

**Title:** Coherent Control and Readout of a 25-level Trapped  $^{137}\text{Ba}^+$  Ion Qudit

**Speakers:** Pei Jiang Low

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**Subject:** Quantum Information

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**Abstract:**

Scaling up the Hilbert space of a quantum computer is essential to demonstrate quantum advantage over classical computing. Other than packing more quantum information carrier in a quantum computing system, one avenue for increasing the quantum computing Hilbert space is by encoding more logical states per quantum information carrier (making it a qudit), rather than a traditional 2-state qubit encoding. Trapped ions typically exhibit multiple (meta)stable electronic energy states that are suitable for logical state encodings and are a suitable platform for qudit applications. In this talk, I present our work on realizing a 25-level trapped  $^{137}\text{Ba}^+$  ion qudit, which is the maximal encoding allowed by the single-shot readout protocol that we employ. We show that the harnessing of these additional logical states per ion can be done by just adding a few more lasers to the control system, without any trapping architecture modification compared to a qubit system. I will also present the experimental challenges and limitations, fundamental and technical, relevant to realizing a 25-level trapped  $^{137}\text{Ba}^+$  ion qudit.



# Coherent Control and Readout of a 25-level Trapped $^{137}\text{Ba}^+$ Ion Qudit

Pei Jiang Low

30<sup>th</sup> September 2024

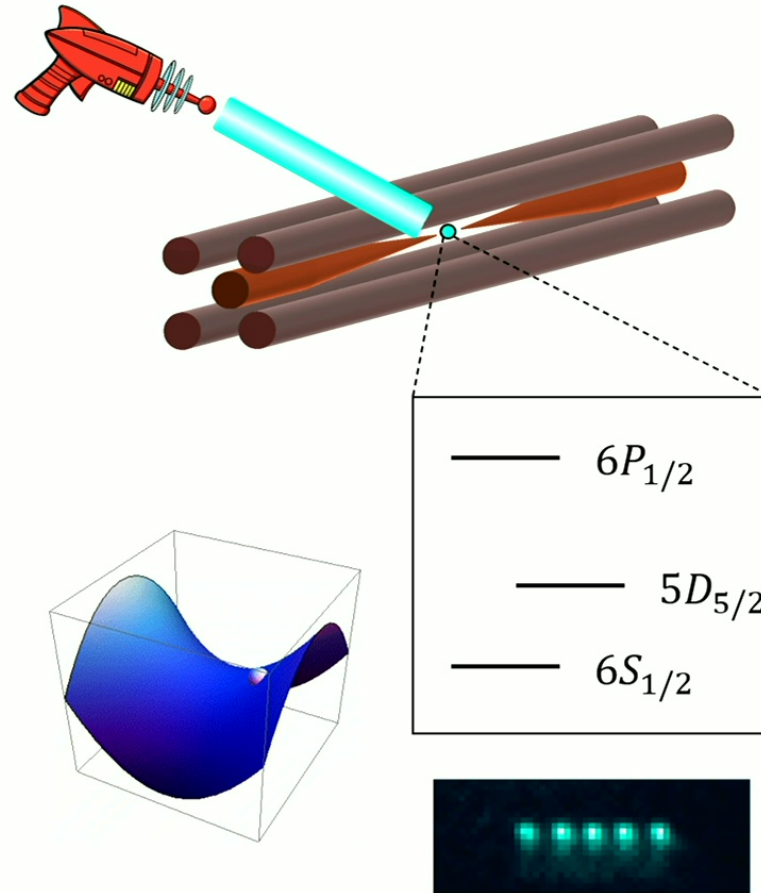


UNIVERSITY OF  
**WATERLOO**

**IQC** Institute for  
Quantum  
Computing

# Brief Summary of Trapped Ion Quantum Computing

- Ion(s) trapped in vacuum with electric field.
- Electronic energy levels of the ion used to encode computational states.
- Lasers used to manipulate the electronic states.
- Chain of ion move with shared phonon modes due to Coulomb interactions, used for entanglement between ions.

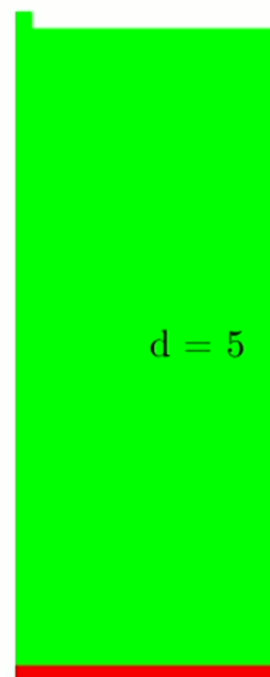


# Why Qudits?

- Quantum advantage over classical computing requires sufficiently large computational space.
- Qudit scales up Hilbert space with the same number of physical information carrier units:  $2^N$  (qubits) vs  $d^N$  (qudits).



Qudits: 2 or more encoding per information carrier.



2D representation of the size of Hilbert space

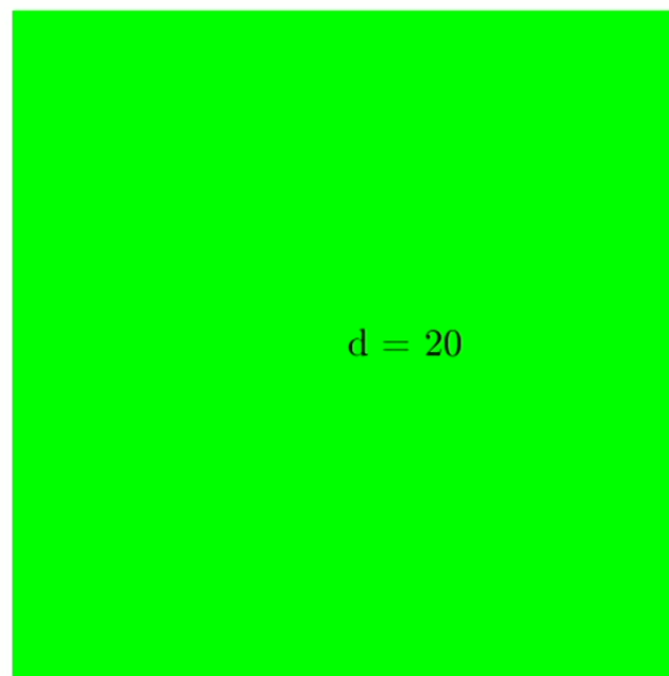
Red – Hilbert space spanned by qubits

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- Qudit scales up Hilbert space with the same number of physical information carrier units:  $2^N$  (qubits) vs  $d^N$  (qudits).
- Larger Hilbert space can be used for simplifying qubit operations.<sup>1,2</sup>
- Higher spin systems' quantum simulation more direct with qudits. E.g. Spin 1 system using qutrits.<sup>3,4</sup>
- More relaxed quantum error correction threshold.<sup>5,6,7</sup>

<sup>1</sup>Ralph, T. C., Resch, K. J. & Gilchrist, A. Phys. Rev. A 75, 022313, DOI: 10.1103/307 PhysRevA.75.022313 (2007).

<sup>2</sup>Lanyon, B. P. et al. Nat. Phys. 5, 134–140, DOI: 305 10.1038/nphys1150 (2008).

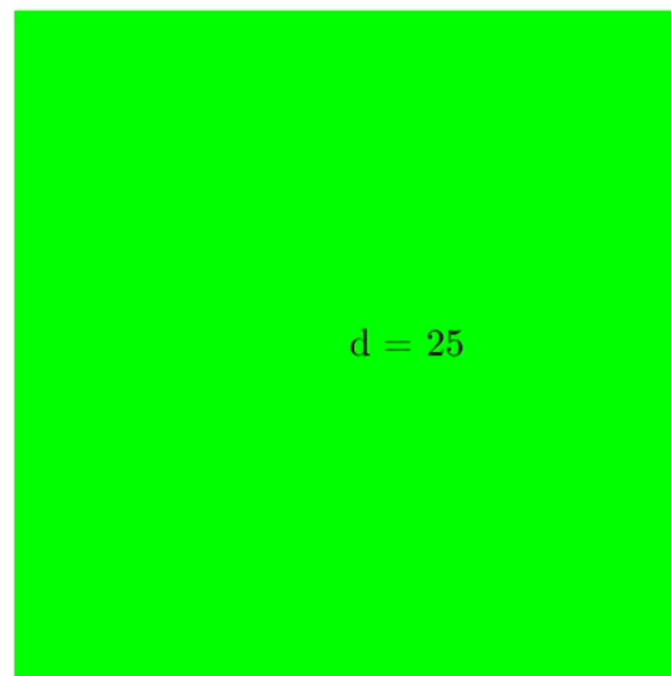
<sup>3</sup>Senko, C. et al. Phys. Rev. X 5, 021026, DOI: 301 10.1103/PhysRevX.5.021026 (2015).

<sup>4</sup>Andrade, B. et al. Quantum Sci. Technol. 7, 034001, DOI: 10.1088/2058-9565/ac5f5b (2022).

<sup>5</sup>Campbell, E. T., Anwar, H. & Browne, D. E. Phys. Rev. X 2, 041021, DOI: 10.1103/PhysRevX.2.041021 (2012).

<sup>6</sup>Andrist, R. S., Wootton, J. R. & Katzgraber, H. G. Phys. Rev. A 91, 042331, DOI: 10.1103/PhysRevA.91.042331 (2015).

<sup>7</sup>Watson, F. H. E., Anwar, H. & Browne, D. E. Phys. Rev. A 92, 299 032309, DOI: 10.1103/PhysRevA.92.032309 (2015).



