

Title: Efficient quantum simulations with quDits

Speakers: Christine Muschik

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Hardware-efficient quantum computing using qudits

Christine Muschik





www.quantum-interactions.com



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Quantum
Computing

PI PERIMETER
INSTITUTE

Simulating 2D lattice gauge theories on a qudit quantum computer

Michael Meth,¹ Jan F. Haase,^{2,3,4} Jinglei Zhang,^{2,3} Claire Edmunds,¹ Lukas Postler,¹ Andrew J. Jena,^{2,3} Alex Steiner,¹ Luca Dellantonio,^{2,3,5} Rainer Blatt,^{1,6,7} Peter Zoller,^{8,6} Thomas Monz,^{1,7} Philipp Schindler,¹ Christine Muschik^{*,2,3,9} and Martin Ringbauer¹

¹Universität Innsbruck, Institut für Experimentalphysik, Technikerstraße 25a, Innsbruck, Austria

²Institute for Quantum Computing, University of Waterloo, Waterloo, Ontario, N2L 3G1, Canada

³Department of Physics & Astronomy, University of Waterloo, Waterloo, Ontario, N2L 3G1, Canada

⁴Institut für Theoretische Physik und IQST, Universität Ulm, Albert-Einstein-Allee 11, D-89069 Ulm, Germany

⁵Department of Physics and Astronomy, University of Exeter, Stocker Road, Exeter EX4 4QL, United Kingdom

⁶Institute for Quantum Optics and Quantum Information of the Austrian Academy of Sciences, Technikerstraße 21a, Innsbruck, Austria

⁷Alpine Quantum Technologies GmbH, Innsbruck, Austria

⁸Universität Innsbruck, Institut für Theoretische Physik, Technikerstraße 21a, Innsbruck, Austria

⁹Perimeter Institute for Theoretical Physics, Waterloo, Ontario, N2L 2Y5, Canada

Particle physics underpins our understanding of the world at a fundamental level by describing the interplay of matter and forces through gauge theories. Yet, despite their unmatched success, the intrinsic quantum mechanical nature of gauge theories makes important problem classes notoriously difficult to address with classical computational techniques. A promising way to overcome these roadblocks is offered by quantum computers, which are based on the same laws that make the classical computations so difficult. Here, we present a quantum computation of the properties of the basic building block of two-dimensional lattice quantum electrodynamics, involving both gauge fields and matter. This computation is made possible by the use of a trapped-ion qudit quantum processor, where quantum information is encoded in d different states per ion, rather than in two states as in qubits. Qudits are ideally suited for describing gauge fields, which are naturally high-dimensional, leading to a dramatic reduction in the quantum register size and circuit complexity. Using a variational quantum eigensolver we find the ground state of the model and observe the interplay between virtual pair creation and quantized magnetic field effects. The qudit approach further allows us to seamlessly observe the effect of different gauge field truncations by controlling the qudit dimension. Our results open the door for hardware-efficient quantum simulations with qudits in near-term quantum devices.

<https://arxiv.org/abs/2310.12110>



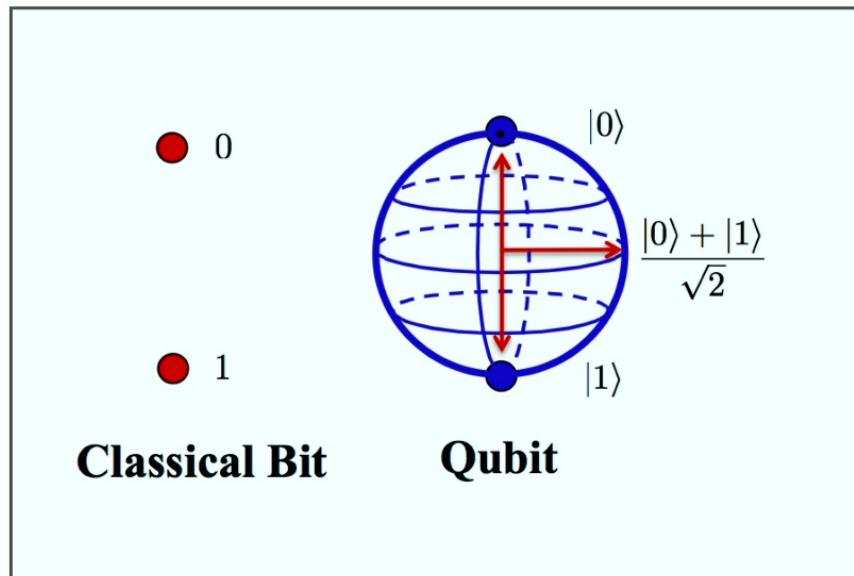
Overview

1. Introduction
2. Using qudits for problems in particle physics
3. Simulating both: gauge fields and matter
4. Increasing d
5. Conclusions

Classical computing today:
Almost exclusively based on binary encoding.



Quantum computing: Qubit encoding currently dominating

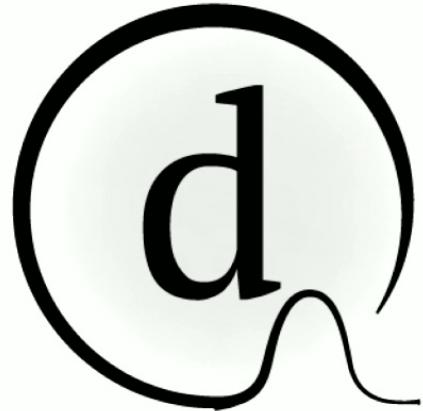


quBelt



CIFAR

Beyond binary

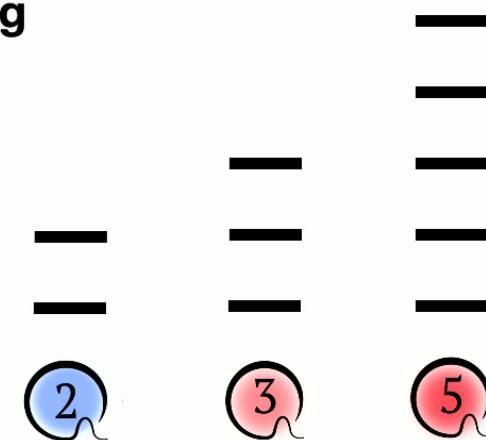


Beyond binary



Today's quantum hardware: **capable of qudit encoding**

- Trapped ions
- Superconducting architectures
- Rydberg atoms in optical tweezers
- Ultracold atoms in optical lattices
- Nuclear spins
- Photonic systems



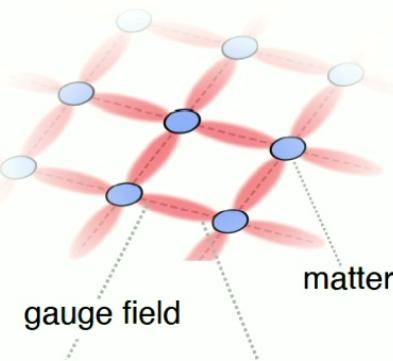
Beyond binary

(d)

