

**Title:** Quantum simulation in the presence of errors

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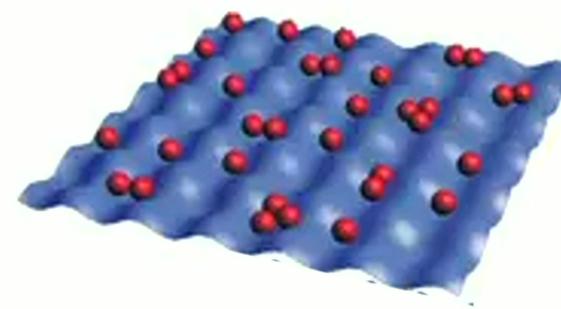
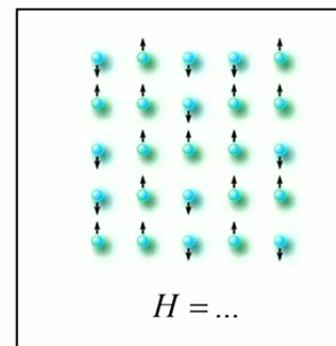
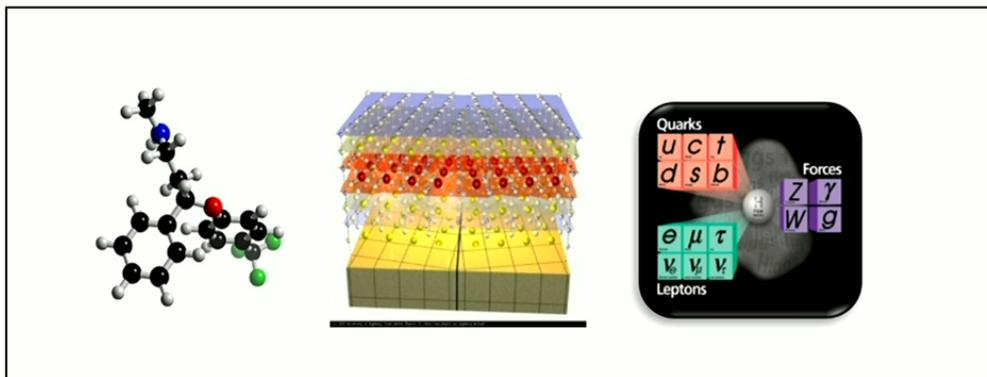
BMZ  
Hans-Kopfermann

# Analog quantum simulation in the presence of errors

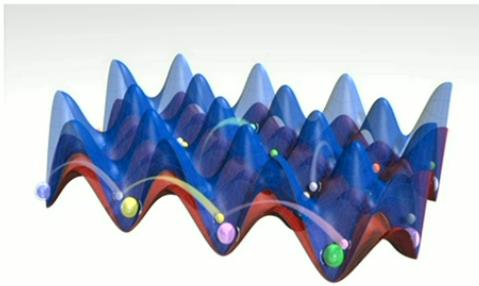
Ignacio Cirac

Waterloo-Munich joint workshop,  
Perimeter Institute, September 30, 2024

# Quantum simulation

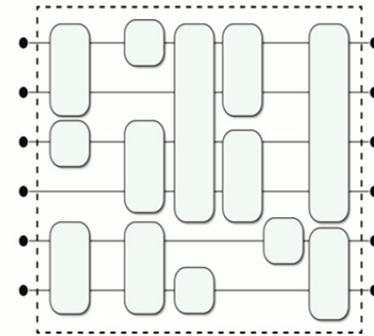


## Analog



- Easier to build
- More natural for fermions

## Digital



- Universal
- Error correction

# Quantum Advantage

**Question:** is quantum advantage possible in analog quantum simulation?

- Wrt best known classical algorithms
- In terms of complexity
  
- Non-standard scenario (eg, analog, thermodynamic limit,...)
- Fair comparison: with „certified“ methods

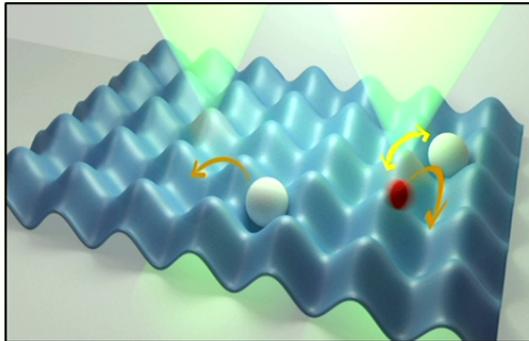
**Question:** is quantum advantage possible in the presence of errors?