2D representation of the size of Hilbert space

Red – Hilbert space spanned by qudits

# Trapped Ions as Qudits

- Common to find more than 2 stable energy states.
- Binary restriction usually for magnetic field insensitivity.
- Magnetic field noise can be overcome<sup>1</sup>, encode more!
- **Main question:** Qudits advantageous in theory, realistically advantageous? Practical challenges vs theoretical gain.
- Universal single qudit operations + one ion-ion entangling gate form universal gate set.<sup>2,3,4</sup>
  - Any 2 encoded states in an ion connected via a series of coherent Givens rotations → ability to perform universal single qudit operations.
- Entangling gate can be the same as qubit entangling gate. Assessing single qudit gates challenges = assessing universal qudit gate challenges.

$$\begin{array}{cccccccccc}
 m_{\tilde{F}} = & -4 & -3 & -2 & -1 & 0 & 1 & 2 & 3 & 4 \\
 \tilde{F} = & & & & & & & & & \\
 1 & & & & & \overline{|22\rangle} & \overline{|23\rangle} & \overline{|24\rangle} & & \\
 2 & & & & \overline{|17\rangle} & \overline{|16\rangle} & \overline{|19\rangle} & \overline{|18\rangle} & \overline{|20\rangle} & 5D_{5/2} \\
 3 & & \overline{|5\rangle} & \overline{|11\rangle} & \overline{|10\rangle} & \overline{|13\rangle} & \overline{|7\rangle} & \overline{|14\rangle} & \overline{|12\rangle} & \\
 4 & \overline{|1\rangle} & \overline{|2\rangle} & \overline{|3\rangle} & \overline{|6\rangle} & \overline{|15\rangle} & \overline{|21\rangle} & \overline{|9\rangle} & \overline{|4\rangle} & \overline{|8\rangle}
 \end{array}$$

$$\begin{array}{ccccccc}
 \tilde{F} = & 2 & \text{---} & \text{---} & \overline{|0\rangle} & \text{---} & \text{---} \\
 & 1 & & \text{---} & \text{---} & \text{---} & \\
 & & & & & & 6S_{1/2}
 \end{array}$$

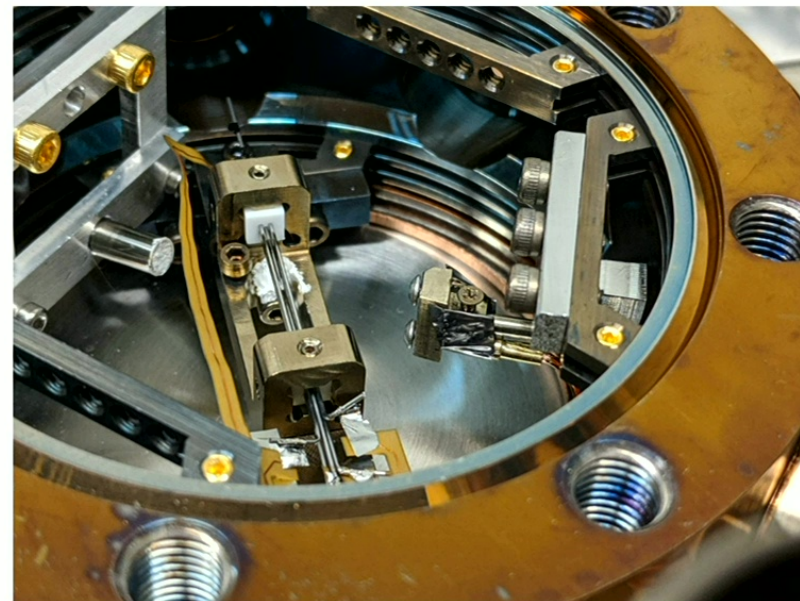
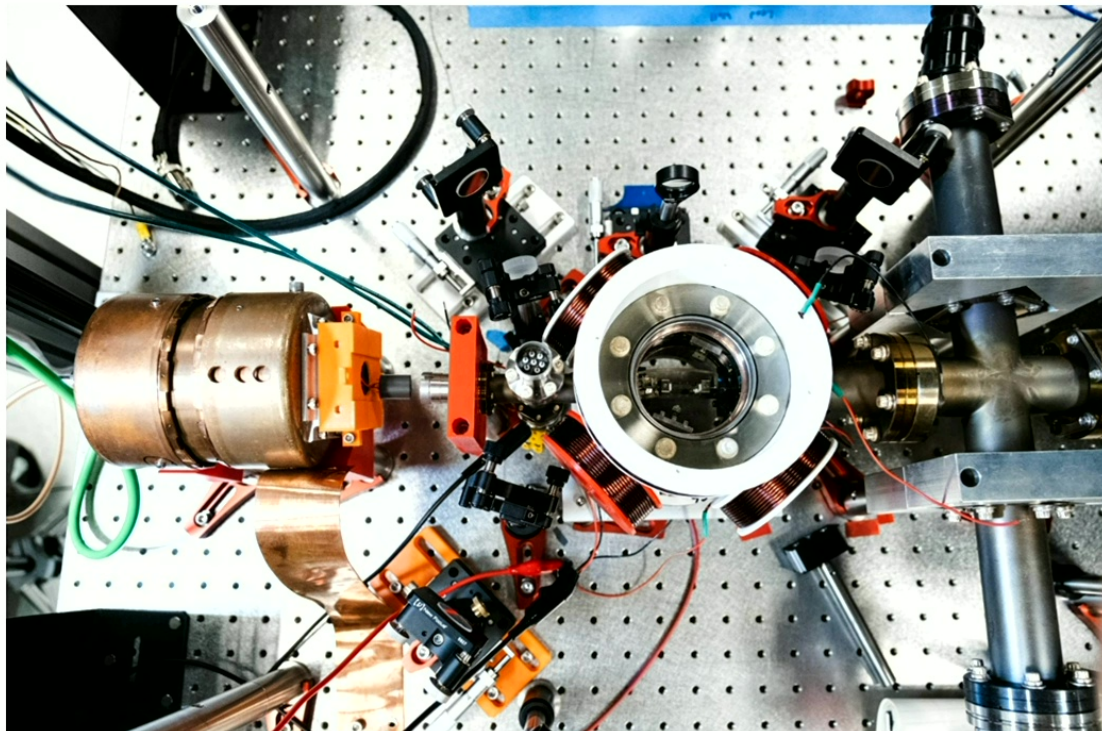
<sup>1</sup>Ruster, T. et al. Appl. Phys. B 122, 254, DOI: 10.1007/s00340-016-6527-4 (2016).

<sup>2</sup>Schirmer, S. G., Greentree, A. D., Ramakrishna, V. & Rabitz, H. J. Phys. A: Math. Gen. 35, 8315, DOI: 10.1088/0305-4470/35/39/313 (2002).

<sup>3</sup>Ivanov, P. A., Kyoseva, E. 332 S. & Vitanov, N. V. Phys. Rev. A 74, 022323, DOI: 10.1103/PhysRevA.74.022323 (2006).

<sup>4</sup>Ringbauer, M. et al. Nat. Phys. 18, 1053–1057, DOI: 10.1038/278 s41567-022-01658-0 (2022).

# Our Ion Trap: Aquamarine



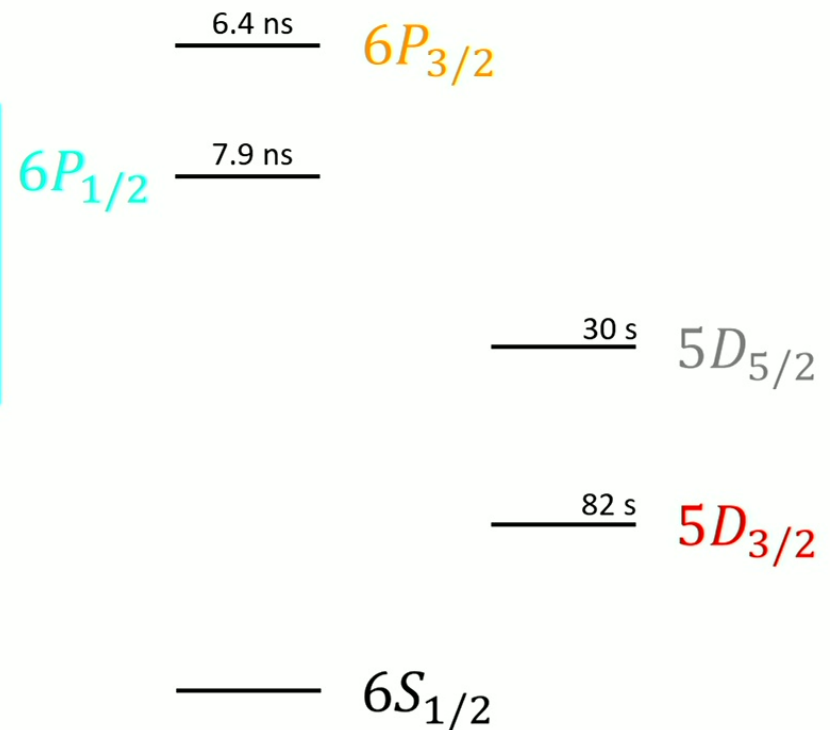


# Barium+ Ion's Energy Levels

- $6S_{1/2}$ : stable ground state.

## Fluorescence group

- $6P_{1/2}$ : drive with 493 nm laser for fluorescence.
- $5D_{3/2}$ : driven with 650 nm laser to repump to  $6P_{1/2}$ .



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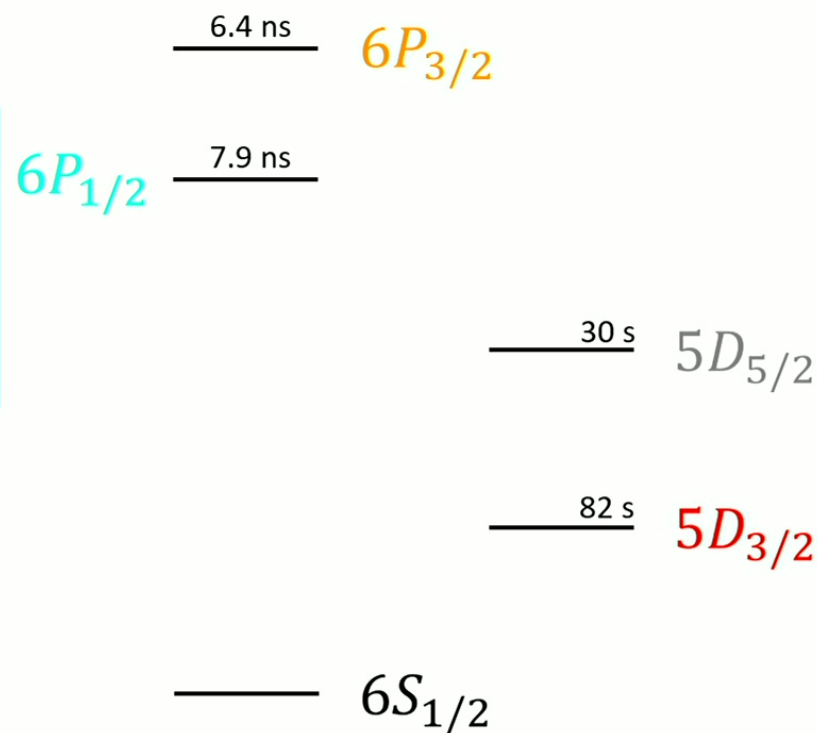
## Fluorescence group

- $6P_{1/2}$ : drive with 493 nm laser for fluorescence.
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## Metastable state control group

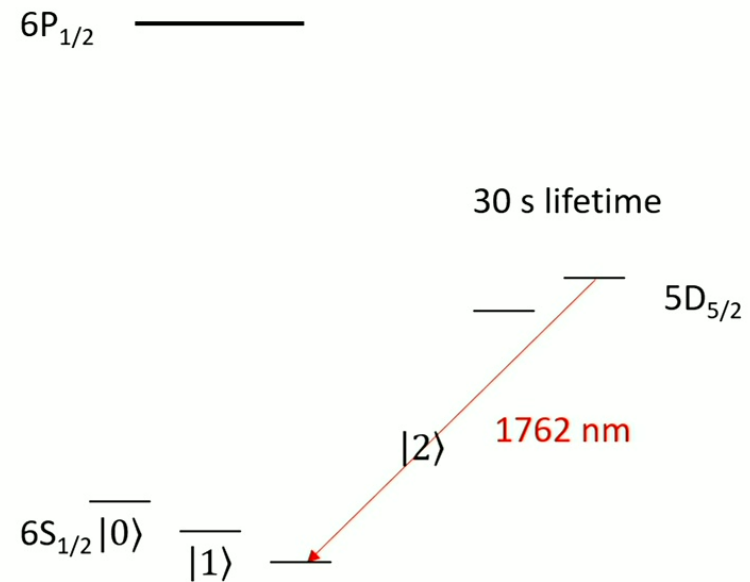
- $5D_{5/2}$ : metastable\* state for QC\*\*. Drive with 1762 nm laser.
- $6P_{3/2}$ : drive with 614 nm to reset  $5D_{5/2}$  experiments.

