\left(\frac{n_n}{n_p}\right)_{EQ} = e^{-\frac{m_n - m_p}{T}}$$

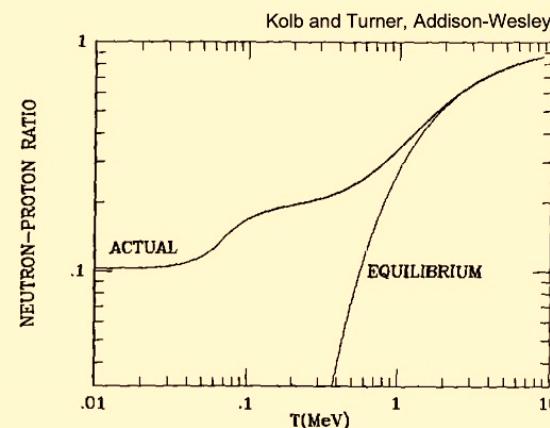


Fig. 4.1: The equilibrium and actual values of the neutron to proton ratio.

BBN is sensitive to the **expansion rate**, the **photon** and **neutrino temperatures**, and the baryon-to-photon ratio.



N_{eff} is impacted by the ratio of photon and neutrino temperatures. Impacts expansion rate.

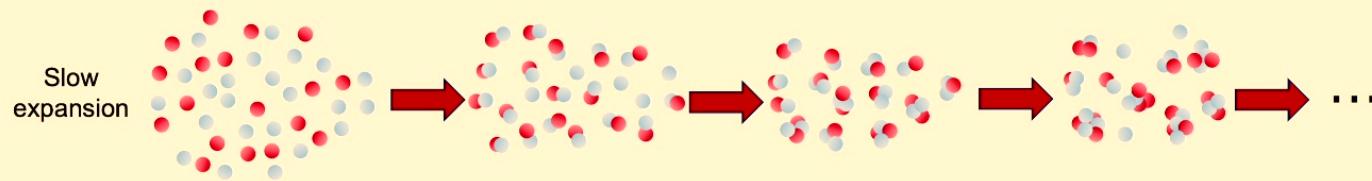
$$N_{\text{eff}} = \left(\frac{\rho_R - \rho_\gamma}{\rho_{\nu, \text{std}}} \right)_0 \rightarrow N_{\text{eff}} \sim \left(\frac{T_\nu}{T_\gamma} \right)_0^4 + \text{dark radiation}$$

BBN is sensitive to the **expansion rate**, the **photon** and **neutrino temperatures**, and the baryon-to-photon ratio.

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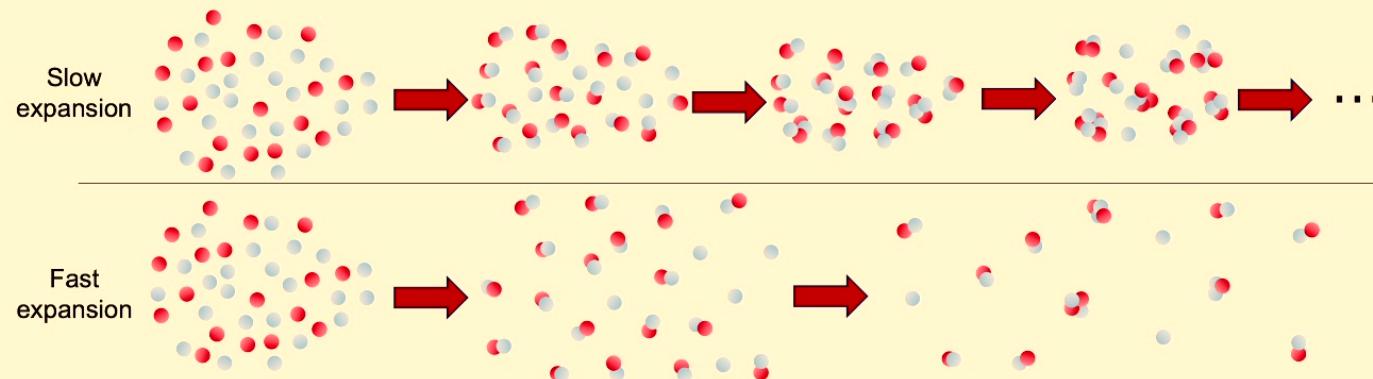


Faster expansion rate means heavy elements can't form. Larger abundance of deuterium.



BBN is sensitive to the **expansion rate**, the photon and neutrino temperatures, and the baryon-to-photon ratio.

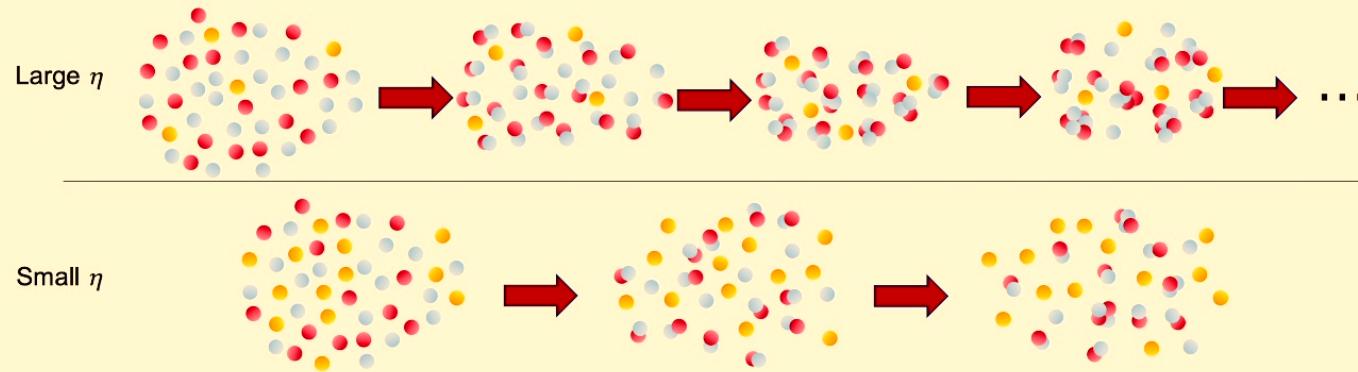
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BBN is sensitive to the **expansion rate**, the photon and neutrino temperatures, and the baryon-to-photon ratio.



Larger η means more frequent interactions between nuclides. Destroys deuterium.



BBN is sensitive to the expansion rate, the photon and neutrino temperatures, and the **baryon-to-photon ratio**.

Why BBN

- Many potential signals
 - Expansion rate (N_{eff})
 - Relative photon and neutrino temperatures
 - Weak rate freeze out
 - N_{eff}
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- Complements CMB analyses
- New tools make rigorous analyses possible

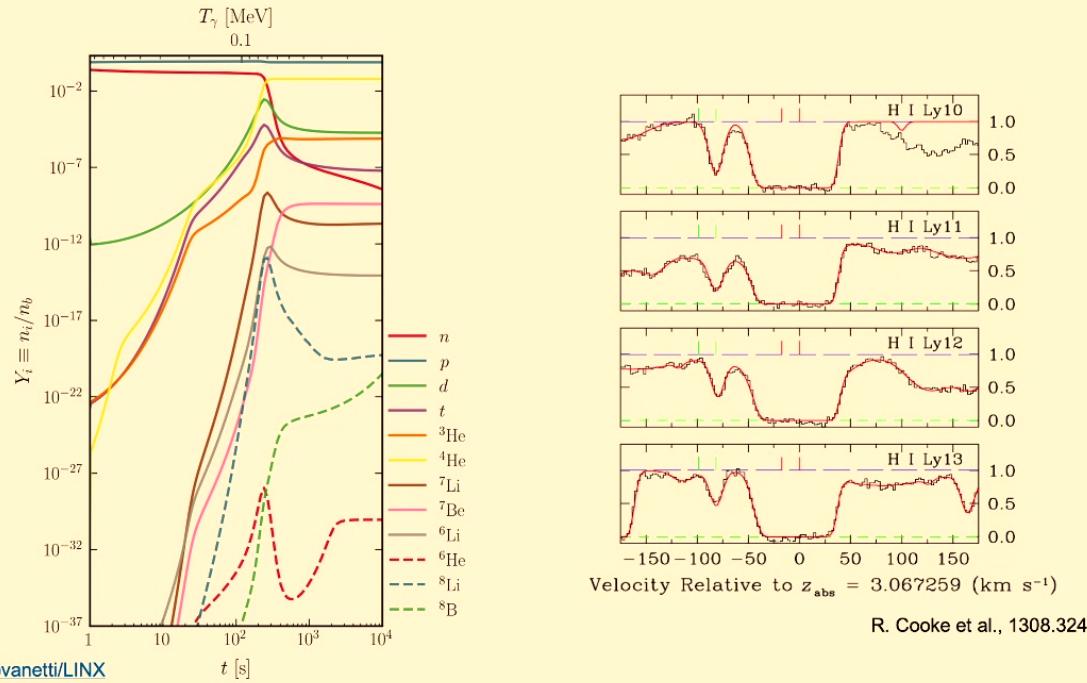
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Prediction meets measurement



