Title: Quantum Permutations as Quantum Coordinate Transformations

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Series: Quantum Foundations, Quantum Information

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Quantum Permutations as Quantum Coordinate Transformations

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Motivation: uniting symmetries

Quantum Mechanics

General Relativity

Unitary Symmetry: invariance to **basis change**

Coordinate Symmetry: invariance to **naming of spacetime points**

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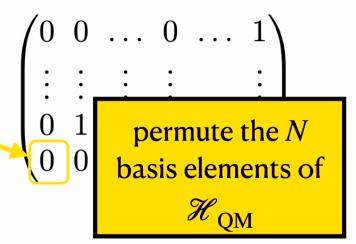
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Quantum Permutations: intuition

Classical Permutations:

 \mathbf{Q} does j become i?

A yes/no =
$$o/1$$



"quantum permute" the N subspaces of

$$\mathcal{H}_{\mathrm{QM}} \otimes \tilde{\mathcal{H}}$$

 $a_{11} \dots a_{1N}$

: :

 $u_{N1} \ldots u_{NN}$

Quantum Permutations:

Q where does j become i?

A a subspace (of $\tilde{\mathcal{H}}$)

Not a group, not closed under composition

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$$u = \begin{pmatrix} |\uparrow\rangle\langle\uparrow| & 0 & |\downarrow\rangle\langle\downarrow| \\ |\downarrow\rangle\langle\downarrow| & |\uparrow\rangle\langle\uparrow| & 0 \\ 0 & |\downarrow\rangle\langle\downarrow| & |\uparrow\rangle\langle\uparrow| \end{pmatrix}$$

An example for a quantity $u = \begin{pmatrix} |\uparrow\rangle\langle\uparrow| & 0 & |\downarrow\rangle\langle\downarrow\downarrow| \\ |\downarrow\rangle\langle\downarrow\downarrow| & |\uparrow\rangle\langle\uparrow| & 0 \\ 0 & |\downarrow\rangle\langle\downarrow\downarrow| & |\uparrow\rangle\langle\uparrow| \end{pmatrix}$ where $u = \begin{pmatrix} |\uparrow\rangle\langle\uparrow| & 0 & |\downarrow\rangle\langle\downarrow\downarrow| \\ |\downarrow\rangle\langle\downarrow\downarrow| & |\uparrow\rangle\langle\uparrow| & 0 \\ 0 & |\downarrow\rangle\langle\downarrow\downarrow| & |\uparrow\rangle\langle\uparrow| \end{pmatrix}$

$$\mathcal{H} = \bigoplus_{x} |x\rangle \otimes \tilde{\mathcal{H}} = \bigoplus_{y} |x\rangle \otimes \tilde{\mathcal{H}} = \bigoplus_{y} |x\rangle \otimes \tilde{\mathcal{H}} = \bigoplus_{x} |x\rangle \otimes \bigoplus_{y} |x\rangle \otimes \bigoplus_{y} |x\rangle \otimes \bigoplus_{y} |x\rangle \otimes \bigoplus_{y} |x\rangle \otimes \bigoplus_{x} |x\rangle \otimes \bigoplus_{x} |x\rangle \otimes \bigoplus_{x} |x\rangle \otimes \bigoplus_{y} |x\rangle \otimes \bigoplus_{x} |x\rangle \otimes \bigoplus_{$$

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Quantum-controlled permutations

The mathematicians call a quantum permutation u 'classical' if all entries u_{xy} **commute**.

This definition is equivalent to a quantum-controlled permutation:

$$u = \sum_{g \in S_N} \sigma_g \otimes \pi_g$$
 c*lassical* permutations complementary orthogonal projections on $\tilde{\mathcal{H}}$

→ The rest of the quantum permutations are richer, we will call them "beyond quantum-controlled" (BQC)

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Example for a BQC permutation

$$u = \begin{pmatrix} |\uparrow\rangle\langle\uparrow| & |\downarrow\rangle\langle\downarrow| & 0 & 0\\ |\downarrow\rangle\langle\downarrow| & |\uparrow\rangle\langle\uparrow| & 0 & 0\\ 0 & 0 & |-\rangle\langle-| & |+\rangle\langle+|\\ 0 & 0 & |+\rangle\langle+| & |-\rangle\langle-| \end{pmatrix}$$