\*Long decay time    \*\*QC: quantum computing



# State Measurement

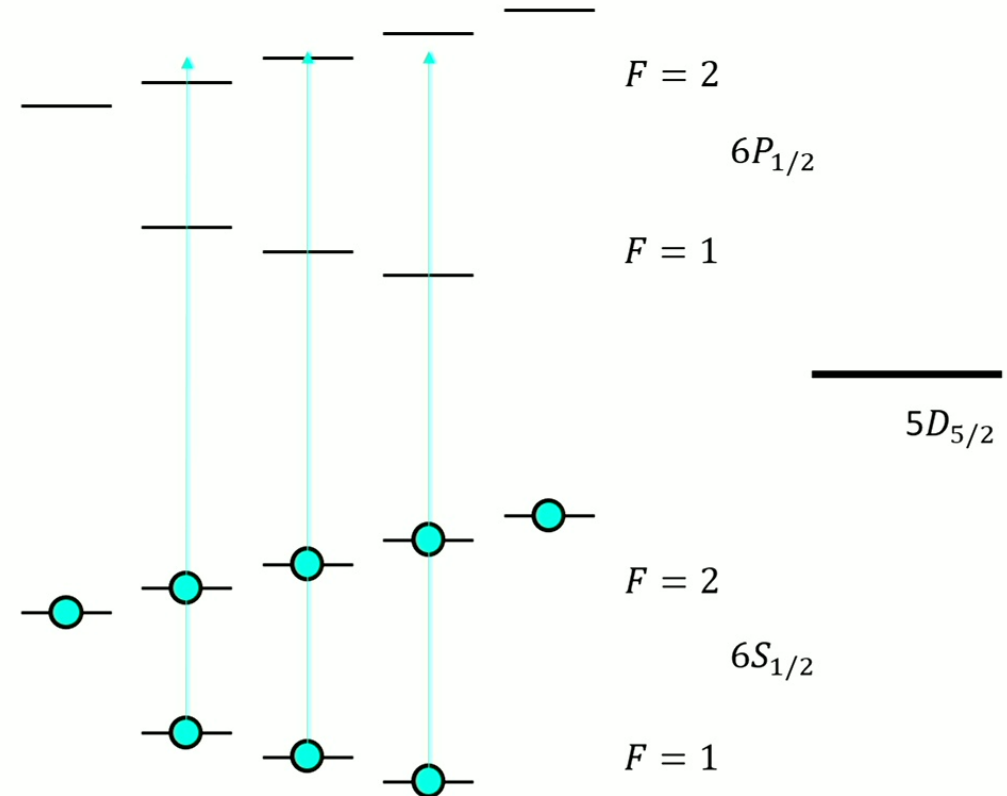
1. Shelve states in  $S_{1/2}$  to corresponding metastable states in  $D_{5/2}$  except for  $|0\rangle$ .
2. Fluoresce all remaining states in  $S_{1/2}$ .
3. Bring back state  $|1\rangle$  to  $S_{1/2}$  level.
4. Repeat steps 2 and 3 for different states.



# State Preparation

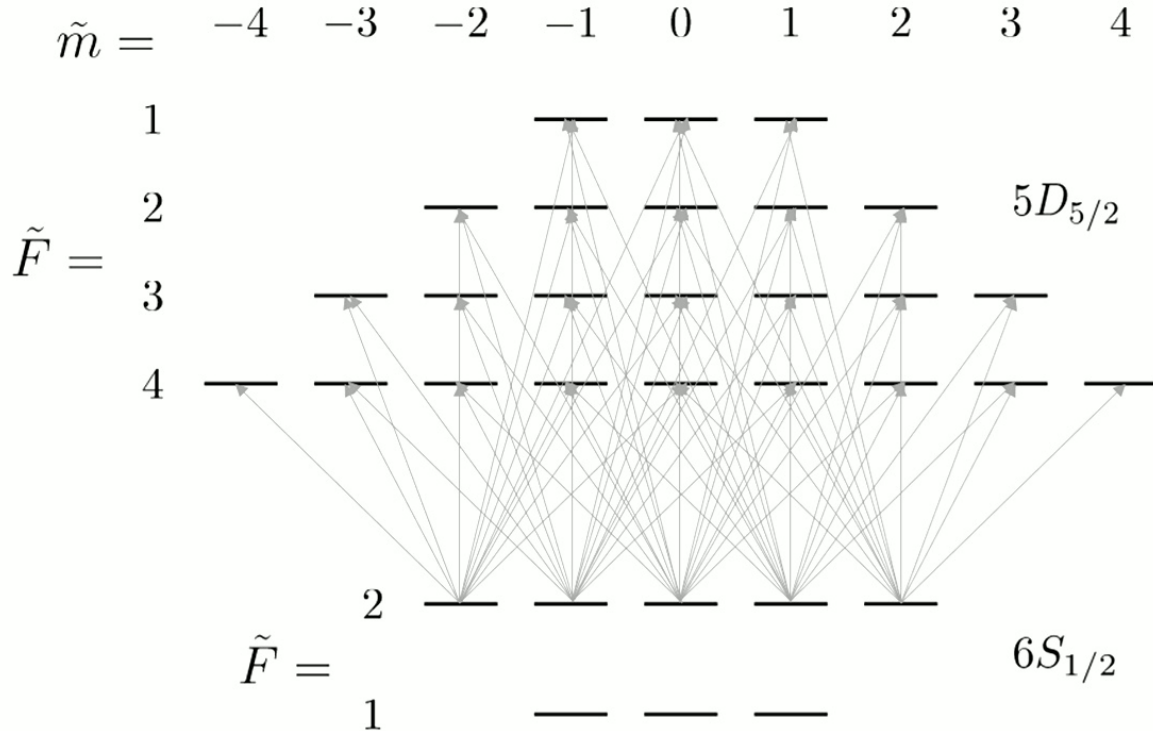
Initializing in  $6S_{1/2}$  state\*:

1. Turn on **493 nm** light with only frequency for pumping out  $6S_{1/2}$ ,  $F = 1$  state populations.
2. Turn on **1762 nm** light with frequencies to drive out population in  $6S_{1/2}$ ,  $F = 2$  states that we do not want to initialize. Each frequency turned on for a  $\pi$ -pulse time.
3. Turn on **614 nm** light to drive out state populations in  $5D_{5/2}$  and allow the state population to decay back to  $6S_{1/2}$ .
4. Repeat steps 1-3 until state population saturates at the desired state to be initialized.



\*An, F. A. et al. Phys. Rev. Lett. 129, 130501, DOI: [10.1103/PhysRevLett.129.130501](https://doi.org/10.1103/PhysRevLett.129.130501) (2022).

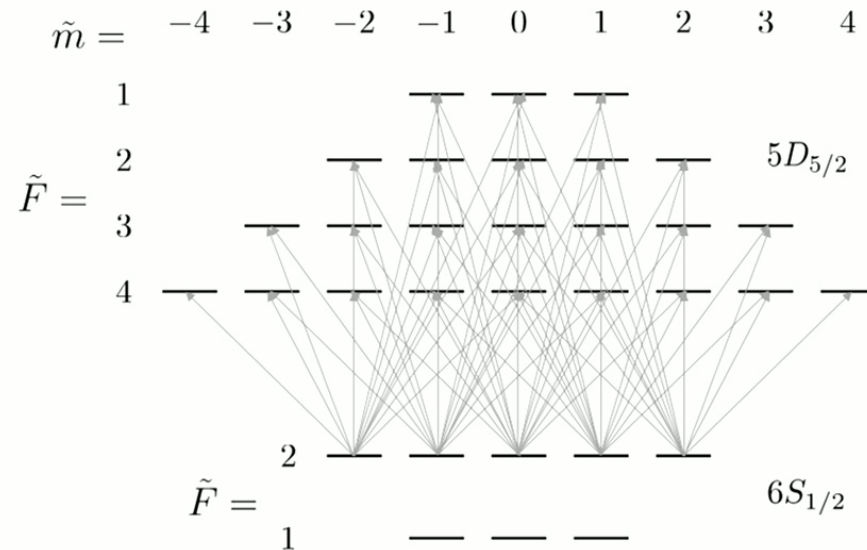
# High Dimensional Qudit Encoding with $^{137}\text{Ba}^+$



- $^{137}\text{Ba}^+$   $6S_{1/2}$  and  $5D_{5/2}$  have 32 non-degenerate states in non-zero magnetic field!
- At most 25-level qudit can be encoded.
- So many transitions! Which ones are useful?