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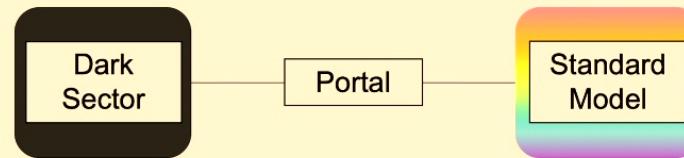
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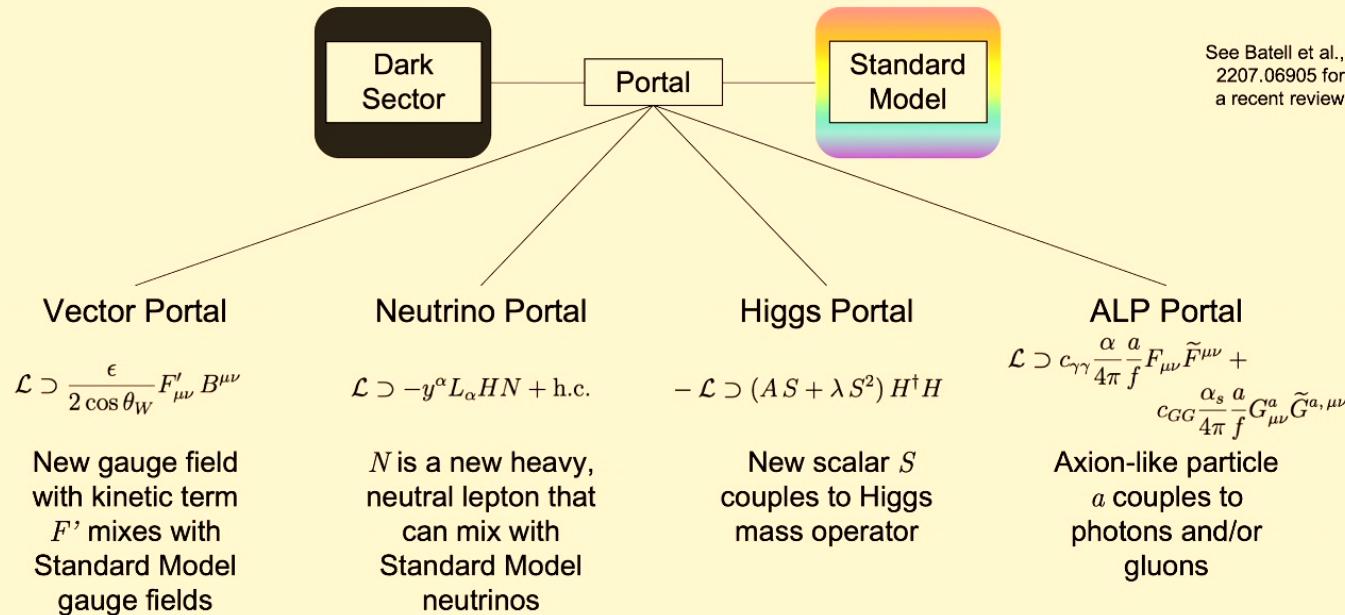
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Minimal Portal Models



See Batell et al.,
2207.06905 for
a recent review

Minimal Portal Models



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BBN constrains electrophilic sub-GeV dark matter

$$\mathcal{L} \supset -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} - \frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2 A'^{\mu}A'_{\mu} + J_{\text{EM}}^{\mu} (A_{\mu} - \epsilon A'_{\mu})$$

+ massive dark matter χ
+ dark radiation

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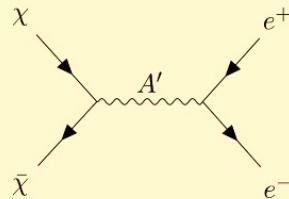
- Vector Portal
- Originally linked to 511 keV Galactic center excess C. Boehm et al., astro-ph/0309686
- Common experimental benchmark

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Small m_{χ} or $m_{A'}$ can contribute directly to N_{eff}



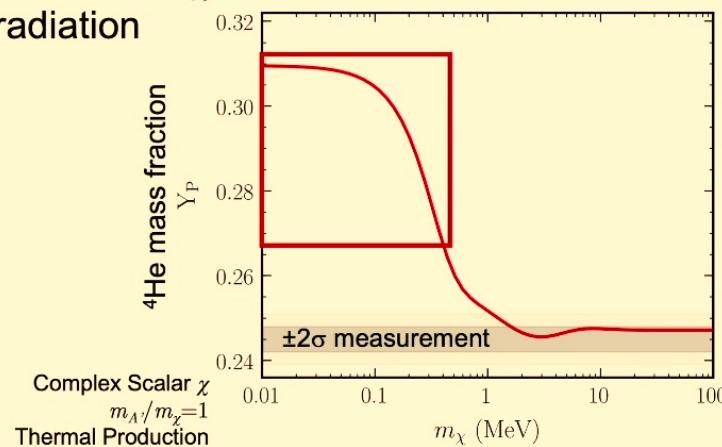
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Measured values from
R.L. Workman *et al.* (PDG), Prog. Theor.
Exp. Phys. 2022, 083C01
R. Cooke *et al.*, 1710.11129

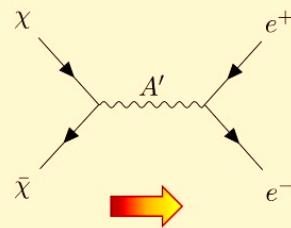


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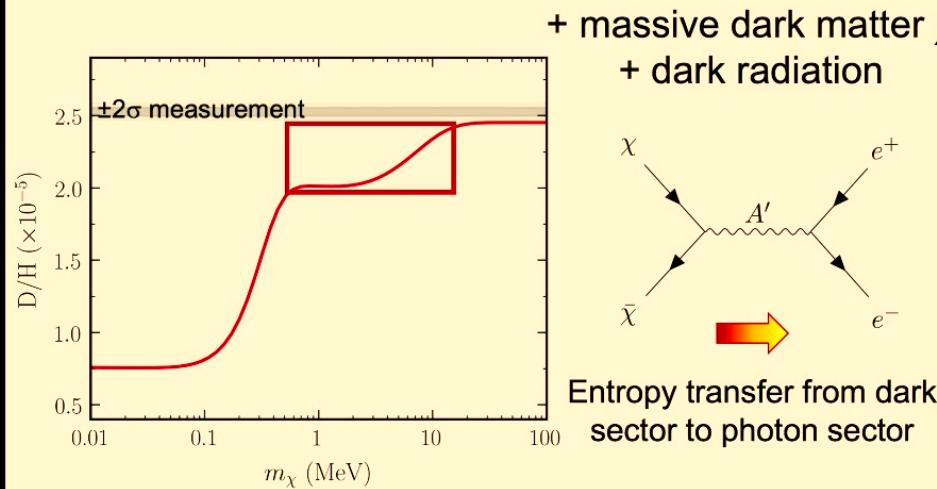
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Entropy transfer from dark sector to photon sector

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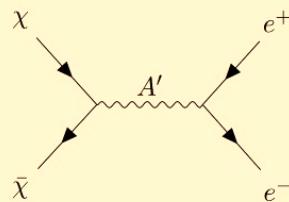
Complex Scalar χ
 $m_{A'}/m_\chi = 1$
 Thermal Production
 Measured values from
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 Exp. Phys. 2022, 083C01
 R. Cooke et al., 1710.11129

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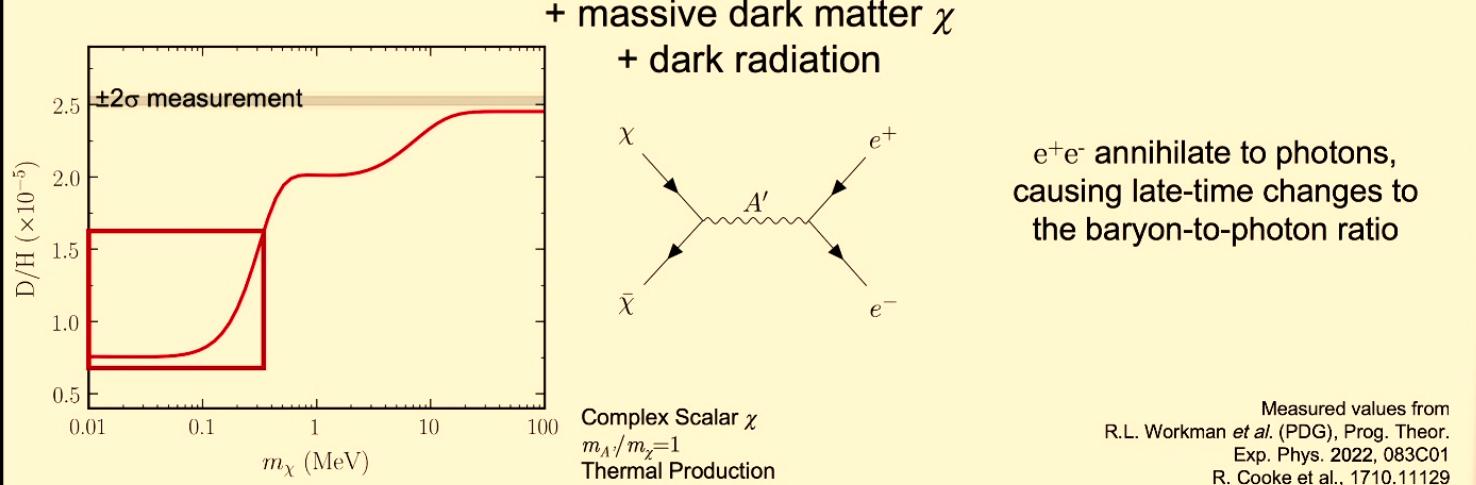
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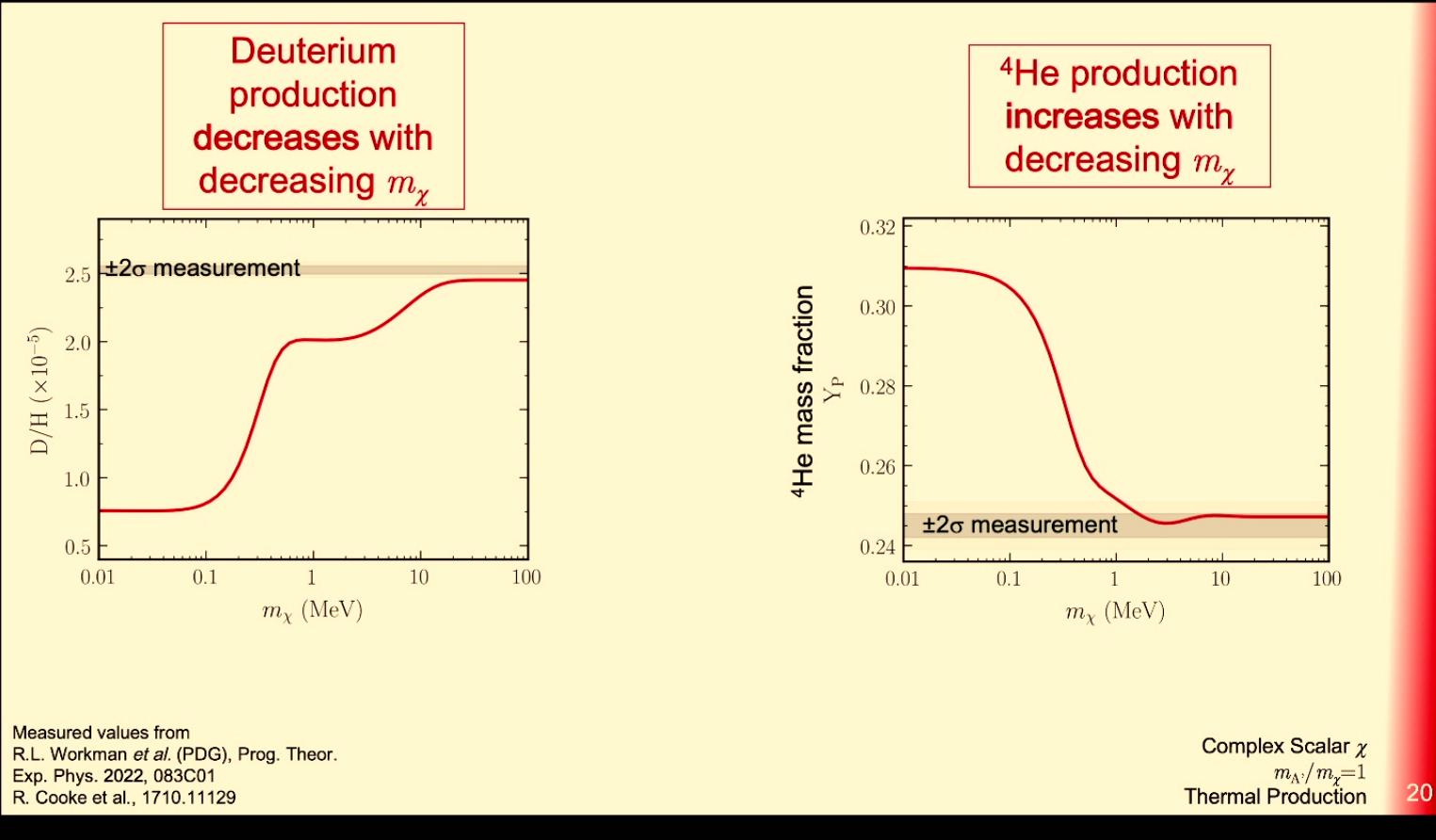
e^+e^- annihilate to photons,
causing late-time changes to
the baryon-to-photon ratio

