

**Title:** Advances in controlling superconducting qubits for quantum computing

**Speakers:** Stefan Filipp

**Collection/Series:** Waterloo-Munich Joint Workshop

**Subject:** Quantum Information

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**URL:** <https://pirsa.org/24090081>



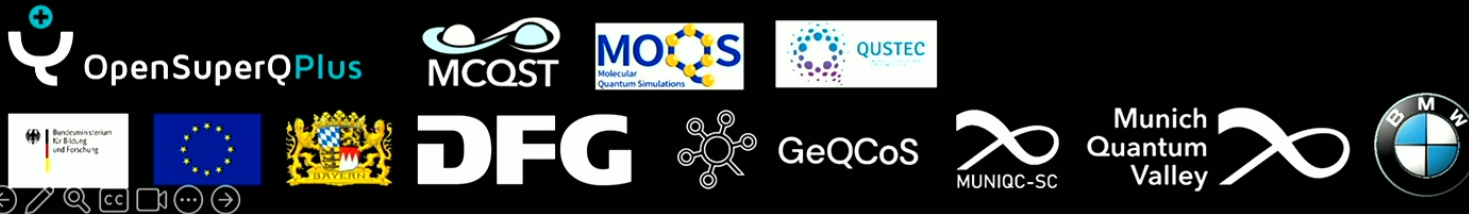
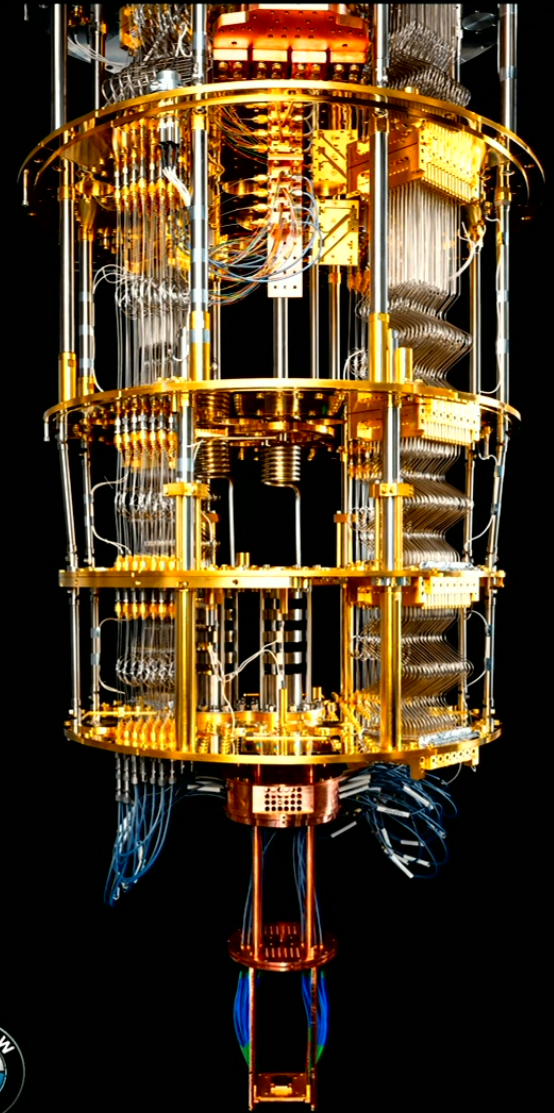
Walther-Meißner-Institut  
DER BAYERISCHEN AKADEMIE DER WISSENSCHAFTEN



# Advances in controlling superconducting qubits for quantum computing

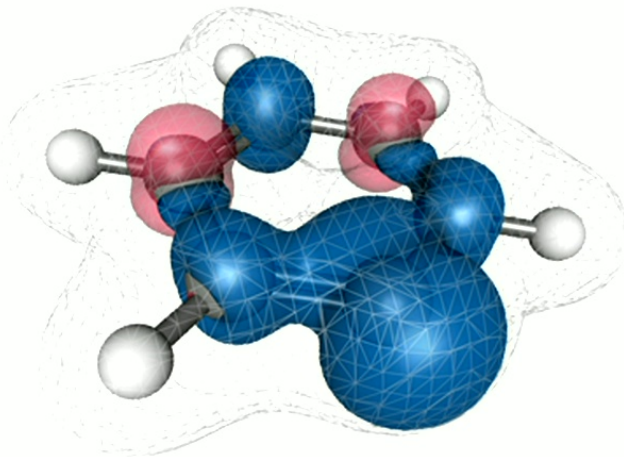
Stefan Filipp  
Walther-Meißner-Institute  
Technical University of Munich

Munich-Waterloo Joint Workshop , Waterloo – Sept 30, 2024

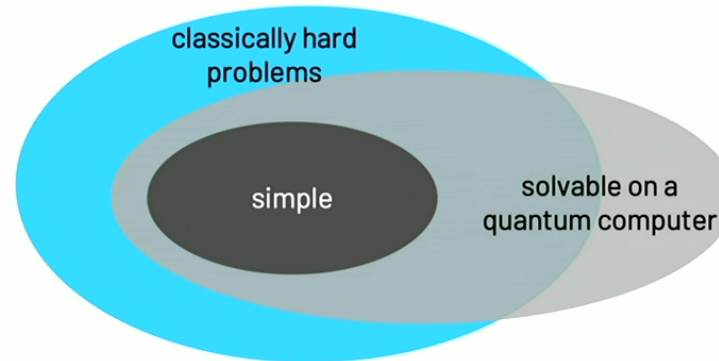


Many problems in science and business are too complex for classical computing systems!

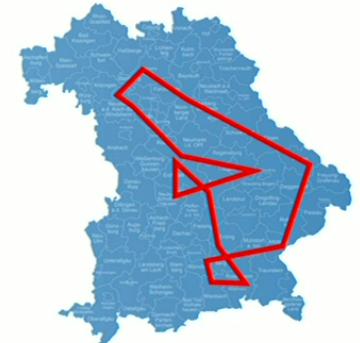
Simulation of  
quantum-mechanical systems:  
Chemistry, material science,...



Algebraic algorithms:  
Machine learning, cryptography, ....

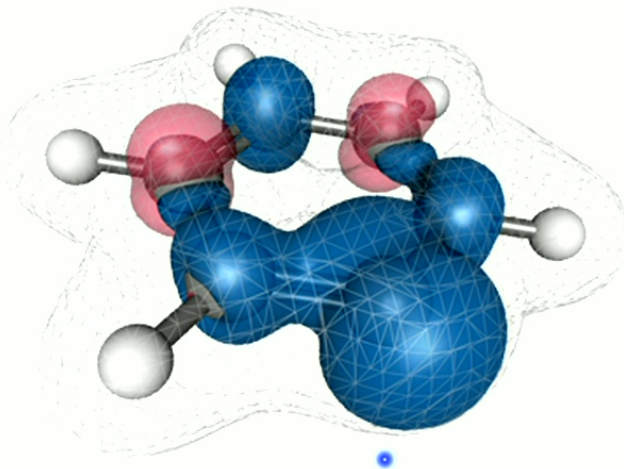


Optimization problems:  
Travelling salesman, logistics, portfolio  
optimization

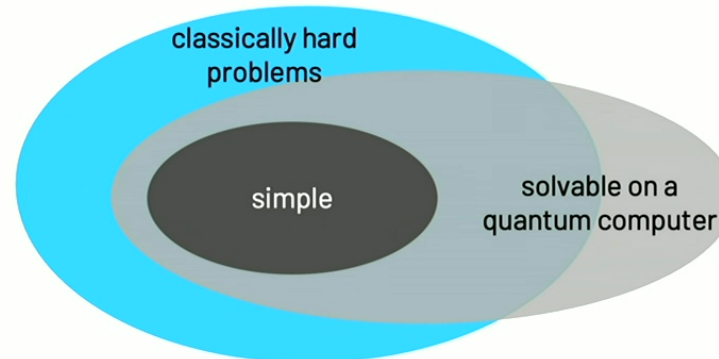


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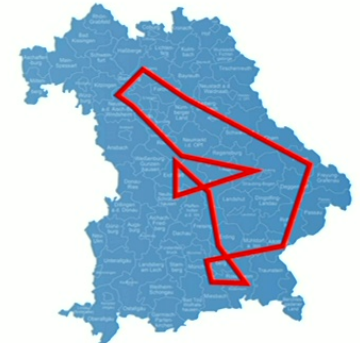
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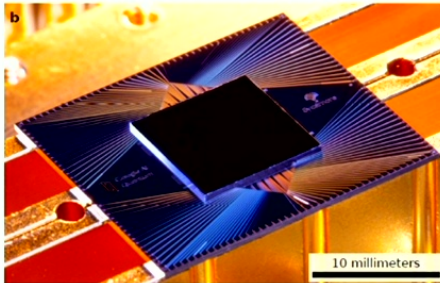


Challenge:

Realization of a quantum computer with many (1.000.000+) **good qubits** and **high-fidelity gate operations**

1<sup>st</sup> quantum computer prototypes are 'on the market' (>100 qubits)

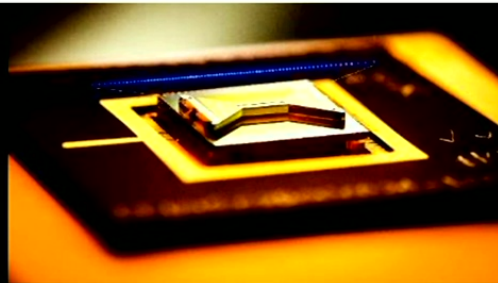
Google AI



IBM



IonQ



AQT



or in development (



OpenSuperQ



GeQCoS



...)