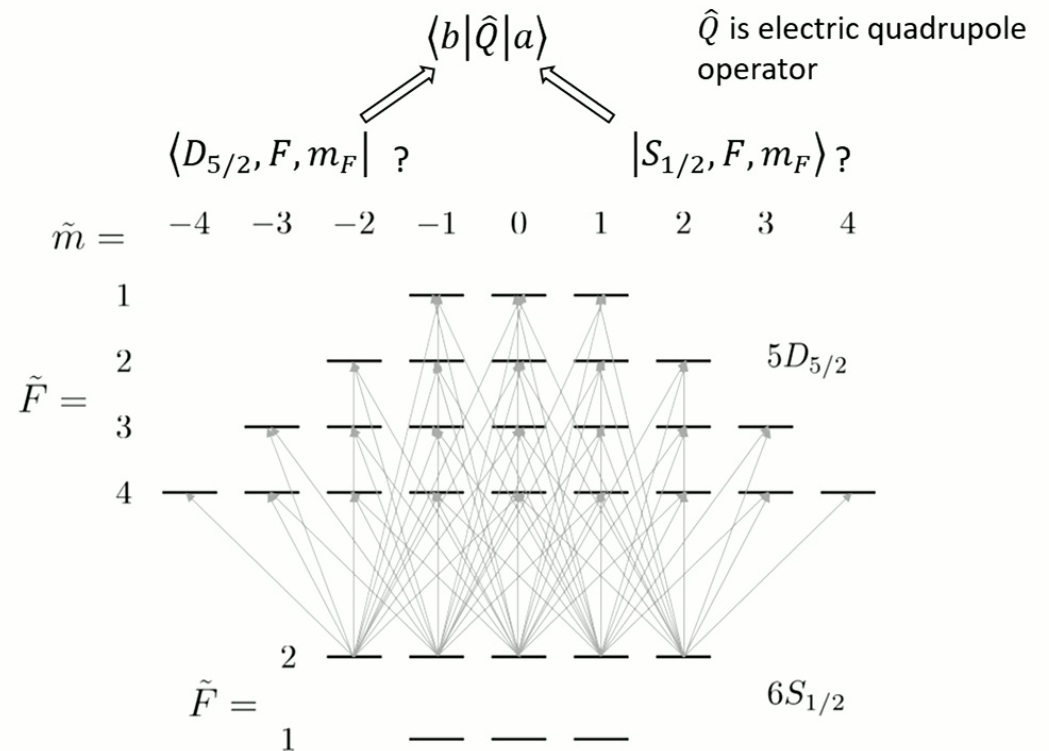
# Choosing Useful Coherent Transitions in $^{137}\text{Ba}^+$

- Calculate transition strengths, pick the ones strong enough for practical use.



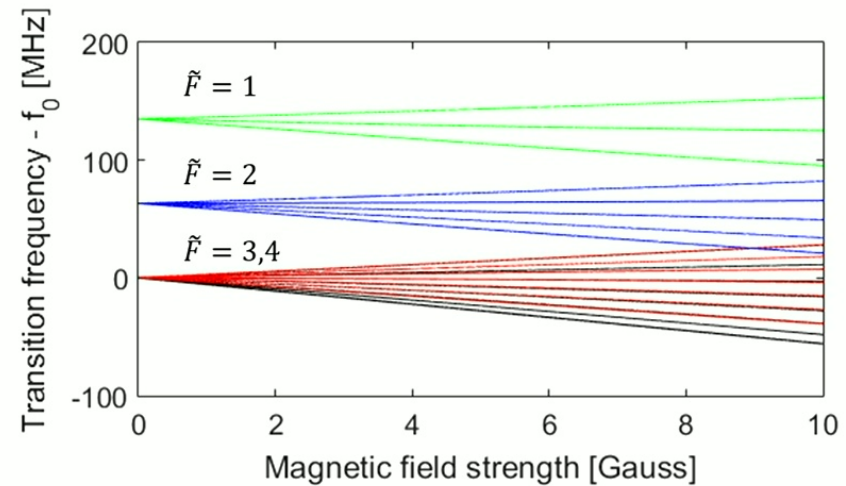
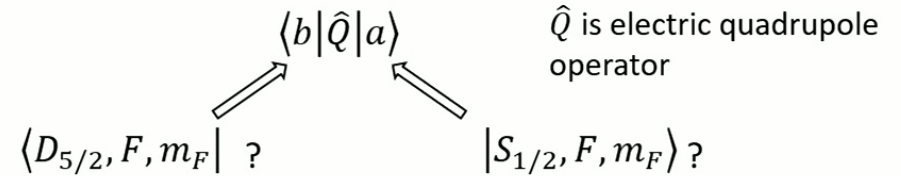
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- $5D_{5/2}$  level has close  $F=3$  and  $F=4$  hyperfine splitting.
- Small B assumption quickly falls apart.



$f_0$  is the energy level of  $\tilde{F} = 4$  at zero magnetic field.



# Choosing Useful Coherent Transitions in $^{137}\text{Ba}^+$

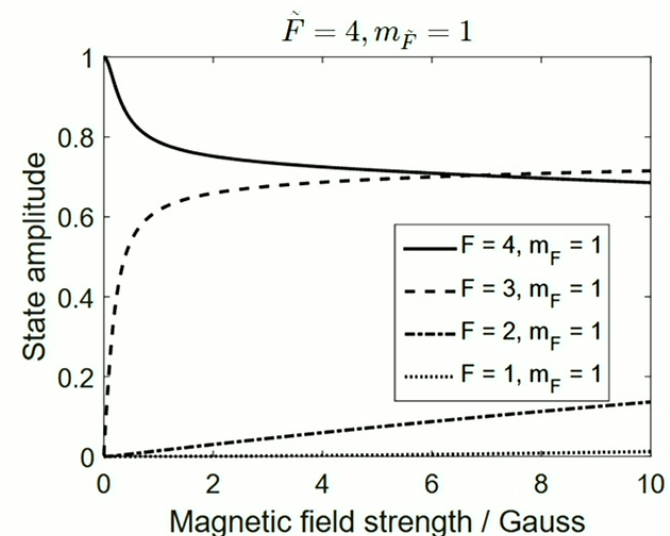
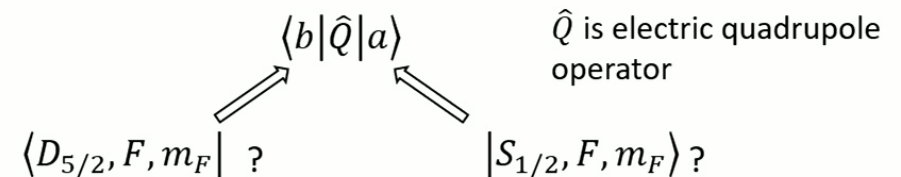
- Calculate transition strengths, pick the ones strong enough for practical use.
- $5D_{5/2}$  level has close  $F=3$  and  $F=4$  hyperfine splitting.
- Small B assumption quickly falls apart.
- Energy eigenstates are not the  $|F, m_F\rangle$  states.
- Obtain energy eigenstates from solving the Hamiltonian.

$$\hat{H}_R = \frac{2\pi A_D}{\hbar} \vec{I} \cdot \vec{J} + hB_Q \frac{\frac{3}{\hbar^4} (\vec{I} \cdot \vec{J})^2 + \frac{3}{2\hbar^2} \vec{I} \cdot \vec{J} - I(I+1)J(J+1)}{2I(2I-1)J(2J-1)} + \frac{B_e \mu_B}{\hbar} (g_J \hat{J}_z + g_I \hat{I}_z)$$