- No error correction
- What is quantum advantage?

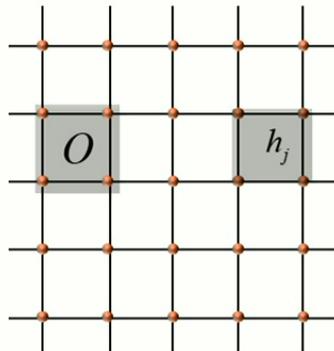
# Quantum Advantage I

R. Trivedi, A. Franco, IC, Nat. Comm. 6507 (2024)

# Quantum Simulation: Goal



$$H = \sum_j h_j$$



Dynamics

$$i d_t |\Psi\rangle = H |\Psi\rangle$$

Ground State

$$H |\Psi\rangle = E_0 |\Psi\rangle$$

Thermal equilibrium

$$\rho = e^{-H/kT}$$

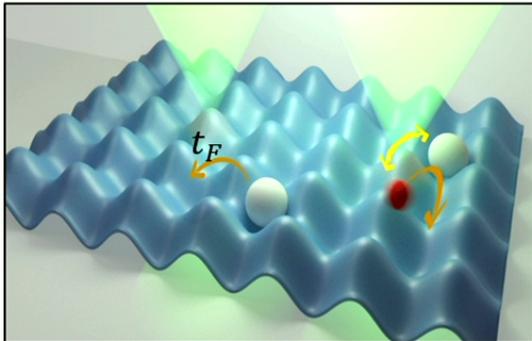


Observable  
 $\langle O \rangle$

# Errors

Errors are extensive

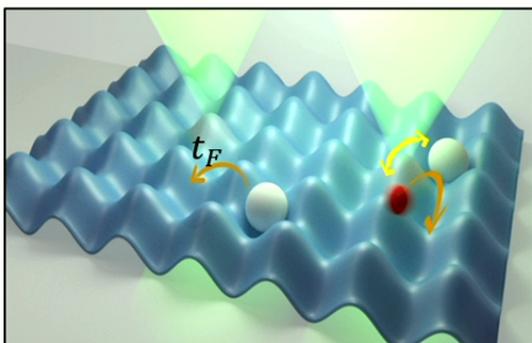
$$H_0 = \sum_j h_j$$



$$H = \sum_j h_j + \underbrace{\varepsilon \sum_j v_j}_{V \sim n \varepsilon}$$

# Errors

$$H_0 = \sum_j h_j$$



Errors are extensive

$$H = \sum_j h_j + \underbrace{\varepsilon \sum_j v_j}_{V \sim n \varepsilon}$$

- Ideal  $|\Psi_0(t)\rangle = e^{-iH_0 t} |\Psi(0)\rangle$
- Real  $|\Psi(t)\rangle = e^{-i(H_0+V)t} |\Psi(0)\rangle$

$$n \varepsilon t \approx 1 \Rightarrow |\Psi\rangle \perp |\Psi_0\rangle$$

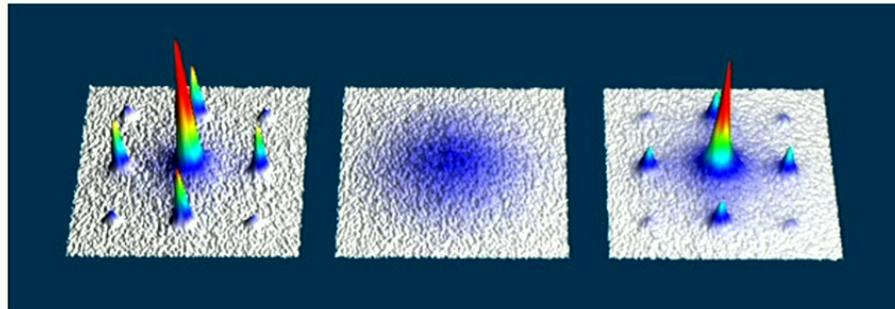
Simulation is limited to  $t \approx \frac{1}{\varepsilon n}$

