

Title: Particle Physics Seminar - TBA

Speakers: David Kaplan

Series: Particle Physics

Date: May 28, 2024 - 1:00 PM

URL: <https://pirsa.org/24050095>

Abstract: Abstract TBA

Zoom link

Non-Linear Quantum Mechanics

David E. Kaplan

w/ S Rajendran — *PRD* (arxiv:2106.10576)

w/ M Raizen and S Rajendran — *PLB* (arxiv:2203.00195)

w/ M Polkovnikov, et al — *PRL* (arxiv:2204.11875)

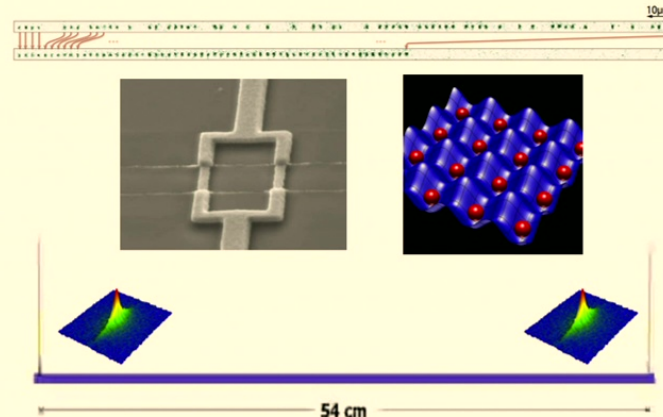
w/ Broz, et al — *PRL* (arxiv:2206.12976)

Modify Quantum Mechanics?

Why(nberg-ian) not?



1927



2024

QM is the only known physical theory that is exactly linear

$$i\partial_t|\chi\rangle = \hat{H}|\chi\rangle$$

Quantum Field Theory

The Schrödinger Equation (Schrödinger picture)

$$i \partial_t |\chi\rangle = \hat{H} |\chi\rangle$$

$$\text{with } \hat{H} = \int d^3x \hat{\mathcal{H}}(\mathbf{x})$$

The time evolution operator

$$\hat{U} = e^{-i\hat{H}t}$$

Already non-linear — how do we modify?
Add state-dependent terms.

Quantum Field Theory

'Non-linear' (state-dependent) terms

$$i \partial_t |\chi\rangle = \int d^3x \left(\hat{\mathcal{H}}(\mathbf{x}) + \epsilon \langle \chi | \hat{\mathcal{O}}_1(\mathbf{x}) | \chi \rangle \hat{\mathcal{O}}_2(\mathbf{x}) \right) |\chi\rangle$$

Time evolution includes term that depends on the state itself

If \mathcal{O}_1 and \mathcal{O}_2 are Hermitian, the norm is constant

$$\partial_t \langle \chi | \chi \rangle = 0$$

Probabilistic interpretation of observables can be maintained

Quantum Field Theory

Path Integral Formulation

Can compute perturbatively — $\langle \chi | \hat{\mathcal{O}}_{1,2} | \chi \rangle$ first at $\epsilon \rightarrow 0$

Then insert as a background field

$$S = \int d^4x \mathcal{L}(\phi(x), \partial_\mu \phi(x)) + \epsilon \left(\langle \chi | \hat{\mathcal{O}}_1 | \chi \rangle(x) \mathcal{O}_2(\phi(x)) + \langle \chi | \hat{\mathcal{O}}_2 | \chi \rangle(x) \mathcal{O}_1(\phi(x)) \right)$$

$$\mathbb{Z} = \int \mathcal{D}\phi e^{i \int d^4x \mathcal{L}(\phi, \partial\phi) + \epsilon (\langle \mathcal{O}_1 \rangle \mathcal{O}_2(\phi) + \langle \mathcal{O}_2 \rangle \mathcal{O}_1(\phi))}$$

QFT Example

YUKAWA COUPLING

$$\hat{\mathcal{H}} \supset y\hat{\phi}\hat{\Psi}\hat{\Psi}$$

Add non-linearity

$$\hat{\mathcal{H}} \supset y\hat{\phi}\hat{\Psi}\hat{\Psi} + y\epsilon \left(\langle \chi | \hat{\phi} | \chi \rangle \hat{\Psi}\hat{\Psi} + \hat{\phi} \langle \chi | \hat{\Psi}\hat{\Psi} | \chi \rangle \right)$$

Non-Linear Quantum Mechanics

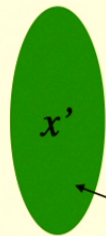
One-particle, non-relativistic limit

$$\mathcal{L} \supset y(\phi + \epsilon \langle \chi | \phi | \chi \rangle) \bar{\Psi} \Psi$$

To get NR theory for fermions Ψ , compute $\langle \phi \rangle$.