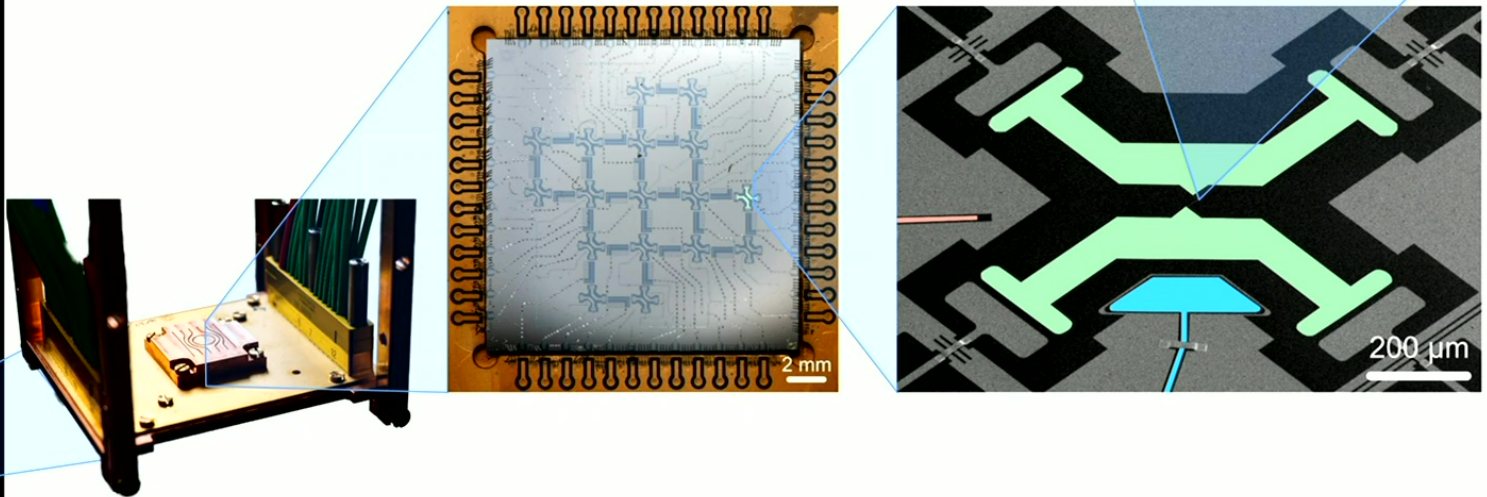
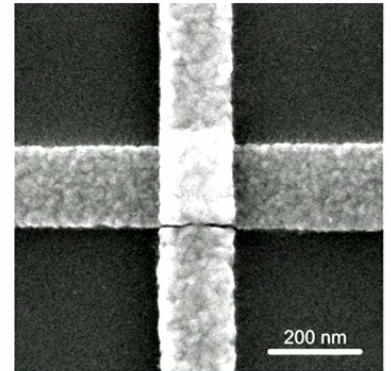
on different platforms (solid state systems, atomic systems, photons, ...)

→ next milestones: [practical quantum algorithm](#), [logical qubits](#)

→ many challenges ahead (scaling, coherence, control & readout, system integration)

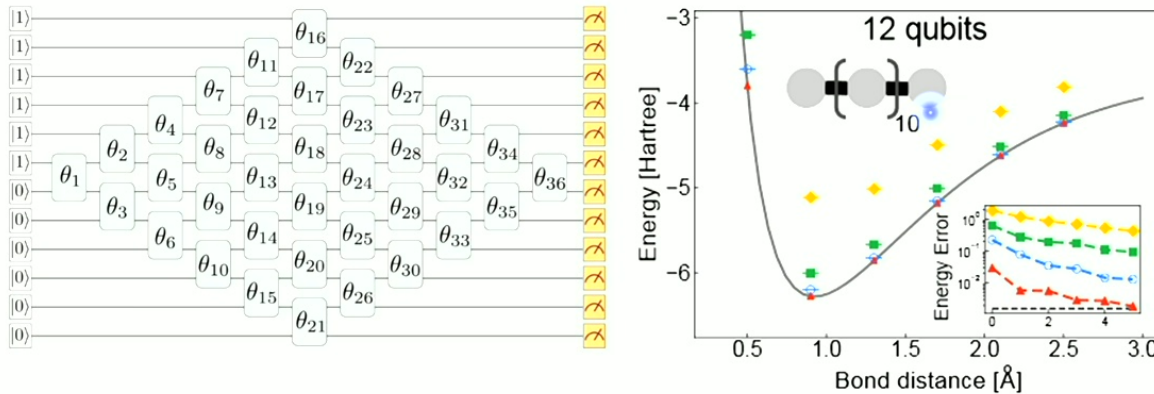
**Features:**

- quantized non-linear superconducting circuits
- typical frequencies: 5 – 10 GHz (microwave range)
- fast gates on ns timescales
- high fidelity gate operations (> 99.9% two-qubit gates)
- scalable fabrication technology



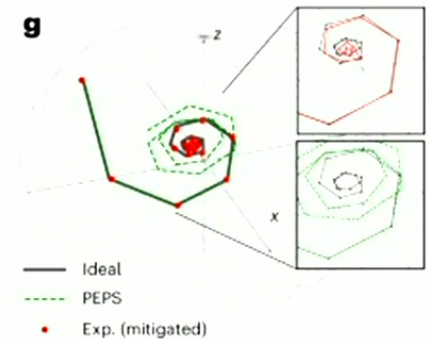
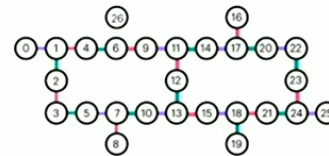
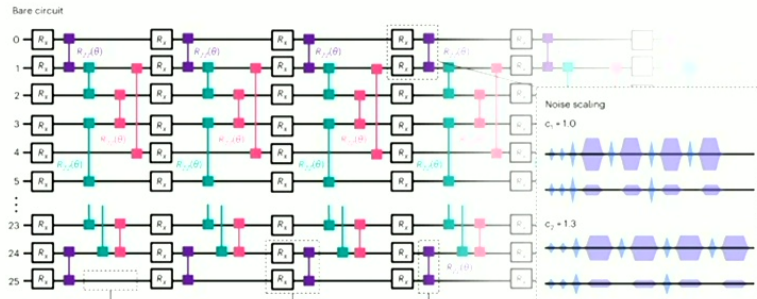
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## Quantum chemistry: calculate ground state of simple molecules (hydrogen chain)



F. Arute, Science 369, 2020;  
A. Kandala, Nature 549 (2017); Nature 567 (2019)

## Material science: dynamics of spin systems (with up to 127 qubits)



Y. Kim et al. (IBM), Nat. Phys 1(2023); Nature (2023)

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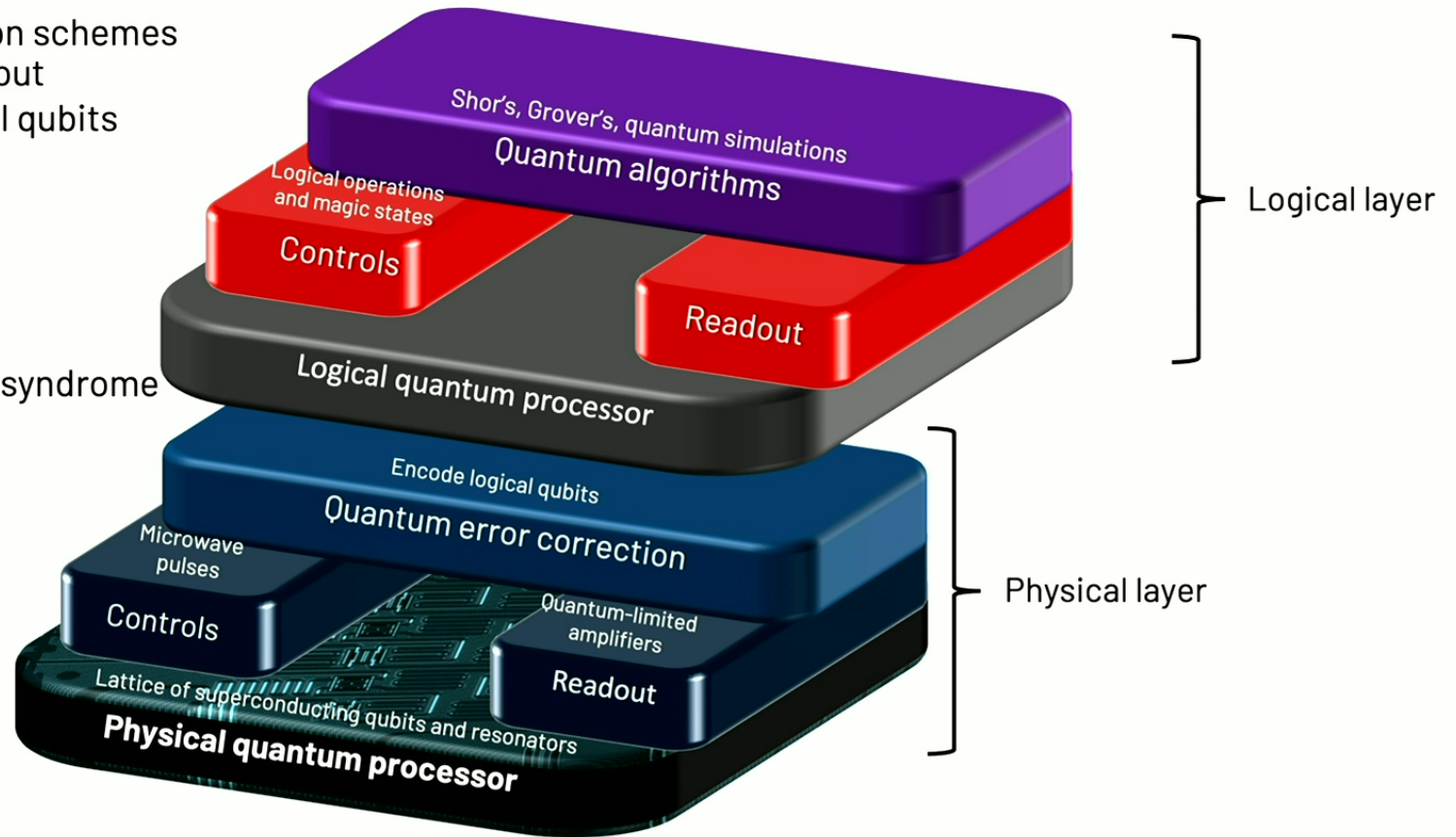
Idea: use error correction schemes to make a 'logical' qubit out of several 'good' physical qubits