Short-depth qudit circuits

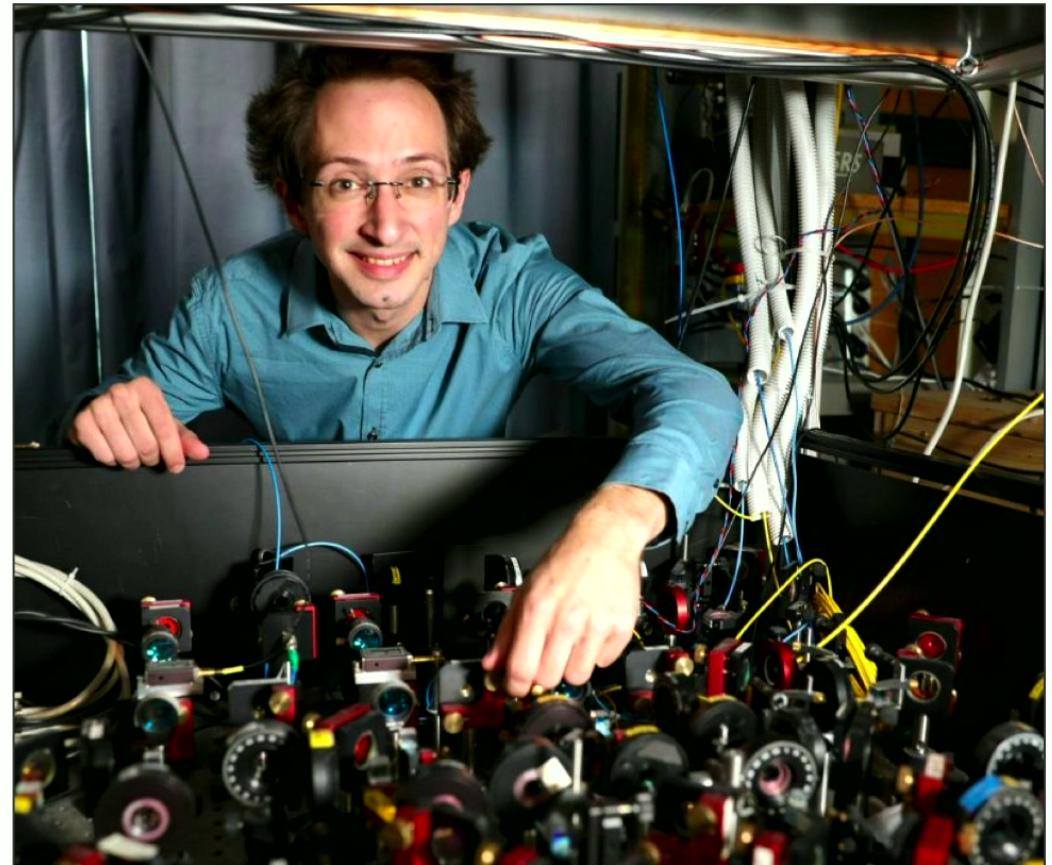


Lattice gauge theories





Mike Meth



Martin Ringbauer

Waterloo & Innsbruck

Theorie

Jan Haase
Jinglei Zhang
Andrew Jena
Luca Dellantonio
Peter Zoller
Christine Muschik

Experiment

Michael Meth
Claire Edmunds
Lukas Postler
Thomas Monz
Philipp Schindler
Rainer Blatt
Martin Ringbauer

arXiv:2310.12110

Gauge theories

Condensed matter systems
(Frustration, topological order)

Gauge theories in particle physics
(Quantum Electrodynamics, Quantum Chromodynamics,...)

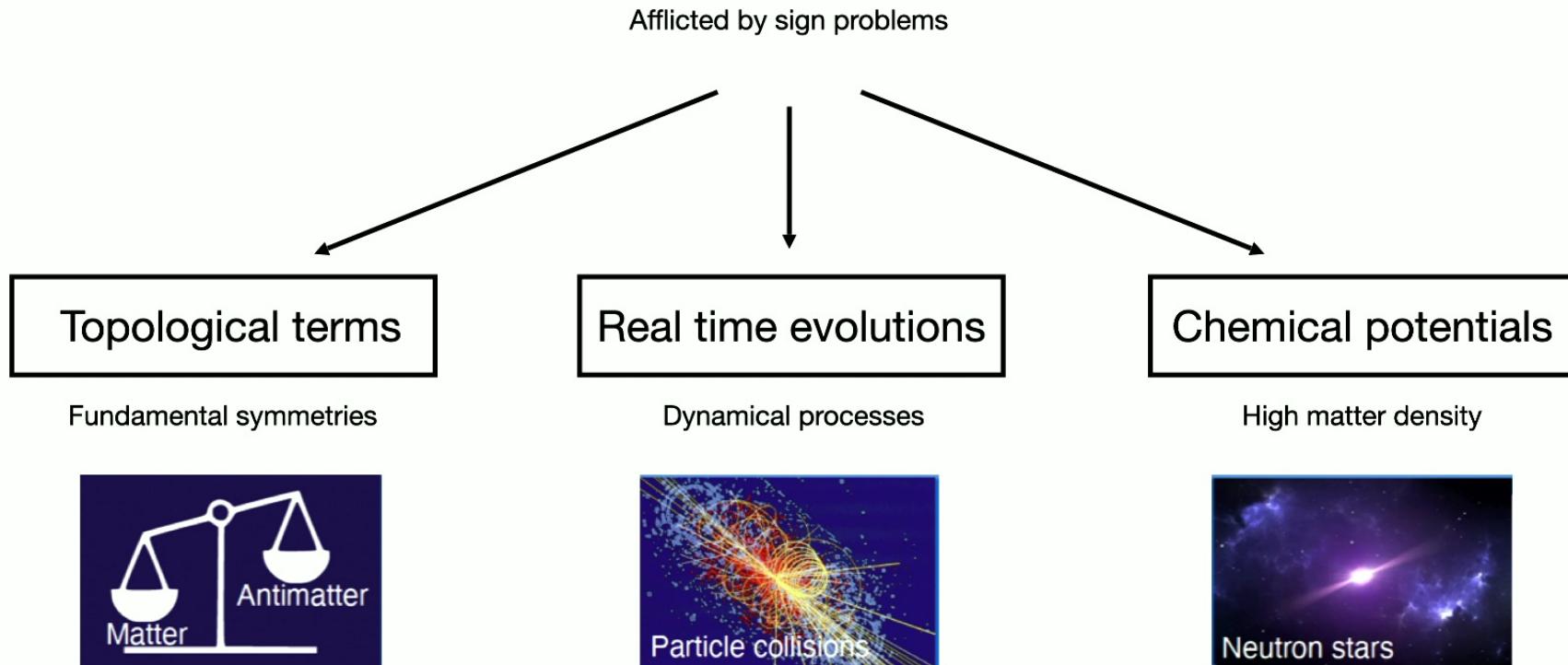
Standard model of
particle physics

FERMIONS (matter particles)			BOSONS (force carriers)	
QUARKS				
u up	c charm	t top	g gluon	H Higgs boson
d down	s strange	b bottom	γ photon	
e electron	μ muon	τ tau	Z^0 Z boson	
ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W^\pm W boson	

Simulations of gauge theories particle physics

1. Understand the physics within the standard model of particle physics.
2. Understand where the standard model fails (observations don't agree with simulations) to discover new laws and to develop new theories beyond the standard model.