Will depend on initial conditions and sources. At zeroth order, Ψ sources ϕ :

ψ wave function for single fermion Ψ



$$\langle \phi \rangle(x) \supset \int d^4x' |\psi(x')|^2 G_R(x - x')$$

Charge density of Ψ

Causal Green's Function

Non-Linear Quantum Mechanics

One-particle, non-relativistic limit - the non-linear Schrödinger Equation

$$i \partial_t \psi(x) = \hat{H}(\mathbf{x}) \psi(x) + \epsilon \int d^4 x' |\psi(x')|^2 G_R(x' - x) \psi(x)$$

Causality guaranteed by the retarded Green's Function:

$$\square G_R(y; x) = \delta^4(x - y)$$

For a massless field:

$$G_R(x - x') = \frac{\delta(t' - (t + |\mathbf{x}' - \mathbf{x}|))}{|\mathbf{x}' - \mathbf{x}|}$$

Entangled Systems and Causality

One-particle, non-relativistic limit

Wave function for two entangled particles:

$$\psi(\mathbf{x}, \mathbf{y}; t) = \sum_{i,j} c_{ij}(t) \psi_i^I(\mathbf{x}) \psi_j^II(\mathbf{y})$$

$$\delta\hat{H}\psi = \int d^3x' d^3y' dt' |\psi(\mathbf{x}', \mathbf{y}'; t')|^2 \left(\boxed{G_R(\mathbf{x}, t; \mathbf{x}', t')} + \underbrace{G_R(\mathbf{y}, t; \mathbf{x}', t') + G_R(\mathbf{x}, t; \mathbf{y}', t')}_{\text{causal}} + \boxed{G_R(\mathbf{y}, t; \mathbf{y}', t')} \right) \psi(\mathbf{x}, \mathbf{y}, t)$$

Can a local action at \mathbf{x} affect the particle at \mathbf{y} ?

$$\delta\hat{H} \subset \int \underbrace{d^3x' d^3y' dt'}_{\text{causal}} |\psi(\mathbf{x}', \mathbf{y}'; t')|^2 G_R(\mathbf{y}, t; \mathbf{y}', t')$$

e.g., need this d^3x' integral unchanged

Entangled Systems and Causality

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$$\delta\hat{H} \subset \int \underbrace{d^3x' d^3y' dt'}_{\text{e.g., need this } d^3x' \text{ integral unchanged}} |\psi(\mathbf{x}', \mathbf{y}'; t')|^2 G_R(\mathbf{y}, t; \mathbf{y}', t')$$

e.g., need this d^3x' integral unchanged

Unchanged if $\psi(\mathbf{x}, \mathbf{y}, t) \rightarrow (\hat{U}_1 \otimes \hat{\mathbf{1}}) \psi(\mathbf{x}, \mathbf{y}, t)$. This is true in linear QM (zeroth order in ϵ)

First order correction (time-dep. p.t.) to particle I potential — use zero-order w.f.s

$$V_1(\mathbf{x}, t) \subset \int d^3x' d^3y' dt' |\psi_0(\mathbf{x}', \mathbf{y}'; t')|^2 G_R(\mathbf{x}, t; \mathbf{x}', t')$$

Similarly factorized — local and unitary. Continue to all orders.

Entangled Systems and Causality

One-particle, non-relativistic limit

Wave function for two entangled particles:

$$\psi(\mathbf{x}, \mathbf{y}; t) = \sum_{i,j} c_{ij}(t) \psi_i^I(\mathbf{x}) \psi_j^II(\mathbf{y})$$

$$\delta \hat{H} \psi = \int d^3x' d^3y' dt' |\psi(\mathbf{x}', \mathbf{y}'; t')|^2 \left(\boxed{G_R(\mathbf{x}, t; \mathbf{x}', t')} + \underbrace{G_R(\mathbf{y}, t; \mathbf{x}', t') + G_R(\mathbf{x}, t; \mathbf{y}', t')}_{\text{causal}} + \boxed{G_R(\mathbf{y}, t; \mathbf{y}', t')} \right) \psi(\mathbf{x}, \mathbf{y}, t)$$

SM Examples

QED

Add non-linearity to vector: $\mathcal{H} \supset e \left(A_\mu + \epsilon \langle \chi | A_\mu | \chi \rangle \right) J^\mu$
(Defined in a gauge, *e.g.*, Weyl)

Or field strength: $F_{\mu\nu} F^{\mu\nu} \rightarrow F_{\mu\nu} F^{\mu\nu} + \epsilon F_{\mu\nu} \langle F^{\mu\nu} \rangle$

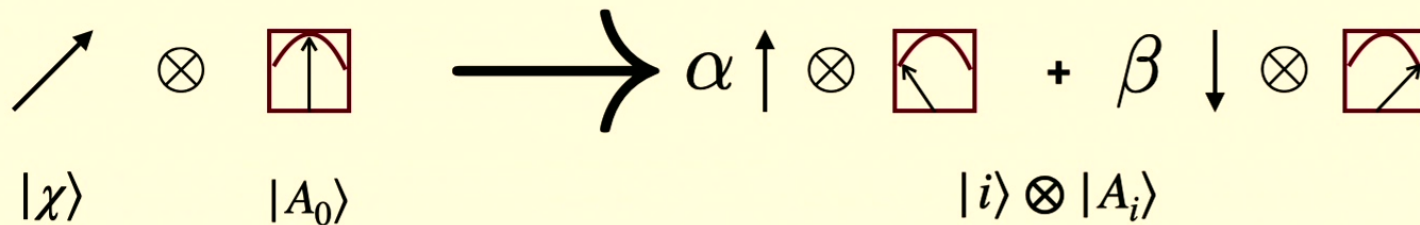
GRAVITY

Replace, *e.g.*, $g_{\mu\nu} \rightarrow \frac{g_{\mu\nu} + \epsilon \langle g_{\mu\nu} \rangle}{1 + \epsilon}$ in interaction terms

Full GCI can be maintained

Measurement in Quantum Mechanics

Time evolution with interaction between the system and measuring device



$$|\chi\rangle \otimes |A_0\rangle \rightarrow \sum_i c_i |i\rangle \otimes |A_i\rangle$$