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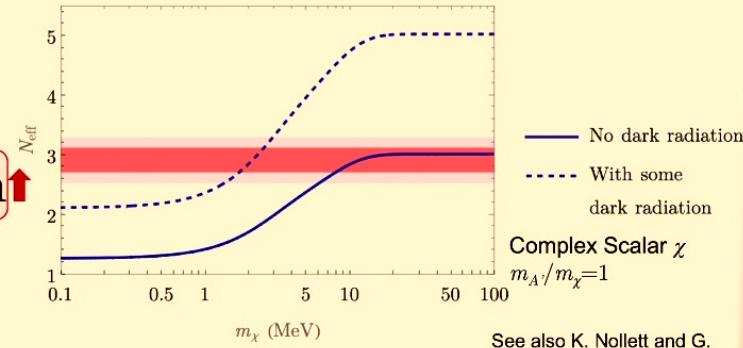


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$$\Delta N_{\text{eff}} \sim \left(\frac{T_{\nu}}{T_{\gamma}} \right)_0^4 + \text{dark radiation}$$



See also K. Nollett and G.
Steigman, 1312.5725

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Many models include dark radiation!

- Dark Higgs Dark Matter

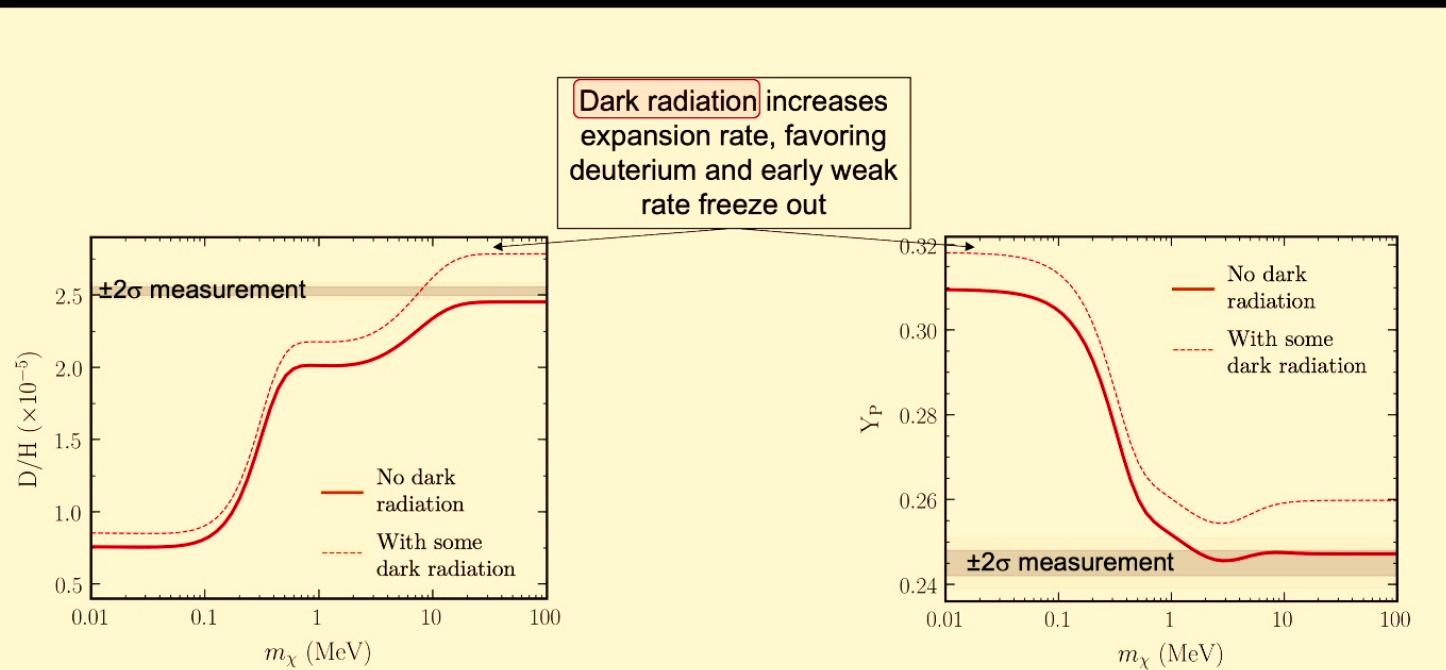
C. Mondino, M. Pospelov,
et al., 2005.02397