Inaccessible to traditional lattice gauge theory



Two major hurdles in quantum LGT calculations

#1: Representing gauge fields

#2: Quantum simulations beyond 1D

Representing gauge fields



Matter: on vertices

Gauge fields: on links

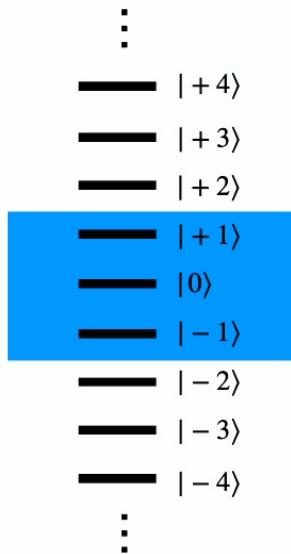
⋮	
—	$ +4\rangle$
—	$ +3\rangle$
—	$ +2\rangle$
—	$ +1\rangle$
—	$ 0\rangle$
—	$ -1\rangle$
—	$ -2\rangle$
—	$ -3\rangle$
—	$ -4\rangle$
⋮	

Gauge fields

Example: lattice QED

$$\hat{E} |E\rangle = E |E\rangle,$$

$$E = 0, \pm 1, \pm 2, \dots$$



Classical simulations and finite dimensional quantum hardware require a **truncation**:

Minimal truncation: quTrit
- field in positive direction
- no flux
- field in negative direction

Truncations for bosonic systems:

- Schwinger boson representation
- Holstein-Primakoff-representation
- Dysen-Maleev transformation
- Highly occupied boson model

Truncations for qubit systems:

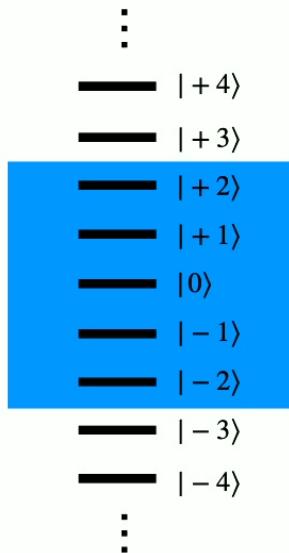
$$\hat{E} \mapsto \hat{S}^z, \quad \hat{V}^- \equiv \begin{bmatrix} 0 & \dots & \dots & 0 \\ 1 & \dots & \dots & 0 \\ 0 & \ddots & \vdots & 0 \\ 0 & \dots & 1 & 0 \end{bmatrix} \quad \text{Alternative:}$$
$$\hat{U} \mapsto \hat{V}^-, \quad \hat{V}^- \equiv \hat{S}^- / |\ell| \quad \hat{V}^- \equiv \hat{S}^- / |\ell|$$
$$\hat{S}^- = \hat{S}^x - i\hat{S}^y$$

Gauge fields

Example: lattice QED

$$\hat{E} |E\rangle = E |E\rangle,$$

$$E = 0, \pm 1, \pm 2, \dots$$



Classical simulations and finite dimensional quantum hardware require a **truncation**:

Better truncation: quQuint
- field in positive direction
- no flux
- field in negative direction

Truncations for bosonic systems:

- Schwinger boson representation
- Holstein-Primakoff-representation
- Dysen-Maleev transformation
- Highly occupied boson model

Truncations for qubit systems:

$$\begin{aligned}\hat{E} &\mapsto \hat{S}^z, & \hat{V}^- &\equiv \begin{bmatrix} 0 & \dots & \dots & 0 \\ 1 & \dots & \dots & 0 \\ 0 & \ddots & \vdots & 0 \\ 0 & \dots & 1 & 0 \end{bmatrix} & \text{Alternative:} \\ \hat{U} &\mapsto \hat{V}^-, & \hat{V}^- &\equiv \hat{S}^- / |l| & \hat{V}^- \equiv \hat{S}^- / |l| \\ && \hat{S}^- &= \hat{S}^x - i\hat{S}^y & \hat{S}^- = \hat{S}^x - i\hat{S}^y\end{aligned}$$

Representing gauge fields

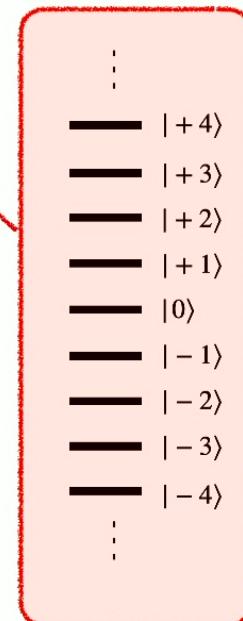


Matter: on vertices

Gauge fields: on links

1D - gauge theory

- Gauge fields can be eliminated



Open boundary conditions

Real-time dynamics of lattice gauge theories with a few-qubit quantum computer

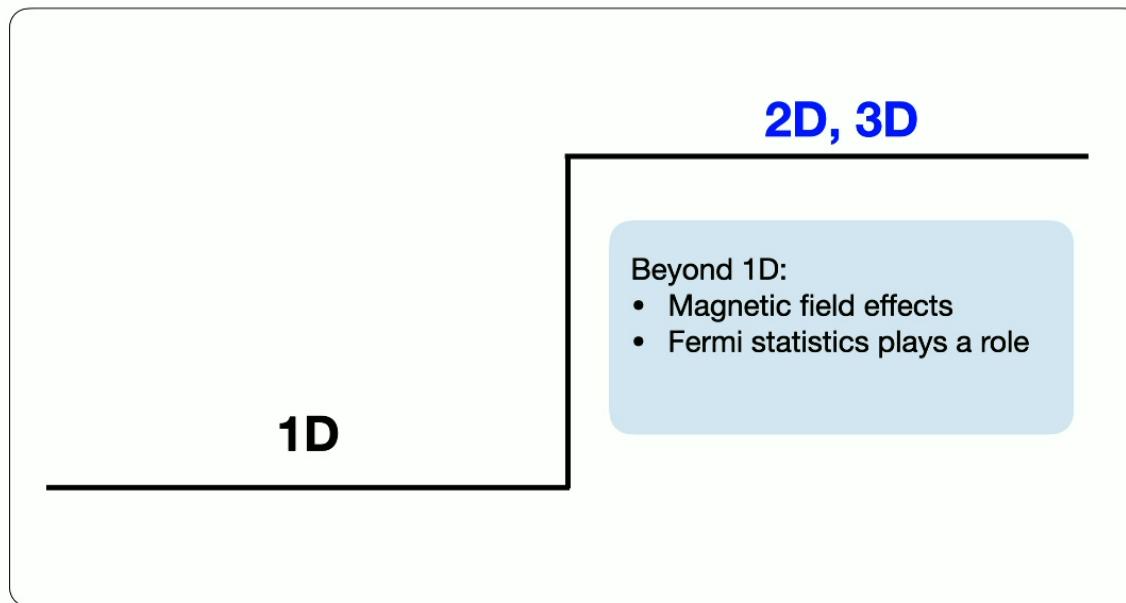
Esteban A. Martinez , Christine A. Muschik , Philipp Schindler, Daniel Nigg, Alexander Erhard,

Markus Heyl, Philipp Hauke, Marcello Dalmonte, Thomas Monz, Peter Zoller & Rainer Blatt

Nature **534**, 516–519 (2016)

- ➡ Trapped ion quantum computer
- ➡ Quantum simulation of 1D-QED
- ➡ Real-time dynamics of electron-positron pair creation.