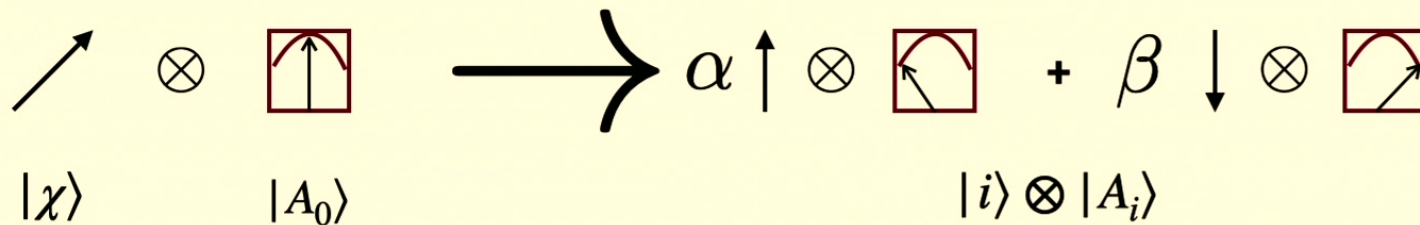
Prediction of Quantum Mechanics (“Many Worlds”)

Pick: $\langle A_j | A_i \rangle = \delta_{ij}$

Trace over device $\longrightarrow \rho_{|\chi\rangle} = \sum_i c_i c_i^* |i\rangle \langle i|$ ‘Probabilistic’ outcome

Measurement in Quantum Mechanics

Time evolution with interaction between the system and measuring device

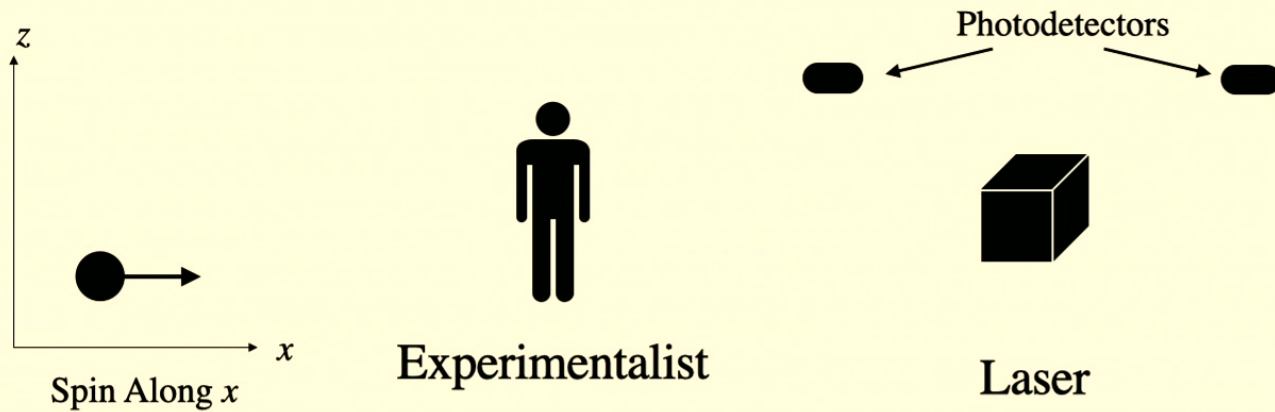


In NLQM, stationary states generally **not** orthogonal —
the effective Hamiltonian depends on the initial state of the system

No Guarantee: $\langle A_j | A_i \rangle = 0$ for $i \neq j$

$$|\Psi\rangle \otimes |A_0\rangle \rightarrow \sum_i c_i |i\rangle \otimes |A_i\rangle + \epsilon \sum_{i,j} d_{i,j} |i\rangle \otimes |A_j\rangle \quad \text{'Measurement noise'}$$