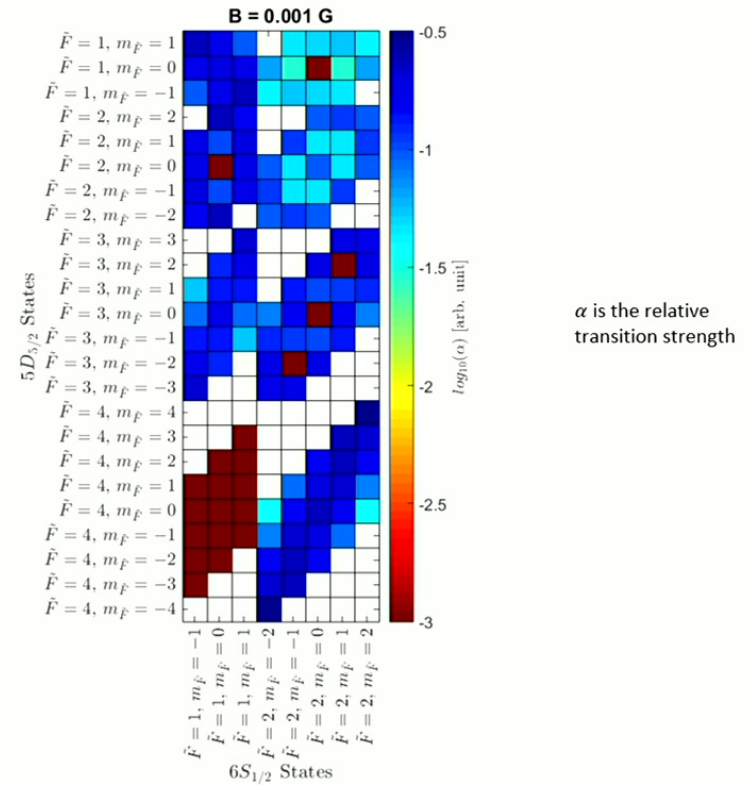
$A_D$  is the dipole hyperfine constant

$B_Q$  is the quadrupole hyperfine constant

$B_e$  is the magnetic field strength

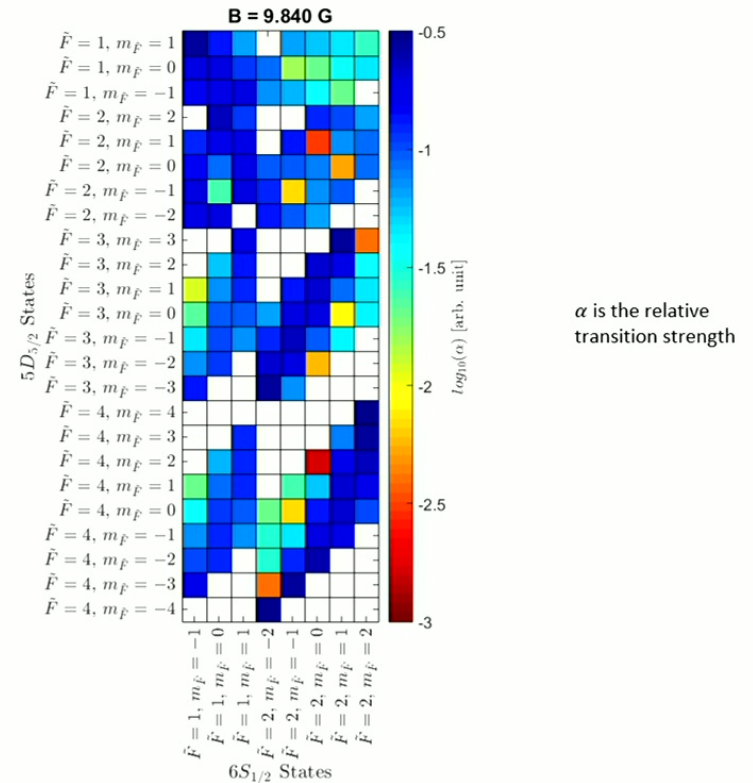


# Evolution of $6S_{1/2}$ to $5D_{5/2}$ Transition Strengths



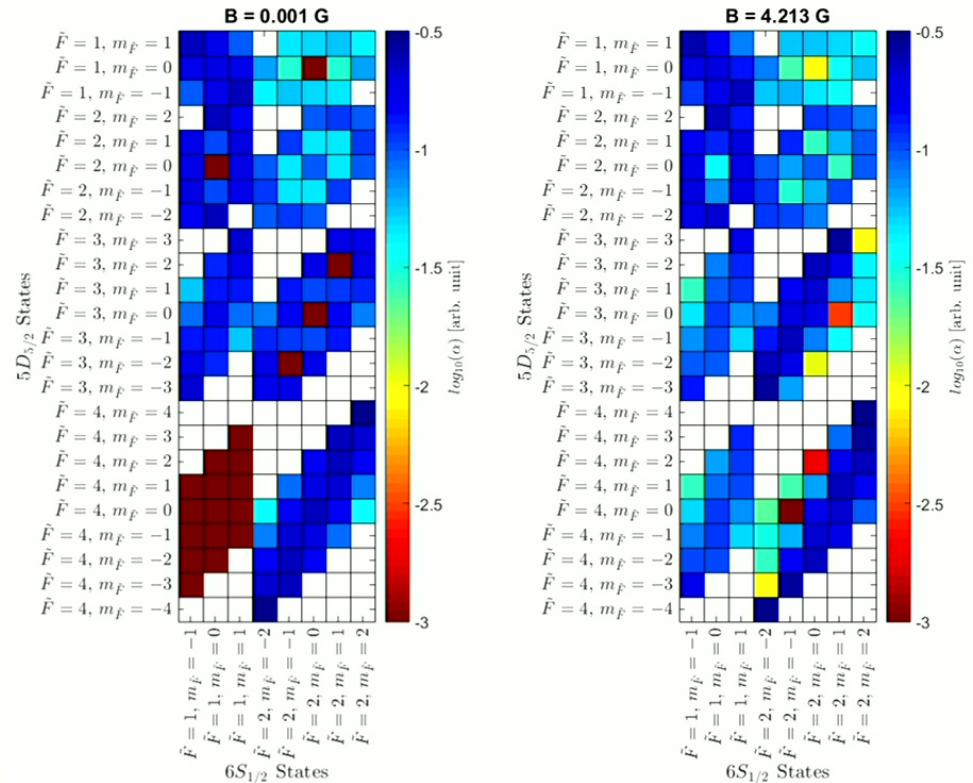
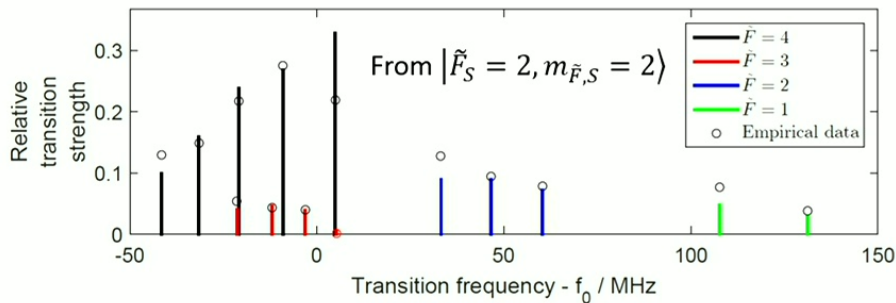
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- Increasing magnetic field quickly changes the transition strengths significantly.



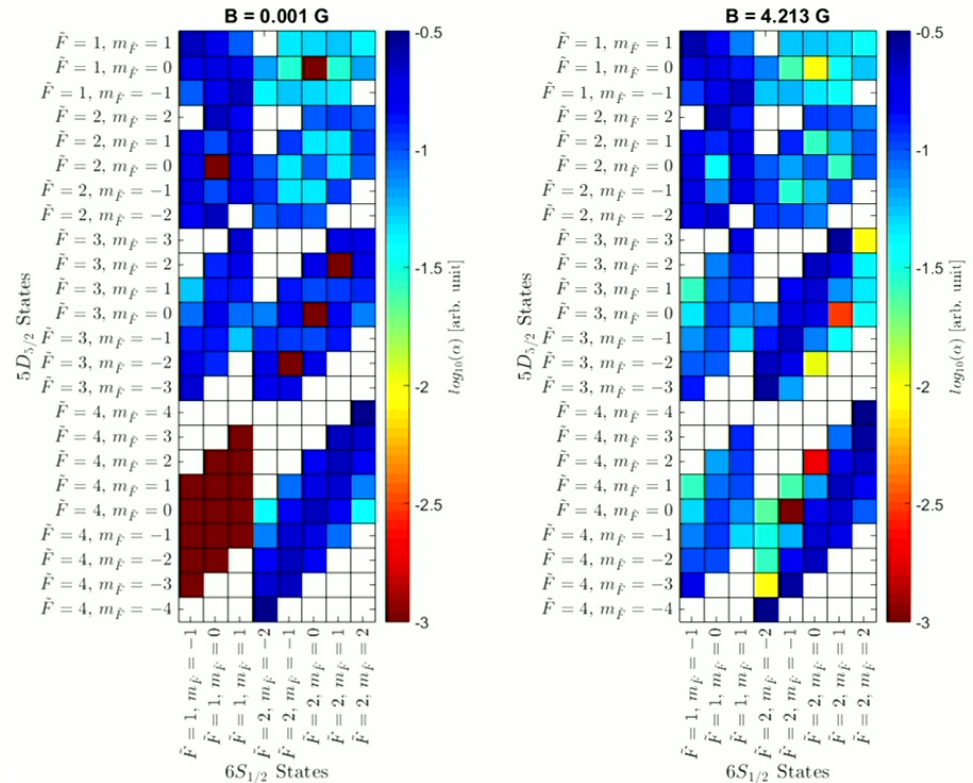
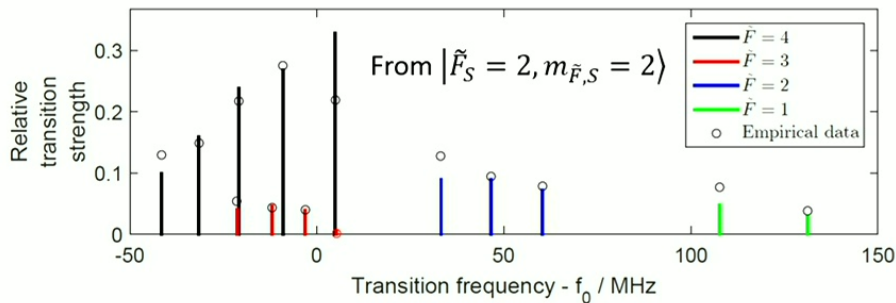
# Evolution of $6S_{1/2}$ to $5D_{5/2}$ Transition Strengths

- Increasing magnetic field quickly changes the transition strengths significantly.
- Allowed hyperfine transitions become practically forbidden and vice versa.
- Empirical data agree well with theoretical calculations.

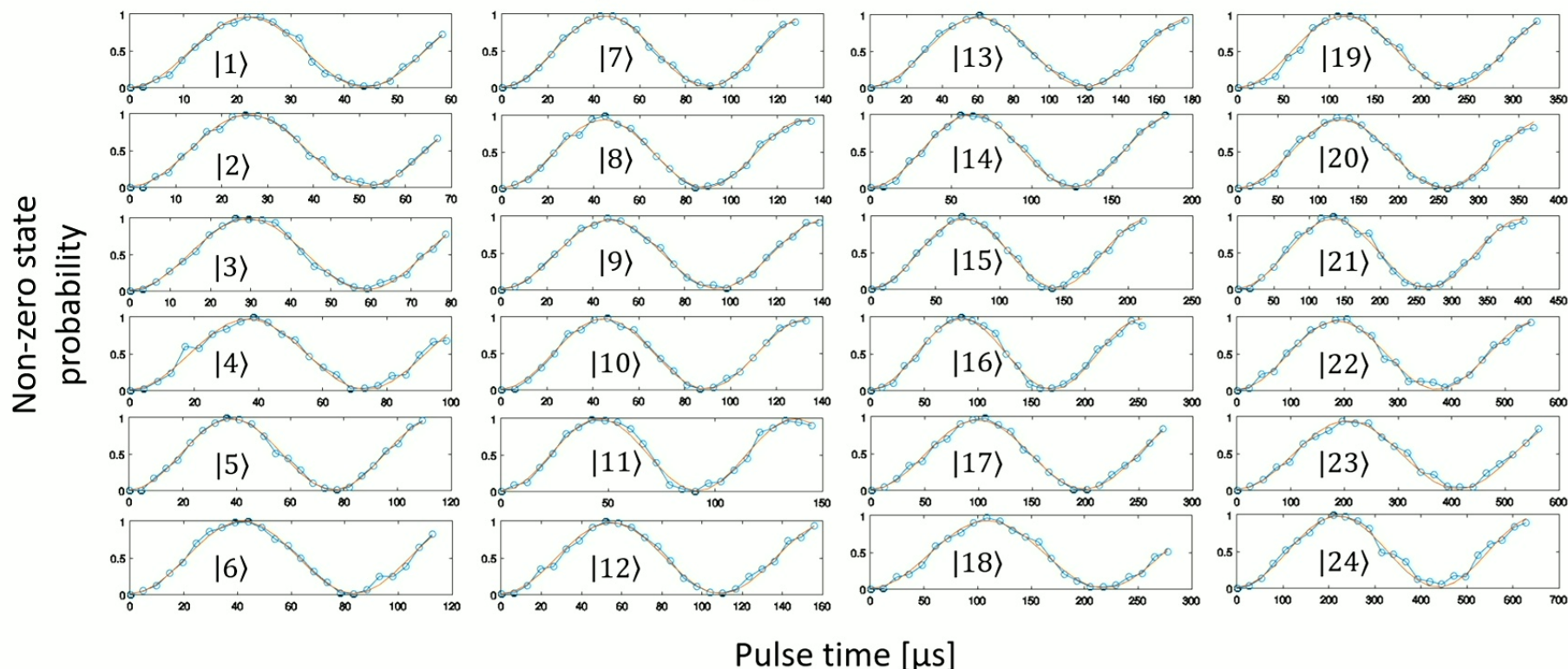


# Evolution of $6S_{1/2}$ to $5D_{5/2}$ Transition Strengths

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- Deviations verified to be from inaccurate assumption of pure linear polarization of 1762 nm laser. Data fits theory perfectly once laser geometric factor is calibrated for (shown later).



# Coherent Rabi Cycling of Qudit States



**Rabi cycling of state population of  $|0\rangle$  (a state in  $S_{1/2}$ ) and other non-zero states.**

Transitions not overcrowded for Rabi frequencies in the order 10 kHz at  $\sim 4$  G magnetic field.  
Sufficient control to perform universal single qudit gates.

