

Title: Energy and speed bound in GPTs - VIRTUAL

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Series: Quantum Foundations

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Abstract: Information-theoretic insights have proven fruitful in many areas of quantum physics. But can the fundamental dynamics of quantum systems be derived from purely information-theoretic principles, without resorting to Hilbert space structures such as unitary evolution and self-adjoint observables? Here we provide a model where the dynamics originates from a condition of informational non-equilibrium, the deviation of the system's state from a reference state associated to a field of identically prepared systems. Combining this idea with three basic information-theoretic principles, we derive a notion of energy that captures the main features of energy in quantum theory: it is observable, bounded from below, invariant under time-evolution, in one-to-one correspondence with the generator of the dynamics, and quantitatively related to the speed of state changes. Our results provide an information-theoretic reconstruction of the Mandelstam-Tamm bound on the speed of quantum evolutions, establishing a bridge between dynamical and information-theoretic notions.

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[Zoom link](#)



# Energy and speed bound in GPTs

L. Giannelli & G. Chiribella

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# Energy in Quantum Theory

- Energy is the expectation value of the Hamiltonian

$$\langle \psi | H | \psi \rangle = E$$

- Hamiltonian  $\equiv$  generator of reversible evolutions

$$i\hbar \frac{d}{dt} |\psi\rangle = H |\psi\rangle$$

- This dual role of the Hamiltonian is not an accident!  
Rather it is intrinsically connected with the quantum formalism [1]

[1] E. Grgin and A. Petersen, Journal of Mathematical Physics 15, 764 (1974)

# Energy in informational terms



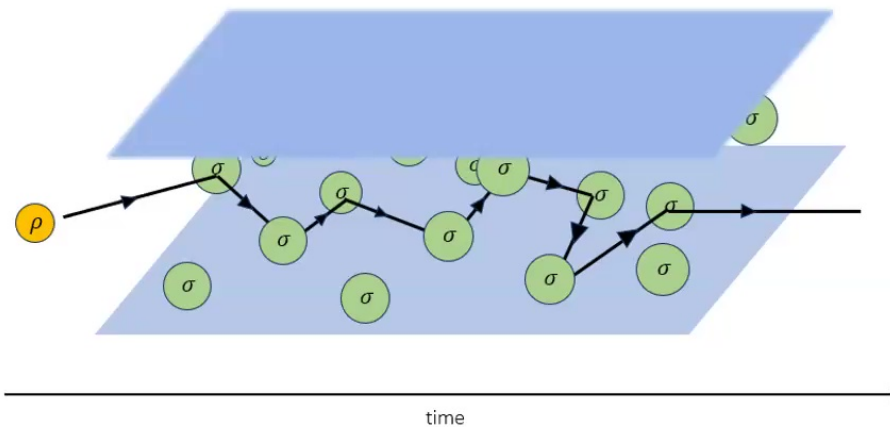
- Despite a central role in quantum mechanics, energy is often neglected in information-theoretic derivation
- Operational theories focus on statistical predictions rather than dynamics

Reconstruct quantum dynamics from  
informational principles



# An informational perspective on the dynamics

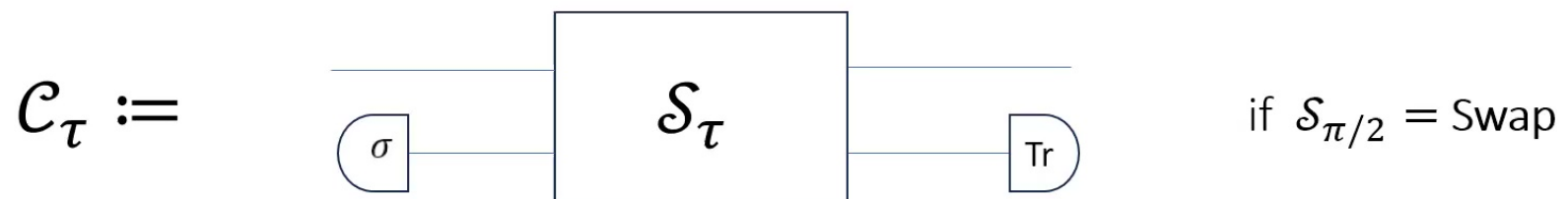
Dynamical evolution is triggered by **information non-equilibrium**.



We can picture a **field**, composed of **identical bosons**.