Linear Quantum Mechanics



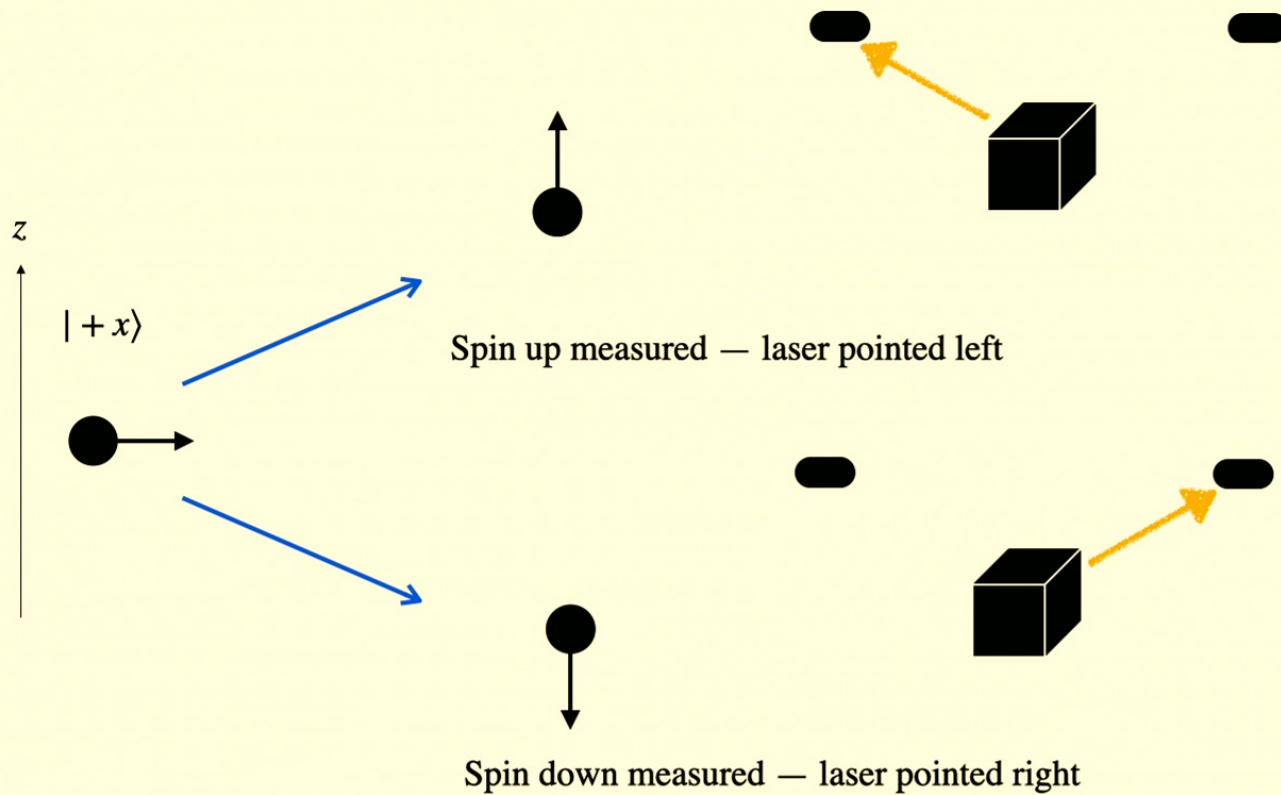
Initial State : $|\chi(0)\rangle$

Represents Full Quantum State (spin, experimentalist...)

Goal: Create Macroscopic Superposition

Method: Measure spin along z .
Depending upon outcome, send laser along different directions

Linear Quantum Mechanics

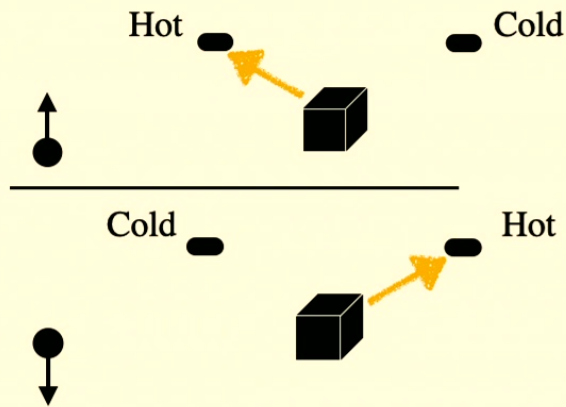


$$|\chi(t)\rangle = \frac{1}{\sqrt{2}} (|+z\rangle|L\rangle|Env_L\rangle + |-z\rangle|R\rangle|Env_R\rangle)$$

Linear Quantum Mechanics

Which photodetectors light up?

$$\mathcal{H} \supset eA_\mu J^\mu$$



Transition Matrix Elements

$$\langle +z | \langle L | \langle Env_L | A_\mu(x_L) J^\mu(x_L) | +z \rangle | L \rangle | Env_L \rangle \neq 0$$

$$\langle +z | \langle L | \langle Env_L | A_\mu(x_R) J^\mu(x_R) | +z \rangle | L \rangle | Env_L \rangle = 0$$



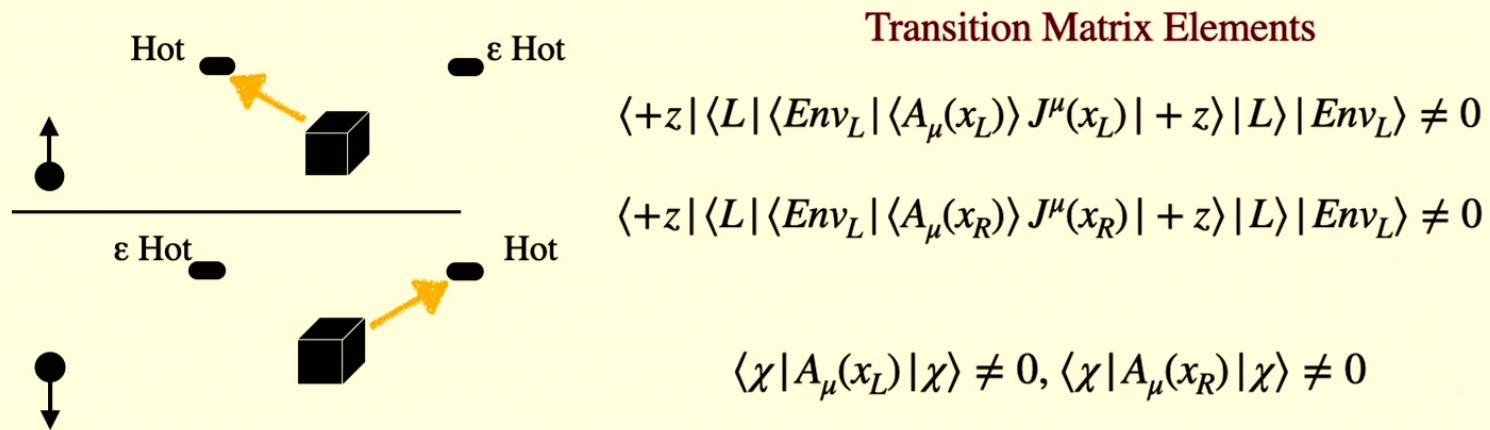
$$\langle L | A_\mu(x_R) | L \rangle = 0$$

$$|\chi(t)\rangle = \frac{1}{\sqrt{2}} (|+z\rangle |L\rangle |Env_L\rangle + |-z\rangle |R\rangle |Env_R\rangle)$$

Non-Linear Quantum Mechanics

Which photodetectors light up?

$$\mathcal{H} \supset eA_\mu J^\mu + e\varepsilon \langle A_\mu \rangle J^\mu$$



Communication between “worlds”

Non-linearity visible despite Environmental De-coherence!