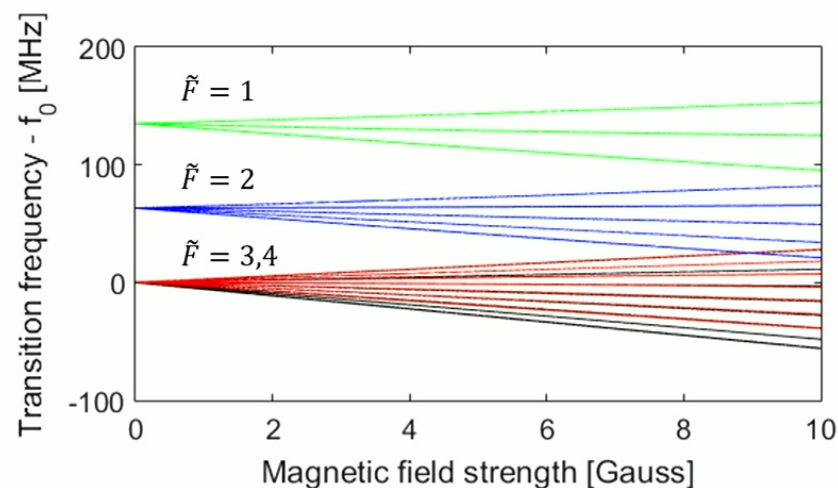
# Laser Parameters Calibrations for Qudits

- Qubits: 1 transition – 1 laser frequency, 1 pulse time calibrations.
- Qudits: many transitions – many calibrations?

# Laser Parameters Calibrations for Qudits

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- Qudits: many transitions – 2 frequencies, 5 pulse times calibrations.
- Frequency calibrations:
  - Measure 2 frequencies,  $f_{offset}$  and  $f_B$ , can compute the difference  $\Delta f = f_B - f_{offset}$  as a proxy for the magnetic field strength.
  - $f_{offset}$  also acts as a calibration measurement for drifts of frequency locking devices (e.g. Fabry–Pérot cavity).
  - Relying on theory alone still gives frequency deviations of order 1 kHz.
  - Do a polynomial fit with historical data, linear fit is sufficient for our magnetic field drifts.

$$f_n = a_{n,1}\Delta f + f_1 + a_{n,2}$$

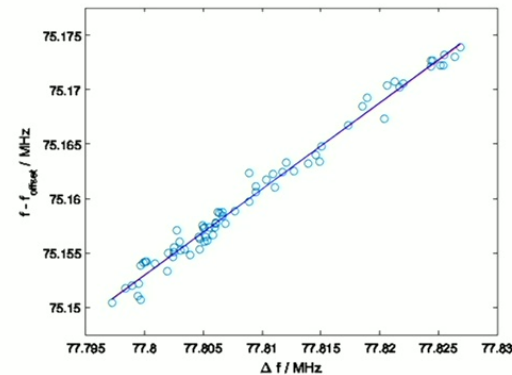
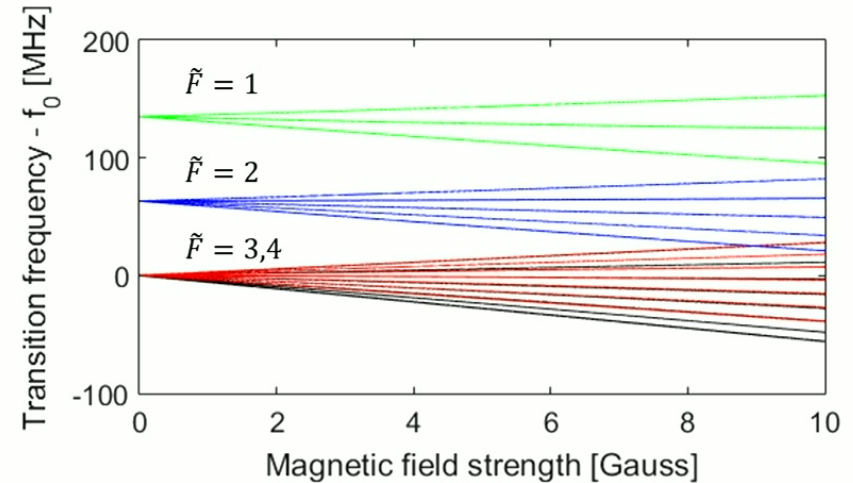




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Red line: Free parameter fit for both gradient,  $a_{n,1}$ , and offset,  $a_{n,2}$ .

Blue line: Free parameter fit for offset only, gradient from theoretical B field sensitivity.

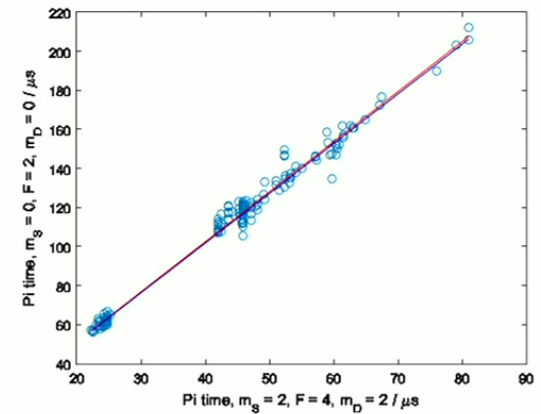
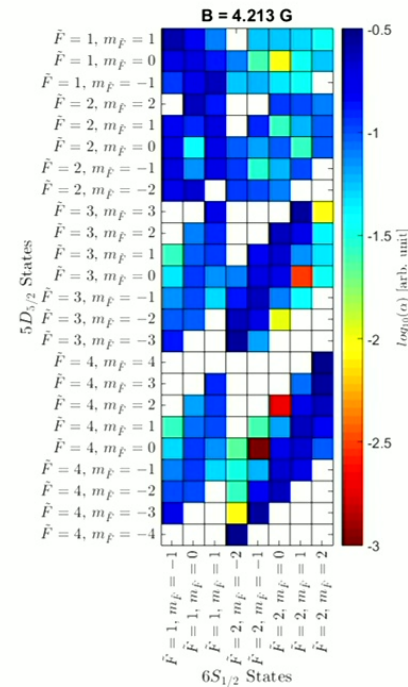
B field sensitivity from theory is accurate enough for practical use!

20

# Laser Parameters Calibrations for Qudits

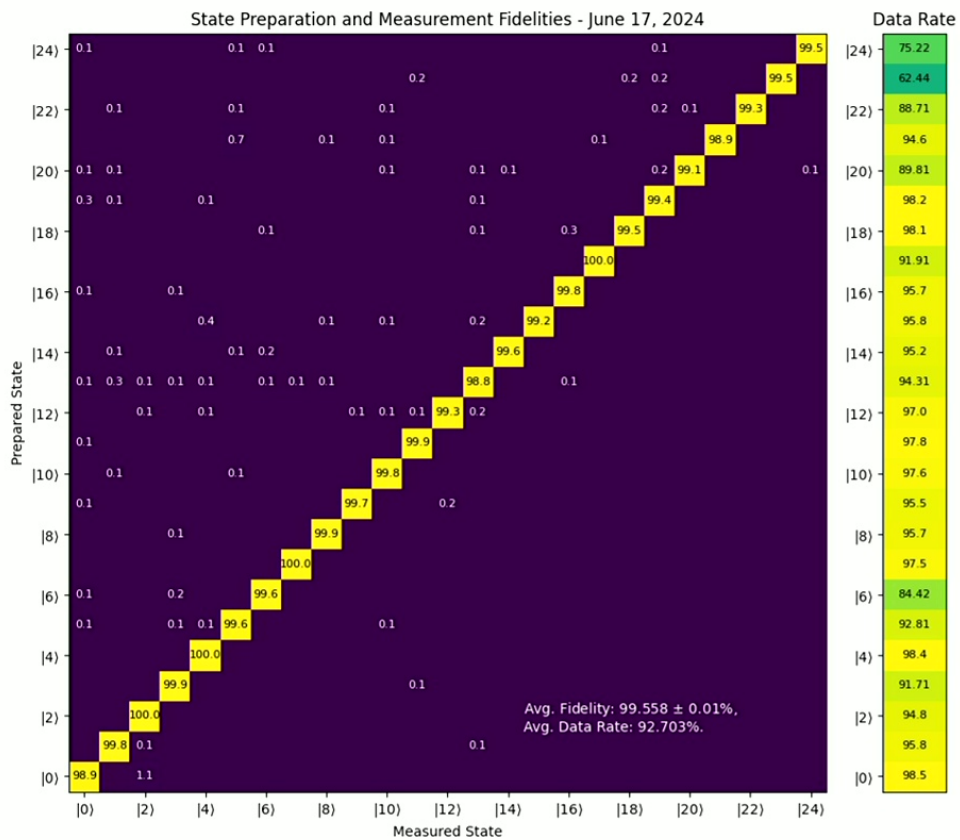
- Qubits: 1 transition – 1 laser frequency, 1 pulse time calibrations.
- Qudits: many transitions – 2 frequencies, 5 pulse times calibrations.
- Pulse time calibrations:
  - Measure 5 reference  $\pi$ -pulse times  $\tau_{\Delta m,0}$ , each corresponding to a  $-2 \leq \Delta m \leq 2$  transition. Having a reference pulse time for each  $\Delta m$  makes it robust to laser polarization drifts.
  - Scale the pulse times of other transitions with a scaling parameter  $b_n$ , either determined from historical data or calculated from theory
- Theory calculations are practically sufficient!

$$\tau_n = b_n \tau_{\Delta m,0}$$



Red line: Free parameter fit for both gradient,  $b_n$ .  
 Blue line: Gradient calculated from theory.

