

**Title:** Non-linear gravitational waves in Horndeski gravity

**Speakers:** Hugo Roussille

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**Abstract:**

The non-linear dynamics of gravitational wave propagation in spacetime can contain drastic new phenomenology that is absent from the linearised theory. In this talk, I will probe the non-linear radiative regime of Horndeski gravity by making use of disformal field redefinition. I will discuss how disformal transformations alter the properties of congruences of geodesics and in particular how they can generate disformal gravitational waves at the fully non-linear level. I will illustrate this effect by presenting a new exact radiative solution in Horndeski gravity describing a scalar pulse. Analysing the non-linear dynamics of this new radiative solution will show that it contains tensorial gravitational waves generated by a purely time-dependent scalar monopole. This intriguing result is made possible by the higher-order nature of Horndeski gravity.

# Non-linear gravitational waves in Horndeski gravity

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Hugo Roussille, Jibril Ben Achour, Mohammad Ali Gorji

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50 years of Horndeski Gravity — Perimeter Institute and University of Waterloo



## Introduction

- Non-linear gravitational wave solutions known in GR (pp-waves, Kundt, Robinson-Trautman) [Robinson, Trautman '60; Ehlers, Kundt '62]
- Not directly useful for comparison to observations but crucial to explore non-linear radiative regime of a theory (ex: non-linear memory effects [Christodoulou '91])
- This work: investigate phenomenology of scalar-tensor mixing on non-linear gravitational waves in modified gravity
- Use an exact solution of a Horndeski theory as a toy model
- Take advantage of disformal transformations as solution-generating techniques

1



## Disformal transformations of scalar-tensor theories

### Disformal transformations

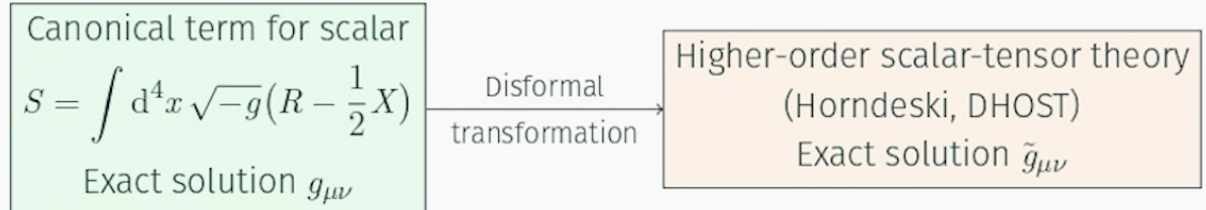
Generalization of conformal transformations [Bekenstein '93]

$$\tilde{g}_{\mu\nu} = A(\phi, X)g_{\mu\nu} + B(\phi, X)\phi_\mu\phi_\nu \quad (X = \phi_\mu\phi^\mu)$$

- In vacuum: field redefinition  $S[g_{\mu\nu}, \phi] \iff \tilde{S}[\tilde{g}_{\mu\nu}, \phi]$
- With matter:  $S[g_{\mu\nu}, \phi] + S_{\text{matter}}[g_{\mu\nu}] \not\iff \tilde{S}[\tilde{g}_{\mu\nu}, \phi] + S_{\text{matter}}[\tilde{g}_{\mu\nu}]$
- Horndeski theories: stable under  $A(\phi), B(\phi)$



## Solution-generating technique



→ use disformal transformations as a solution-generating technique [Anson, Babichev+ '21; Ben Achour, Liu+ '20; Ben Achour, Liu+ '20; Faraoni, Leblanc '21]

- Simplest transformation:  $\tilde{g}_{\mu\nu} = g_{\mu\nu} + B_0 \phi_\mu \phi_\nu$
- Obtain a solution of Horndeski with  $\tilde{F}_2(\tilde{X}) = \frac{1}{\sqrt{1 - B_0 \tilde{X}}}$



## Non-linear GW in GR

### Usual linearised setup

- Add a perturbation to a background:  $g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$
- Obtain a wave propagation equation  $\bar{\square} h_{\mu\nu} = 0$

→ sufficient for most cases but missing complex dynamics of the theory:  
non-linear memory effects [Christodoulou '91], colliding waves [Szekeres '70], solitons...

### Simplest non-linear example: pp-wave

[Brinkmann '25; Bondi '57; Robinson, Trautman '60; Penrose '65]

$$ds^2 = -H_{ab} x^a x^b du^2 + 2 du dv + \delta_{ab} dx^a dx^b$$

- describes propagation of a plane wave in vacuum along  $\partial_v$
- property close to linearised waves: profiles  $H_{ab}$  can be added
- can define polarisations through components of  $H_{ab}$





## Metric element

Solution of Einstein-scalar theory [Tahamtan, Svitek '15; Tahamtan, Svitek '16]

$$S = \int d^4x \sqrt{-g} \left( R - \frac{1}{2} \nabla_\mu \phi \nabla^\mu \phi \right)$$

$$ds^2 = -K(x, y) dw^2 - 2 dw d\rho + \frac{\rho^2 - \chi(w)^2}{P(x, y)^2} (dx^2 + dy^2)$$

lightlike  
coordinate

$$\phi = \frac{1}{\sqrt{2}} \log \left( \frac{\rho - \chi(w)}{\rho + \chi(w)} \right)$$