Errors are extensive

$$H = \sum_j h_j + \varepsilon \sum_j v_j$$

Observables are intensive

$$m = \frac{1}{n} \sum_{j=1}^n \langle s_j^z \rangle$$

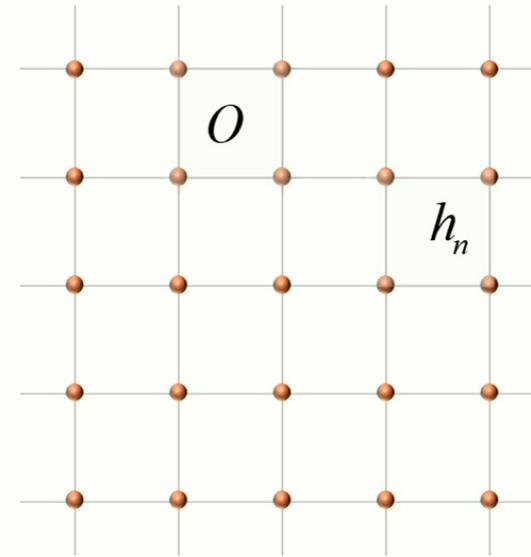


Bloch, Esslinger, Greiner, Hänsch, Nature (2002)

- **Model:**  $H_0 = \sum_j h_j$
- **Q. Simulator:**  $H = \sum_j h_j + \varepsilon \sum_j v_j$
- **Observable:**  $o = \frac{1}{n} \sum_{j=1}^n \langle o_j \rangle$

**Stable:** if  $|o_{H_0} - o_H| \leq f(\varepsilon)$  independent of  $n$

- **Stable problems:**
  - Dynamics
  - Ground states of gapped Hamiltonians\*
  - Thermal states (non critical)
  - Critical Gaussian states



# Quantum Simulation without Errors

- Q. Simulator:  $H = \sum_j h_j$
- Thermodynamic limit:  $o^* = \lim_{n \rightarrow \infty} \langle o_n \rangle$

Does not fit the standard scenario of complexity

Aharonov and Irani (2022)  
Wathson and Cubitt (2022)

We have a:

- Classical algorithm:  $o^{\text{cl}}$
- Quantum algorithm:  $o^{\text{q}}$
- **Question:** Given  $\delta > 0$ , what is computational time to obtain  $|o^{\text{cl,q}} - o^*| < \delta$

$$T^{\text{q}} = \text{polynomial in } 1/\delta \begin{cases} \text{Dynamics} \\ \text{Ground state (1/poly}(N)\text{ gap)} \end{cases}$$

$$T^{\text{cl}} = \text{superpolynomial (dynamics) or exponential (ground state) in } 1/\delta$$

# Quantum Simulation with Errors

- Q. Simulator:  $H = \sum_j h_j + \varepsilon \sum_j v_j$

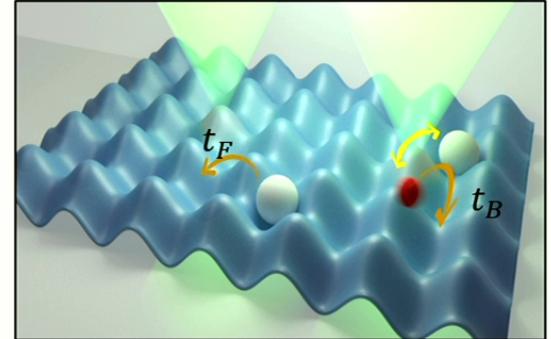
- Quantum advantage:

- With a quantum simulator, the error will be  $\delta = f(\varepsilon)$
- What is the time in a classical computer to reach an error  $\delta = f(\varepsilon)$  ?
- How does that time scale with  $\varepsilon$ ?