- Heavy Particle Decay models

As considered in K. Bleau,
J. Bramante, and C.
Cappiello, 2309.06482

- Dark QED avoids N_{eff} constraint on dark photons

A. Arvanitaki, S.
Dimopoulos, M. Galanis, D.
Racco et al., 2108.04823

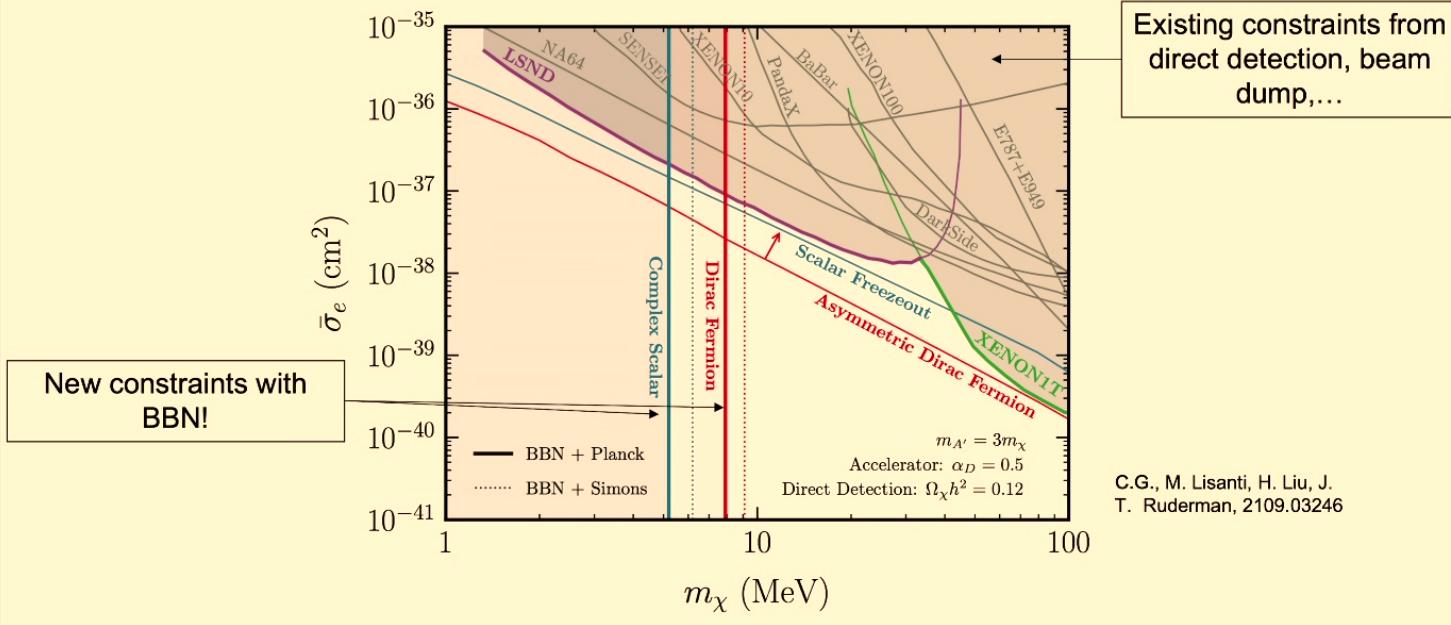


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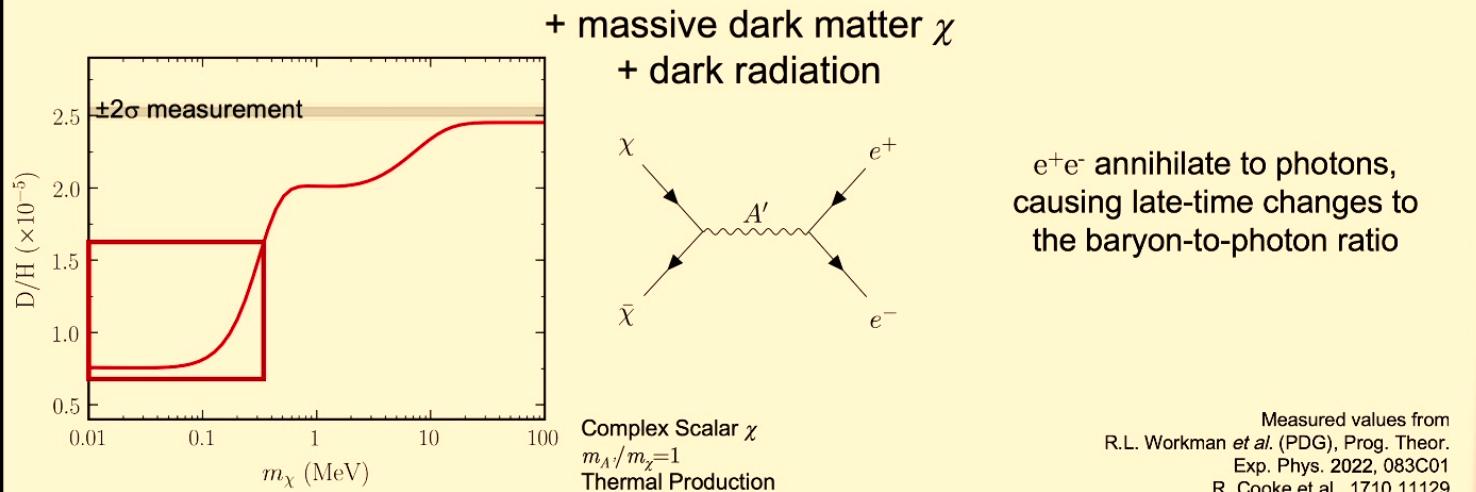
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BBN constrains electrophilic sub-GeV dark matter



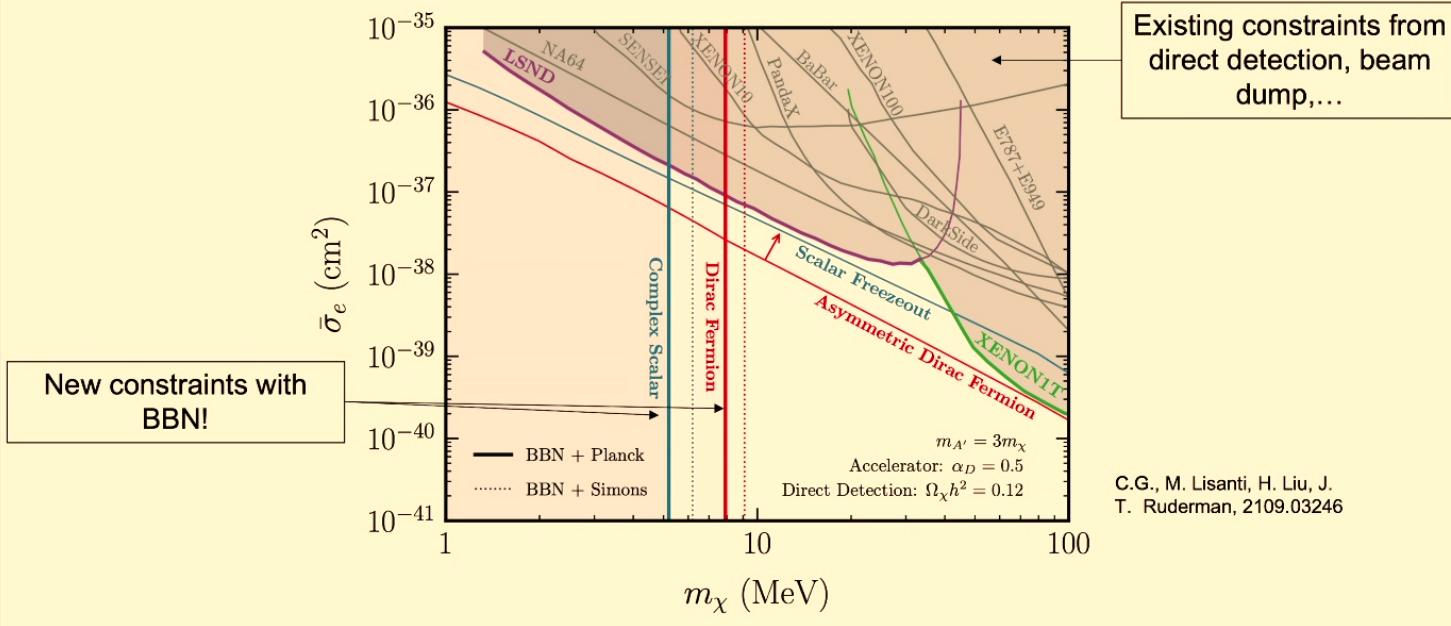
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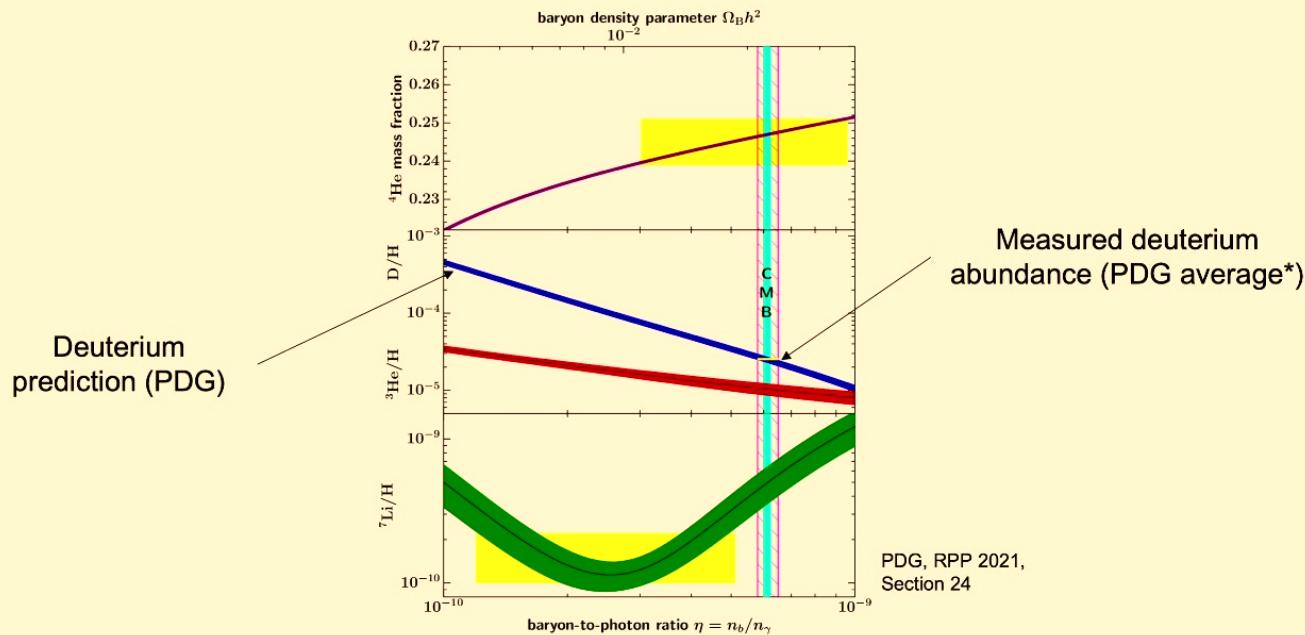
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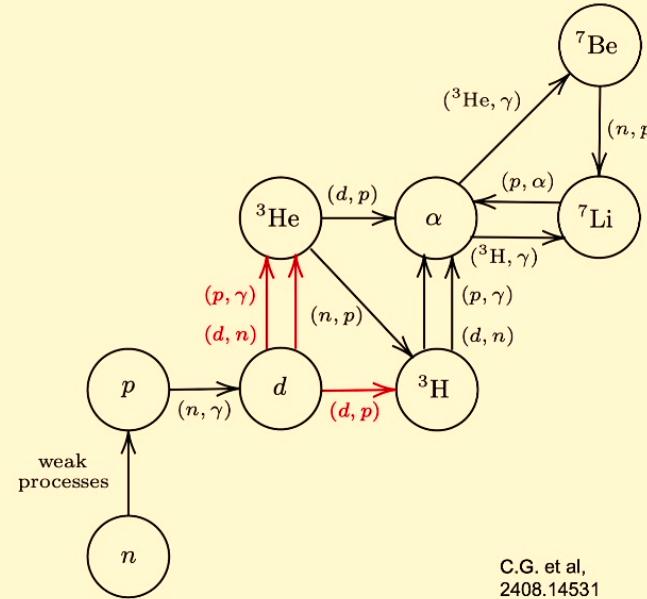
Tension between CMB and BBN η