Gauge theories for particle physics beyond 1D



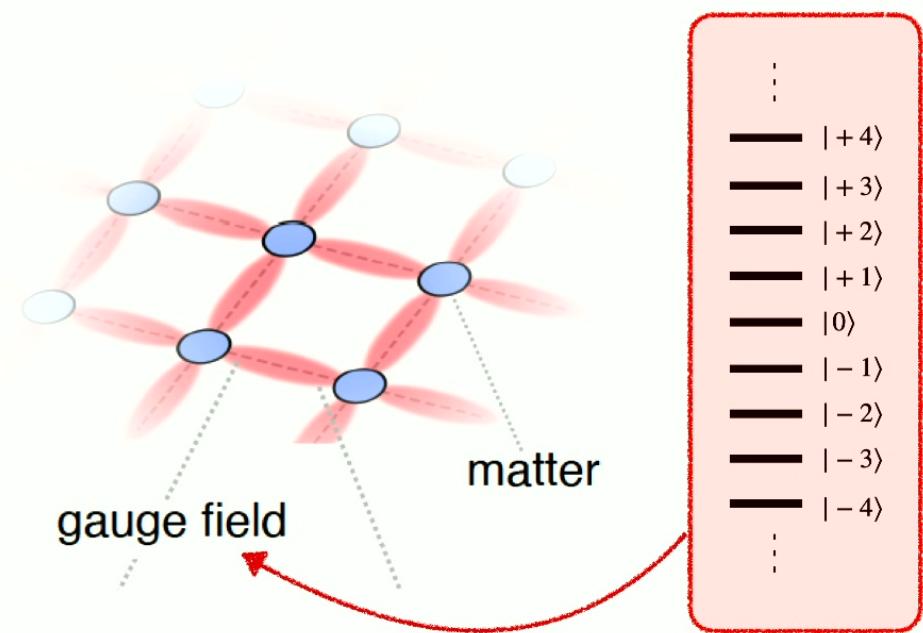
Gauge theories beyond 1D

2D gauge theory

- Not all gauge fields can be eliminated
- i.e. "dynamical" gauge fields

→ Magnetic field effects

$$\begin{array}{c} \hat{U}_{\mathbf{n}+\mathbf{e}_y, \mathbf{e}_x}^\dagger \\ \square \\ \hat{U}_{\mathbf{n}, \mathbf{e}_y}^\dagger \quad \hat{U}_{\mathbf{n}+\mathbf{e}_x, \mathbf{e}_y} \\ \hat{U}_{\mathbf{n}, \mathbf{e}_x} \end{array}$$



SCIENTIFIC
AMERICAN

In a First, Quantum Computer Simulates High-Energy Physics

The technique could allow quantum computers to address otherwise-intractable problems in particle physics

By Davide Castelvecchi, Nature magazine on June 23, 2016

Muschik says that her team is already making plans to use two-dimensional configurations

Teaching quantum computers to address problems from particle physics

Review article (QTFLAG collaboration):

Simulating Lattice Gauge Theories within Quantum Technologies: arXiv:1911.00003

Experimental demonstrations

Gauge theories for particle physics beyond 1D

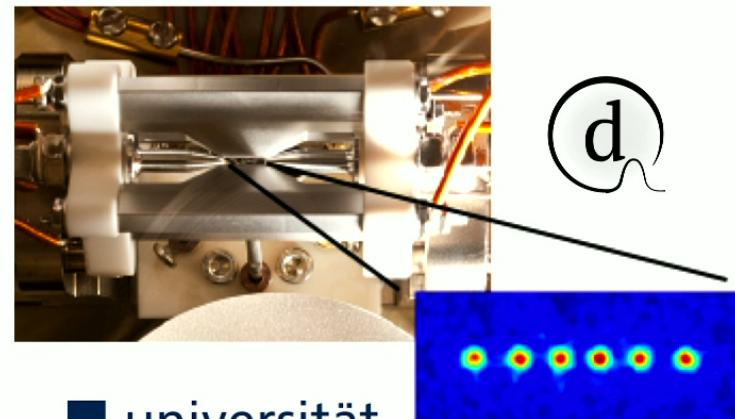
- ➡ N. Klco, M. J. Savage, and J. R. Stryker, Phys. Rev. D **101**, 074512 (2020).
- ➡ A. Ciavarella, N. Klco, and M. J. Savage, Phys. Rev. D **103**, 094501 (2021).
- ➡ S. A Rahman, R. Lewis, E. Mendicelli, and S. Powell, Phys. Rev. D **104**, 034501 (2021).
- ➡ A. N. Ciavarella and I. A. Chernyshev, Phys. Rev. D **105**, 074504 (2022).
- ➡ S. A Rahman, R. Lewis, E. Mendicelli, and S. Powell, Phys. Rev. D **106**, 074502 (2022).
- ➡ S. A. Rahman, R. Lewis, E. Mendicelli, and S. Powell, “Real time evolution and a traveling excitation in SU(2) pure gauge theory on a quantum computer,” (2022), arXiv:2210.11606 [hep-lat].
- ➡ A. N. Ciavarella, “Quantum Simulation of Lattice QCD with Improved Hamiltonians,” (2023), arXiv:2307.05593 [hep-lat].

→ Impressive Advances!
(but gauge fields or matter fields are trivial)

Experimental qudit system

Linear ion-trap quantum processor
with all-to-all connectivity

Extending qubit entangling gates
to mixed-dimensional qudit systems

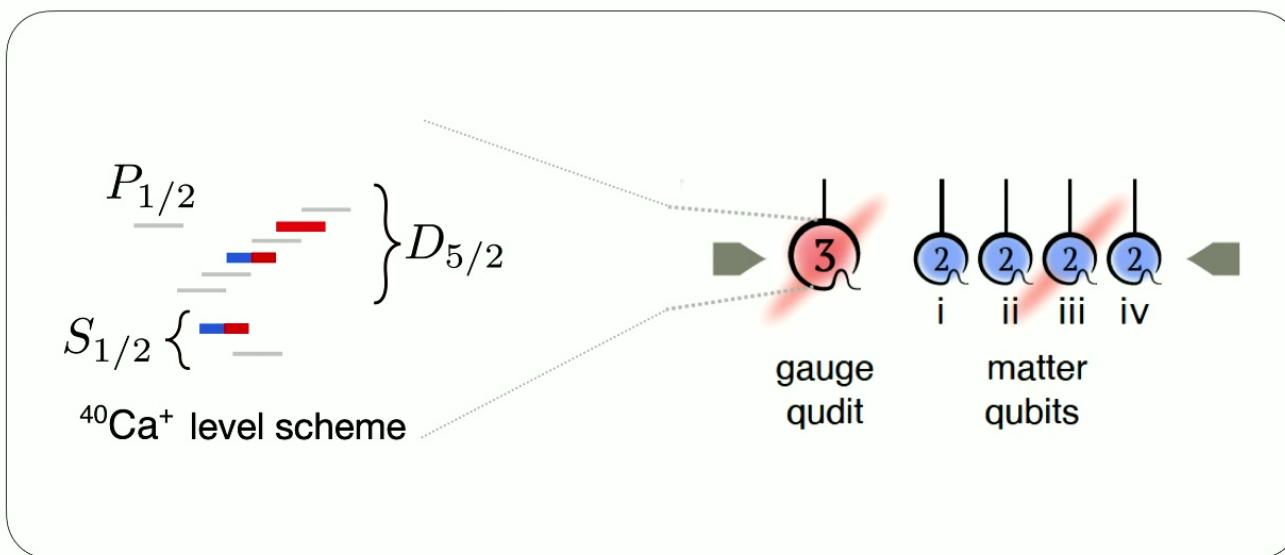


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innsbruck

Martin Ringbauer

M. Ringbauer, M. Meth, L. Postler, R. Stricker, R. Blatt, P. Schindler, T. Monz, Nature Physics 18, 1053 (2022)

Encoding qudits in trapped ions



Experimental demonstration

Gauge theories for particle physics beyond 1D



Including both - dynamical gauge and matter fields

Necessary theory steps

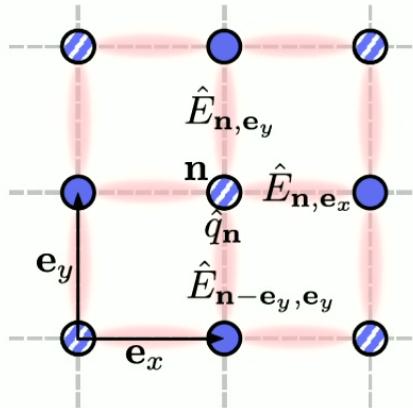
Step A: Efficient Hamiltonian formulation

A resource efficient approach for quantum and classical simulations of gauge theories in particle physics
J. Haase, L. Dellantonio, A. Celi, D. Paulson, A. Kan, K. Jansen, C. Muschik
Quantum 5, 393 (2021).

