

**Title:** Higher-Spin Charges in Gravity

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**Abstract:**

I will describe shortly what the higher-spin charges are and why they are relevant from a holographic point of view. Besides, I will emphasize the recent result according to which they are realized as Noether charges in a non-linear regime.

# CANONICAL REALIZATION OF HIGHER SPIN SYMMETRIES

Lightning Talk  
Celestial Holography Summer School 2024

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## MOTIVATIONS

- ▶ What are the *symmetries* of the  $\mathcal{S}$ -matrix? [Strominger, Narayanan, Himwich]
- ▶ To which extent do they constraint scattering amplitudes with soft particles? *Soft Theorems*. [Strominger, Pate, Raclariu, Pasterski, Ball]
- ▶  $w_{1+\infty}$  as symmetry of truncated EE. [Freidel, Pranzetti, Raclariu, Geiller]
- ▶  $w_{1+\infty}$  as residual gauge symmetry of twistor space. [Mason, Sharma, Adamo, Ruzziconi]
- ▶ (How) is the  $w_{1+\infty}$  symmetry deformed by *interactions* (higher  $G_N$  corrections)?

### Our Main Result

There is a non-perturbative canonical realization of a 'beyond the wedge'  $w_{1+\infty}$  algebroid onto the Ashtekar-Streubel phase space  $\mathcal{P}$ .

- ▶ AS symplectic form:

$$\Omega_{\text{AS}} = \frac{1}{4\pi G} \int_{\mathcal{I}} \delta \bar{N} \wedge \delta C$$

## WHAT'S KNOWN

- ▶  $w_{1+\infty}$  was identified as a symmetry algebra constraining *OPE* between soft gravitons. [Strominger and al.]
- ▶ A collection of Soft and Hard *higher spin charges* were then constructed on  $\mathcal{P}$ . [Freidel, Pranzetti, Raclariu, Geiller] Components in a  $1/r$  exp. of Weyl tensor
- ▶ They satisfy the SDGR Bianchi id.:  $\partial_u Q_s = DQ_{s-1} + (s+1)CQ_{s-2}$
- ▶ Their Hard *action* on soft gravitons was matched with OPE.
- ▶ These charges form a  $w_{1+\infty}$  *algebra* defined by the bracket

$$\{Q(T_s), Q(T'_{s'})\}^{(1)} = Q^{(1)} \left( \underbrace{(s+1)T_s DT'_{s'} - (s'+1)T'_{s'} DT_s}_{[T_s, T'_{s'}]_{s+s'-1}^W} \right) \quad (1)$$

under 2 restrictions:

- Truncation in  $G_N$
- Wedge condition:  $D^{s+2}T_s = 0$  Think about an extension of gbm's to all spins

- ▶ The W-bracket is *not* a Lie bracket for all  $s$  *outside* of the wedge

## BEYOND THE WEDGE

- ▶ Action of *time* and *field dependent* symmetry parameters  $\tau_s(u, z, \bar{z})$  on  $\mathcal{P}$ :

$$\delta_\tau C = -D^2\tau_0 + N\tau_0 + 2DC\tau_1 + 3CD\tau_1 - 3C^2\tau_2 \quad (2)$$

↳ *Only 3* parameters!

- ▶  $\delta_\tau$  is a *Lie algebroid action*,

$$[\delta_\tau, \delta_{\tau'}]C = -\delta_{[[\tau, \tau']]}C \quad (3)$$

when  $\tau_s$  satisfy a set of dual EOM:

$$\partial_u \tau_s = D\tau_{s+1} - (s+3)C\tau_{s+2}$$

- ▶ *New* bracket

$$\tau_s|_{u=0} = T_s(z, \bar{z})$$

$$[[\tau, \tau']]_s := [\tau, \tau']_s + \underbrace{(\delta_{\tau'}\tau)_s - (\delta_\tau\tau')_s}_{\text{algebroid/Barnich-Troessaert}} \quad s \geq -1, \quad (4)$$

$$[\tau, \tau']_s = \underbrace{\sum_{n=0}^{s+1} (n+1)(\tau_n D\tau'_{s+1-n} - \tau'_n D\tau_{s+1-n})}_{\text{W-bracket structure}} - \underbrace{(s+3)C(\tau_0\tau'_{s+2} - \tau'_0\tau_{s+2})}_{\text{field dependency}}$$

## NOETHER REPRESENTATION

► Defining

$$Q_\tau = \sum_{s=-1}^{\infty} \int_S Q_s \tau_s \quad (5)$$

↳ Plays the role of a *generating* functional!

we showed that

$$Q_\tau = \frac{1}{4\pi G} \int_{\mathcal{I}} \bar{N} \delta_\tau C \quad (6)$$

### Result

$Q_\tau$  is a *Noether* charge and generates a *beyond* the *wedge* and *non-perturbative*  $w_{1+\infty}$  algebroid on  $\mathcal{P}$ .

$$\begin{aligned} I_{\delta_\tau} \Omega_{AS} &= -\delta Q_\tau & \{Q_\tau, O\} &= \delta_\tau O \\ Q_\tau &= I_{\delta_\tau} \Theta_{AS} & \{Q_\tau, Q_{\tau'}\} &= -Q_{[[\tau, \tau']]} \end{aligned} \quad (7)$$

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