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Reaction network

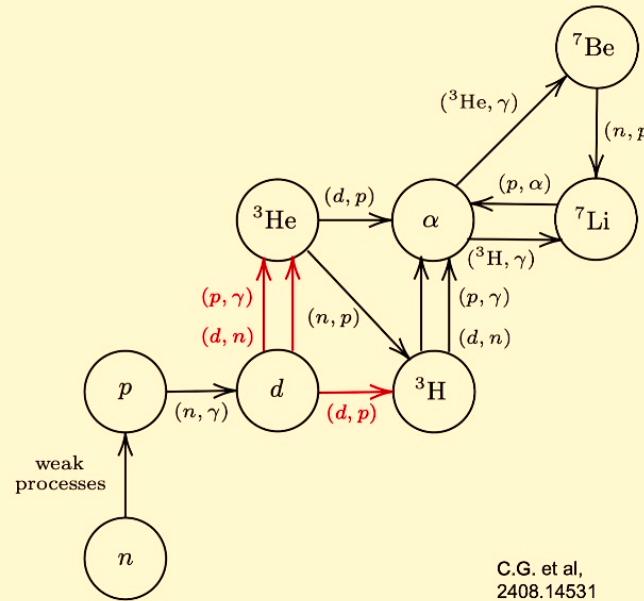
- Main input to BBN codes is a reaction network.
 - Includes rates for all reactions in the network



C.G. et al,
2408.14531

Reaction network

- Main input to BBN codes is a reaction network.
 - Includes rates for all reactions in the network
- Treatment of nuclear physics data differs between major BBN codes.
- **Absence or presence of tension is due entirely to choice of reaction network.**



C.G. et al,
2408.14531

The model

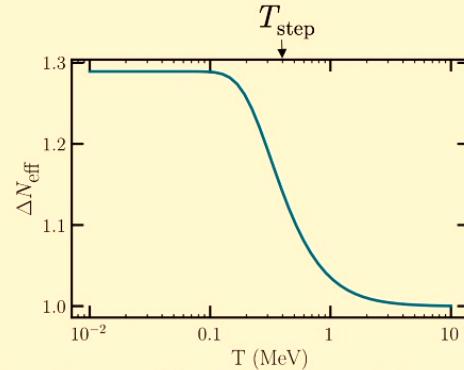
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- Dark sector equilibrates with SM neutrinos via repeated oscillation and scattering
- Freeze out at T_{step}
- Dominant effect is a “step” in N_{eff}

D. Aloni et al, 2111.00014

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The model



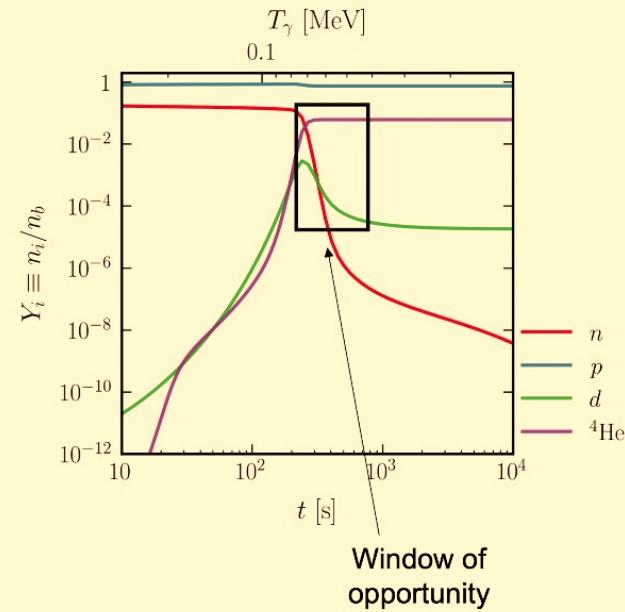
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A step in understanding the deuterium “tension”

- Increasing N_{eff} increases D/H, but also increases ${}^4\text{He}$
- Step in N_{eff} modifies D/H without modifying ${}^4\text{He}$

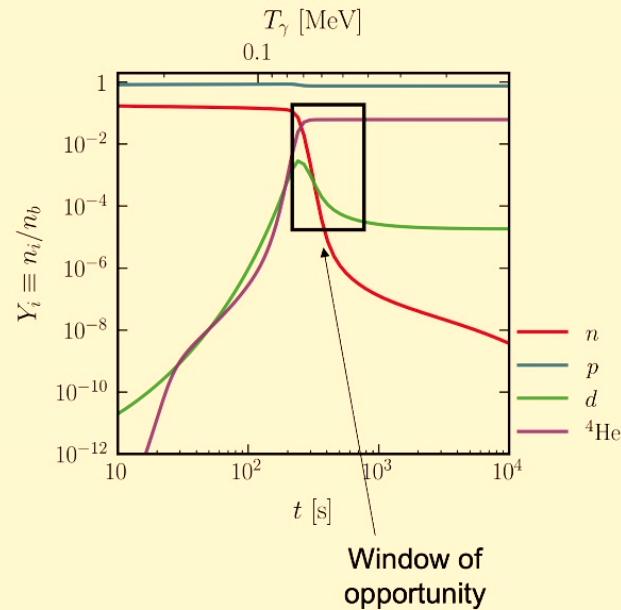


See also A. Berlin et al., 1904.04256

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A step in understanding the deuterium “tension”

- Increasing N_{eff} increases D/H, but also increases ${}^4\text{He}$
- Step in N_{eff} modifies D/H without modifying ${}^4\text{He}$
- This works across a wide range of step temperatures and changes to N_{eff}



See also A. Berlin et al., 1904.04256

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The model

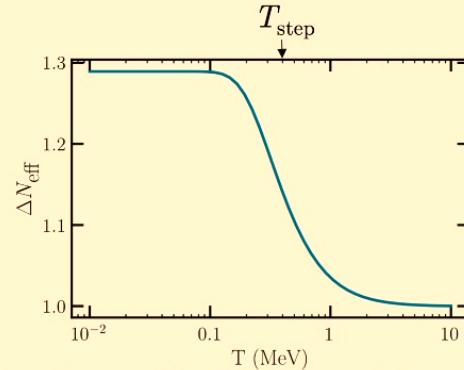
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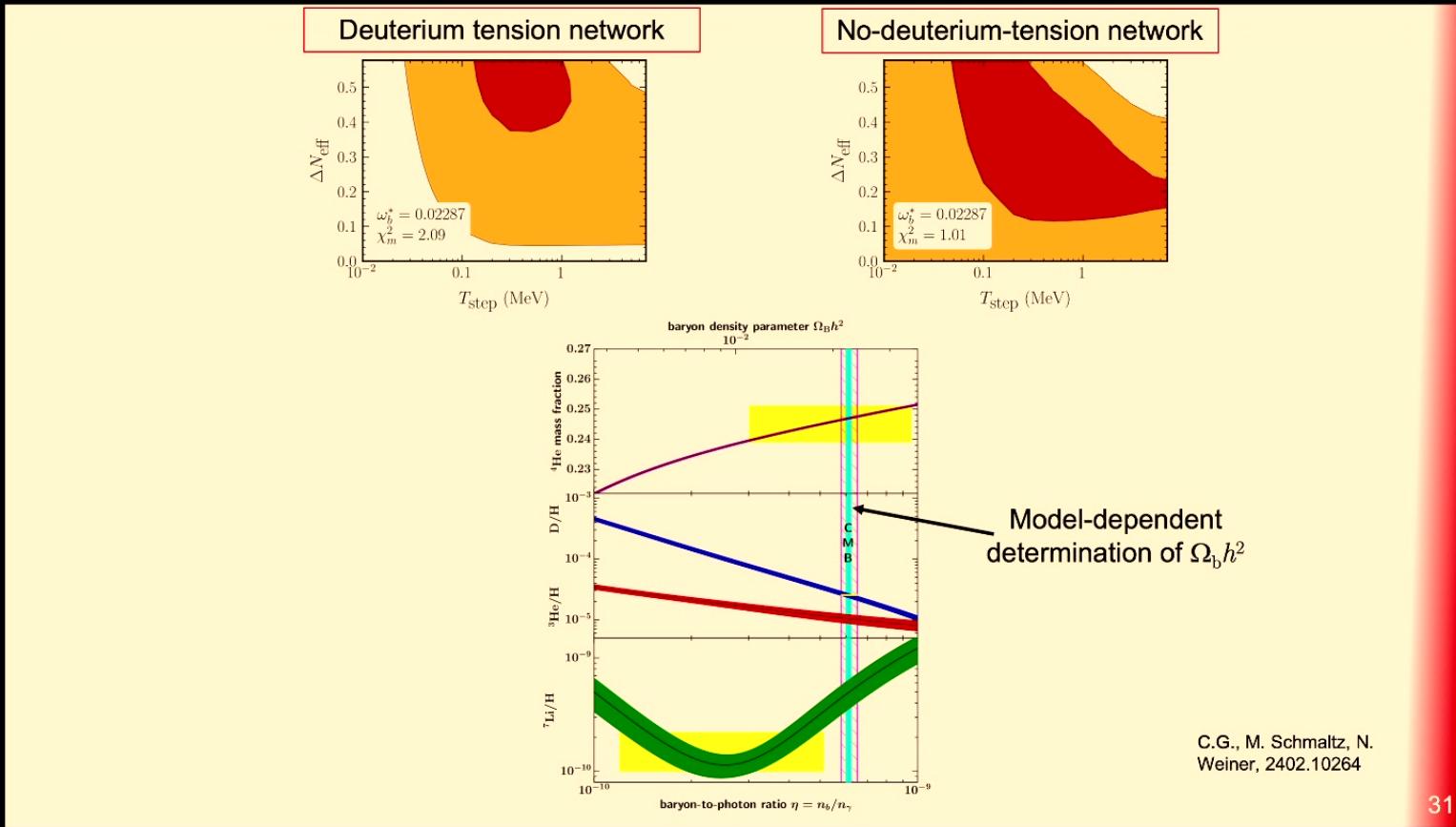
The model

$$\mathcal{L} \supset \underbrace{m_{\text{dark}} \nu_d \nu_d}_{\text{Dark fermion, MeV-scale mass}} + \underbrace{m_{\text{mix}} \nu_d \nu_{\text{SM}}}_{\text{Dark sector mixing}} + \underbrace{\lambda \phi \nu_d \nu_d}_{\text{Light scalar}}$$