Polchinski: “Everett Phone”

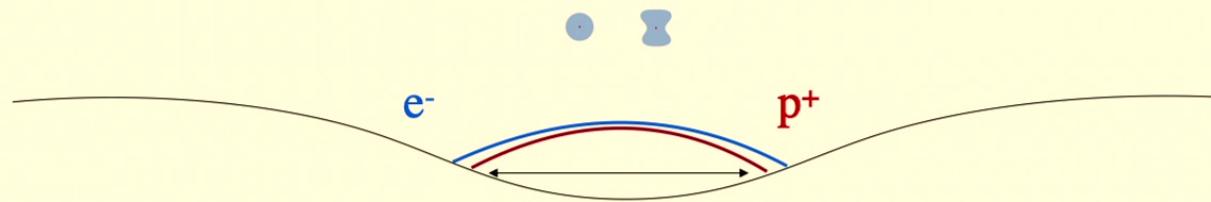
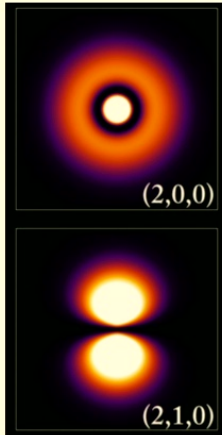
Constraints and Tests: Part I

Constraints: Atomic Levels

What does this do to the Lamb Shift?

Say charged particles see their own w.f.:

$$i \frac{\partial}{\partial t} \psi(x) = H(\mathbf{x}) \psi(x) + \epsilon_\gamma \alpha \int d^4 x' |\psi(x')|^2 G_R(x' - x) \psi(x)$$



Electron spread over the trap (micron) dilutes the electric field and thus the level splitting

Proton's wave function also produces a field that nearly cancels the electron wave function.

Key — center of mass coordinate cannot be separated from relative coordinate due to locality.

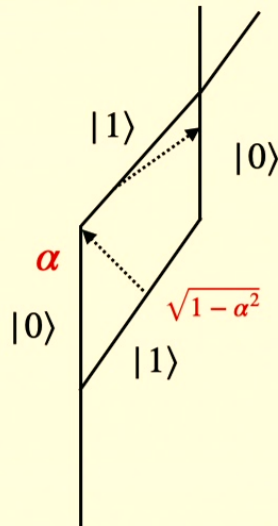
$$\epsilon_\gamma < 10^{-2}$$

Experimental Tests: Local

Interferometry - interaction between paths

Take an ion - split its wave-function

Coulomb Field of one path interacts with the other path



$$|\psi\rangle = \alpha |0\rangle + \sqrt{1 - \alpha^2} |1\rangle$$

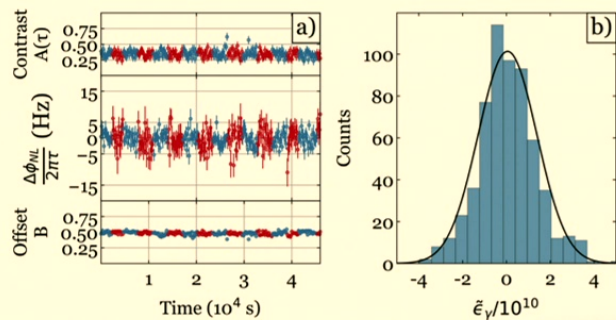
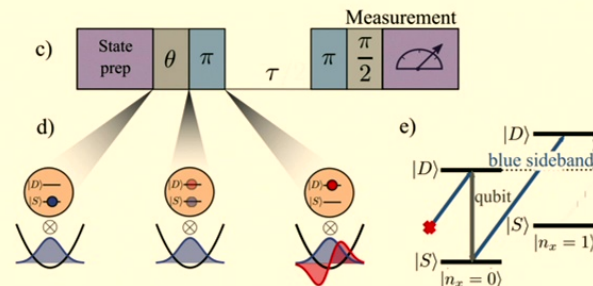
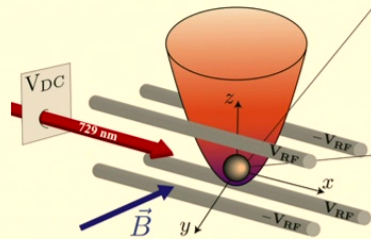
Gives rise to phase shift that depends on the intensity α^2 of the split