Step B: Efficient Quantum Simulation scheme

Towards simulating 2D effects in lattice gauge theories on a quantum computer
D. Paulson, L. Dellantonio, J. Haase, A. Celi, A. Kan, A. Jena, C. Kokail, R. van Bijnen, K. Jansen, P. Zoller, C. Muschik
PRX Quantum 2, 030334 (2021).

2D lattice QED



Even $n_x + n_y$ lattice sites

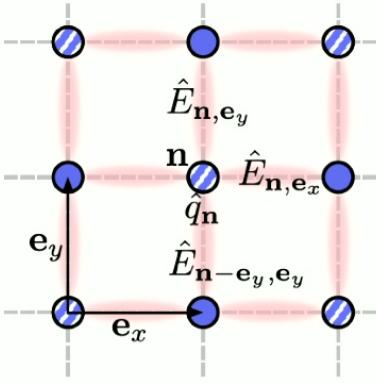
$$\circlearrowleft \cong |\downarrow\rangle \cong |e\rangle, \quad \hat{q}_{\mathbf{n}} = -1$$

$$\bullet \cong |\uparrow\rangle \cong |v\rangle, \quad \hat{q}_{\mathbf{n}} = 0$$

Odd $n_x + n_y$ lattice sites

$$\circlearrowright \cong |\downarrow\rangle \cong |v\rangle, \quad \hat{q}_{\mathbf{n}} = 0$$

$$\circlearrowright \cong |\uparrow\rangle \cong |p\rangle, \quad \hat{q}_{\mathbf{n}} = +1$$



2D lattice QED

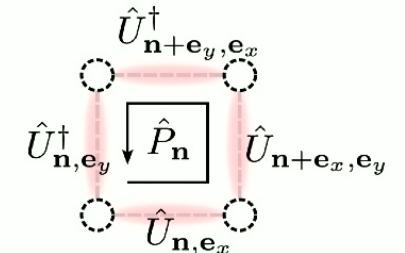
$$\hat{H} = g^2 \hat{H}_E + \frac{1}{g^2} \hat{H}_B + m \hat{H}_m + \Omega \hat{H}_k$$

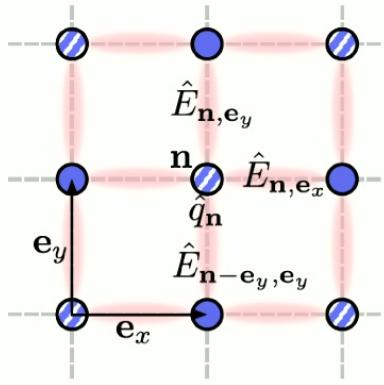
Electric field energy $\longrightarrow \hat{H}_E = \frac{1}{2} \sum_{\mathbf{n}} \left(\hat{E}_{\mathbf{n}, \mathbf{e}_x}^2 + \hat{E}_{\mathbf{n}, \mathbf{e}_y}^2 \right),$

Magnetic field energy $\longrightarrow \hat{H}_B = -\frac{1}{2} \sum_{\mathbf{n}} \left(\hat{P}_{\mathbf{n}} + \hat{P}_{\mathbf{n}}^\dagger \right), \quad \hat{P}_{\mathbf{n}} = \hat{U}_{\mathbf{n}, \mathbf{e}_x} \hat{U}_{\mathbf{n}+\mathbf{e}_x, \mathbf{e}_y} \hat{U}_{\mathbf{n}+\mathbf{e}_y, \mathbf{e}_x}^\dagger \hat{U}_{\mathbf{n}, \mathbf{e}_y}^\dagger.$

Particle rest mass $\longrightarrow \hat{H}_m = \sum_{\mathbf{n}} (-1)^{n_x+n_y} \hat{\phi}_{\mathbf{n}}^\dagger \hat{\phi}_{\mathbf{n}},$

Kinetic energy term
(pair creation) $\longrightarrow \hat{H}_k = \sum_{\mathbf{n}} \sum_{\mu=x,y} \left(\hat{\phi}_{\mathbf{n}} \hat{U}_{\mathbf{n}, \mathbf{e}_\mu}^\dagger \hat{\phi}_{\mathbf{n}+\mathbf{e}_\mu}^\dagger + \text{H.c.} \right).$





2D lattice QED

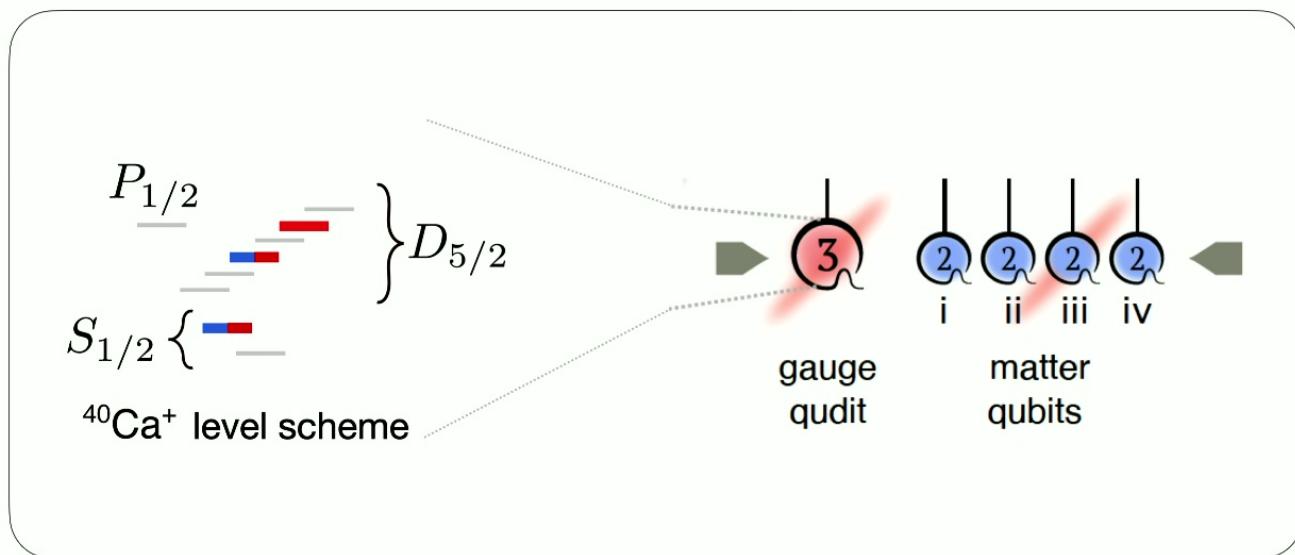
$$\hat{H} = g^2 \hat{H}_E + \frac{1}{g^2} \hat{H}_B + m \hat{H}_m + \Omega \hat{H}_k$$

