

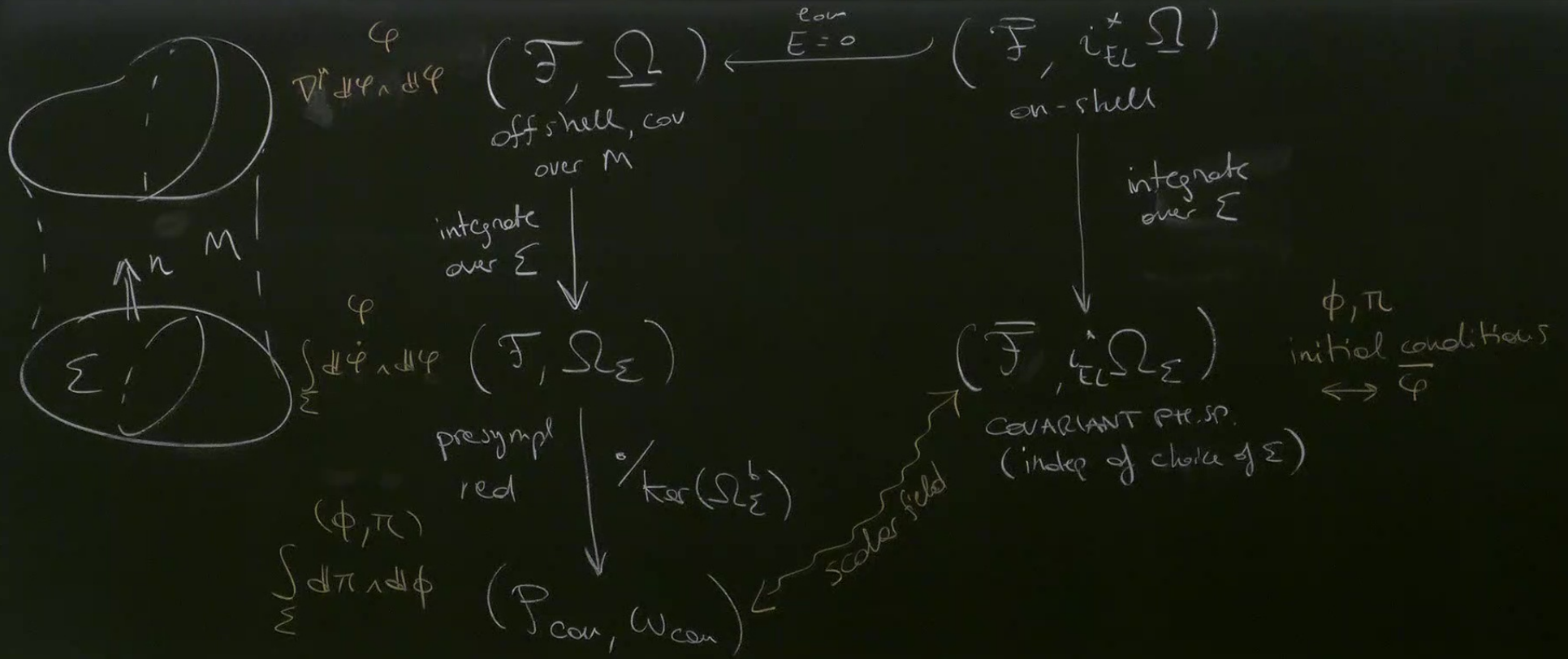
Title: Quantum Gravity Lecture

Speakers: Aldo Riello

Collection: Quantum Gravity 2023/24

Date: April 29, 2024 - 9:00 AM

URL: <https://pirsa.org/24040032>



Maxwell

$$(A_\mu, \int_{\Sigma} F^{\mu\nu} \wedge dA_\nu)$$

(sol. of Maxwell eqs)

$$\Sigma = \{t=0\}$$

$$(A_\mu, \int_{\Sigma} F^{0i} \wedge dA_i)$$

$$(\vec{J}, \int_{\Sigma} \vec{E} \cdot d\vec{\Omega}_\Sigma)$$

$$a_i = A_i|_{\Sigma}$$
$$E^i = F^{0i}|_{\Sigma}$$

$$(a_i, E^i), \int_{\Sigma} E^i \wedge da_i$$

off shell canonical  
ph. space.