$$|0\rangle \rightarrow |00 \dots 0\rangle$$

$$|1\rangle \rightarrow |11 \dots 1\rangle$$

100-1000 x

measure parity as error syndrome

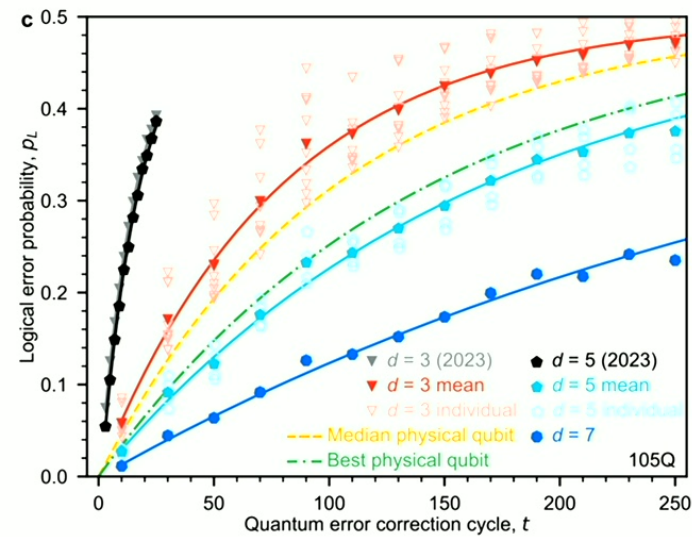
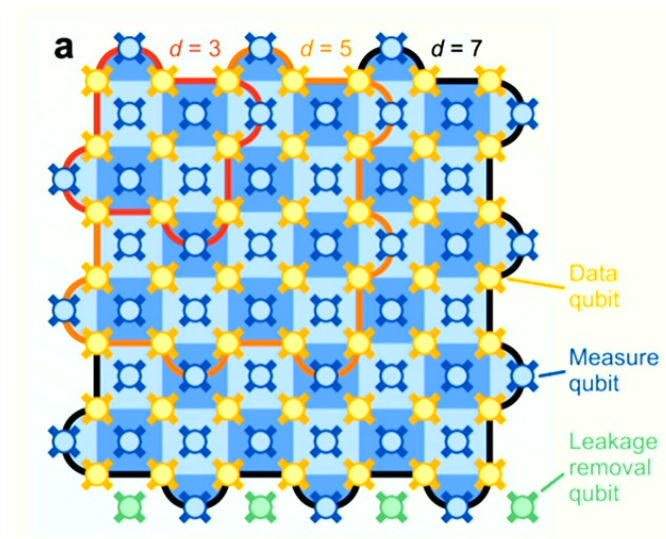


[Gambetta, Chow, Steffen, npj Quantum Information 3, 2 (2017)]



Surface Code error threshold:

Logical error should decrease for increasing number of physical qubits per logical qubit



[R. Acharya et al., Nature (2023);  
Google AI, Nature (2024)]

*Scaling: guarantee performance at scale*

cross-coupling and cross-talk, uniformity & reproducibility, scalable control, I/O, size of qubits, thermal budget

*System: guarantee stable operation conditions*

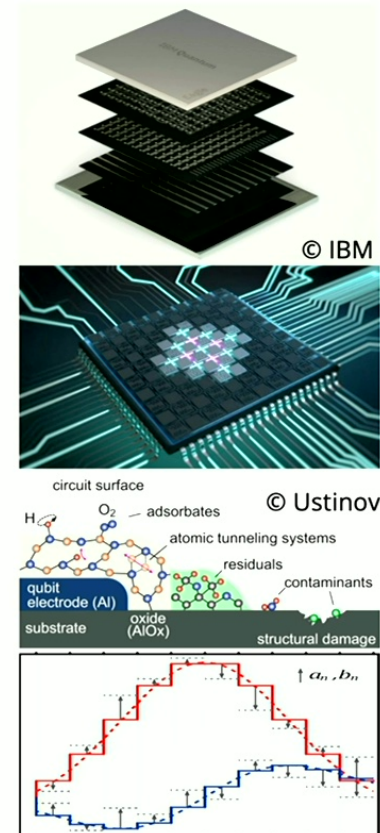
automated calibration & bring-up, run-time environment, characterization & verification, quantum/classical integration, (cryogenic) electronics,...

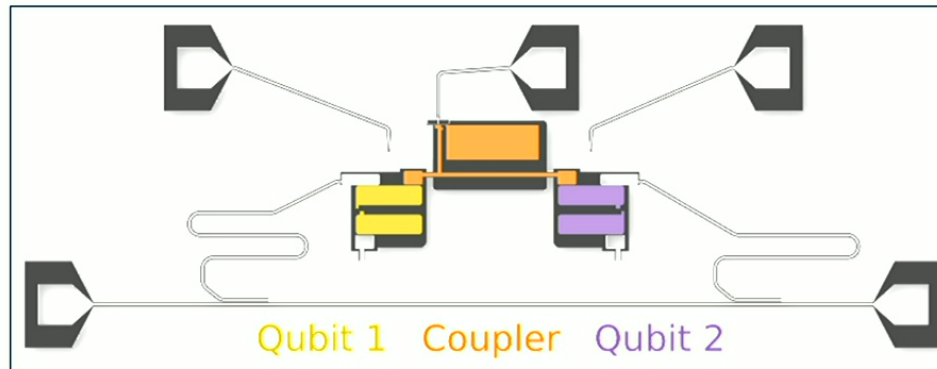
*Coherence: maximize lifetime of quantum states*

identification of loss channels and noise sources (two-level fluctuators, quasi-particles, B-field fluctuations), mitigation (by design, by choice of materials, by fabrication)

*Control & readout: coherence-limited high-fidelity gates*

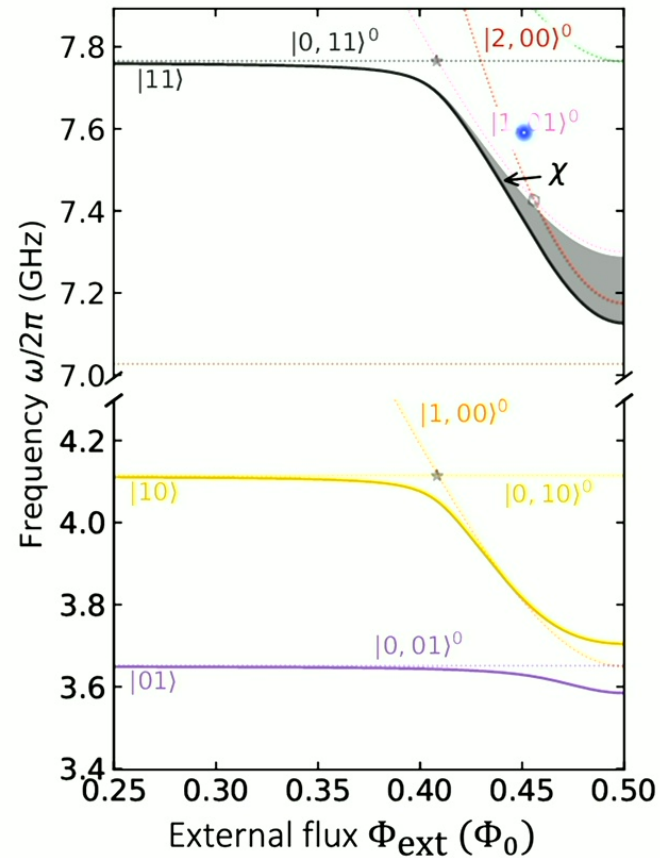
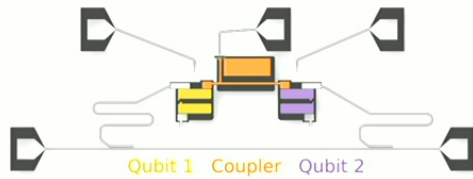
pulse optimization, benchmarking sequences, multi-qubit operations & extended gate sets,...