Classical Gauss law: $\nabla \mathbf{E}(\mathbf{r}) - \rho(\mathbf{r}) = 0$

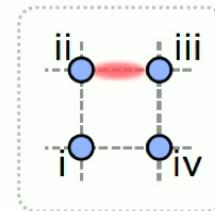
QED Gauss law: $\hat{G}_{\mathbf{n}} |\Psi_{\text{phys}}\rangle = 0,$

$$\hat{G}_{\mathbf{n}} = \sum_{\mu} (\hat{E}_{\mathbf{n}, \mathbf{e}_{\mu}} - \hat{E}_{\mathbf{n}-\mathbf{e}_{\mu}, \mathbf{e}_{\mu}}) - \hat{q}_{\mathbf{n}}$$

Experiment



Basic building block: Plaquette



Experiment

Variational Quantum Eigensolver



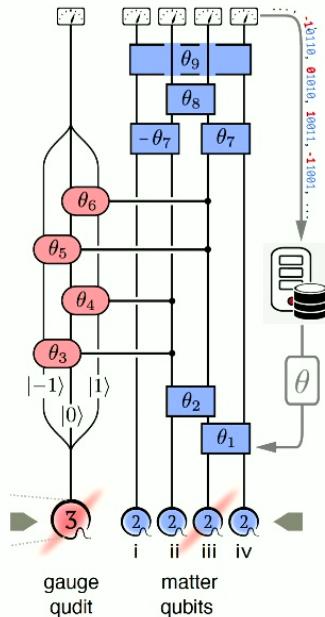
Ground state preparation



Hybrid qubit-quTrit variational circuit

Measure plaquette expectation value

$$\langle \hat{\square} \rangle \propto \langle \hat{H}_B \rangle$$



Experiment

Variational Quantum Eigensolver



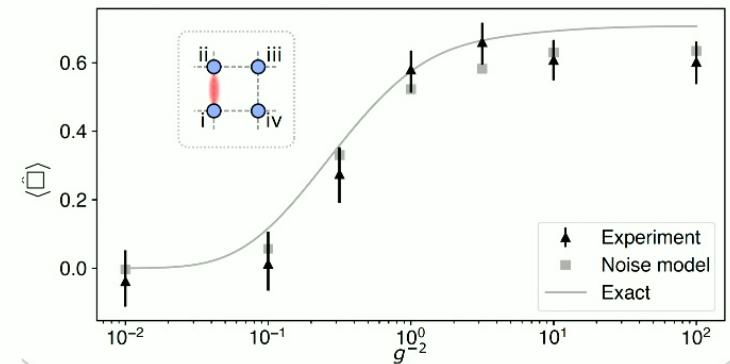
Ground state preparation



Hybrid qubit-quTrit variational circuit

Measure plaquette expectation value

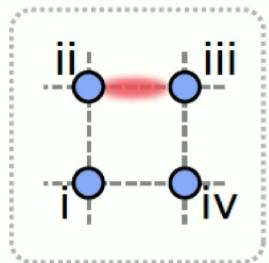
$$\langle \hat{\square} \rangle \propto \langle \hat{H}_B \rangle$$



Qudit encoding: short circuits



Basic building block: Single plaquette:



	Qudit enc			Qubit enc.		
dimension d	3	5	7	3	5	7
register size	5	5	5	7	9	11
CNOT count	26	34	42	96	174	252
CNOT fidelity	99%			99.5%		
approx. circ. fid.	77%	71%	66%	38%	17%	8%
CNOT fidelity	99.5%			99.5%		
approx. circ. fid.	88%	84%	81%	62%	42%	28%

QuDit review article:

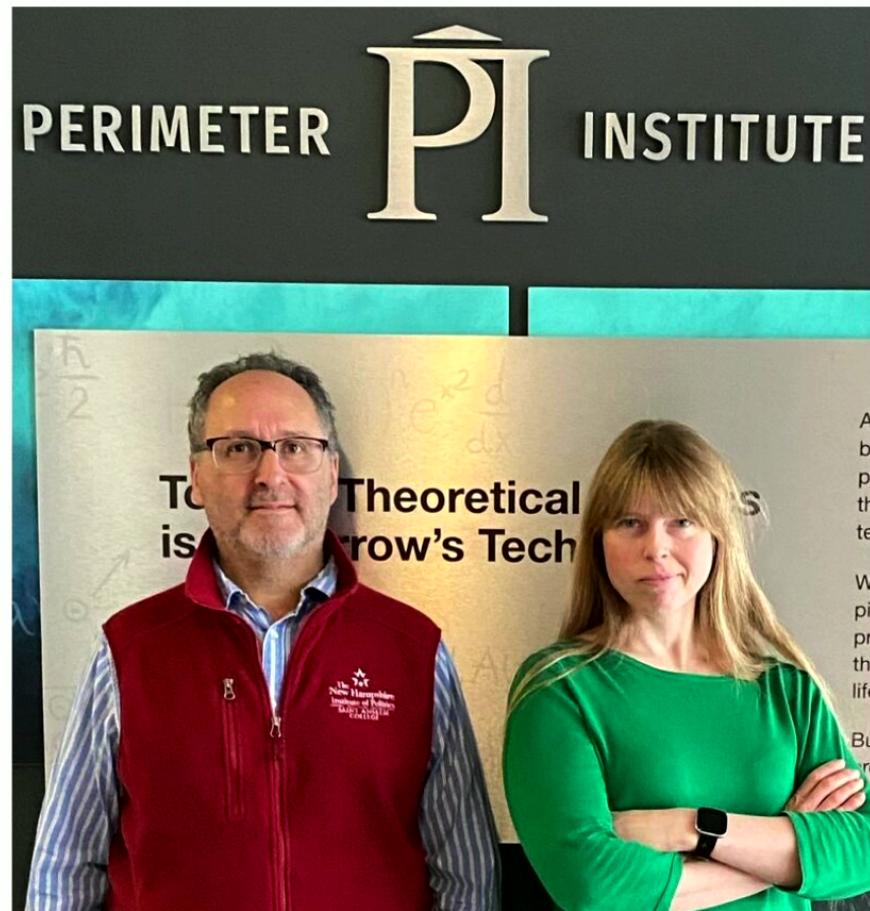
Y. Wang, Z. Hu, B. C. Sanders,
and S. Kais, Front. Phys. 8 (2020).

3.14 reasons why **Barry Sanders** was the worst guest to spend PI day with:

- During his entire stay, I have not heard him say anything remotely irrational.
- He repeated himself several times!
- His visit has, unfortunately, definitely ended.