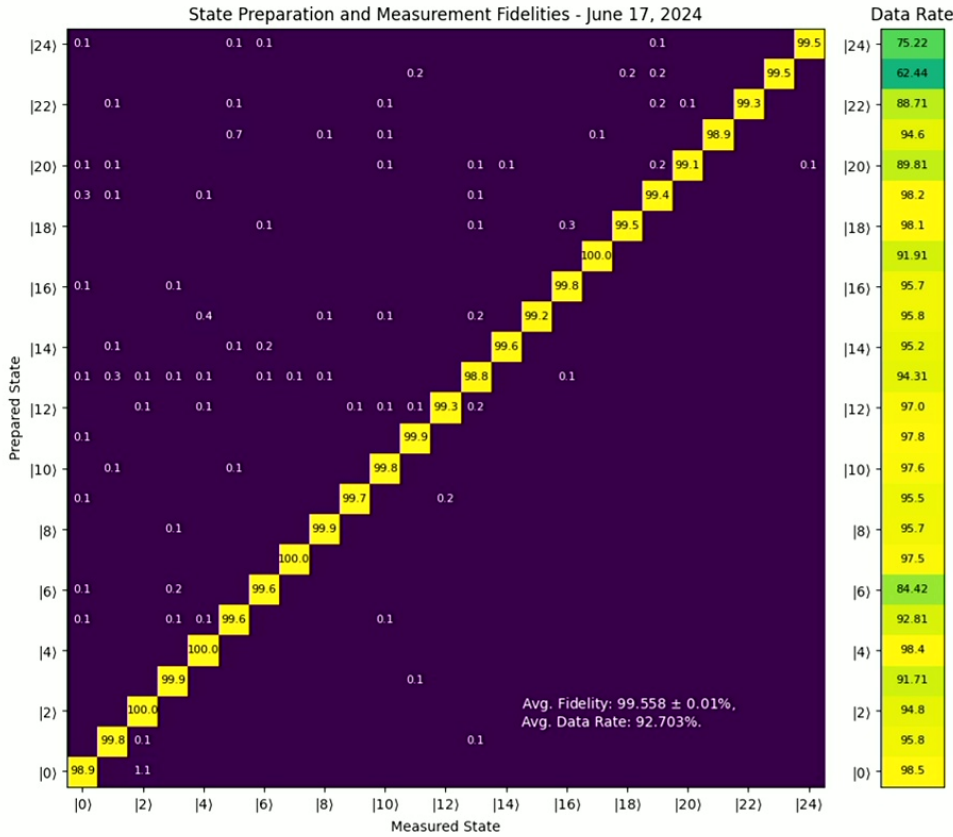
# 25-Level Qudit SPAM Results



$$m_{\tilde{F}} = \begin{matrix} -4 & -3 & -2 & -1 & 0 & 1 & 2 & 3 & 4 \\ 1 & & & \overline{|22\rangle} & \overline{|23\rangle} & \overline{|24\rangle} & & & \\ 2 & & & \overline{|17\rangle} & \overline{|16\rangle} & \overline{|19\rangle} & \overline{|18\rangle} & \overline{|20\rangle} & 5D_{5/2} \\ 3 & & \overline{|5\rangle} & \overline{|11\rangle} & \overline{|10\rangle} & \overline{|13\rangle} & \overline{|7\rangle} & \overline{|14\rangle} & \overline{|12\rangle} \\ 4 & \overline{|1\rangle} & \overline{|2\rangle} & \overline{|3\rangle} & \overline{|6\rangle} & \overline{|15\rangle} & \overline{|21\rangle} & \overline{|9\rangle} & \overline{|4\rangle} & \overline{|8\rangle} \end{matrix}$$

$$\tilde{F} = \begin{matrix} 2 & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} \\ 1 & & \text{---} & \text{---} & \text{---} & \text{---} \end{matrix} \quad 6S_{1/2}$$

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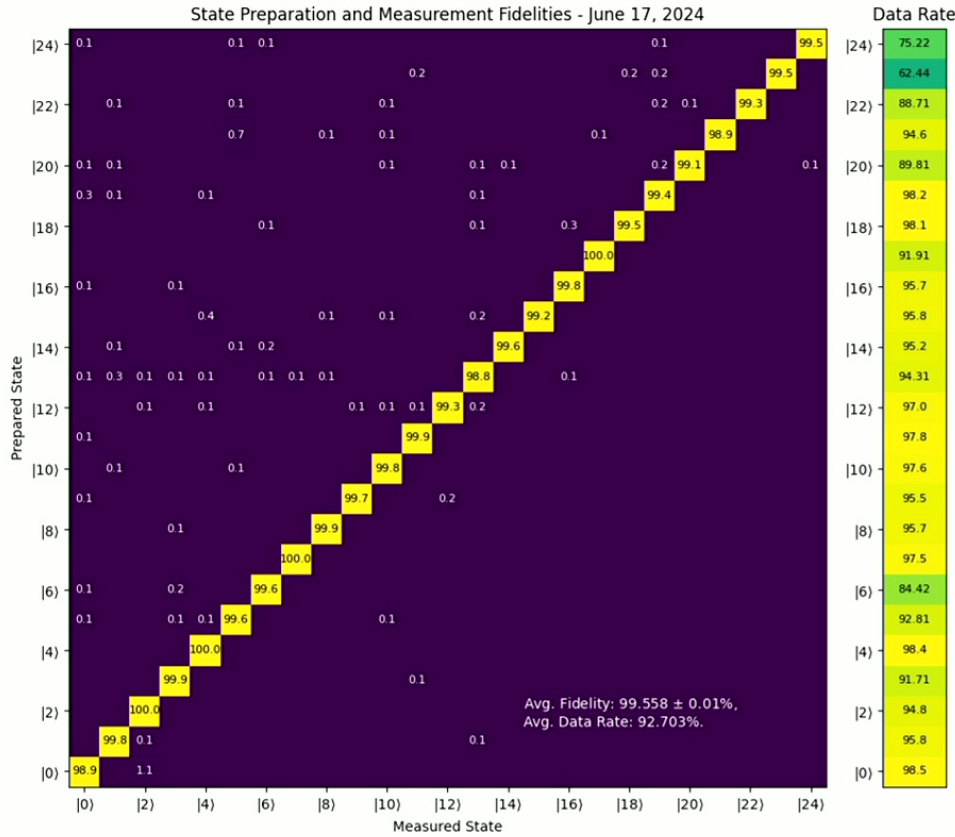


- Heralded state preparation (thanks to Oxford's group\*) – repeat experiment if first detect bright measurement for  $|n \neq 0\rangle$ .
- Average 25-level SPAM fidelity of  $99.56 \pm 0.1\%$ . For reference, 4 qubits (16 states) at 99.9% error  $\sim 99.6\%$  error.

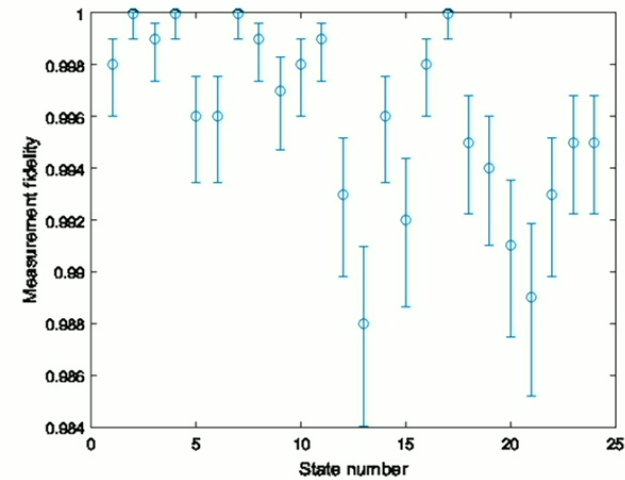
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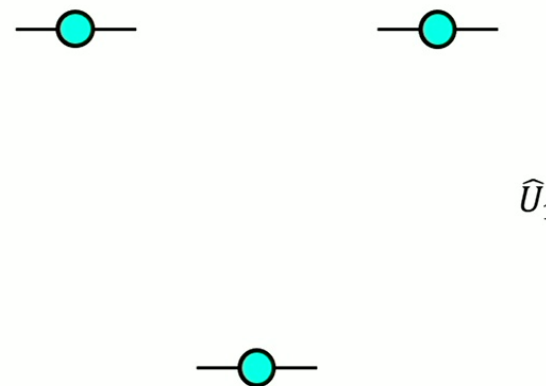


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- Average 25-level SPAM fidelity of  $99.56 \pm 0.1\%$ . For reference, 4 qubits (16 states) at 99.9% error  $\sim$ 99.6% error.
- Errors mainly from spontaneous  $5D_{5/2}$  decay. Predicted average error of 0.2%.



# Qudit Coherence

- Demonstrate coherence among all the encoded states with a generalized qudit Ramsey experiment.
- Prepare a state with equal superposition with  $\hat{U}_1$ .
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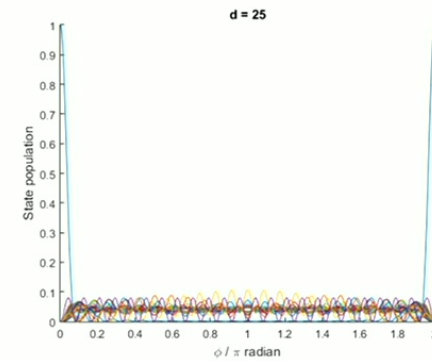
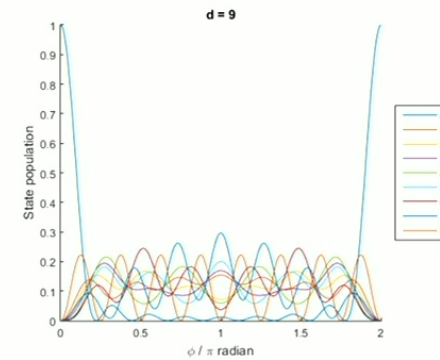
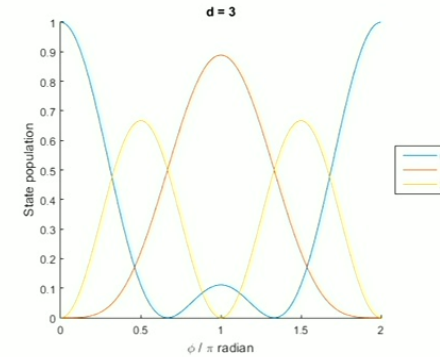


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$$|\langle 0|\psi\rangle|^2 = \frac{1}{d} + \frac{2}{d^2} \sum_{m=1}^{d-1} (d-m) \cos(m\phi)$$

- $|\langle 0|\psi\rangle|^2 = 1$  when  $\phi$  is zero or integer multiples of  $2\pi$ .