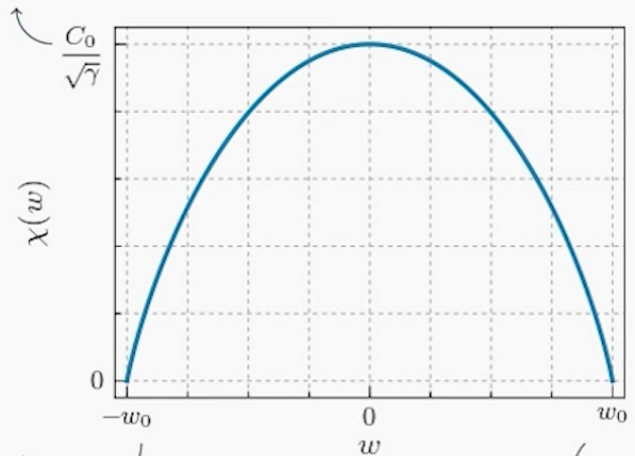
wave  
pulse

- Wave propagation towards outgoing  $\rho$
- Not spherically symmetric
- Presence of an apparent horizon
- Fully non-linear solution
- Petrov type II



## Representing the wave pulse

scalar charge



null infinity  
(past)

null infinity  
(future)

- Curvature of 2D space  $(\partial_x, \partial_y)$ :

$$\mathcal{K} = \frac{\chi^2(w)}{C_0^2} K(x, y)$$

- Scalar pulse  $\chi$  goes from 0 to max and back to 0
- Longitudinal wave generated by scalar field monopole
- Empty spacetime at remote past and future

6





## Description of the scalar-tensor solution

$$ds^2 = (-K(x, y) + B_0 \phi_w^2) dw^2 - 2(1 - B_0 \phi_w \phi_\rho) dw d\rho + B_0 \phi_\rho^2 d\rho^2 + \frac{\rho^2 - \chi(w)^2}{P(x, y)^2} (dx^2 + dy^2)$$

- Wave pulse  $\chi$  unchanged: scalar monopole
- Apparent horizon and singularities unchanged qualitatively
- Remote past and future ( $w \rightarrow \pm w_0$ ): empty non-spherical spacetime
- Petrov classification: type I while seed was type II  $\rightarrow$  loss of algebraic speciality



## GW content of spacetime

How can one read the polarisation content of a GW?

Linearised GW

Read off from the components of  $h_{\mu\nu}$

→

Non-linear GW

No extraction of wave profile!

- **Main idea:** tidal effects experienced by a photon around its worldline  $\bar{\gamma}$  [Penrose '76], with parallel transported null tetrad  $E_A^\mu$  and parameter  $W$
- In this setup, always recover pp-wave geometry:

$$ds^2 = 2dWdV + \delta_{AB}dX^A dX^B - H_{AB}X^A X^B dW^2$$

- Read off polarisation from components of  $H_{AB} = \bar{R}_{\mu\nu\rho\sigma} E_W^\mu E_A^\nu E_W^\rho E_B^\sigma$   
↪ evaluated on  $\bar{\gamma}$

8

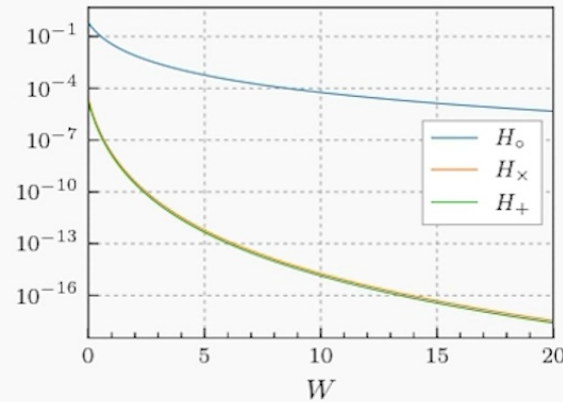


## Waves in the transformed solution

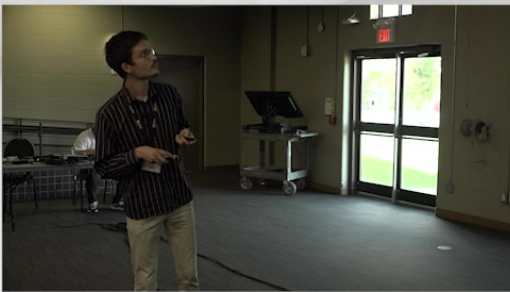
$$H_{AB} = \begin{pmatrix} a_0 + B_0 a_1 + B_0^2 a_2 & 0 \\ 0 & a_0 + B_0 a_1 + B_0^2 a_2 \end{pmatrix} + B_0^2 \begin{pmatrix} +b_2 & c_2 \\ c_2 & -b_2 \end{pmatrix} + \mathcal{O}(B_0^3)$$

$\downarrow$   
 scalar waves  $H_0$

$\rightarrow H_\times$   
 $\leftarrow H_+$



- the disformal transformation sources *tensorial* gravitational waves
- keep the complete non-perturbative character of the metric



## Conclusion

- First **exact radiative non-linear solution** in Horndeski beyond plane waves
- Contains non-linear superposition of shear and breathing modes generated by a scalar monopole
- Probe new effects in strong regime scalar-tensor gravity: **additional contribution to GWs**
- Open question: consequence for GWs in the case of scalar-tensor cosmology?

10



Thank you for your attention!