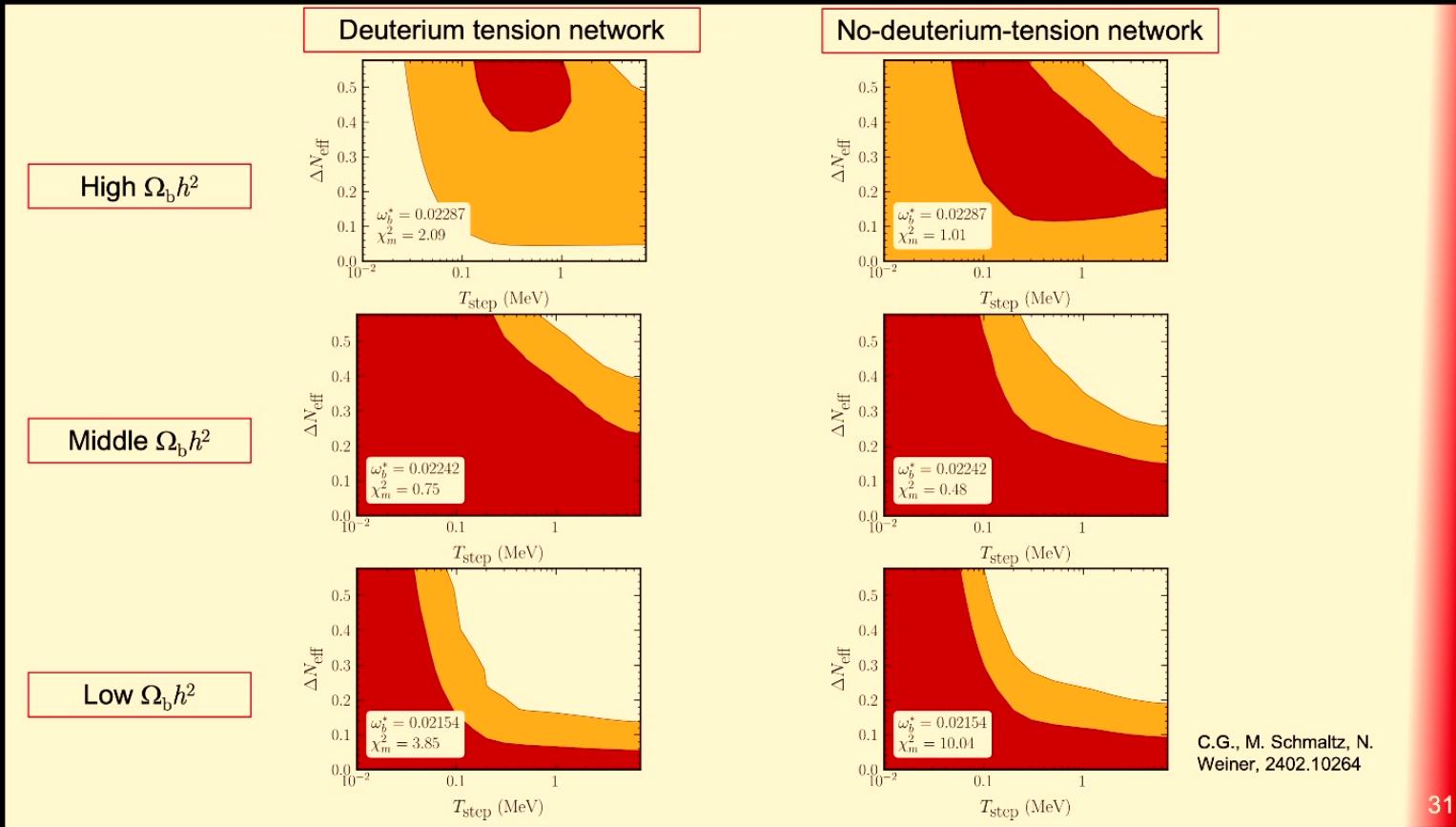
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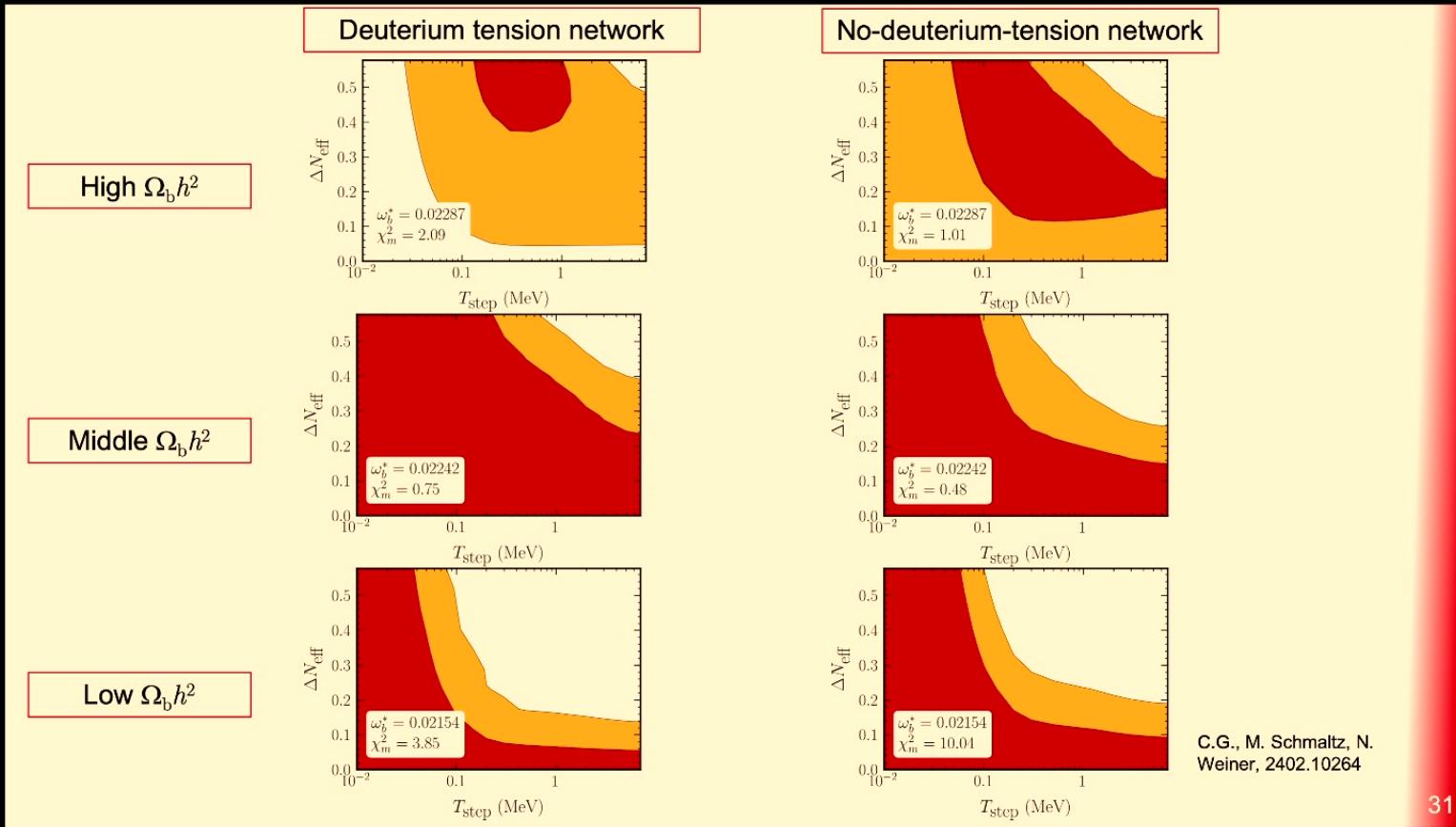
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Major hurdles in these analyses

- Switching reaction networks is hard, **but important**.
- Scans take **several days** to run.
- Calculating uncertainties is **hard**.



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Goals:

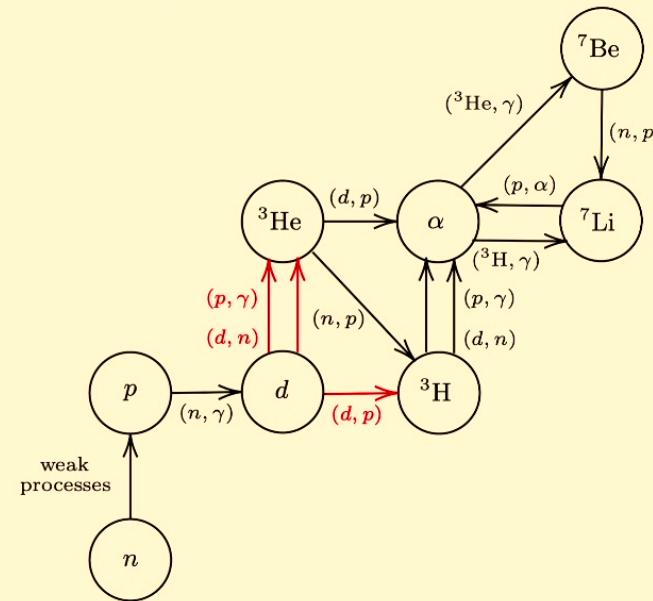
- 1) A fast BBN code
- 2) Easy calculation of uncertainties
- 3) Easy switching between reaction networks

We want to put BBN analyses on the same footing as CMB analyses when performing *parameter estimation*.