A **particle** crossing the field evolves if and only if it is in a different state

# Quantum collision models



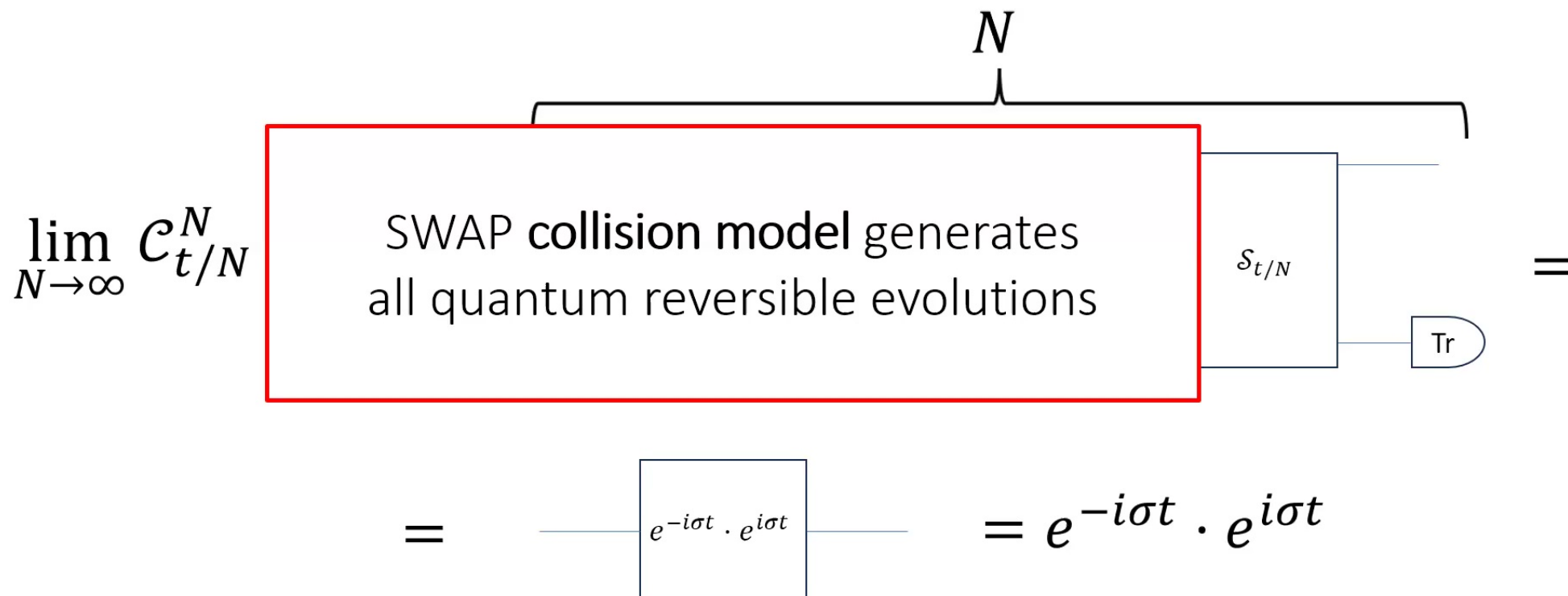
$\mathcal{C}_\tau \circ \dots \circ \mathcal{C}_\tau$  approximate the unitary evolution  $e^{-i\sigma t} \cdot e^{i\sigma t}$  by applying in sequence the operation  $\mathcal{C}_\tau$

[1] S. Lloyd, M. Mohseni, and P. Rebentrost, Nature physics **10**, 631 (2014)

# Collision models in the continuous-time limit

$$\lim_{N \rightarrow \infty} \mathcal{C}_{t/N}^N = \lim_{N \rightarrow \infty} \left[ \overbrace{\left[ \begin{array}{c} \sigma \text{---} \mathcal{S}_{t/N} \text{---} \text{Tr} \\ \sigma \text{---} \mathcal{S}_{t/N} \text{---} \text{Tr} \end{array} \right] \cdots \left[ \begin{array}{c} \mathcal{S}_{t/N} \\ \sigma \text{---} \text{Tr} \end{array} \right]}^N \right] = \\
 = \left[ \begin{array}{c} e^{-i\sigma t} \cdot e^{i\sigma t} \end{array} \right] = e^{-i\sigma t} \cdot e^{i\sigma t}$$