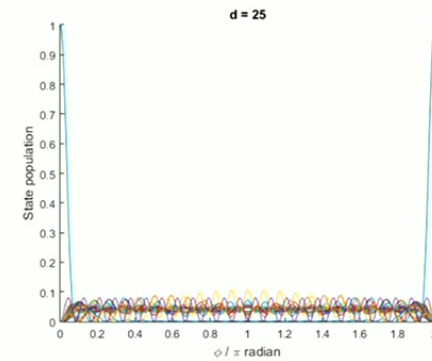
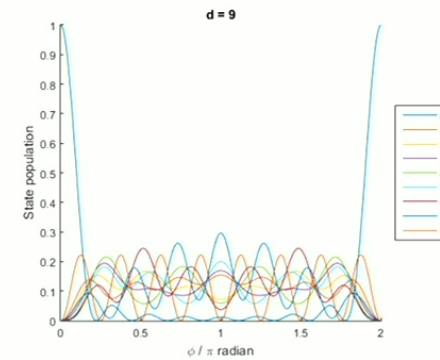
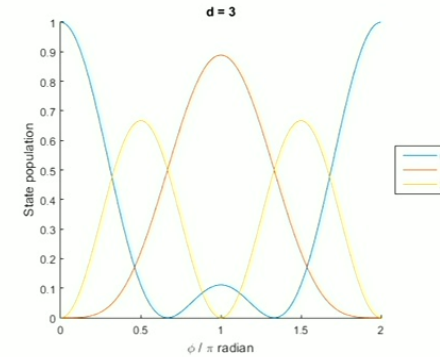


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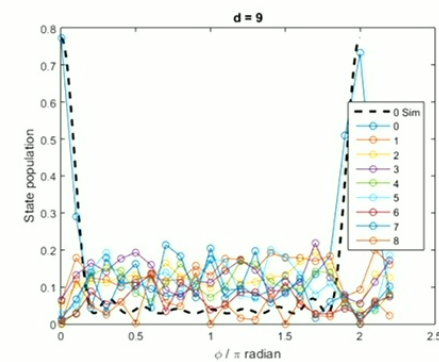
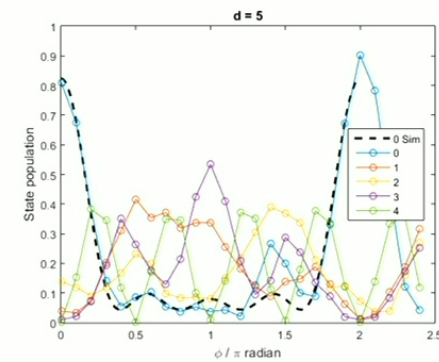
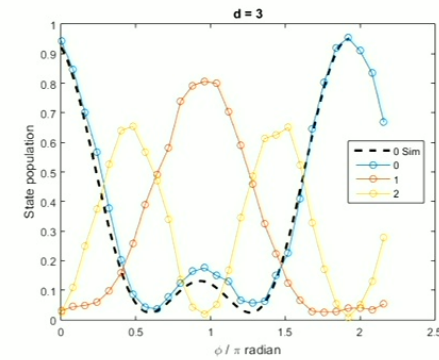
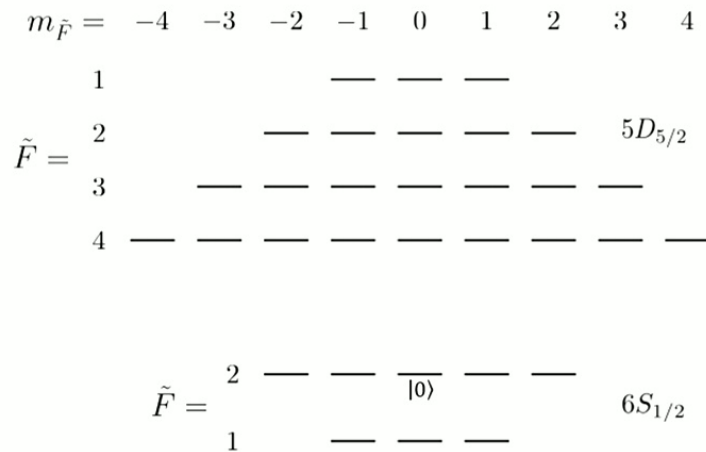
- $|\langle 0|\psi\rangle|^2 = 1$  when  $\phi$  is zero or integer multiples of  $2\pi$ .
- $|\langle 0|\psi\rangle|^2 = 0$  when  $\phi$  is integer multiples of  $\frac{2\pi}{d}$ .
- Cosine terms average out to zero if  $\phi$  is not coherently controlled in experiments,  $|\langle 0|\psi\rangle|^2$  converge to  $\frac{1}{d}$ .





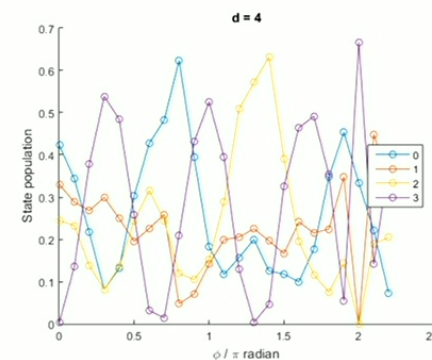
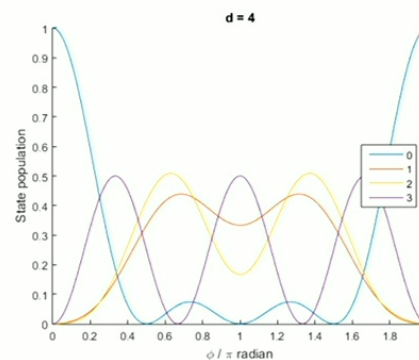
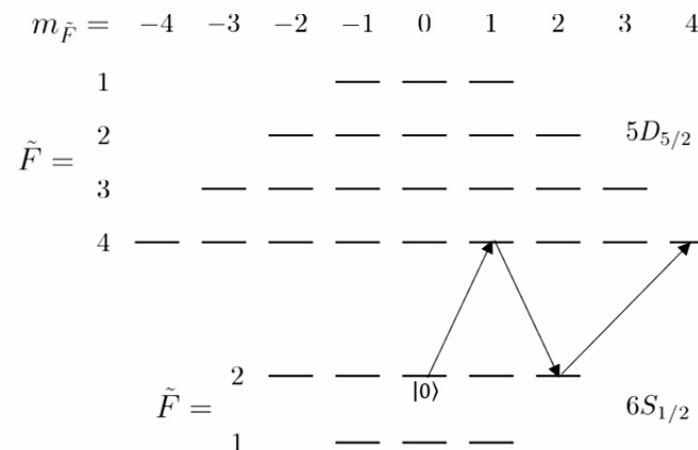
# Qudit Coherence

- $|0\rangle$  encoded in the  $F = 2, m = 0$  ground state.
- States reachable from  $|0\rangle$  with only one 1762 nm transition allows encoding up to 9 states with our calibrated data.
- Successful demonstration of qudit coherence up to  $d = 9$ .



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# Qudit Coherence

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- Tried 4-level qudit Ramsey phase scan with bussing, unexpected results obtained.
- Allowing frequency detunings in the simulations explains the discrepancy.

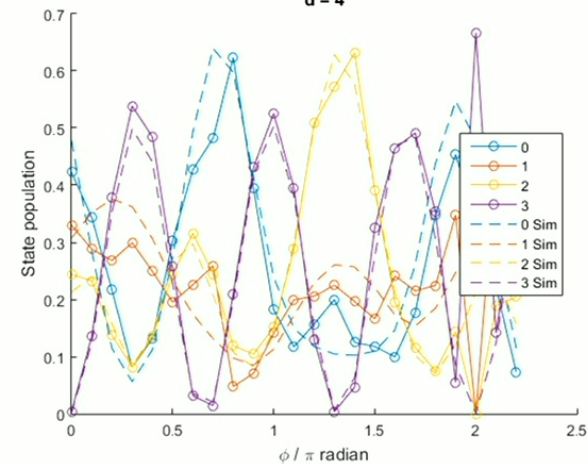
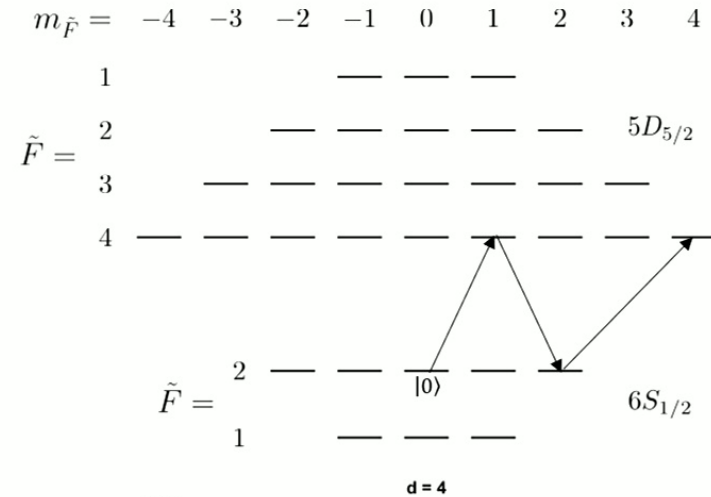
$$| \langle 0 | \psi \rangle |^2 = \frac{1}{d} + \frac{2}{d^2} \sum_{m=1}^{d-1} \cos(m\phi + \Delta_m t) + \frac{2}{d^2} \sum_{m=1}^{d-1} \sum_{n=1}^{d-1} \cos((m-n)\phi + (\Delta_m - \Delta_n)t)$$

- Two states with distance  $D$  ( $D$  2-level transitions) compounds  $D$  number of frequency detunings:

$$\Delta_{net} = \Delta_1 + \Delta_2 + \dots + \Delta_D$$

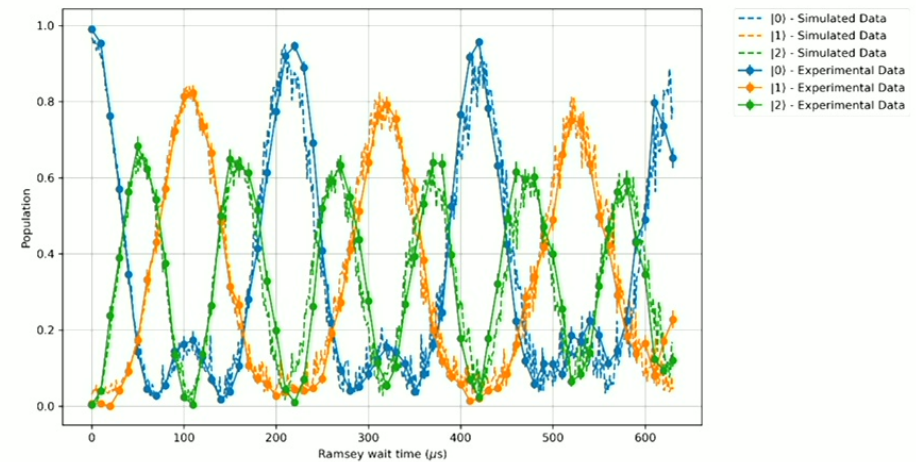
$$\text{std}(\Delta_{net}) = \sqrt{\text{Var}(\Delta_1) + \text{Var}(\Delta_2) + \dots + \text{Var}(\Delta_D)}$$

- Qudit encoding with distances  $D$  needs calibration parameters that are overall  $\sqrt{D}$  more accurate.



# Qudit Coherence

- Can add wait time between  $\hat{U}_1$  and  $\hat{U}_2$  to study decoherence time, vary  $\phi$  with time.
- Investigation ongoing.



# Summary

- **Main question:** Qudits advantageous in theory, realistically advantageous? Practical challenges vs theoretical gain.
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  - Unique to  $^{137}\text{Ba}^+$  or  $^{137}\text{Ba}^+$ -like ions – non-trivially different energy level structures.
  - Calibration time for parameters do not scale with qudit dimension.
  - Transition frequencies are not too crowded for order 10 kHz Rabi frequencies.
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