BBN Nuisance Parameters

Nuclear physics gives us an **uncertain** rate for each reaction.

Need to sample reaction rates as **nuisance parameters!**



Public BBN Codes

Name	Language	Time Per Solve	Comments
AlterBBN	C	< 1s	Incomplete implementation of neutrino decoupling, weak rates; old nuclear rates.
PRIMAT	Mathematica	$O(1 \text{ min})$	Extremely accurate, but very slow.
PArthENoPE	Fortran	< 1s	Fast, but challenging to modify for parameter estimation.
PRyMordial	Python	$O(10 \text{ s})$	Accurate. Full parameter estimation possible, but slow. Written with new physics in mind.

All current BBN codes have to **make compromises** when it comes to parameter estimation.

O. Pisanti et al., 0705.0290
A. Arbey, 1106.1363
R. Consiglio et al., 1712.04378
A. Arbey et al., 1806.11095
C. Pitrou et al., 1801.08023
S. Gariazzo et al., 2103.05027
A.K. Burns et al., 2307.07061

Slide courtesy of Hongwan Liu

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Public BBN Codes

Name	Language	Time Per Solve	Comments
AlterBBN	C	< 1s	Incomplete implementation of neutrino decoupling, weak rates; old nuclear rates.
PRIMAT	Mathematica	O(1 min)	Extremely accurate, but very slow.
PArthENoPE	Fortran	< 1s	Fast, but challenging to modify for parameter estimation.
PRyMordial	Python	O(10 s)	Accurate. Full parameter estimation possible, but slow. Written with new physics in mind.
LINX	Python+JAX	<0.1s	As accurate as PRyMordial. Fast enough for MCMC methods.

*LINX: Light Isotope Nucleosynthesis with JAX

C.G. et al., 2408.14538

No new approximations, not an emulator

Publicly available at <https://github.com/cgiovannetti/LINX>

Slide courtesy of Hongwan Liu

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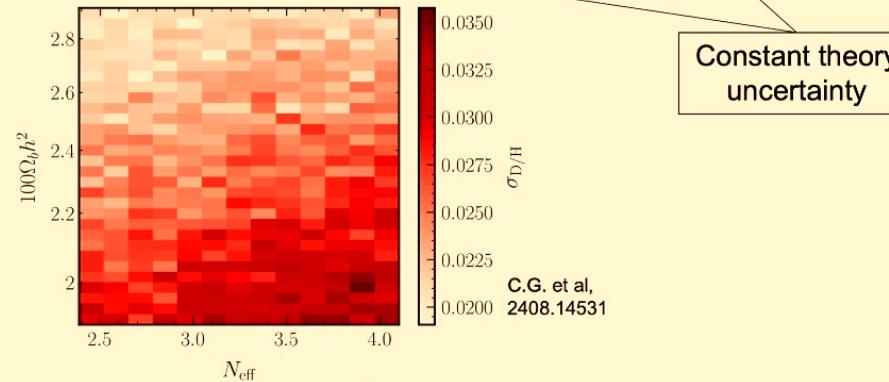
BBN Nuisance Parameters

$$\log \mathcal{L}_{\text{BBN,PArth.}} = -\frac{1}{2} \left[\frac{\left(Y_P(\Omega_b h^2, N_{\text{eff}}) - Y_P^{\text{obs}} \right)^2}{\sigma_{Y_P^{\text{obs}}}^2 + \sigma_{Y_P}^{\text{th2}}} + \frac{\left(D/H(\Omega_b h^2, N_{\text{eff}}) - D/H^{\text{obs}} \right)^2}{\sigma_{D/H^{\text{obs}}}^2 + \sigma_{D/H}^{\text{th2}}} \right]$$

Constant theory
uncertainty

BBN Nuisance Parameters

$$\log \mathcal{L}_{\text{BBN,PArth.}} = -\frac{1}{2} \left[\frac{\left(Y_P(\Omega_b h^2, N_{\text{eff}}) - Y_P^{\text{obs}} \right)^2}{\sigma_{Y_P^{\text{obs}}}^2 + \sigma_{Y_P}^{\text{th2}}} + \frac{\left(D/H(\Omega_b h^2, N_{\text{eff}}) - D/H^{\text{obs}} \right)^2}{\sigma_{D/H^{\text{obs}}}^2 + \sigma_{D/H}^{\text{th2}}} \right]$$



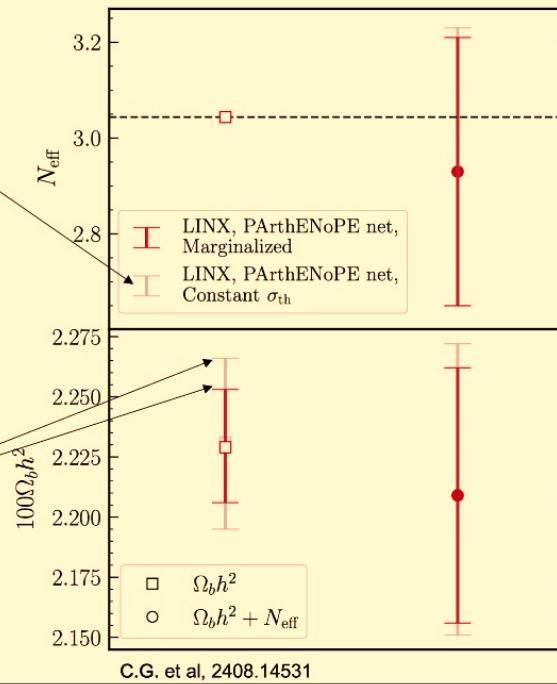
$$\log \mathcal{L}_{\text{BBN,LINX}} = -\frac{1}{2} \left[\left(\frac{Y_P(\Omega_b h^2, N_{\text{eff}}, \vec{\nu}_{\text{BBN}}) - Y_P^{\text{obs}}}{\sigma_{Y_P^{\text{obs}}}^2} \right)^2 + \left(\frac{D/H(\Omega_b h^2, N_{\text{eff}}, \vec{\nu}_{\text{BBN}}) - D/H^{\text{obs}}}{\sigma_{D/H^{\text{obs}}}^2} \right)^2 \right]$$

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BBN Nuisance Parameters

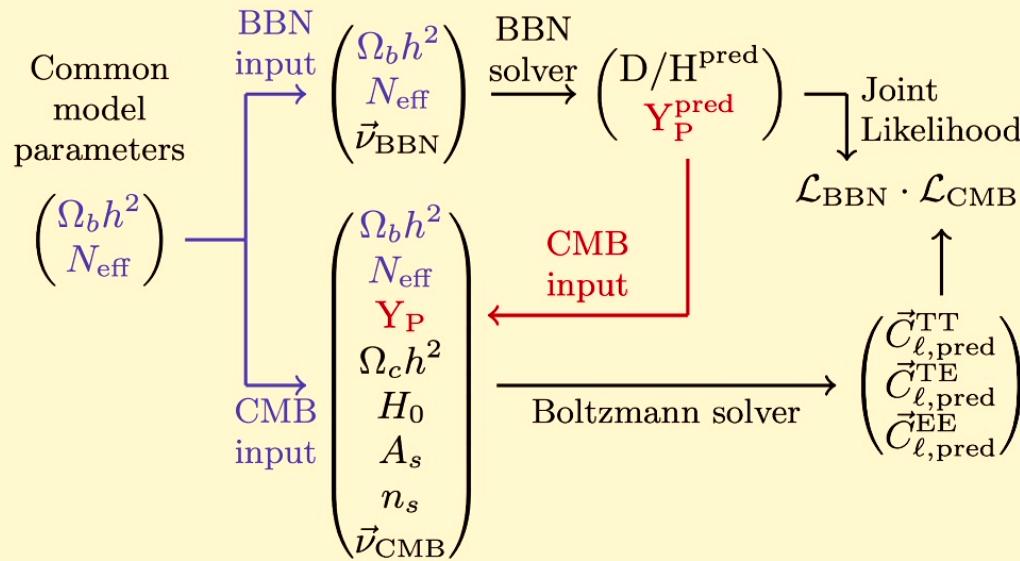
Old analyses (see e.g. Schöneberg 2401.15054, Planck 1807.06209) assume a constant theory uncertainty

LINX performs marginalization to better estimate error bars—in some cases, a 30% reduction from previous work!