Not the same b.c.  
there leftover  
from constraining  
 $E^i$  (at "fixed time")

ie. Gauss constraint  
 $\nabla_i E^i \approx 0$

$v)$   $\leftarrow$  (sol. of Maxwell eqs)

$(\vec{J}, \vec{E}, \Omega_\Sigma)$

Not the same b.c. there left over from constraining  $E^i$  (at "fixed time")

iii. Gauss constraint  $\nabla_i E^i \approx 0$

RMK

• Noether 2  $((D_\alpha^I)^+ E_I \equiv 0)$   
 there are less eom than fields  
 $\rightarrow$  no 1-to-1 corresp. between on shell histories and initial data

• Noether 2'  $(\epsilon_\Sigma^* \underline{J} = C_\alpha \underline{\dot{x}}^\alpha + d \underline{j})$   
 $\mathcal{P}_\Sigma(\vec{r}) \approx 0$   
 Gauss constraint

Diagnosis: need to quotient out gauge sym!

Maxwell

$$(A_\mu, \int_{\Sigma} dF^{\mu\nu} \wedge dA_\nu)$$

(sol. of Maxwell eqs)

$$\Sigma = \{t=0\}$$

$$(A_\mu, \int_{\Sigma} dF^{0i} \wedge dA_i)$$

$$(\bar{J}, \int_{\Sigma} \star E \wedge \Omega_{\Sigma})$$

$$a_i = A_i|_{\Sigma}$$

$$E^i = F^{0i}|_{\Sigma}$$

$$((a_i, E^i), \int_{\Sigma} dE^i \wedge da_i)$$

off shell canonical ph. space.

Not the same b.c. there left over from constraining  $E^i$  (at "fixed time")

inc. Gauss constraint  $\nabla_i E^i \approx 0$

$$(\mathcal{P}^{red}, \mathcal{W}^{red})$$

$$(\bar{J}/g, \Omega^{red})$$

RMK

- Noether there or
- no cons data
- Noether

Diagnosis

Symplectic reduction

$$(\mathcal{P}, \omega) : (a, E^i) \quad \omega = \int_{\Sigma} \underbrace{dE^i \wedge dA_i}$$

$$\text{constraint : } C = D_i E^i = 0$$

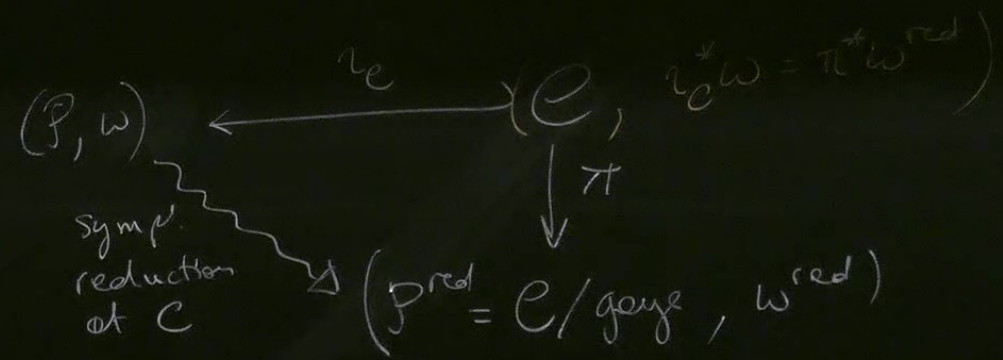
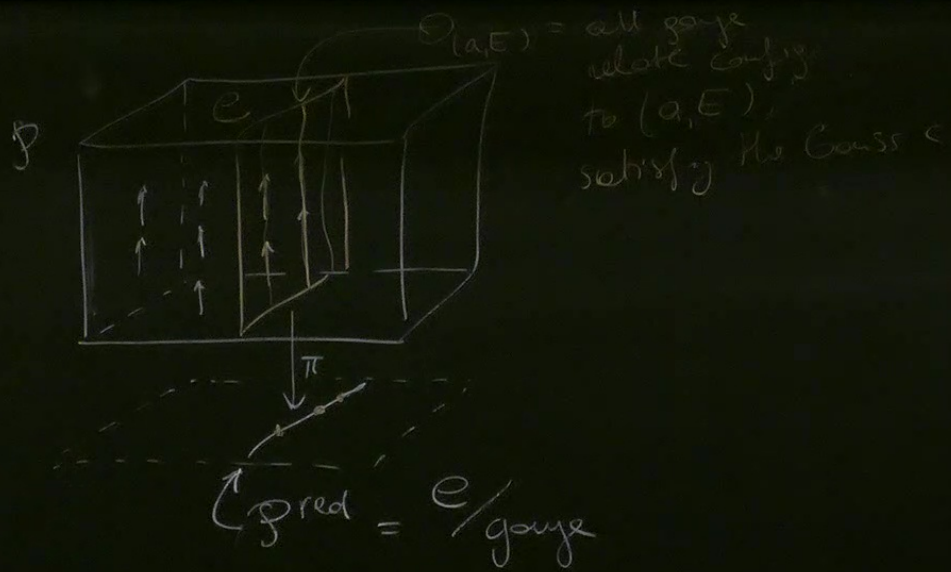
$$\mathcal{C} = \{ (a, E) : \nabla_i E^i = 0 \} \hookrightarrow \mathcal{P}$$