# Collision models in the continuous-time limit





# Collision models in the continuous-time limit

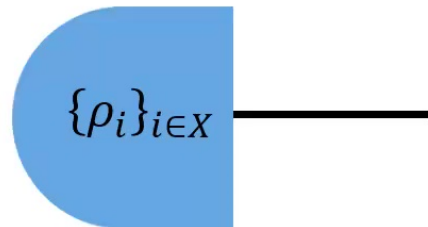
$$\lim_{N \rightarrow \infty} \mathcal{C}_{t/N}^N = \lim_{N \rightarrow \infty} \left[ \overbrace{\left( \begin{array}{c} \sigma \text{---} \mathcal{S}_{t/N} \text{---} \text{Tr} \\ \sigma \text{---} \mathcal{S}_{t/N} \text{---} \text{Tr} \\ \dots \\ \sigma \text{---} \mathcal{S}_{t/N} \text{---} \text{Tr} \end{array} \right)}^N \right] = \\
 = \begin{array}{c} \uparrow \\ \boxed{e^{-i\sigma t} \cdot e^{i\sigma t}} \\ \text{---} \end{array} = e^{-i\sigma t} \cdot e^{i\sigma t}$$



# Outline

1. Introduction of the framework:
  - a) GPTs
  - b) generalized collision models
  - c) our assumptions
  
2. Derivation of a **observable – generator correspondence**:
  - a) state – generator
  - b) observable
  
3. Derivation of an operational speed bound

# Operational framework...

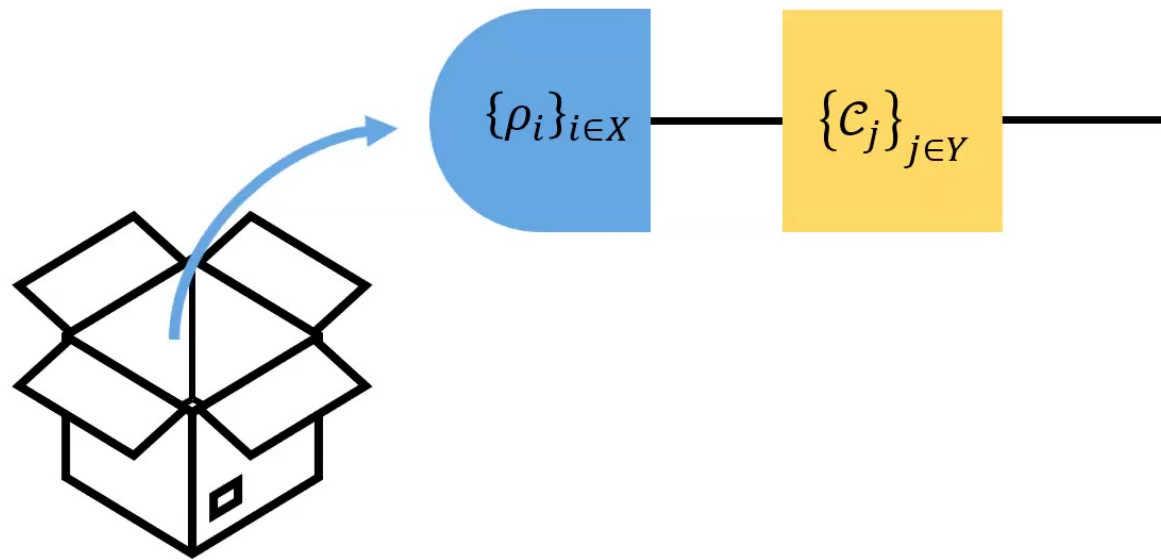


L. Hardy, *Quantum theory from five reasonable axioms* (2001), arxiv:quant-ph/0101012

J. Barrett, *Phys. Rev. A* **75**, 032304 (2007)

H. Barnum, J. Barrett, M. Leifer, and A. Wilce, *Phys. Rev. Lett.* **99**, 240501 (2007)

# Operational framework...

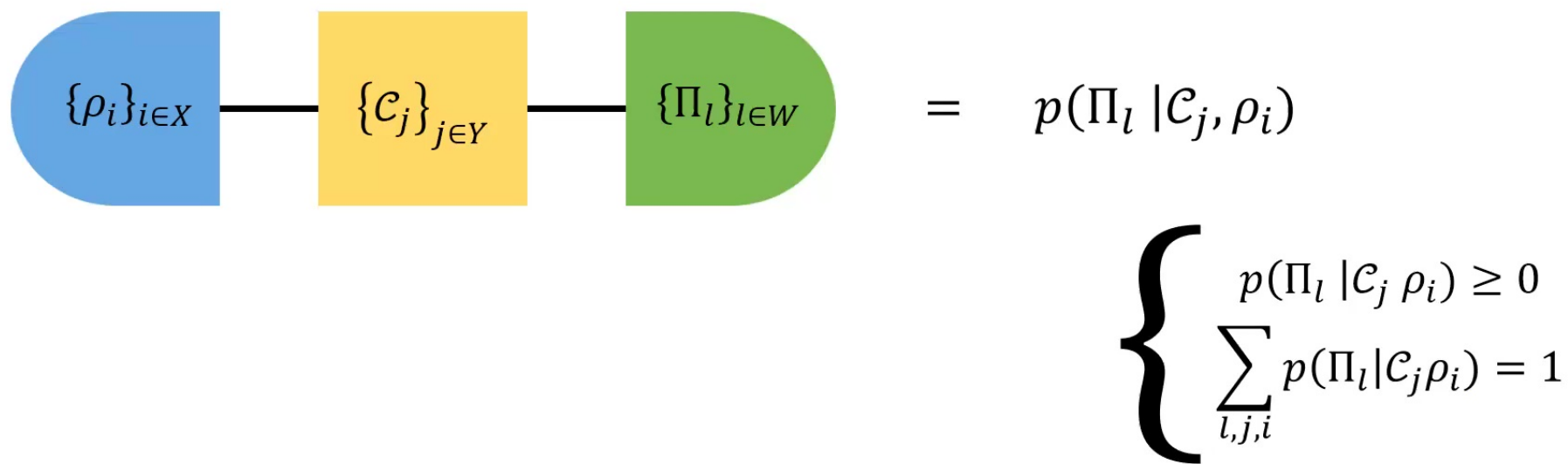


L. Hardy, *Quantum theory from five reasonable axioms* (2001), arxiv:quant-ph/0101012

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... combined with probability



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... combined with probability



$$\rho_i = \begin{pmatrix} p_1^i \\ p_2^i \\ \vdots \\ p_n^i \end{pmatrix}$$

L. Hardy, *Quantum theory from five reasonable axioms* (2001), arxiv:quant-ph/0101012

J. Barrett, *Phys. Rev. A* **75**, 032304 (2007)

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# GPT framework

- A **system**  $\mathcal{S}$  is a finite dimension real vector space
- A **state**  $\rho$  is an element of a convex set  $\subset \mathcal{S}$