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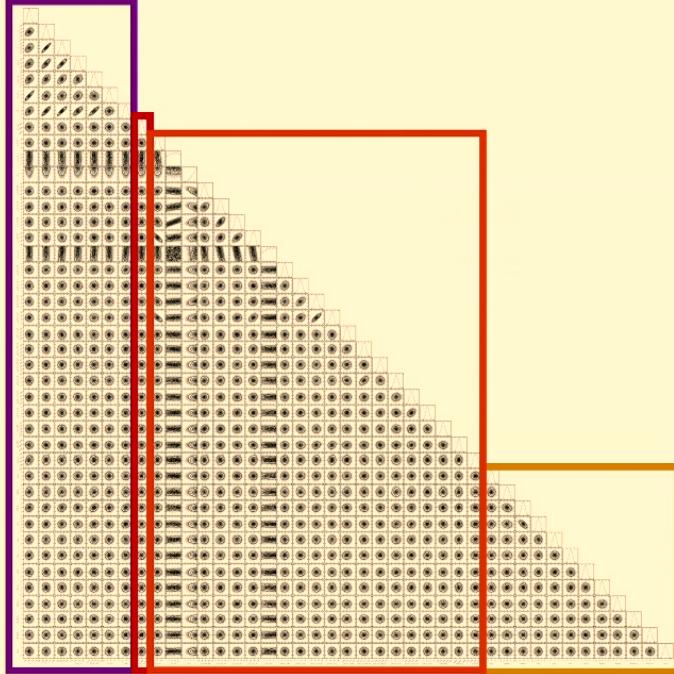
CMB + BBN Joint Fit



C.G. et al,
2408.14531

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CMB + BBN Joint Fit

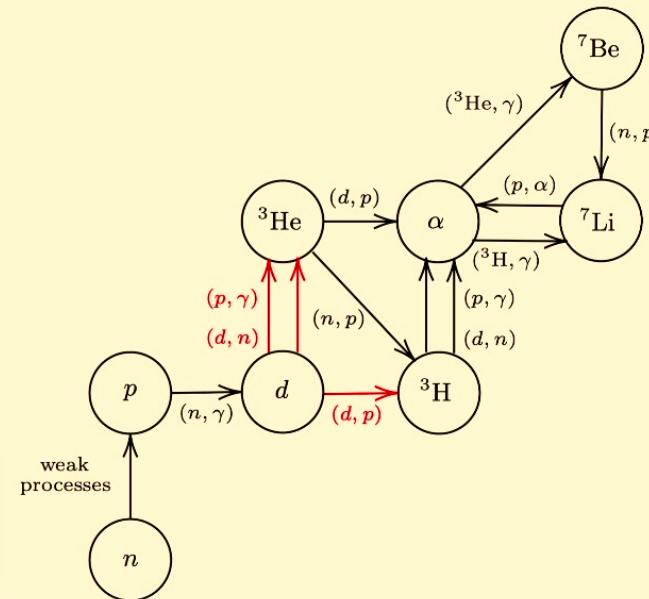
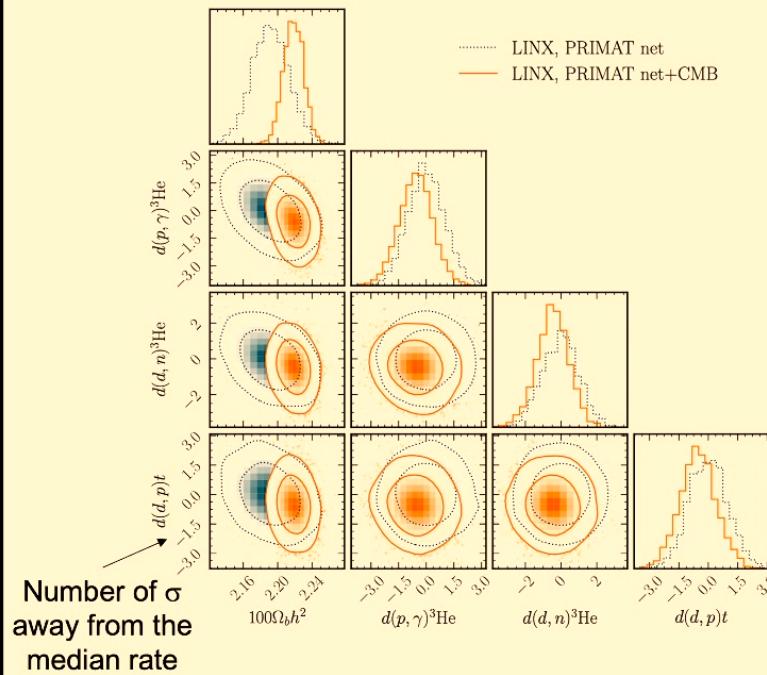


Can vary **cosmological parameters**, **neutron lifetime**, **CMB nuisance parameters**, and **individual reaction rates**

Runtime of ~2 days on 192 standard memory cores using nested sampling.

Long runtime dominated by CLASS

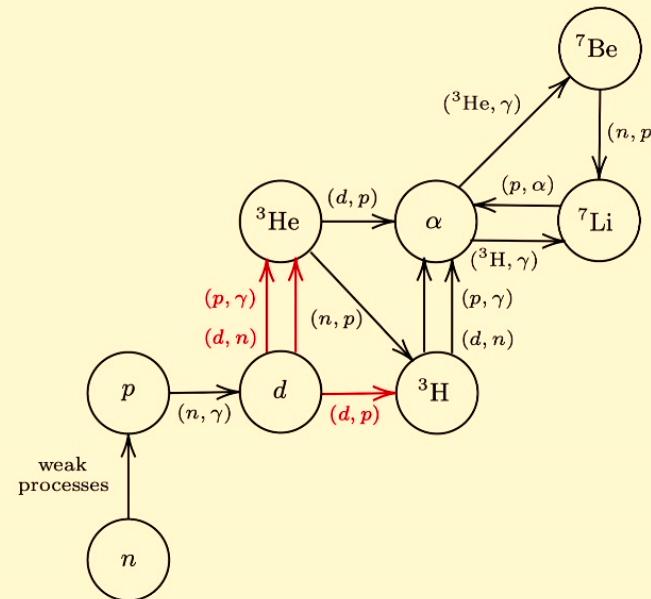
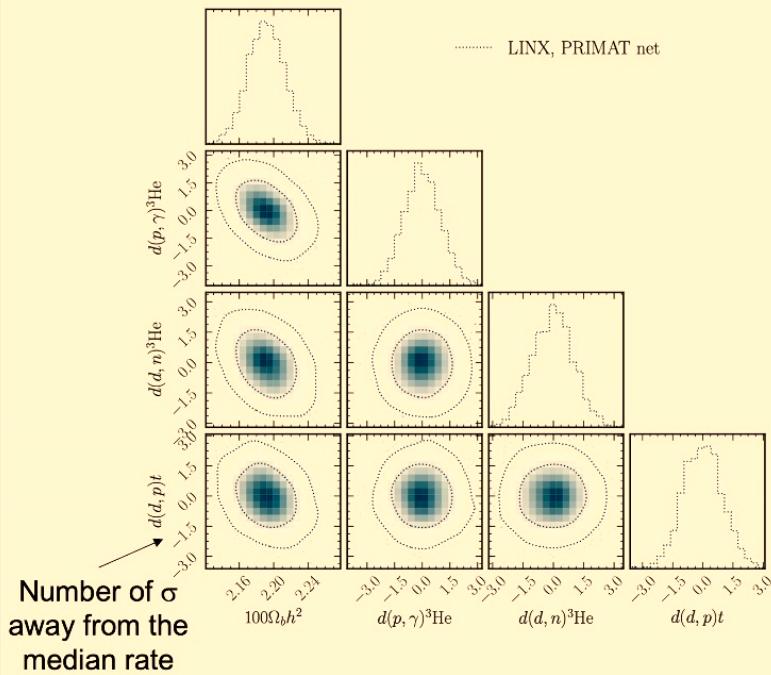
Why track nuisance parameters?



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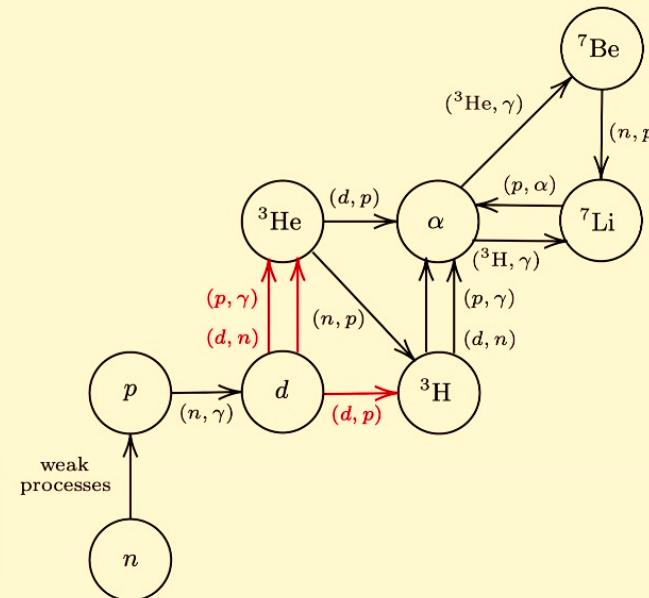
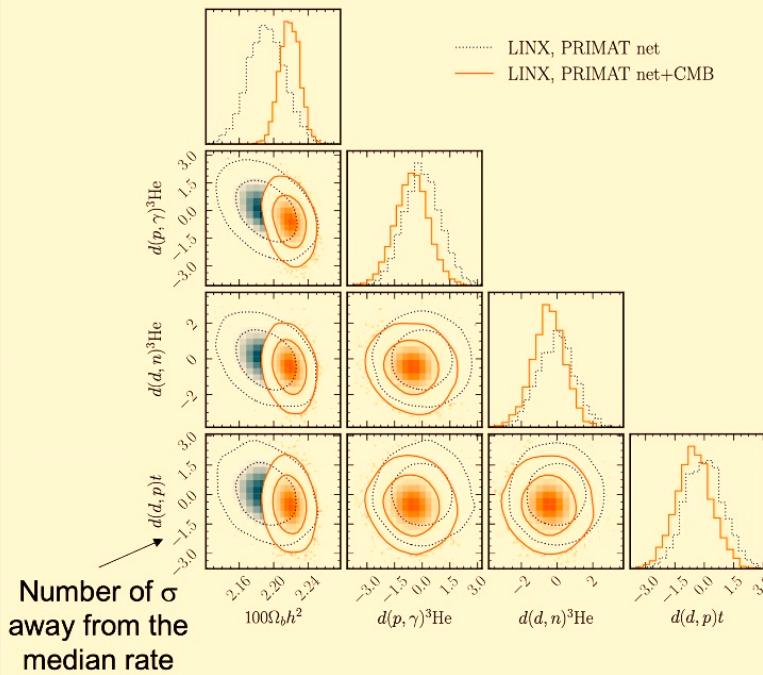
Why track nuisance parameters?



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Why track nuisance parameters?



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