Use intensity dependence to combat systematics

Test ... Using a Vibrational Mode of a Trapped Ion

J. Broz, et al. (2022)

State: $|\psi\rangle = a|0\rangle + b|1\rangle$

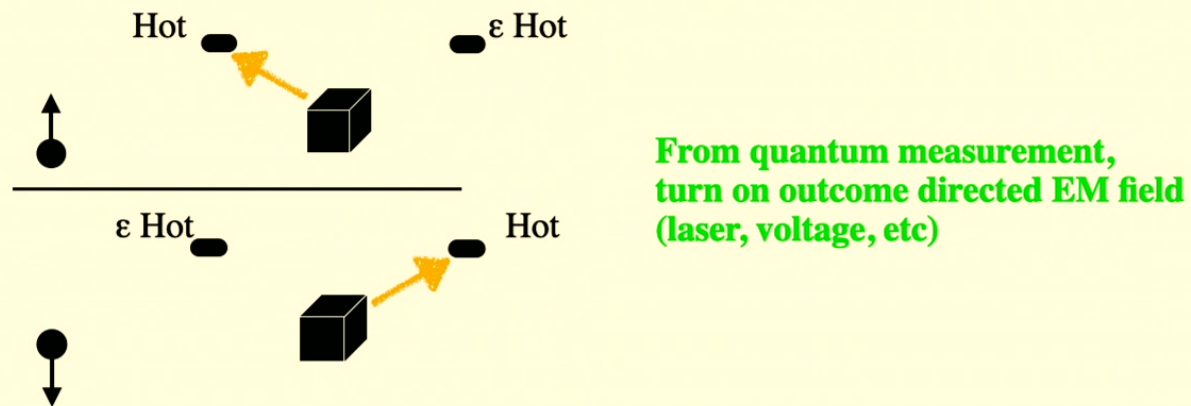


$$\text{Non-linear phase: } \phi_{NL} = \epsilon_{\gamma} \alpha \frac{10a^2 + b^2}{30\sqrt{2\pi\hbar}x_0} \tau$$

$$\epsilon_{\gamma} = 5 \pm 5.4 \times 10^{-12}$$

Experimental Tests: Macroscopic

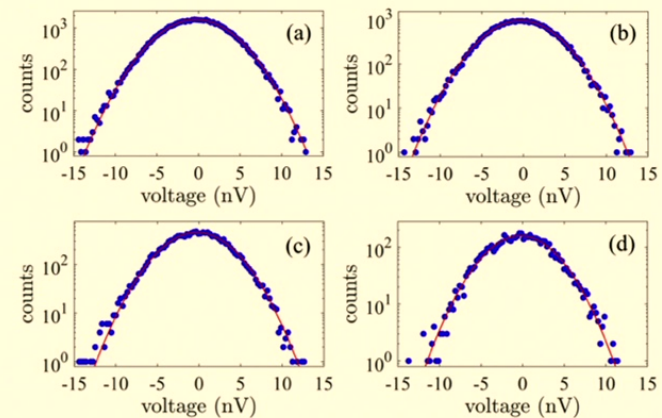
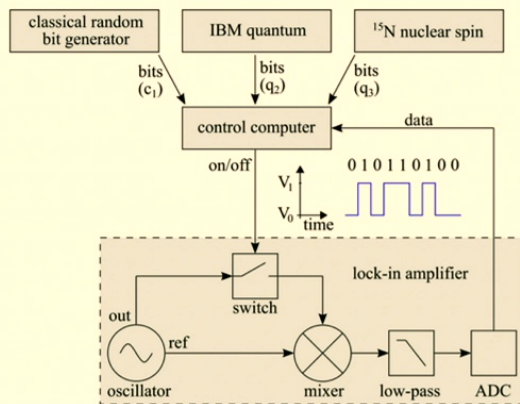
Create macroscopic superposition
Create Expectation value of EM/Gravity
Search for Expectation value



Similar tests for gravity

Experimental limit on non-linear QM using a voltmeter and quantum bits

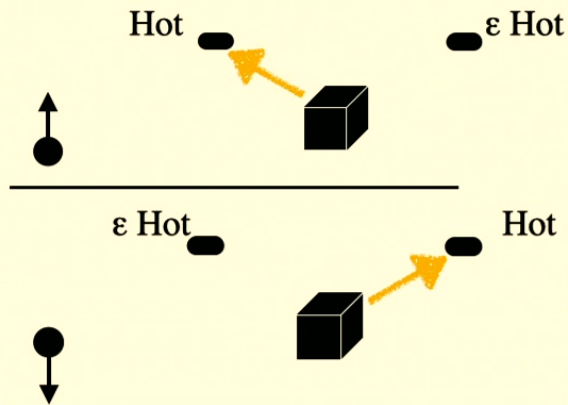
M. Polkovnikov, et al (2022)



$$|\epsilon_\gamma| \leq 4.7 \times 10^{-11}$$

Delicate Non-Linearity

O performs the laser experiment Today
 - discovers non-linear quantum mechanics!



$$|\chi\rangle = \frac{1}{\sqrt{2}} (|L\rangle|O_L\rangle + |R\rangle|O_R\rangle)$$

Now O wants to repeat experiment

Suppose $|O_U\rangle$ decides to run experiment at 9am on October 2
 But $|O_D\rangle$ runs experiment on 9am on October 12

State at 9am on October 2

$$|\chi\rangle = \frac{1}{\sqrt{2}} \left(|L\rangle|O_L\rangle \frac{|L'\rangle|O'_L\rangle + |R'\rangle|O'_D\rangle}{\sqrt{2}} + |R\rangle|O_R\rangle \right)$$

Delicate Non-Linearity

State on 9am on October 2

Compare with State today

$$|\chi\rangle = \frac{1}{\sqrt{2}} \left(|L\rangle|O_L\rangle \frac{|L'\rangle|O'_L\rangle + |R'\rangle|O'_D\rangle}{\sqrt{2}} + |R\rangle|O_R\rangle \right)$$

$$|\chi\rangle = \frac{1}{\sqrt{2}} (|L\rangle|O_L\rangle + |R\rangle|O_R\rangle)$$

$$\langle L|\langle O_L|\langle L'|\langle O'_L|\langle A_\mu(x_R)\rangle J^\mu(x_R)|\chi(t = \text{Oct 2})\rangle = \frac{1}{2}\langle L|\langle O_L|\langle A_\mu(x_R)\rangle J^\mu(x_R)|\chi(t = \text{today})\rangle$$

Effect is 1/2 of prior effect!