$T^q = \text{polynomial in } 1/\varepsilon$

$T^{\text{cl}} = \text{superpolynomial (dynamics) or exponential (ground state) in } 1/\varepsilon$

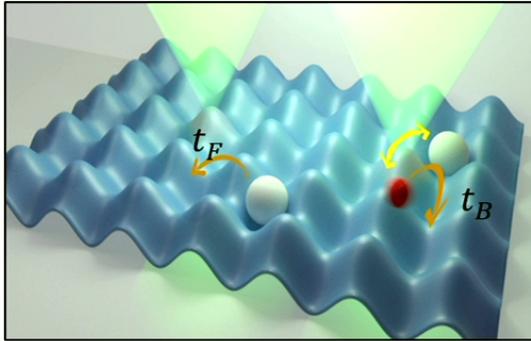
- **Example:** if hardware error is reduced by a factor of 10, i.e.  $\varepsilon \rightarrow \varepsilon/10$   
the classical computational time (depth)  $T^{\text{cl}} \rightarrow (T^{\text{cl}})^{10}$



# Quantum Advantage II

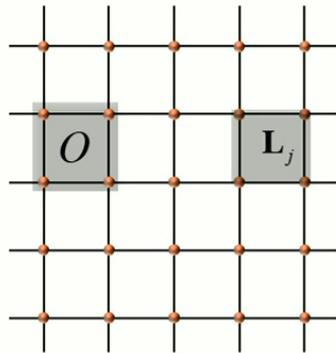
V. Kashyap, G. Styliaris, S. Mouradian, IC, R. Trivedi, arXiv:2404.11081

# Simulating dissipative Quantum Systems: Goal



~~$$H = \sum_j H_j$$~~

$$\mathbf{L} = \sum_j \mathbf{L}_j$$



$$d_t \rho(t) = \mathbf{L} \rho(t)$$

$$\mathbf{L} \rho = 0$$

Dynamics

Steady state  
(Fixed point)

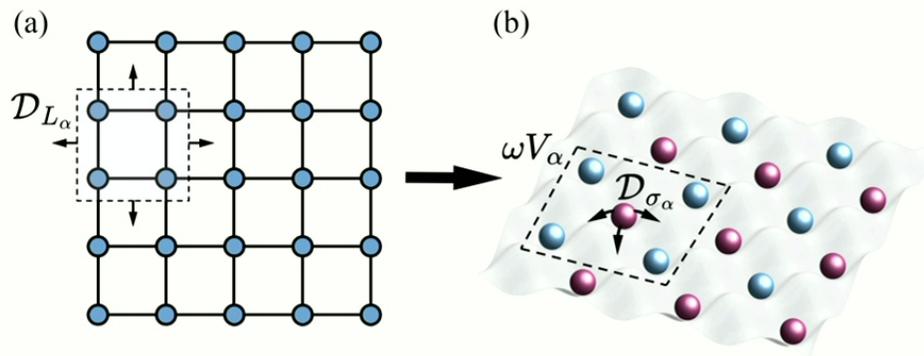


Observable

$$\langle O \rangle$$

# Quantum simulation of dissipative systems

Simulating dissipative systems:



Pastawski, Clemente, IC, Phys. Rev. A (2011)

# Quantum Simulation without Errors

**Rapid Mixing:** A local observable  $O$  (with  $\|O\| < 1$ ) in a spatially local Lindbladian  $L$  with unique steady state,  $\sigma$ , is rapidly mixing if it converges exponentially fast to  $\text{Tr}(O\sigma)$

$$\left| \text{Tr} \left[ O e^{Lt} \rho(0) \right] - \text{Tr}(O\sigma) \right| < k e^{-\gamma t}$$

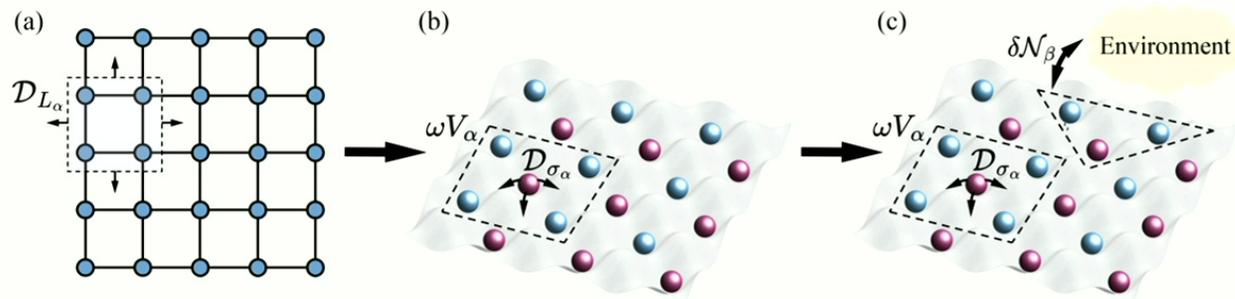
Toby Cubit, Angelo Lucia, Spyridon Michelakis, David Perez-Garcia, CMP (2015)