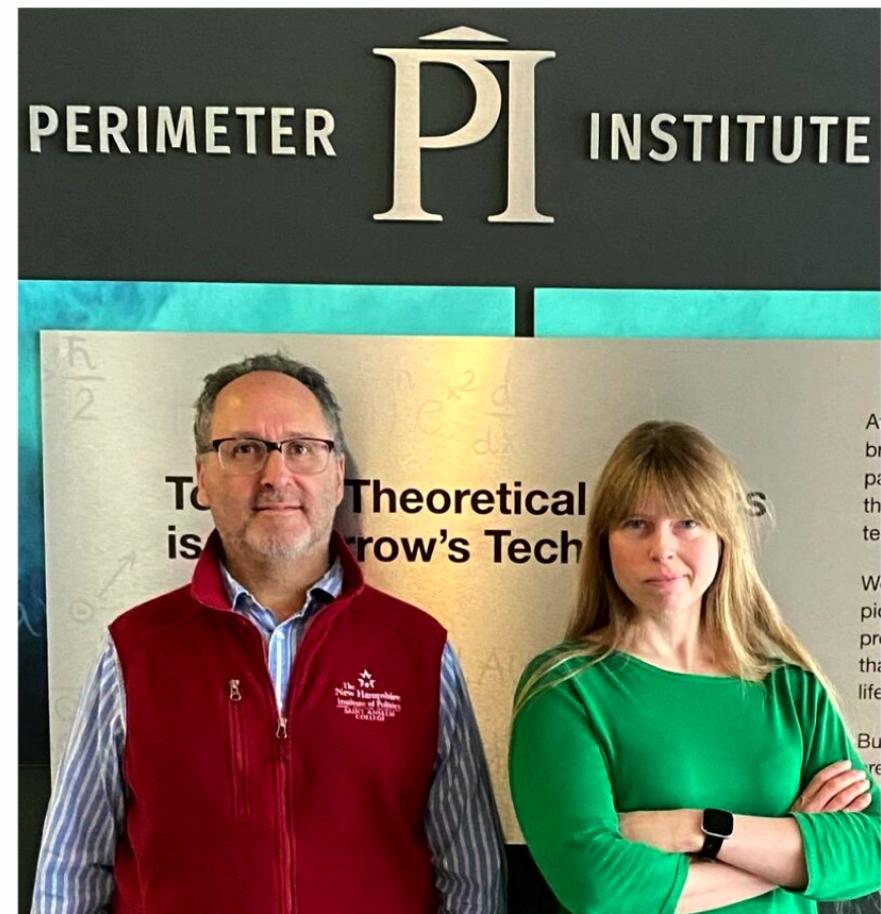
(and if you watched him moderating the Quantum Computing panel on the Quantum Days 2024, you know that he will never allow a conversation to go in a circle.)

On the bright side: We spent PI day at PI!



QuDit review article:

Y. Wang, Z. Hu, B. C. Sanders,
and S. Kais, Front. Phys. 8 (2020)



Increasing d

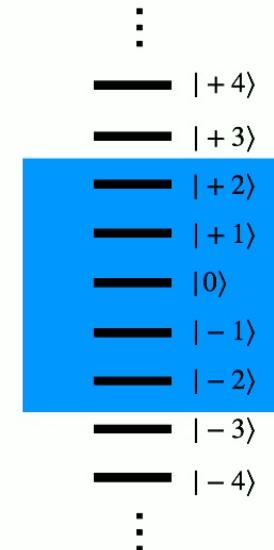
So far:

- Relevant physics could be simulated.
- Using the minimal truncation $d = 3$.

Next step:

- Study how we can improve the truncation.
- Go to $d = 5$.

Gauge fields

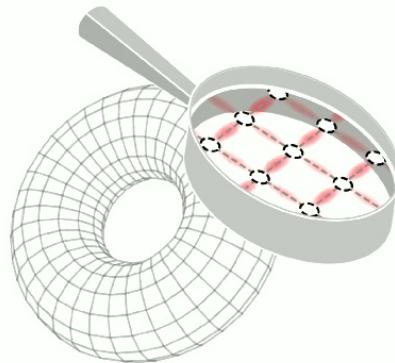


Increasing d

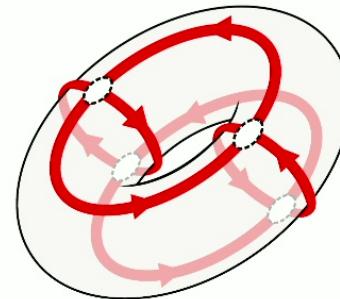
Improving the gauge-field truncation

Concrete example:

- Study the dependence of $\langle \hat{\square} \rangle$ on the bare coupling at different discretizations.
- Consider QED on a 2D lattice with periodic boundary conditions.
- No matter fields.



2D-QED with periodic boundary conditions



One plaquette with periodic boundary conditions

Increasing d

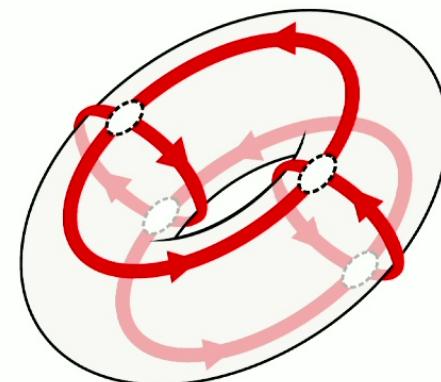
Improving the gauge-field truncation

Pure gauge theory:

$$\hat{H} = g^2 \hat{H}_E + (1/g^2) \hat{H}_B$$

Single plaquette:

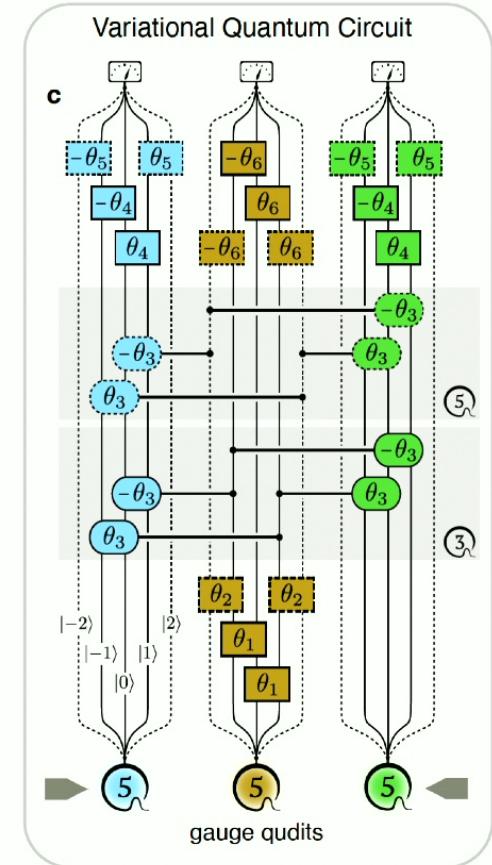
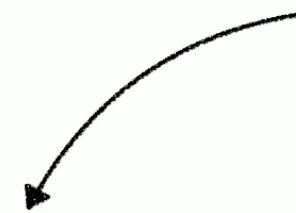
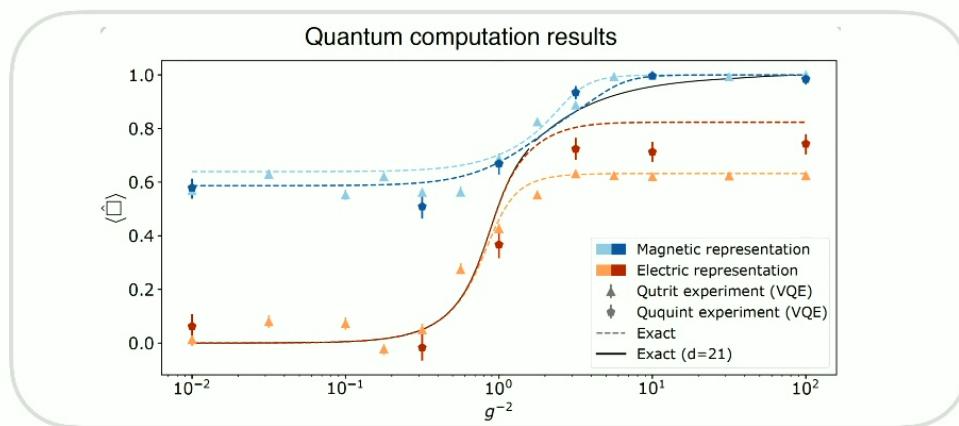
- 4 vertices.
- 8 links = 8 gauge fields.
- Gauss law: 5 independent gauge fields.
- For ground state properties **3 gauge fields**.