But, full effect if O_U and O_D perform experiment at same time!

Quantum Pollution: Without adequate care, superpositions may diverge wildly, preventing exploitability. Not automatic
- but need careful protocols!

Delicate Non-Linearity

State on 9am on October 2

Compare with State today

$$|\chi\rangle = \frac{1}{\sqrt{2}} \left(|L\rangle|O_L\rangle \frac{|L'\rangle|O'_L\rangle + |R'\rangle|O'_D\rangle}{\sqrt{2}} + |R\rangle|O_R\rangle \right)$$

$$|\chi\rangle = \frac{1}{\sqrt{2}} (|L\rangle|O_L\rangle + |R\rangle|O_R\rangle)$$

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Effect is 1/2 of prior effect!

But, full effect if O_U and O_D perform experiment at same time!

Quantum Pollution: Without adequate care, superpositions may diverge wildly, preventing exploitability. Not automatic
- but need careful protocols!

Particles have been scattering for 13 billion years. Cosmological dilution?

What part of the wave function...

$$|\chi\rangle = \alpha |Us\rangle + \beta |Them\rangle$$

$$\mathcal{H} \supset eA_\mu J^\mu + e\epsilon \langle A_\mu \rangle J^\mu$$

$$|\chi\rangle = \alpha |Us\rangle + \beta |Them\rangle \rightarrow \langle \chi | A_\mu | \chi \rangle = |\alpha|^2 \langle U | A_\mu | U \rangle + |\beta|^2 \langle T | A_\mu | T \rangle$$

$$\langle \chi | A_\mu | \chi \rangle \rightarrow |\alpha|^2 \langle U | A_\mu | U \rangle$$

$$\epsilon \rightarrow \epsilon_{\text{eff}} \equiv \epsilon |\alpha|^2$$

For $\alpha \ll \beta$, the wave function is dominated by something we are not a part of.
Can't turn on coherent fields over there.

Local exploitability completely determined by unchangeable
initial conditions dramatic difference from linear QM

Classical World?

$$|U(t)\rangle = |\text{Earth}\rangle + \delta|\text{Black}\rangle \quad \text{Or} \quad |U(t)\rangle = \delta|\text{Earth}\rangle + |\text{Black}\rangle$$

Are there natural quantum amplifiers, for e.g. in chaotic systems?



$$\Delta x \sim 100\text{nm}$$

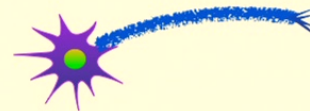
Changing classical evolution of a system requires coherent motion of N atoms

Probability that N atoms coherently move in some way: p^N

With $p \sim O(1)$ scattering probability

What about the weather?

What about my brain??



100's + ions to get one neuron to fire

$$|U(t)\rangle = |\text{Earth}\rangle + \delta|\text{Black}\rangle \quad \text{Reasonable}$$

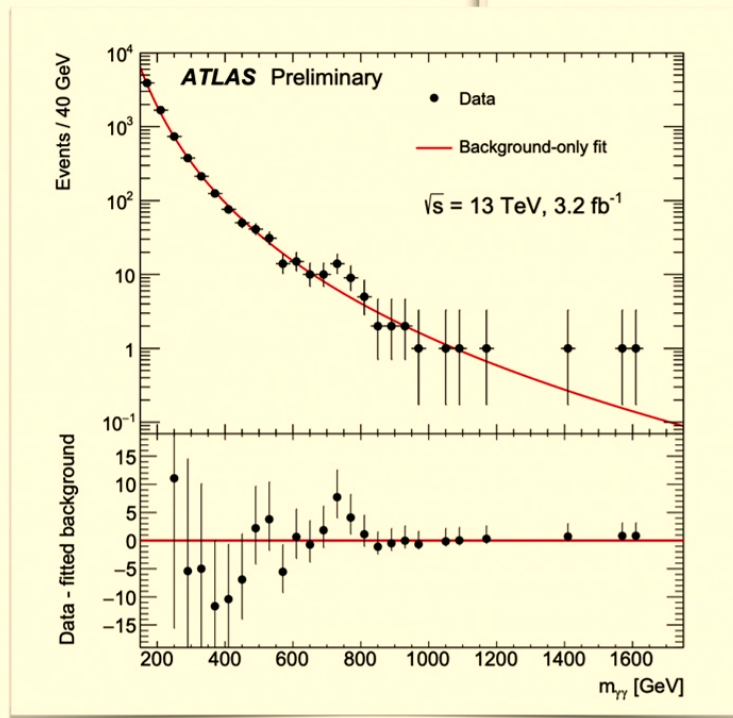
Quantum Amplifiers are Hard!

Natural Quantum Dilution

The '750 GeV' resonance!