**Simulation:** To achieve additive error  $\delta$  in  $\langle O \rangle$  in steady state requires an analog Quantum Simulator with

$$T^q = \Theta \left[ \gamma^{-k} \delta^{-1} \log(\delta^{-1}) \right]$$

**Advantage:** No classical algorithm in 2D that achieves that in time  $\text{poly}(\gamma^{-1}, \delta^{-1})$  unless BQP=BPP

# Errors



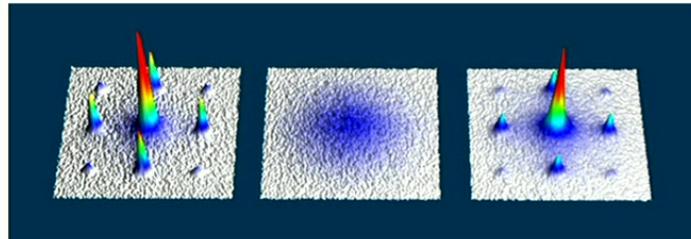
Errors are extensive

$$\mathbf{L} = \sum_j \mathbf{L}_j + \varepsilon \sum_j \mathbf{N}_j$$

$$\left\| \sum_j \mathbf{N}_j \right\| = O(n)$$

Observables are intensive

$$m = \frac{1}{n} \sum_{j=1}^n \langle s_j^z \rangle$$



Bloch, Esslinger, Greiner, Hänsch (2002)

# Quantum Simulation with Errors

Q. Simulator:  $\mathbf{L} = \sum_j \mathbf{L}_j + \varepsilon \sum_j \mathbf{N}_j$

Simulation: In the presence of error rate  $\varepsilon$ , one can obtain  $\langle O \rangle$  in steady state with an error  $\delta$  and time  $t_{\text{sim}}$

$$\delta = O\left[\sqrt{\varepsilon}\right] \quad T^Q = O\left[\varepsilon^{-1/2} \log(\varepsilon^{-1/2})\right]$$

Advantage: There cannot exist a randomized classical algorithm in 2D that achieves that with the same precision in time  $T^{\text{cl}} = \text{poly}(\varepsilon^{-1})$  unless BQP=BPP

Example: if hardware error is reduced by a factor of 10,  $\delta \rightarrow \delta/10$   
the classical computational time (depth)  $T^{\text{cl}} \rightarrow (T^{\text{cl}})^{10}$

# Techniques

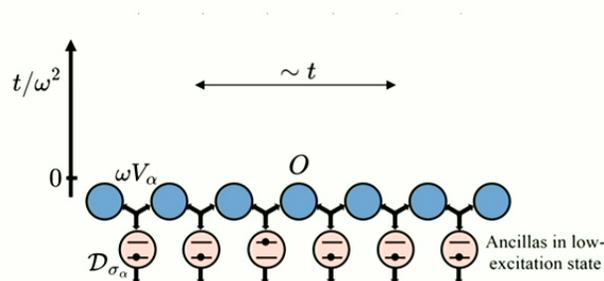
Rahul Trivedi, Adrian Franco-Rubio, IC, Nat. Comm. 6507 (2024)

Vikram Kashyap, Georgios Styliaris, Sara Mouradian, IC, Rahul Trivedi, arXiv:2404.11081



R. Trivedi

## Engineered dissipation



Pastawski, Clemente, Cirac (2011)

## Quantum advantage

### Adapt to a 2D geometry

Verstraete, Wolf, Cirac (2009)

Aharonov, van Dam, Kempe, Landau, Lloyd, Regev (2008)

Stability with respect to errors in all qubits  
both in time and steady state

Quantum **simulation** is (arguably) the most suitable application for quantum computers

Currently, **analog** and digital quantum computers are very well suited for quantum simulation **despite errors**