- A **measurement** is a collection of linear functional  $\{e_j\}$  such that  $\sum_j e_j(\rho) = 1$ ;
- An **effect** is an element in a measurement: a linear functional over the state space with value in  $[0,1]$
- A **(reversible) transformation** is an (invertible) linear map between state spaces

# GPT framework

- A **system**  $\mathcal{S}$  is a finite dimension real vector space

- A **state**  $\rho$  is an

We assume

- A **measurement**

- the state and effect spaces,
- the set of transformations, to be **CONVEX** and **CLOSED**

$$\sum_j e_j(\rho) = 1;$$

- An **effect** is an  
space with va

the state

- A **(reversible) transformation** is an (invertible) linear map between state spaces



# What is already there about the dynamics?

From the GPT framework ONLY, we can go as far as identifying

- Reversible transformations as a subgroup of  $SO(N)$
- The generators of reversible transformations as skew-symmetric matrices, or a subalgebra thereof

Analogous to unitary and skew-Hermitian matrices in quantum theory

# Collision models in GPT

- Same architecture of the quantum version, this time

$$\mathcal{C}_\tau :=$$



$$\lim_{N \rightarrow \infty} \mathcal{C}_{t/N}^N = e^{G_\sigma t} \quad \text{where } G_\sigma := (I \otimes u)G_{tot}(\cdot \otimes \sigma)$$



# Operational assumptions

3. **Diagonalization [4]:** every state  $\rho$  can be written as a convex sum of perfectly distinguishable pure states
4. **Purity Preservation [5]:** the parallel or sequential composition of pure states (transformations, effects) is a pure state (transformation, effect)

[4] H. Barnum, M. P. Müller, and C. Ududec, *New Journal of Physics* 16, 123029 (2014)

[5] G. Chiribella and C. M. Scandolo, *Entanglement as an axiomatic foundation for mechanics* (2016), arXiv:1608.04459

# State-generator duality

For every collision model, the correspondence  $\sigma \mapsto G_\sigma$  between states and generators is injective:

- if two states generate the same collisional dynamics, then they are the same state
- the maximally mixed state is the only state generating the trivial dynamics

# How to define an observable

We take inspiration from [5,6]

- Observables are linear combinations of co-existing pure effects:

$$X = \sum_i c_i e_i$$

where  $c_i \in \mathbb{R}$  and

$\{e_i\}_{i=1}^d$  is a valid measurement of the theory composed by pure effects

[5] G. Chiribella and C. M. Scandolo, *Entanglement as an axiomatic foundation for mechanics* (2016), arXiv:1608.04459

[6] G. Chiribella, C. M. Scandolo, L. Giannelli, *Sharp theories with purification*, in preparation

# How to define an observable

We take inspiration from [5,6]

- Observables are linear combinations of co-existing pure effects:

$$X = \sum_i c_i e_i$$

- Expectation value:  $\langle X \rangle_\rho := X(\rho) = \sum_i c_i e_i(\rho)$  for every state  $\rho$

[5] G. Chiribella and C. M. Scandolo, *Entanglement as an axiomatic foundation for mechanics* (2016), arXiv:1608.04459

[6] G. Chiribella, C. M. Scandolo, L. Giannelli, *Sharp theories with purification*, in preparation

# The generator-observable correspondence

- $G_\sigma \leftrightarrow \sigma$ , for every state  $\sigma = \sum_i p_i \psi_i$
- $\psi_i \overset{[3,4]}{\leftrightarrow} e_{\psi_i}$  such that  $\{e_{\psi_i}\}$  form a measurement

[3] M. P. Müller and C. Ududec, Physical Review Letters 108, 10.1103/physrevlett.108.130401 (2012)

[4] H. Barnum, M. P. Müller, and C. Ududec, New Journal of Physics 16, 123029 (2014)

# The generator-observable correspondence

- $G_\sigma \leftrightarrow \sigma$ , for every

- $\psi_i \leftrightarrow e_{\psi_i}$  such that

$e^{G_\sigma t}$  and  $e^{\alpha G_\sigma t}$  are different dynamics, for any  $\alpha \neq 1$

- ENERGY observable:  $H := \sum_i p_i e_{\psi_i}$



## The generator-observable correspondence

- $G_\sigma \leftrightarrow \sigma$ , for every state  $\sigma = \sum_i p_i \psi_i$
- $\psi_i \leftrightarrow e_{\psi_i}$  such that  $\{ e_{\psi_i} \}$  form a measurement
- ENERGY observable:

$$H := \lambda_{max} \sum_i p_i e_{\psi_i}$$

where  $\lambda_{max}$  is the maximum singular value of the generator  $G_\sigma$

Energy as  $\langle H \rangle_\rho := \lambda_{max} \sum_i p_i e_{\psi_i}(\rho)$

- $\langle H \rangle_\rho \geq 0$  for every state  $\rho$
- Performing  $\{e_{\psi_i}\}$  estimates the expectation value of the energy
- Energy is invariant under time evolution [Informational equilibrium]:

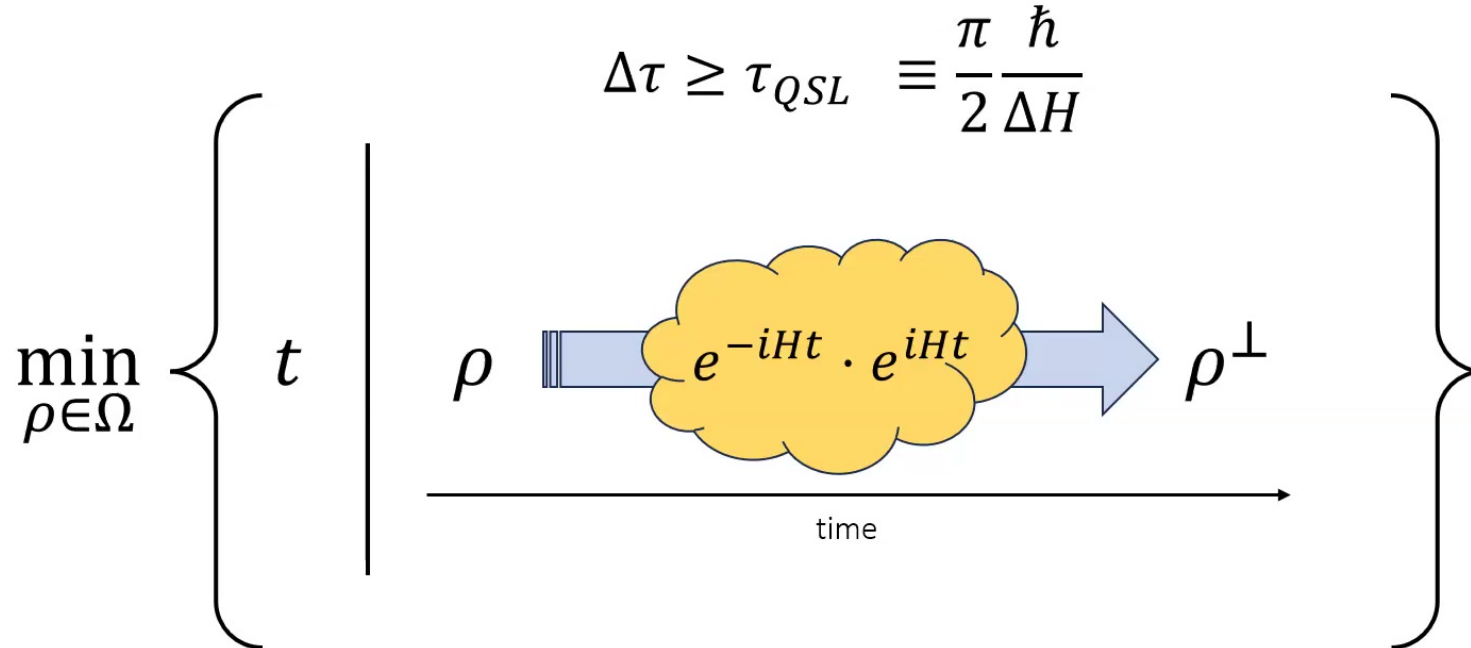
$$\text{for } \rho_t = e^{G\sigma t} \rho, \quad \langle H \rangle_{\rho_t} = \langle H \rangle_\rho \text{ for every } t$$

- The whole probability distribution of the ideal energy measurement is invariant:

$$\text{for } \rho_t = e^{G\sigma t} \rho, \quad e_{\psi_i}(\rho_t) = e_{\psi_i}(\rho) \text{ for every } t, i$$

# Quantum speed bound

Lower bound on the time necessary for a system to evolve between every two **orthogonal** states, according to a given Hamiltonian  $H$



# An improper uncertainty relation

- Time is not an observable [7]
- Time has to be interpreted as the internal time of the system [8-10]
- the relation expresses the minimum time necessary to evolve accordingly to a certain dynamics -> speed bound

[7] W. Pauli, General principles of quantum mechanics

[8] L. Mandelstam and I. Tamm, J. Phys. **9** (1945)

[9] Y. Aharonov and D. Bohm, Phys. Rev. **122**, 1649 (1961)

[10] J. Uffink, American journal of physics **61**, 935 (1993)

# Let's prepare the ground in GTPs

Let  $D_t$  be a dynamic. For any state  $\rho$  we define the quantity

$$v_\rho(t, t_0) := \frac{\|D_t\rho - D_{t_0}\rho\|}{t - t_0} ,$$

as the *evolution speed* of  $\rho$  from  $t_0$  to  $t$

# Bound #1

Let  $U_t = e^{At}$  be a reversible dynamics.

For any state  $\rho$

$$v_\rho(t, t_0) \leq \|A\rho\| = v_\rho(t)$$

the average speed of the evolution is upper bounded by

the **instantaneous speed**  $v_\rho(t) = \lim_{h \rightarrow 0} v_\rho(t+h, t)$

## Rearranging the terms...

- $\Delta t = t - t_0$ ,
  - $D(\rho_t, \rho_{t_0}) := \|\rho_t - \rho_{t_0}\|/\sqrt{2}$
- then the previous equation gives

$$\Delta t \geq \frac{D(\rho_t, \rho_{t_0})}{\Delta H}$$

Quantum speed limit

$$\Delta \tau \geq \frac{\pi \hbar}{2 \Delta H}$$

# Take-home points

Characterized the dynamics in informational terms:

- Introduced the **informational equilibrium** assumption
- Derived a **generator-observable duality** in GPTs
- Derived an operational **speed limit**



# Future works

- Compare toy theories with our informational-dynamical assumptions
- Derive tighter speed limits
- Can we give up strong symmetry? (sharp theories with purification)

Thank you for the attention!