- **Fixed frequency** qubits with control via dedicated **drive lines**
- Two-qubit **CPHASE** gate by adiabatic flux control of **tunable coupler**

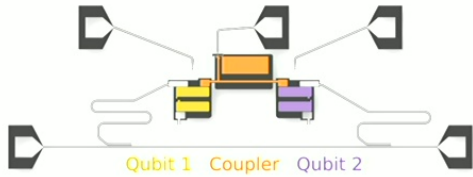
|         | Frequency $\omega$ | Anharmonicity $\alpha$ | $T_1$           | $T_2^{Echo}$    | $T_2^*$        |
|---------|--------------------|------------------------|-----------------|-----------------|----------------|
| Qubit 1 | 4.115 GHz          | -261(1) MHz            | 52(4) $\mu$ s   | 111(41) $\mu$ s | 23(4) $\mu$ s  |
| Qubit 2 | 3.651 GHz          | -275(1) MHz            | 103(20) $\mu$ s | 54(11) $\mu$ s  | 5.6(5) $\mu$ s |
| Coupler | 3.7 - 6.3 GHz      | -124 (1) MHz           | 16(5) $\mu$ s   | —               | —              |



$$\chi = \omega_{11} - (\omega_{01} + \omega_{10})$$

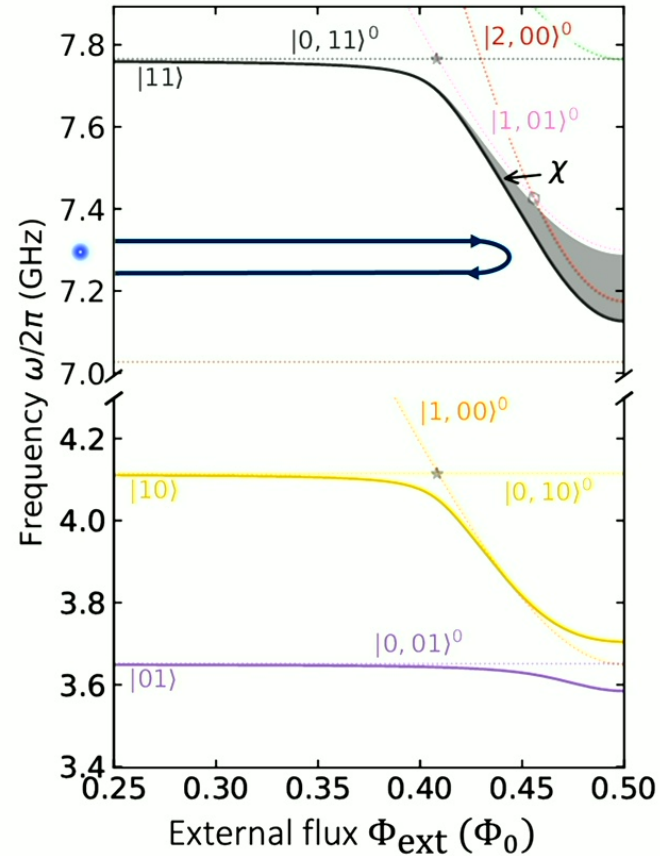
adiabatic states:  $|n_1 n_2\rangle$   
bare states:  $|n_c, n_1 n_2\rangle^0$

[M. C. Collodo et al., Phys. Rev. Lett. 125, 240502 (2020)  
Y. Xu et al., Phys. Rev. Lett. 125, 240503 (2020)  
J. Chu & F. Yan, Physical Review Applied, 16(5), 054020 (2021).  
Stehlik, J. et al., Physical Review Letters, 127(8), 080505 (2021)]



$$U = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & e^{i\alpha} & 0 & 0 \\ 0 & 0 & e^{i\beta} & 0 \\ 0 & 0 & 0 & e^{i(\varphi+\alpha+\beta)} \end{pmatrix}$$

[M. C. Collodo et al., Phys. Rev. Lett. 125, 240502 (2020)  
 Y. Xu et al., Phys. Rev. Lett. 125, 240503 (2020)  
 J. Chu & F. Yan, Physical Review Applied, 16(5), 054020 (2021).  
 Stehlik, J. et al., Physical Review Letters, 127(8), 080505 (2021)]






$$\chi = \omega_{11} - (\omega_{01} + \omega_{10})$$

$$\varphi = \int_0^{t_p} \chi(\Phi_{\text{ext}}(\tau)) d\tau$$

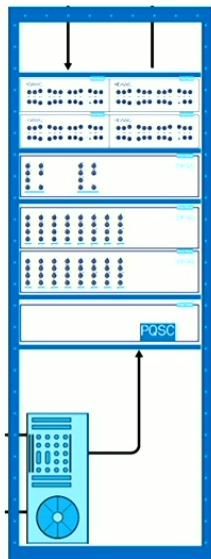
adiabatic states:  $|n_1 n_2\rangle$   
 bare states:  $|n_c, n_1 n_2\rangle^0$

Different pulse shapes: optimize parameters for highest fidelity (numerical simulation)

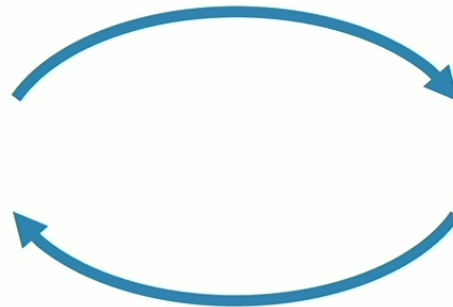
|   | #Parameters | Pulse shape   | $t_p$ ( $\mathcal{F} > 99.9\%$ ) |
|---|-------------|---|----------------------------------|
| Gaussian-square                             | 3           |  | 51 ns                            |
| Fourier Series                              | 5           |  | 15 ns                            |
| Piecewise-constant<br>Using AWG limitations | 22          |  | 9 ns                             |

Trade-off between parameter count and gate time / fidelity

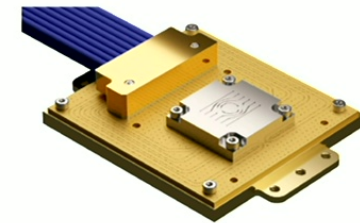
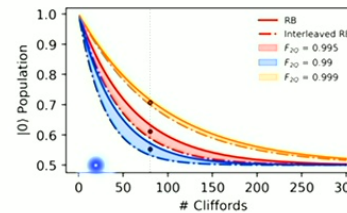
C<sup>3</sup>-toolset: N. Wittler, *et al.*, *Phys. Rev. Applied* 15, 034080 (2021)



adjust pulse  
parameters  
gradient free  
optimization



measure fidelity  
randomized  
benchmarking



## Cost function:

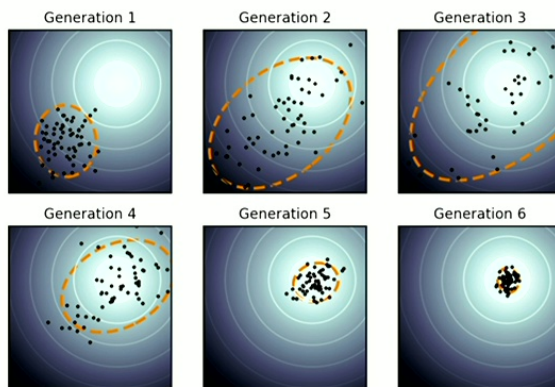
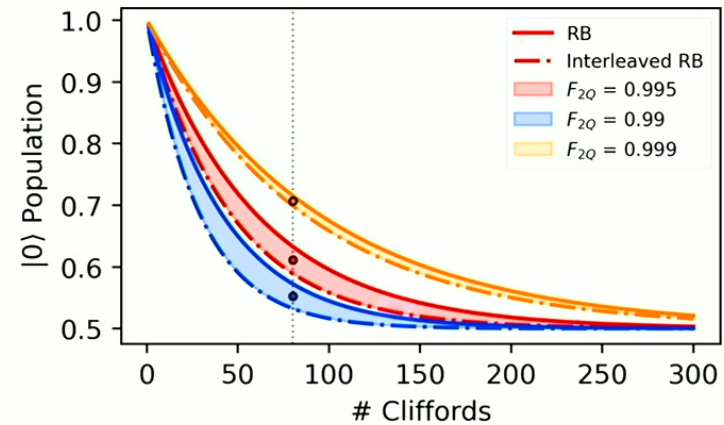
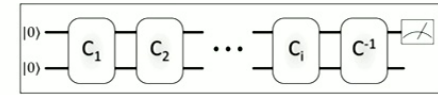
$N$  random pseudo-identity clifford sequences of length  $M$

## Optimization:

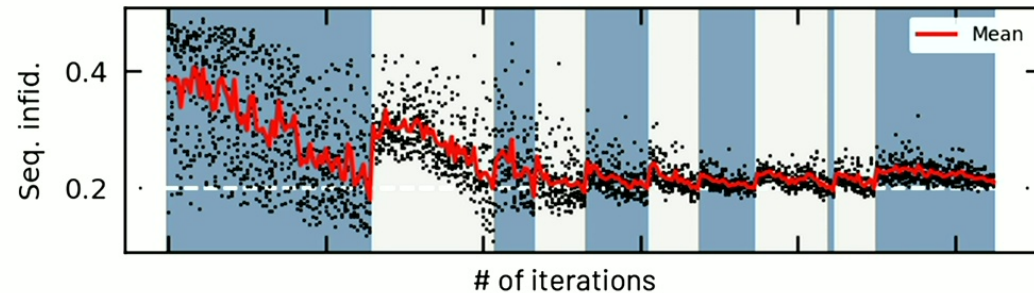
CMA-ES Algorithm (Noise resilient, point-cloud-based evolutionary algorithm)