Search for resonances decaying to photon pairs in 3.2 fb^{-1} of pp collisions at $\sqrt{s} = 13 \text{ TeV}$ with the ATLAS detector

Dec 15, 2015



INF-2015-081

ATLAS
server

claim

reference search 562 citations

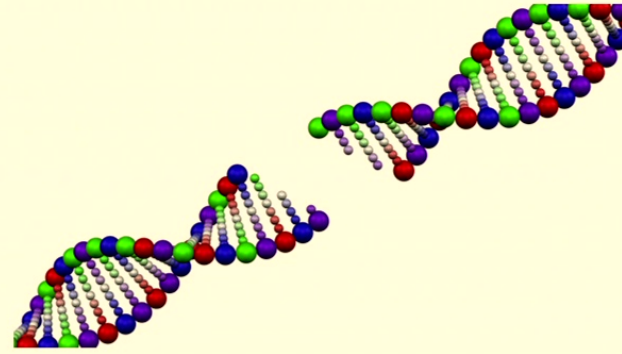
Have we been diluting our wave function on Earth for the past 100 years?

Evolutionary Dilution?

Is $N \sim O(\text{few})$ for evolution?

Maybe for RNA/DNA?

RNA formation?



Is Evolution a quantum amplifier?!?

$$|U(t=0)\rangle = \text{[Image of a rocky landscape]}$$



$$|U(t)\rangle = |\text{Earth}\rangle(|\text{Human}\rangle + |\text{No life}\rangle + \dots)$$

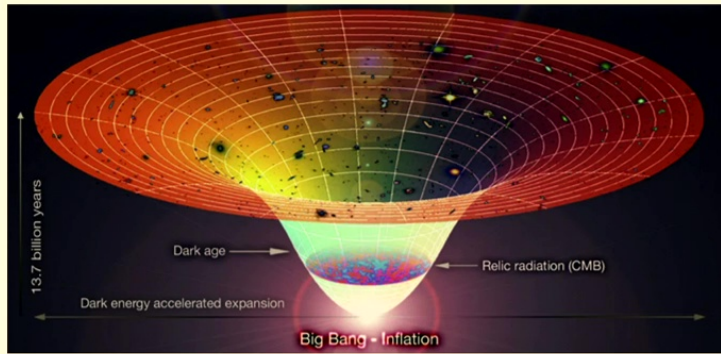
$$|U(t=0)\rangle = \text{[Image of a rocky landscape]}$$



$$|U(t)\rangle = |\text{Earth}\rangle(|\text{Human}\rangle + |\text{Alien}\rangle + \dots)$$

Cosmological Quantum Amplifier: Inflation

Standard cosmic inflation:
rapidly places quantum state in a homogenous and isotropic state
(Bunch-Davies Vacuum)



Homogeneous state becomes massive
superposition of statistically similar
Universes!

$$|\chi\rangle = \sum_i c_i |U_i\rangle, \quad c_i \sim e^{-N}$$

Most of the Universe: The space-time point where the
Earth is located is in intergalactic space

Probes of Non-Linear Quantum Mechanics

Non-Linear effects deeply tied to unchangeable initial and evolution of total quantum state

Probe not just non-linear quantum mechanics but full evolutionary history of quantum state!

Three Scenarios

Classical Universe (human dilution)

$$|U(t)\rangle = |\text{🌍}\rangle + \delta|\text{■}\rangle$$

Canonical Inflation

$$|U(t)\rangle = \delta|\text{🌍}\rangle + \delta|\text{■}\rangle$$

Quantum Biology

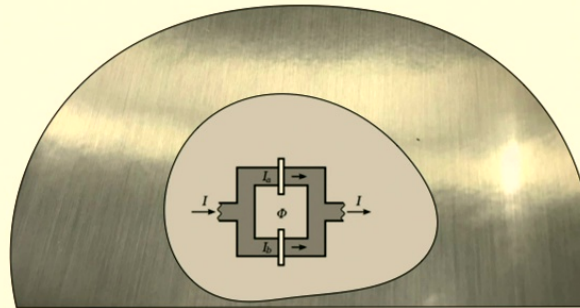
$$|U(t)\rangle = |\text{🌍}\rangle(\delta|\text{👤}\rangle + |\text{👽}\rangle + \dots)$$

Tests for a Quantum-Diluted Earth

Look for coherent fields turned on in all parts of the wavefunction

$$eJ^\mu(A_\mu + \epsilon_\gamma \langle A_\mu \rangle)$$

Remove an ancient magnetic source in our part of the wave function and measure the field in that region.



$$|U(t)\rangle = | \text{Earth} \rangle (\alpha | \text{missing magnetic material} \rangle + \beta | \text{magnetic material} \rangle)$$

Could do the same with gravitational sources on Earth (*e.g.*, Bikini Atoll)

Future Directions

1. Gravity for non-classical states
2. Black Hole horizon with interfering metrics
3. How does non-linear QM affect inflation?
4. Model building with $\langle h \rangle$, $\langle R \rangle$, ...

Conclusions

Remarkably easy to modify QM in a causal consistent way

Locality makes many past tests insensitive — **new probes required**

NL effects can be experimentally tested by
amplifying quantum measurements

Quantum amplification in the history of the universe
could suppress local non-linearities

If locally diluted, non-zero fields across the wave function could be
detected (ancient magnetic/gravitational fields, cosmological metric)

If NLQM is locally accessible, it will radically change
what we can do technologically

Thank you!