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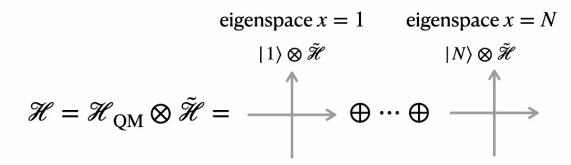
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| | Usual QM, with ${\mathscr H}_{\mathrm{QM}}$ | Here, with $\mathscr{H}_{\mathrm{QM}} \otimes 	ilde{\mathscr{H}}$ |
|--|---|--|
| coordinate system | choice of 1-d eigenspaces of position | choice of multi-d eigenspaces of position |
| (discrete) coordinate transformations | classical permutations | quantum permutations = (discrete) quantum coordinate transformations |

$$\mathcal{H} = \bigoplus_{x} |x\rangle \otimes \tilde{\mathcal{H}} = \bigoplus_{y} |x\rangle \otimes \tilde{\mathcal{H}} = \bigoplus_{y} |x\rangle \otimes \bigoplus_{y} u_{xy} \tilde{\mathcal{H}} = \bigoplus_{y} |x\rangle \otimes \bigoplus_{y} |x\rangle$$

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A particle on a line, but with degeneracy in the position operator *X*.



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The action of a magic unitary on a state:

$$u | \psi \rangle = \bigoplus_{x} \psi(x) | x \rangle \otimes \sum_{y} u_{xy} | \phi_{y} \rangle = \begin{pmatrix} u_{11} & \dots & u_{1N} \\ \vdots & & \vdots \\ u_{N1} & \dots & u_{NN} \end{pmatrix} \begin{pmatrix} | 1 \rangle | \phi_{1} \rangle \\ \vdots \\ | N \rangle | \phi_{N} \rangle \end{pmatrix}$$

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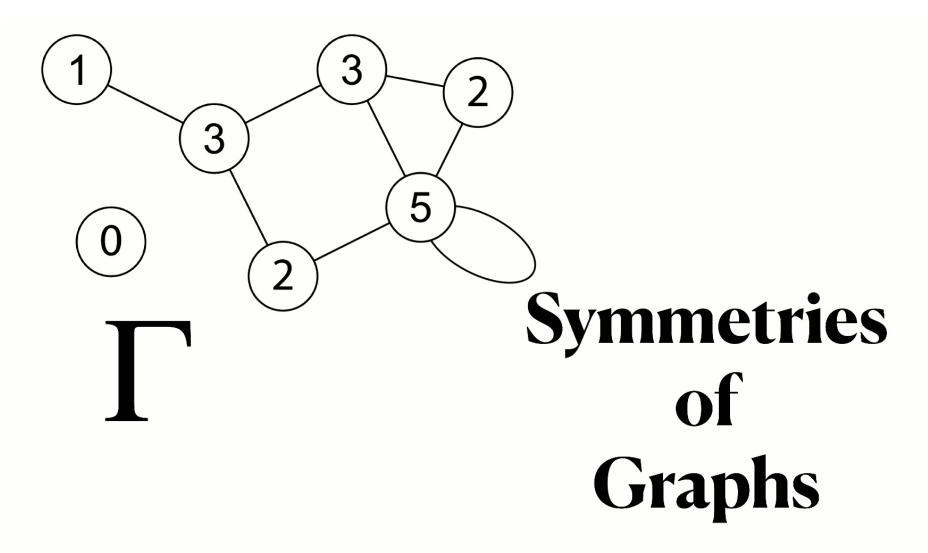
This way, one can "transfer" properties from the state into the descriptive redundancy.

$$|\psi\rangle_{\mathrm{QM}} = \sum_{x} \psi(x) |x\rangle$$
Embed in enlarged space
$$|\psi\rangle = \sum_{x} \psi(x) |x\rangle |e_{x}\rangle$$
ONB of $\tilde{\mathcal{H}}$

$$\exists u, u | \psi \rangle = |x = 1\rangle | \phi \rangle$$
 For some $|\phi \rangle \in \tilde{\mathcal{H}}$

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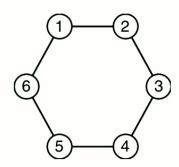
Symmetries of graphs

The classical symmetries of a graph are permutations which preserve the edge structure:

$$\sigma^{\dagger} A_{\Gamma} \sigma = A_{\Gamma}$$
 adjacency matrix

Called automorphisms of the graph, Aut Γ .

e.g. if the graph is a cycle — symmetries include rotations.



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Symmetries of graphs

If several permutations σ_i exist, any quantum controlled unitary

$$u = \sum_{i} \sigma_{i} \otimes \pi_{i}$$

is a quantum-controlled permutation which is also symmetry of the graph.