## Sensitivity:

Gradually increment sequence length  $M$  for sensitivity improvement



Wikipedia: CMA-ES



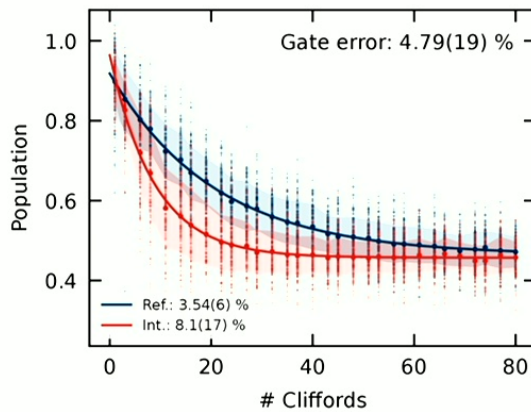
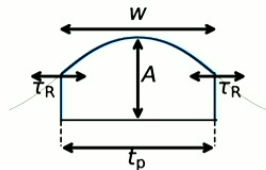
[Hansen N., arXiv:1604.00772 (2016); J. Kelly et al., Phys. Rev. Lett. 112, 240504 (2014); N. Glaser et al. In preparation (2024)]



## Gaussian Square

- 3x Shape parameters
- 2x Z-Rotation

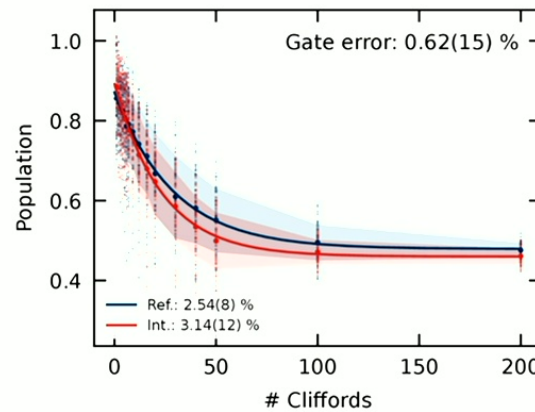
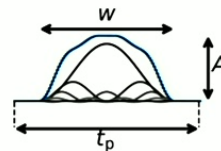
Convergence:  
40 evolutions (40 min)



## Fourier Series

- 5x Fourier parameter
- 1x Width
- 2x Z-Rotation

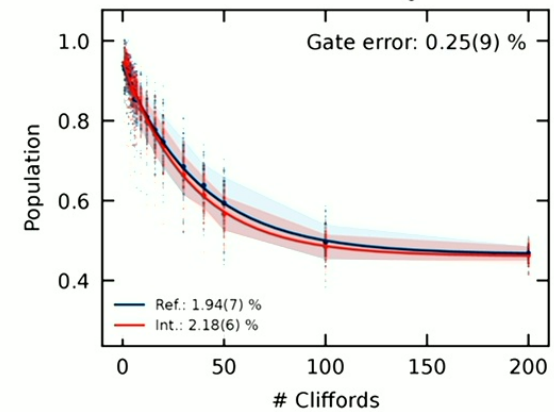
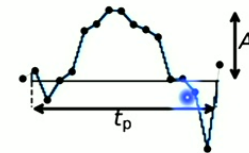
Convergence:  
80 evolutions (80 min)



## Linear Interpolation

- 17x Nodes
- 1x Amplitude, 1x Width
- 2x Z-Rotation

Convergence:  
260 evolutions (250 min)

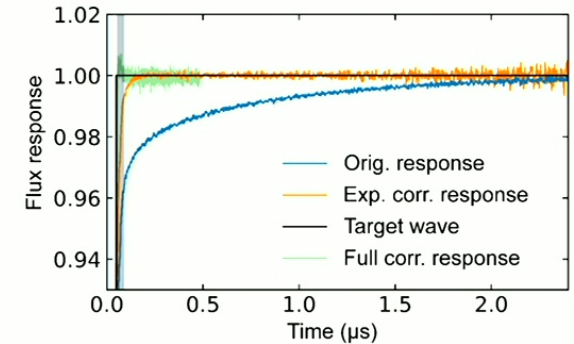


### Correct effect of long term flux distortions

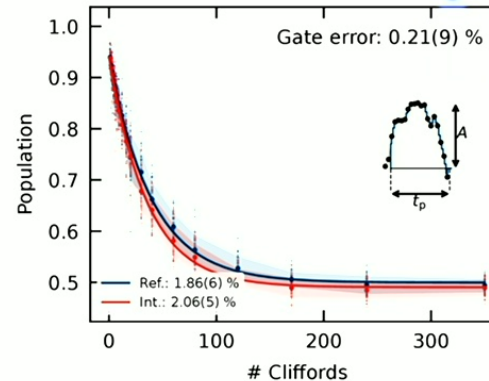
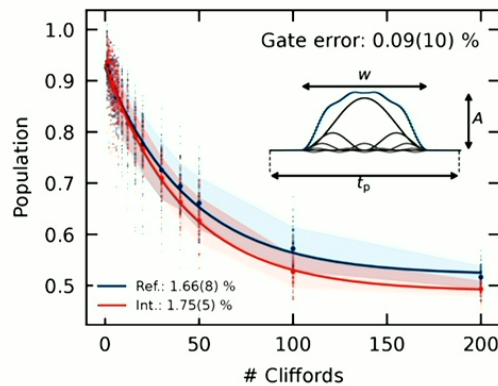
- cryoscope measurement
- correction with HDAWG real-time exponential filters + FIR filter

→ error for 60ns Fourier series reduced by a factor of six:  $\epsilon = 0.09 \pm 0.1\%$

→ fast 20 ns gate with linear interpolation:  $\epsilon = 0.21 \pm 0.09\%$



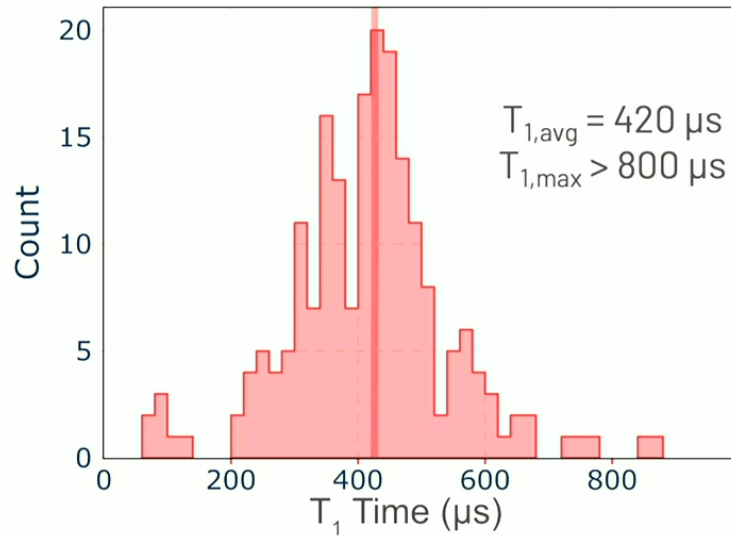
M. Rol et al., *Applied Physics Letters* 116.5 (2020).



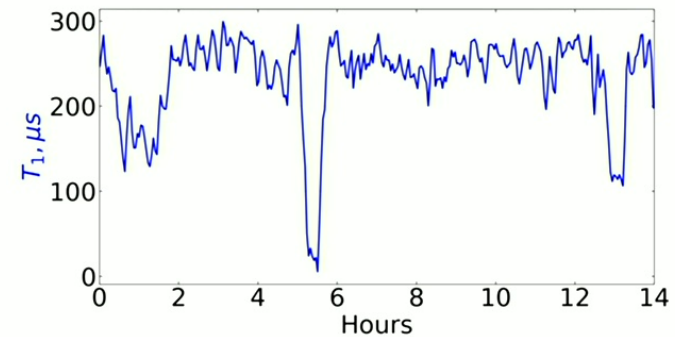
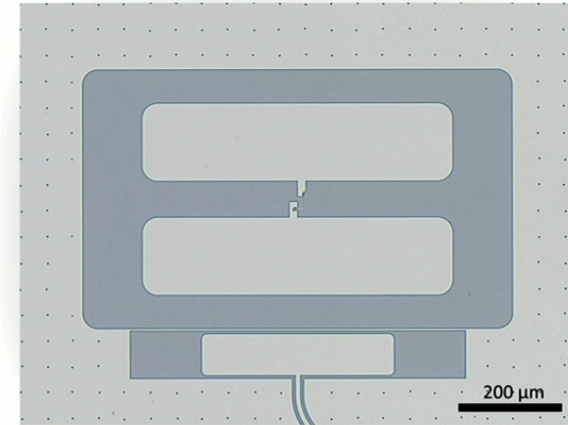
[N. Glaser et al., in preparation (2024)]

- Fixed-frequency single transmon qubits (frequency around 4.4 GHz)
- Al/AIO<sub>x</sub>/Al Josephson junctions with Niobium ground plane

[L. Koch, N. Bruckmoser, in preparation (2024)]



$$\rightarrow T_{1,2}/T_{gate} > 10000$$



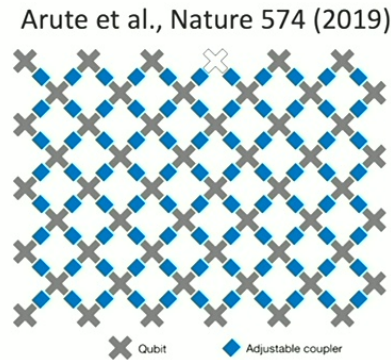
→ fluctuating environment (TLSs)