J. F. Haase, L. Dellantonio, A. Celi, D. Paulson, A. Kan, K. Jansen, and C. A. Muschik, Quantum **5**, 393 (2021).

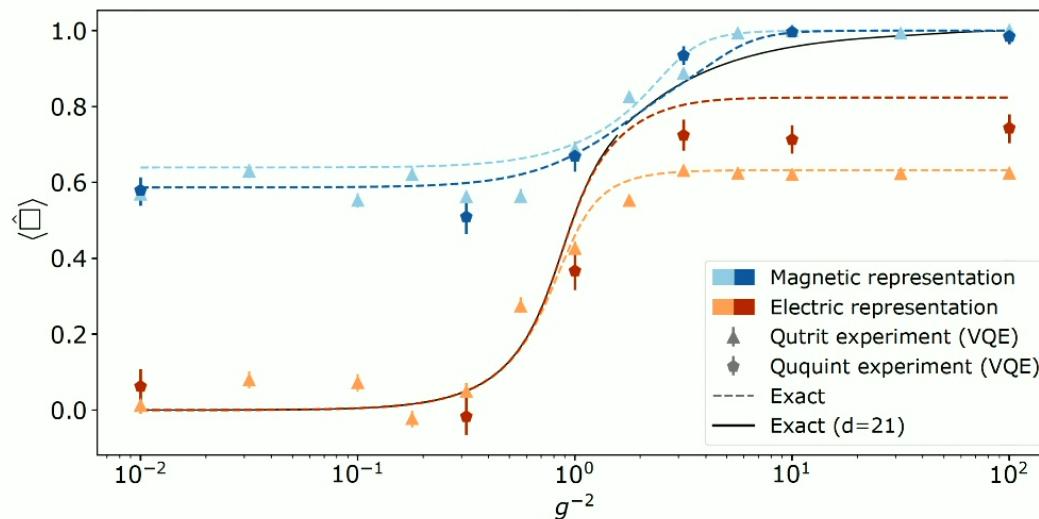
Experiment:

- VQE protocol to prepare the ground state.
- Measurement of $\langle \hat{\square} \rangle$ in the ground state
- for different values of the bare coupling g .



The closing of the gap between the values in the E-representation (red) and the B-representation (blue) corresponds to a better representation of the ground state.

Also: avoid “freezing” for large values of g^{-2} .

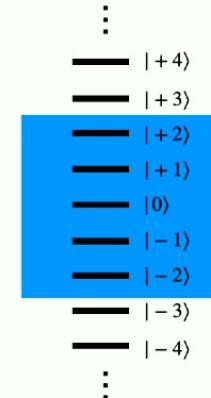


Take home messages:

- Full VQE ground state search for 3-level qutrits and 5-level ququints.
- Qudits allows to seamlessly observe the effect of different gauge field truncations by controlling the qudit dimension.

Experimental demonstration

Gauge theories for particle physics beyond 1D

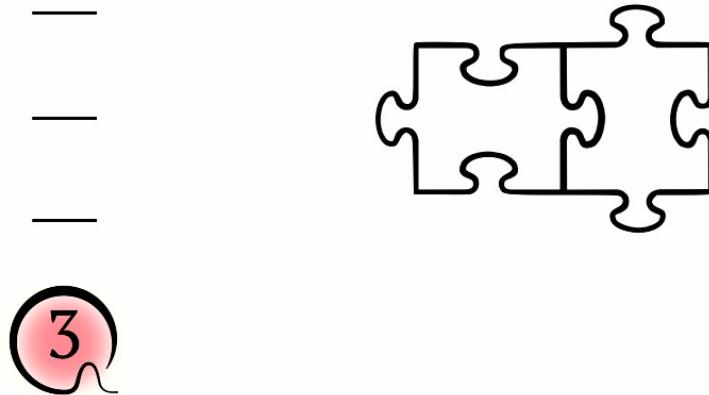


Including both - dynamical gauge and matter fields

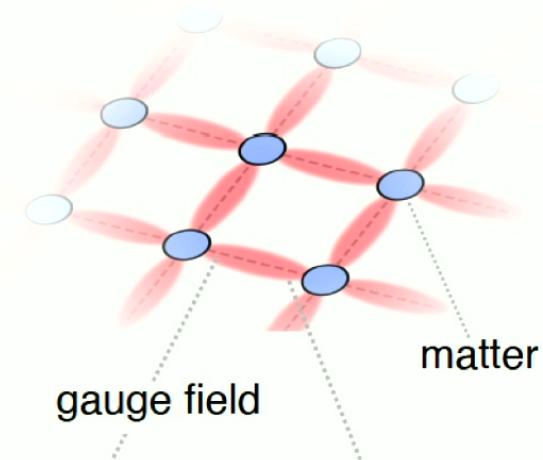
Adjustable discretisation of the gauge fields

Beyond binary

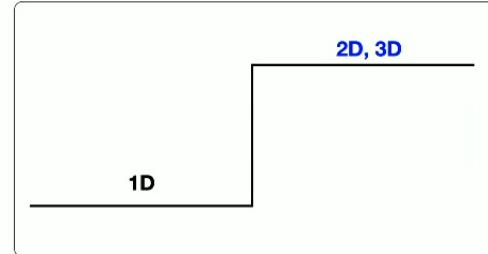
short-depth qudit circuits



Lattice Gauge theories for particle physics



Outlook



Long-term vision:

- Simulate nature in 3 spatial dimensions
- Scaling up the system sizes
- Both make an efficient gauge field representation even more important

Qudit techniques:

- Can be implemented on many quantum computing platforms
- Are directly adaptable to digital quantum simulations of real-time dynamics
Exciting perspective: quantum LGT calculations in regimes that are classically intractable due to sign problems
- Applications in chemistry and materials science
e.g. in the form of exotic large-spin models: see e.g. work at Harvard (N. Maskara et. al.)

Qudit workshop

9-10 July 2024, Waterloo

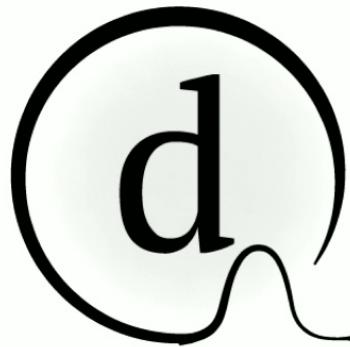
Quantum Algorithms

Quantum Simulations

Quantum Networks

Quantum Sensing

Theory + Experimental



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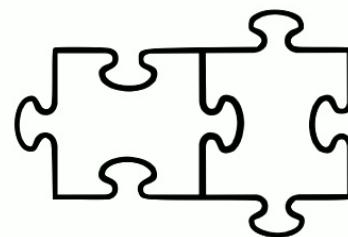
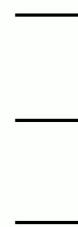
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arXiv:2310.12110

Beyond binary



short-depth qudit circuits



Lattice Gauge theories for particle physics