But this is a trivial extension.

For some graphs there exist symmetries *u* which are **beyond quantum-controlled.**

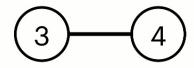
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Symmetries of graphs

E.g. the example we gave earlier is a BQC symmetry of the following graph:

$$u = \begin{pmatrix} |\uparrow\rangle\langle\uparrow| & |\downarrow\rangle\langle\downarrow| & 0 & 0\\ |\downarrow\rangle\langle\downarrow| & |\uparrow\rangle\langle\uparrow| & 0 & 0\\ 0 & 0 & |+\rangle\langle+| & |-\rangle\langle-|\\ 0 & 0 & |-\rangle\langle-| & |+\rangle\langle+| \end{pmatrix}$$



$$u^{\dagger}A_{\Gamma}u = A_{\Gamma}$$

 $u \in QAut$

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Example: Ising model

Spins on a graph Γ

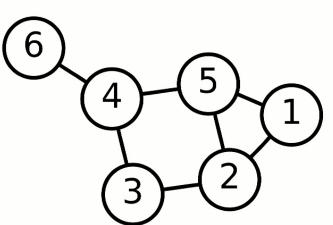
$$H = \sum_{i \sim_{\Gamma} j} S_i S_j$$

Symmetries: Aut Γ , classical automorphisms of the graph Γ

$$\sigma H \sigma^{\dagger} = H$$

Introducing an external space $\tilde{\mathcal{H}}$, the symmetries can be extended to QAut (*)

$$uHu^{\dagger} = H$$



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Application: A particle on a graph

Toy example: A d dimensional spin in one of N nodes of a graph Γ .

State:
$$|\psi\rangle\in\mathcal{H}_N\otimes\mathcal{H}_d,\ |\psi\rangle=\sum_{i\in\Gamma}|i\rangle\,|\psi_i\rangle$$
 Location

Hamiltonian: amount of "alignment" along edges

$$H = \sum_{i \sim_{\Gamma} j} |i\rangle \langle j| \otimes 1_d = A_{\Gamma} \otimes 1_d \quad , \quad \langle \psi | H | \psi \rangle = \frac{1}{2} \sum_{i \sim_{\Gamma} j} \operatorname{Re} \langle \psi_i | \psi_j \rangle$$

The symmetries of this physical system are both the classical **and** quantum automorphisms of Γ .

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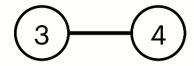
Another example in more detail: A particle on a graph

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Application: A particle on a graph

$$u = \begin{pmatrix} |\uparrow\rangle\langle\uparrow| & |\downarrow\rangle\langle\downarrow| & 0 & 0\\ |\downarrow\rangle\langle\downarrow| & |\uparrow\rangle\langle\uparrow| & 0 & 0\\ 0 & 0 & |+\rangle\langle+| & |-\rangle\langle-|\\ 0 & 0 & |-\rangle\langle-| & |+\rangle\langle+| \end{pmatrix}$$



$$u^{\dagger}Hu = H$$

Classical symmetry is non-trivially extended into the "beyond quantum-controlled" realm

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Can we use this in QRF transformations?

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Can we use this in QRF transformations?

$$S_{A \to C} = \text{SWAP}_{AC} \quad \sum_{g \in G} |g^{-1}\rangle\langle g|_C \quad u_{BC} \quad \text{on a tripartite system } \mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_C$$

$$u_{BC} = \sum_{g \in G} u_g \otimes |g\rangle\langle g| \quad \text{up to a swap and a reflection: a quantum-controlled magic unitary.}$$

What if we use a BQC *u*? E.g. is entanglement + coherence still conserved?

Cepollaro et al, arXiv:2406.19448

e.g. incorporating two distinct symmetry groups G_1 , G_2 to create the following BQC permutation:

$$u_{BC} = \sum_{g_1 \in G_1} u_{g_1} \otimes |g_1\rangle \langle g_1| + \sum_{g_2 \in G_2} v_{g_2} \otimes |g_2\rangle \langle g_2|$$

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Thanks for listening.

- Magic unitaries/quantum permutations are an interesting generalisation of the concept of permutations into the quantum realm.
- They can be interpreted as *changes into quantum coordinate systems*
- There are non-trivial, beyond quantum-controlled permutations which *extend the symmetry groups* of some Hamiltonians.
- Usual QRF transformations are only quantum-controlled.
- Extending them with BQC permutations might be insightful or useful.
- Might allow to incorporate more than one observable/symmetry group.



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