Quantum processors increase in size

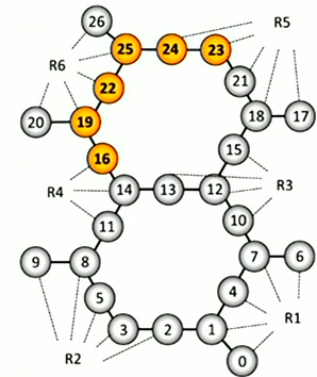
low connectivity and local operations lead to large gate overhead

control multiple local interactions simultaneously

multi-qubit and non-local operations



Petar et al., QS&T 6 (2021)



Gu et al., PRX Quantum 2 (2021)  
 Burkhart et al., PRX Quantum 2 (2021)  
 Zhang et al., PRL 128 (2022)  
 Warren et al., arXiv 2207.02938 (2022)

**Glaser et al., PR Applied 19 (2023)**  
 Kim et al., Nature Physics 1 (2022)  
 Baker et al., Appl PL 120 (2022)  
 Lu et al., PRX Quantum 3 (2022)

efficient state transfer along chains

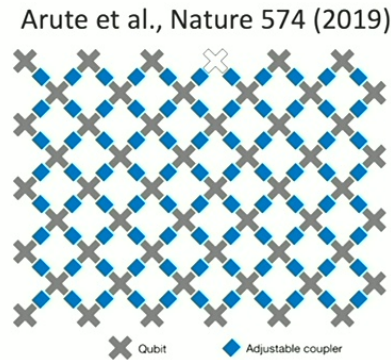
Christandl et al., PRL 92 (2004)  
 Li et al., PRApplied 10 (2018)  
 Genest et al., Annals of Physics 371 (2016)  
 Lemay et al., JPA 49 (2016)  
**Nägele et al., PR Research 4 (2022)**

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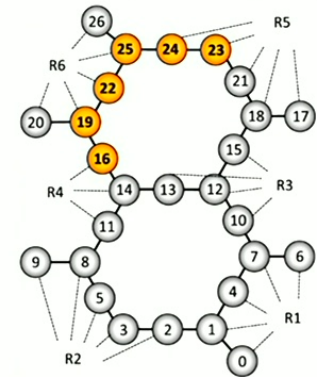
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Petar et al., QS&T 6 (2021)

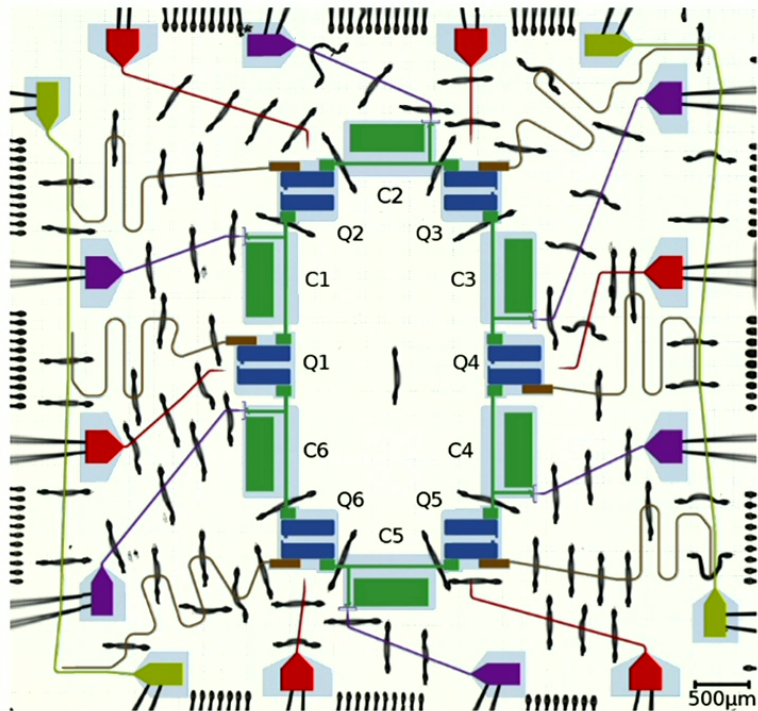


Gu et al., PRX Quantum 2 (2021)  
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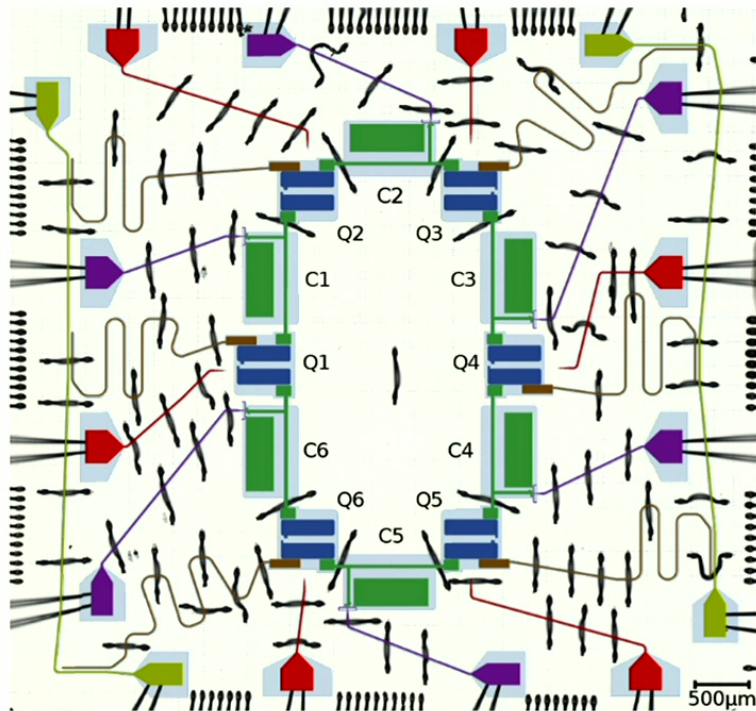
**Glaser et al., PR Applied 19 (2023)**  
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## efficient state transfer along chains

Christandl et al., PRL 92 (2004)  
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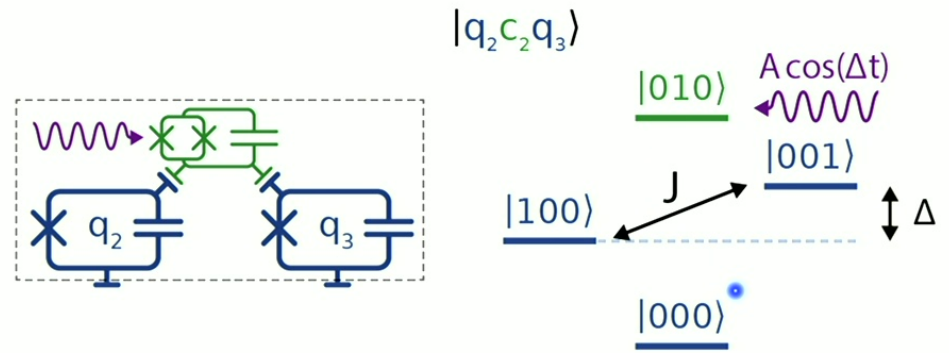


- 6 fixed-frequency **transmons**, individual **drive lines**.
- 6 **resonators**, 2 **feedlines**.
- 6 **tunable-couplers**, individual **flux lines**.

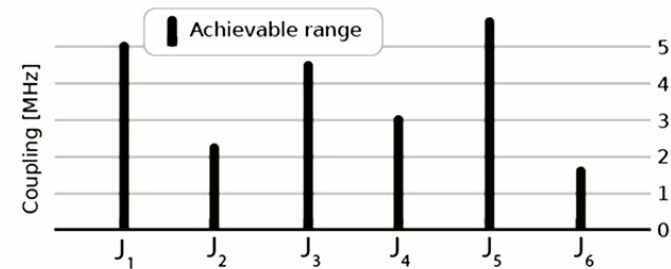


- 6 fixed-frequency **transmons**, individual **drive lines**.
- 6 **resonators**, 2 **feedlines**.
- 6 **tunable-couplers**, individual **flux lines**.

Two-qubit interactions by **parametric drive** of the coupler.



coupling strength **J** is controllable via amplitude **A**.