$i_{\mathcal{C}}^* \omega$  will be degenerate, in fact:

$$i_{\rho(\vec{\zeta})}^* \omega = -dQ_{\Sigma}(\vec{\zeta}) = -d \int_{\Sigma} \nabla_i E^i \zeta + (0)$$

pullback to  $\mathcal{C} = 0 \Rightarrow \rho(\vec{\zeta})$  is in the kernel of  $i_{\mathcal{C}}^* \omega$

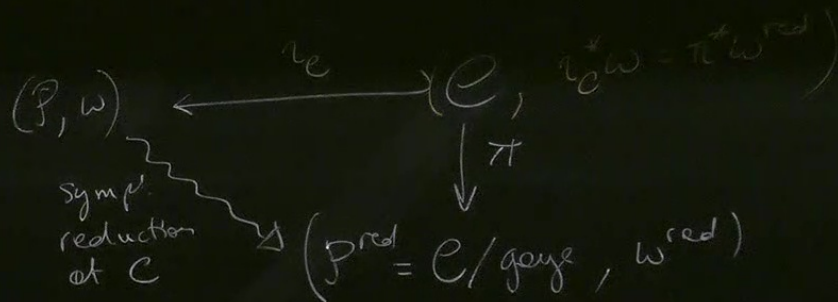
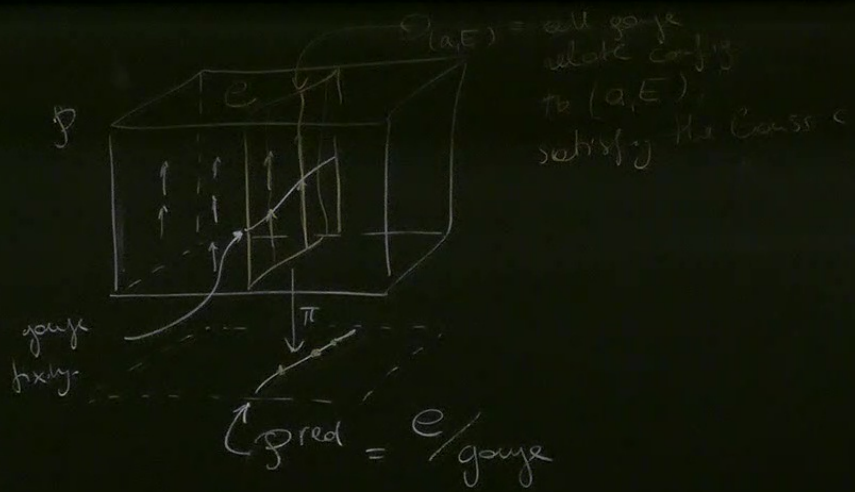
$\uparrow$   
 $\partial \Sigma = \emptyset$



$w^{\text{red}}$  defined by demand:

$$r_e^* w = \pi^* w^{\text{red}}$$

$$(\mathcal{F}^{\text{red}}, \omega^{\text{red}}) \longleftrightarrow (\mathcal{F}/g, \Omega^{\text{red}})$$