# Quantum Computing

## Efficient Simulation of Quantum Chemistry Problems in an Enlarged Basis Set

Maxine Luo<sup>1,2</sup> and J. Ignacio Cirac<sup>1,2</sup>

<sup>1</sup>Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Str. 1, D-85748 Garching, Germany

<sup>2</sup>Munich Center for Quantum Science and Technology (MCQST), Schellingstr. 4, D-80799 München, Germany

(Dated: July 8, 2024)

Standard Hamiltonian:

$$H = h + V$$

$$= \sum_{ij} h_{ij} a_i^\dagger a_j + \frac{1}{2} \sum_{ijkl} V_{ijkl} a_i^\dagger a_k^\dagger a_l a_j,$$



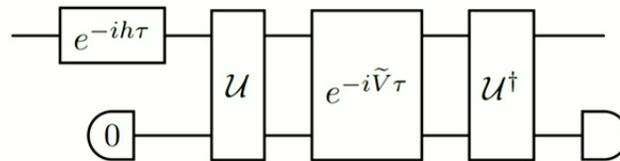
add N-M fictitious modes:

$$c_\alpha = \sum_i u_{i\alpha} a_i + \sum_m v_{m\alpha} b_m$$

New interaction Hamiltonian:

$$\tilde{V} = \sum_{\alpha \neq \beta} \tilde{V}_{\alpha\beta} n_\alpha n_\beta$$

$$V = {}_b \langle 0 | \tilde{V} | 0 \rangle_b$$



# Efficient preparation of MPS

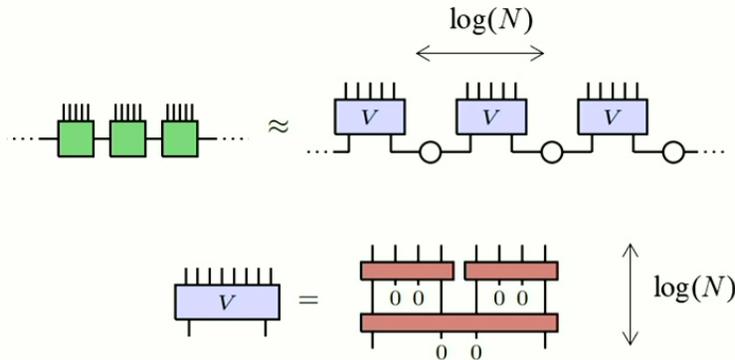
## Preparation of matrix product states with log-depth quantum circuits

Daniel Malz,<sup>1,\*</sup> Georgios Styliaris,<sup>2,3,\*</sup> Zhi-Yuan Wei,<sup>2,3,\*</sup> and J. Ignacio Cirac<sup>2,3</sup>

<sup>1</sup>Department of Mathematical Sciences, University of Copenhagen,  
Universitetsparken 5, 2200 Copenhagen, Denmark

<sup>2</sup>Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Str. 1, 85748 Garching, Germany

<sup>3</sup>Munich Center for Quantum Science and Technology (MCQST), Schellingstr. 4, 80799 München, Germany  
(Dated: July 6, 2023)



Circuit depth:  $T = O(\log N)$

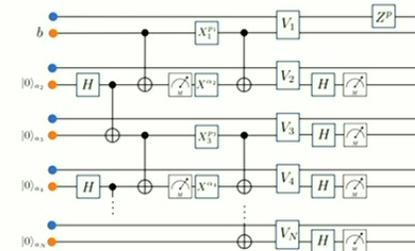
## Approximating many-body quantum states with quantum circuits and measurements

Lorenzo Piroli,<sup>1</sup> Georgios Styliaris,<sup>2,3</sup> and J. Ignacio Cirac<sup>2,3</sup>

<sup>1</sup>Dipartimento di Fisica e Astronomia, Università di Bologna and INFN,  
Sezione di Bologna, via Irnerio 46, I-40126 Bologna, Italy

<sup>2</sup>Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Str. 1, 85748 Garching, Germany

<sup>3</sup>Munich Center for Quantum Science and Technology (MCQST), Schellingstr. 4, D-80799 München, Germany