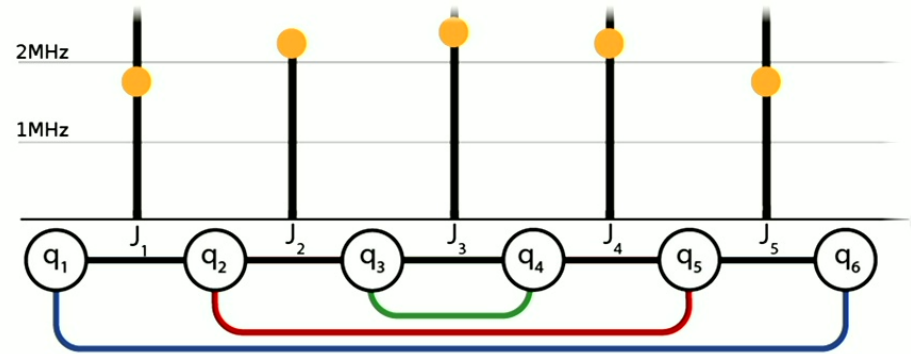
Spin chain Hamiltonian

$$H_{\text{chain}} = \sum_{n=1}^{N-1} J_n (\hat{\sigma}_n^- \hat{\sigma}_{n+1}^+ + \hat{\sigma}_n^+ \hat{\sigma}_{n+1}^-)$$

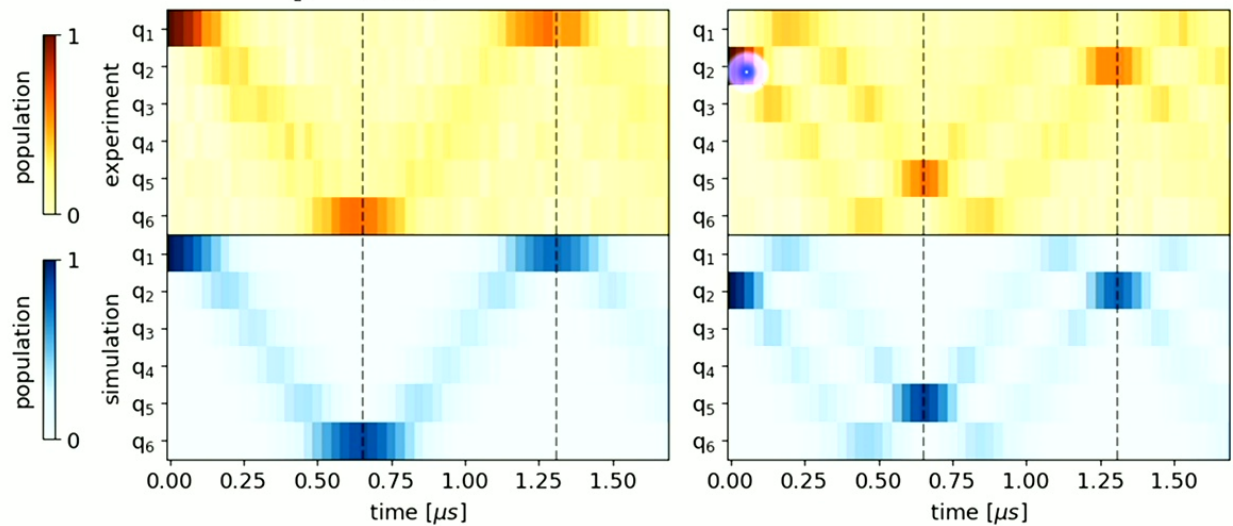
Given  $J_n = \pi \sqrt{n(N-n)}/2\tau$   
For a transfer time  $\tau$

$$H_{PST} = \sum_{n=1}^{\lfloor N/2 \rfloor} \frac{\pi}{2\tau} (\hat{\sigma}_n^- \hat{\sigma}_{N+1-n}^+ + \hat{\sigma}_n^+ \hat{\sigma}_{N+1-n}^-)$$

Effective PST Hamiltonian (single excitation)



$[\tau = 640\text{ns}]$





In the presence of multiple excitations:

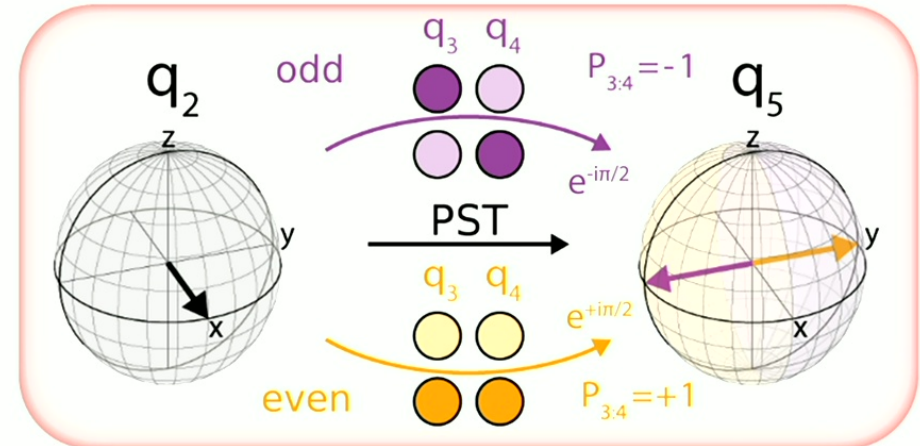
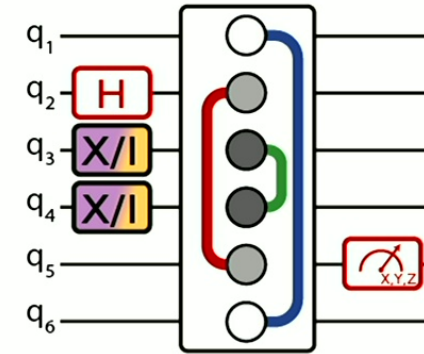
## Effective PST Hamiltonian

$$H_{PST} = \sum_{n=1}^{\lfloor N/2 \rfloor} \underbrace{\left( \prod_{k=n+1}^{N-n} \hat{\sigma}_k^z \right)}_{\hat{P}_{n+1;N-n}} \frac{\pi}{2\tau} (\hat{\sigma}_n^- \hat{\sigma}_{N+1-n}^+ + \hat{\sigma}_n^+ \hat{\sigma}_{N+1-n}^-)$$

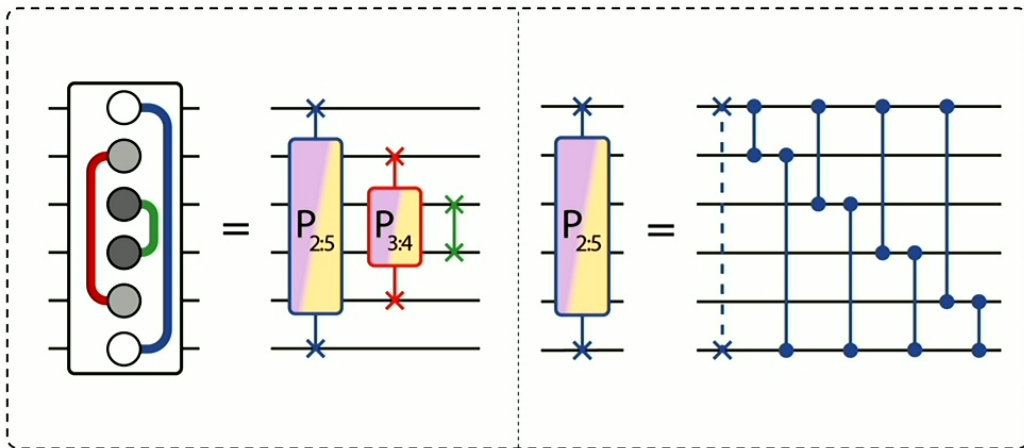
Nägele et al., PRR 4 (2022)

**Parity** controls phase of transfer

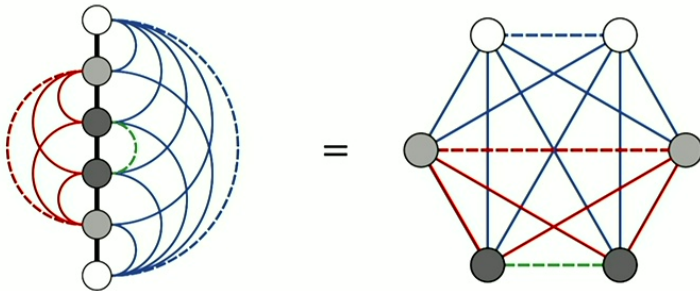
$$iSWAP_{\hat{P}_{3:4}} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & \pm i & 0 \\ 0 & \pm i & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$



PST's decomposition into 2-qubit gates:



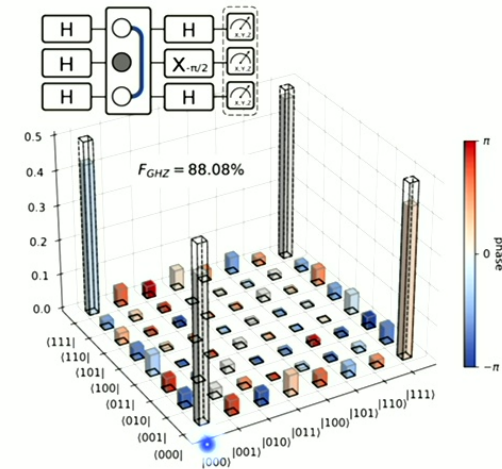
Graph states representation:



Qubits → nodes  
Ising interactions → edges

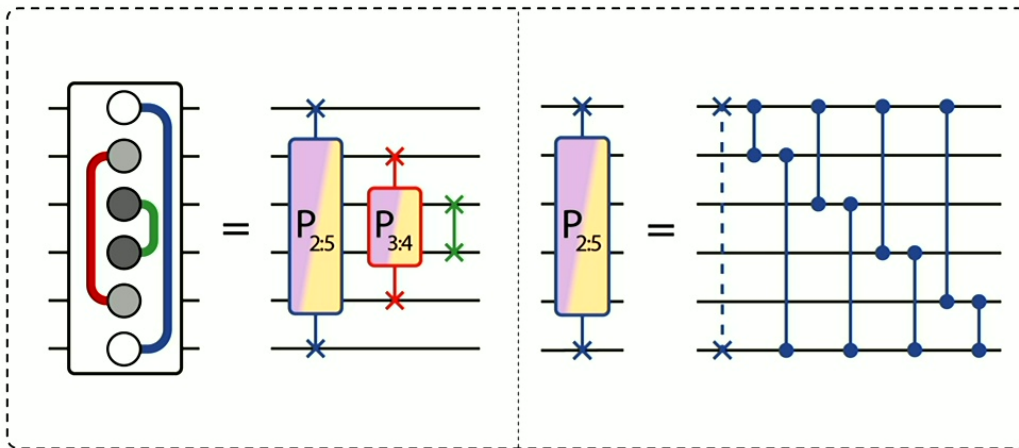
**GHZ-state preparation:** based on single multi-qubit unitary transformation

$$|GHZ\rangle = \frac{|0\rangle^{\otimes N} + |1\rangle^{\otimes N}}{\sqrt{2}}$$

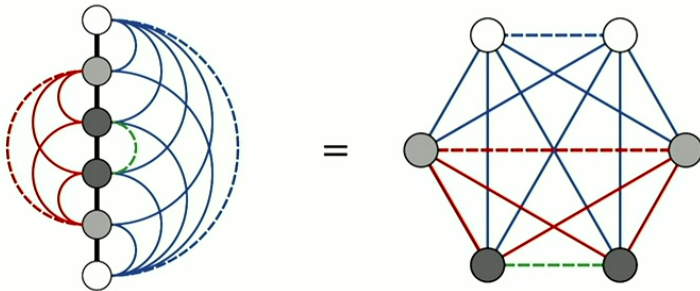


[F. Roy, J. Romeiro et al. arXiv:2405.19408 (2024)]