$\omega^{\text{red}}$  defined by demand:

$$z_e^* \omega = \pi^* \omega^{\text{red}}$$

In practice, working on  $\mathcal{P}^{\text{red}}$  is hard (locality), so one introduces

a "gauge fixing" in  $\mathcal{e}$ . Eg. on each gauge orbit of  $(q, E)$  there is 1! conf such that  $\nabla^i q_i = 0$

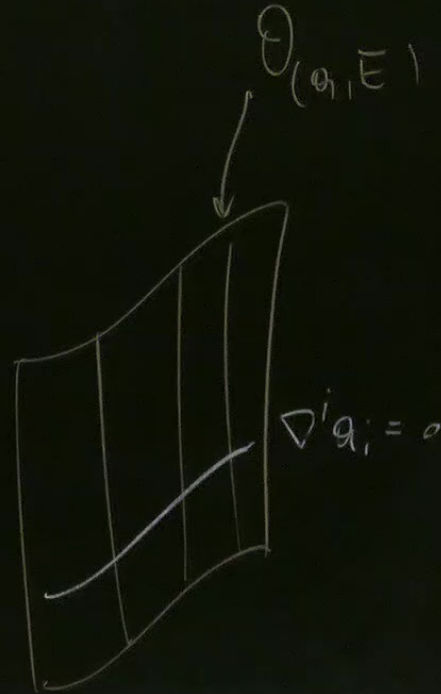


(gauge,  $w^{\text{red}}$ )

demand:

$w^{\text{red}}$

∃! conf such that  $\nabla^i a_i = 0$



pull back to  $e = 0 \Rightarrow p(\bar{z})$  is in the kernel of  $i_e^* \omega$   $d\Sigma = \phi$  a "gauge"

In YM

$\xi(x) \in C^\infty(M)$   
Maxwell

$\xi^\alpha(x) \in C^\infty(M, \mathfrak{g})$   
in YM

$\alpha = \text{basis index}$   
in  $\mathfrak{g}$

$$\mathcal{P} = (a_i^\alpha, E_\alpha^i)$$

$$\omega = \int \sum_\alpha dE_\alpha^i \wedge da_i^\alpha$$

Gauss constraint:  $D_i E_\alpha^i \equiv \nabla_i E_\alpha^i + f_{\alpha\beta}^\gamma a_i^\beta E_\gamma^i \approx 0$

$$i_{p(\bar{z})} \omega = - dQ_\Sigma(\xi) \approx 0$$

a "gauge fixing" in  $\mathcal{C}$ . Eg. on each gauge orbit of  $(q, E)$  there is 1! conf. structure

$$\int_{\rho(\eta)} \int_{\rho(\xi)} \omega = -\int_{\rho(\eta)} Q_{\Sigma}(\xi) \stackrel{\text{equivariance}}{=} -Q_{\Sigma}([\eta, \xi])$$

||

$$- \{Q_{\Sigma}(\eta), Q_{\Sigma}(\xi)\}$$

→ in YM the momentum map / Noether charge is equiv.

$$\{Q_{\Sigma}(\eta), Q_{\Sigma}(\xi)\} = Q_{\Sigma}([\xi, \eta])$$

i.o. the charge (Poisson) algebra is a representation of the gauge alg.

≈ 0

$\eta$  in  $\mathcal{E}$ . E.g. on each page orbit of  $(Q, E)$  there is 1!  $\hookrightarrow$   $\mathcal{E} \rightarrow \mathcal{E}/G$

$$e(\xi) \omega = -\mathbb{L}_{\rho(\eta)} Q_{\mathcal{E}}(\xi) \stackrel{\text{equivariance}}{=} -Q_{\mathcal{E}}([\eta, \xi])$$

$$\eta, Q_{\mathcal{E}}(\xi)\}$$

VM the momentum/Noether charge is equiv.

$$\{Q_{\mathcal{E}}(\eta), Q_{\mathcal{E}}(\xi)\} = Q_{\mathcal{E}}([\xi, \eta])$$

the charge (Poisson) algebra is a representation of the gauge alg.

This guarantees that the reduction procedure is self-consistent  
 (gauge orbits are sub-mfolds of  $\mathcal{E}$  by Frobenius thm)  
 $[e(\xi), e(\eta)] = e([\xi, \eta])$

$\rightarrow$  Rank: Watch out in GR!  
 this will not quite hold  
 [structural feature of timelike diffeo]