# Tensor Network Theory

## Matrix-product unitaries: Beyond quantum cellular automata

Georgios Styliaris,<sup>1,2</sup> Rahul Trivedi,<sup>3,1,2</sup> David Pérez-García,<sup>4,5</sup> and J. Ignacio Cirac<sup>1,2</sup>

<sup>1</sup>Max Planck Institute of Quantum Optics, Hans-Kopfermann-Str. 1, Garching 85748, Germany

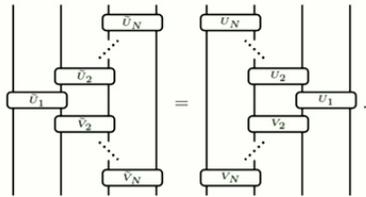
<sup>2</sup>Munich Center for Quantum Science and Technology (MCQST), Schellingstr. 4, 80799 München, Germany

<sup>3</sup>Electrical and Computer Engineering, University of Washington, Seattle, Washington 98195, USA

<sup>4</sup>Departamento de Análisis Matemático, Universidad Complutense de Madrid, 28040 Madrid, Spain

<sup>5</sup>Instituto de Ciencias Matemáticas (CSIC-UAM-UC3M-UCM), 28049 Madrid, Spain

(Dated: June 17, 2024)

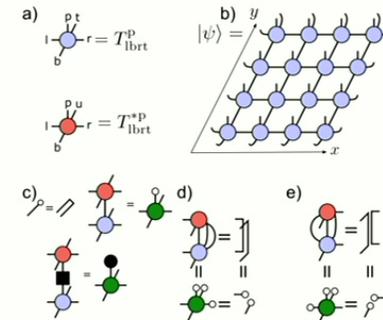


## Dual-isometric Projected Entangled Pair States

Xie-Hang Yu,<sup>1,2</sup> J. Ignacio Cirac,<sup>1,2</sup> Pavel Kos,<sup>1,2,\*</sup> and Georgios Styliaris<sup>1,2,\*</sup>

<sup>1</sup>Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Str. 1, 85748 Garching, Germany

<sup>2</sup>Munich Center for Quantum Science and Technology (MCQST), Schellingstr. 4, 80799 München, Germany



## Regular language quantum states

Marta Florido-Llinàs,<sup>1,2,\*</sup> Álvaro M. Alhambra,<sup>3,†</sup> David Pérez-García,<sup>4,‡</sup> and J. Ignacio Cirac<sup>1,2,§</sup>

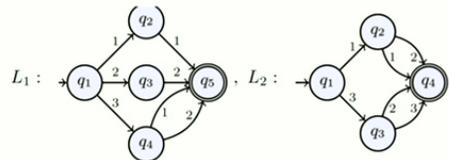
<sup>1</sup>Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Str. 1, 85748 Garching, Germany

<sup>2</sup>Munich Center for Quantum Science and Technology (MCQST), Schellingstr. 4, 80799 München, Germany

<sup>3</sup>Instituto de Física Teórica UAM/CSIC, C/ Nicolás Cabrera 13-15, Cantoblanco, 28049 Madrid, Spain

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(Dated: July 26, 2024)



$$|L_N\rangle := \left( \bigcirc_{l_1} \right) \left[ \begin{array}{c} | \\ \text{A} \\ | \end{array} \right] \left[ \begin{array}{c} | \\ \text{A} \\ | \end{array} \right] \dots \left[ \begin{array}{c} | \\ \text{A} \\ | \end{array} \right] \left( \bigcirc_{r_1} \right),$$

where the bond dimension is  $D = |Q|$ , and

$$\left\{ \begin{array}{l} \bigcirc_{l_1} := \sum_{i \in I} |i\rangle, \\ \bigcirc_{r_1} := \sum_{f \in F} |f\rangle, \end{array} \right. \quad \left[ \begin{array}{c} x \\ | \\ \text{A} \\ | \\ j \end{array} \right] = \begin{cases} 1 & \text{if } j \in \delta(i, x), \\ 0 & \text{otherwise.} \end{cases}$$



A. Franco



G. Styliaris



R. Trivedi

**Collaborators:** Demler (Zürich), Lukin (Harvard), Perez (Madrid), Polzik (Copenhagen), Schuch (Vienna), Shi (Beijing), Vuckovic (Stanford), Verstraete (Cambridge), Zoller (Innsbruck)