PST's decomposition into 2-qubit gates:



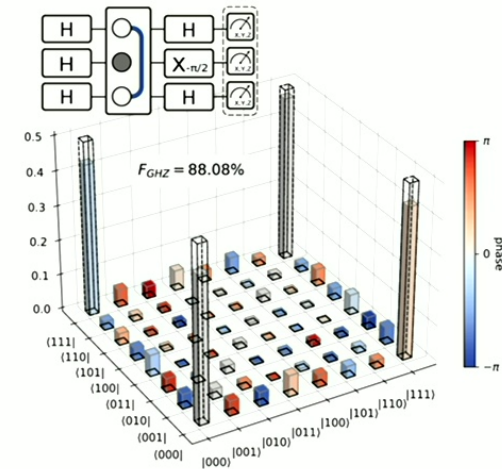
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Qubits → nodes  
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**GHZ-state preparation:** based on single multi-qubit unitary transformation

$$|GHZ\rangle = \frac{|0\rangle^{\otimes N} + |1\rangle^{\otimes N}}{\sqrt{2}}$$



→ reduction of gate overhead, parity-check codes, quantum simulation of fermionic many-body systems, generation of different classes of graph states

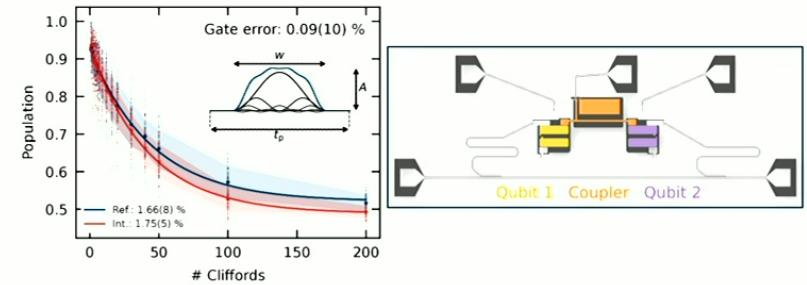
[F. Roy, J. Romeiro et al. arXiv:2405.19408 (2024)]

## Entangling operations based on tunable couplers

High-fidelity two-qubit gates >99.9%

→ extension to full architecture

[N. Glaser et al., in preparation (2024)]

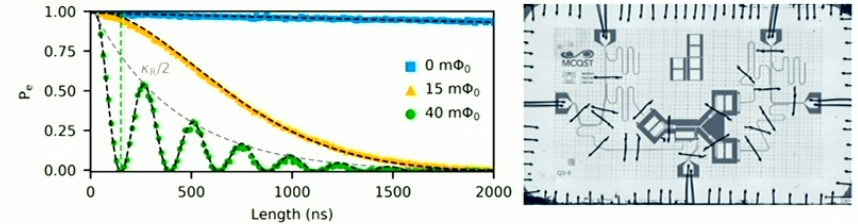


## Architecture for reset, leakage reduction and tunable readout

→ multi-qubit couplers

→ extension to full architecture

[G. Huber et al, arXiv:2403.02203 (2023)]



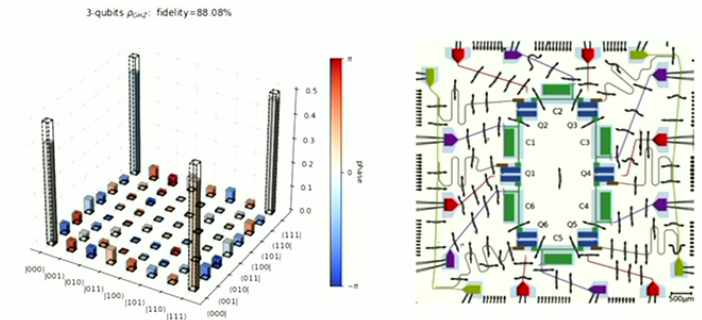
## State transfer and multi-qubit entanglement generation

fractional state transfer and entanglement generation demonstrated

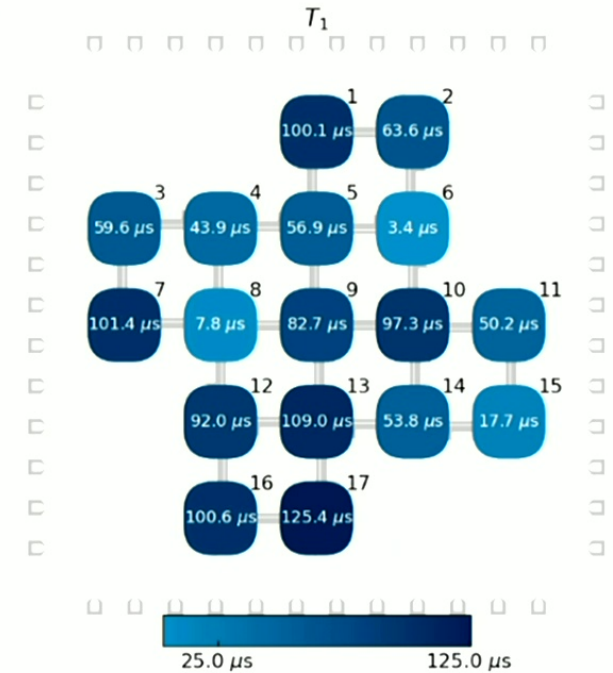
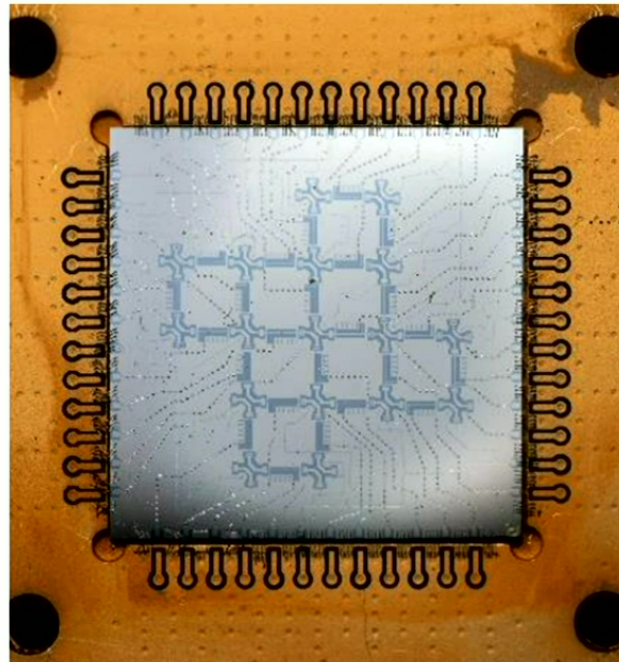
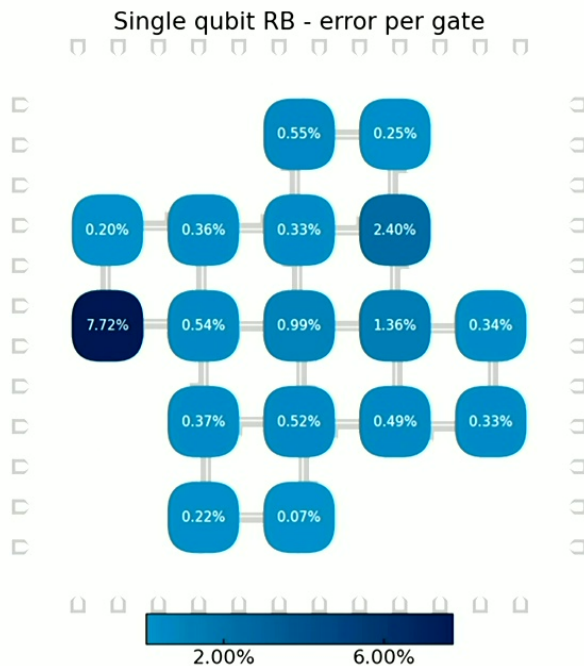
→ extended gate set, multi-qubit gates

→ simulation of fermionic systems, parity measurements

[M. Nägele et al., PRR(2022) F. Roy, J. Romeiro et al., arXiv:2405.19408 (2024)]



## Operation of a 17 qubit processor (transmon-based)





**PI**

Stefan Filipp

**Project Management**

Klaus Liegener

**Researchers**

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Christian Schneider

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Noelia Fernandez  
Gleb Krylov  
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Kevin Kiener  
Leon Koch  
Frederik Pfeiffer  
Lea Richard

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**We also acknowledge support from the other WMI groups.**



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