

Title: Quantum Gravity Lecture

Speakers: Aldo Riello

Collection: Quantum Gravity 2023/24

Date: April 19, 2024 - 9:00 AM

URL: <https://pirsa.org/24040028>

RECAP

Local sym

$$\mathbb{L}_{\rho(\xi)} \underline{\Omega} = d\underline{R}(\xi)$$

NOETHER 1

$$\underline{J}(\xi) := i_{\rho(\xi)} \underline{\Omega} - \underline{R}(\xi)$$

$$d\underline{J}(\xi) = \delta_{\xi} \varphi^I E_I \approx 0$$

$$i_{\rho(\xi)} \underline{\Omega} = -d\underline{J}(\xi) + \underline{S}(\xi)$$

↑
top-1

↑
top-2

"Best case scenario" $\underline{R}(\xi) = 0$, $\mathbb{L}_{\rho(\xi)} \underline{\Omega} = 0$

$$\rightarrow i_{\rho(\xi)} \underline{\Omega} = -d\underline{J}(\xi), \quad \underline{J}(\xi) = i_{\rho(\xi)} \underline{\Omega}$$

Local sym

Using N1

\leadsto NOETHER

(\mathbb{D}^+)

\rightarrow n
b.c.

\rightarrow un
 \rightarrow t

$$dR(\xi)$$

$$i_{\rho(\xi)} \underline{\omega} - R(\xi)$$

$$\approx 0$$

$$) + S(\xi)$$

$$\uparrow \approx d\underline{L}(\xi)$$

\uparrow top-2

$$R(\xi) = 0, \quad \mathbb{L}_{\rho(\xi)} \underline{\omega} = 0$$

$$(\xi), \quad \underline{J}(\xi) = i_{\rho(\xi)} \underline{\omega}$$

Local sym $\xi(x) \in \mathcal{G} = \begin{cases} C^\infty(M, \mathfrak{g}) \\ \mathcal{X}(M) \end{cases}$

Using N1

\leadsto NOETHER 2

$$\delta_\xi \varphi^I = D_\alpha^I \xi^\alpha$$

$$(D^+)_\alpha^I \xi_I \equiv 0$$

\rightarrow not all eqns are indep.

b.c. # α relations among #I eqns

\rightarrow underdetermined dynamics

\rightarrow to save determinism, need to identify
sols

Ex
YM

Ex

YM

$$\varphi^I = A_a^\alpha, \quad I = (\alpha, a)$$

$$D = d + A$$

$$E_I = D_b^a F_{ba}^\alpha$$

$$d_\xi \varphi^I = (D^\xi)^\alpha_a$$

$$D_\alpha^I = D_{a\alpha}^\beta = \delta_\alpha^\beta \partial_a + f_{\alpha\gamma}^\beta A_a^\gamma = D_a^\beta$$

$$(D^\dagger)^\alpha_a = \text{div}_A$$

$$(D^\dagger)^\alpha_a E_I = D^\alpha D^b F_{ab}^\alpha \equiv 0$$

GR

$$\varphi^I = g_{ab}, \quad I = (a, b)$$

$$E_I = G_{ab} \equiv R_{ab} - \frac{1}{2} R g_{ab}$$

$$d_\xi g_{ab} = L_\xi g_{ab} = \nabla_a \xi_b + \nabla_b \xi_a$$

$$(D^\dagger)^\alpha_a E_I = \nabla^a G_{ab} \equiv P$$

"Bianchi"

$$Q_\Sigma(\xi) = \int_\Sigma J(\xi) \approx 0$$

↑
 $\partial\Sigma = \emptyset$

Ex

YM

$$\varphi^I = A_a^\alpha, \quad I = (\alpha, a)$$

$$E_I = D_a^b F_{ba}^\alpha, \quad D = d + A$$

$$d_\xi \varphi^I = (D \xi)^\alpha_a$$

$$D_\alpha^I = D_a \alpha^\beta = \delta_\alpha^\beta \partial_a + f_{\alpha\gamma}^\beta A_a^\gamma = D_a$$

$$(D^+)^\alpha_a = \text{div}_A$$

$$(D^+)^\alpha_a E_I = D_a^e D^b F_{eb}^\alpha \equiv 0$$

GR

$$\varphi^I = g_{ab}, \quad I = (a, b)$$

$$E_I = G_{ab} \equiv R_{ab} - \frac{1}{2} R g_{ab}$$

$$d_\xi g_{ab} = L_\xi g_{ab} = \nabla_a \xi_b + \nabla_b \xi_a$$

$$(D^+)^\alpha_a E_I = \nabla^a G_{ab} \equiv 0$$

"Bianchi"

$$Q_\Sigma(\xi) = \int_\Sigma J(\xi) \approx 0$$

$$\uparrow$$
$$\partial \Sigma = \emptyset$$

↑ holds
 $\forall \xi, \forall \Sigma$

THM

$$\hookrightarrow \underline{J}(\vec{\xi}) \approx d_j(\vec{\xi})$$

Generally: $\underline{J}(\vec{\xi}) = \xi^{\alpha} \left[\begin{array}{c} C_{\alpha} \\ \mathbb{Z} \\ 0 \end{array} \right] + d_j(\vec{\xi})$

$\partial \xi = \phi, \xi = 0$ (YM)

constraints
(Gauss in YM)

$$\int_{\Sigma} \delta p(\xi) \Omega_{\Sigma} = - \int_{\Sigma} d Q_{\Sigma}(\xi) = - \int_{\Sigma} \xi^{\alpha} d C_{\alpha}$$

→ if all goes for the "best", the constraints generates gauge transf. (off shell)

AMBIGUITIES

$$\bullet \underline{f} \mapsto \underline{f} + d\underline{e}$$

$$\bullet \underline{E} \mapsto \underline{E}$$

$$\bullet \underline{H} \mapsto \underline{H} + d\underline{e} + d\underline{t}$$

$$\bullet \underline{R}(\xi) \mapsto \underline{R}(\xi) + \mathbb{L}_{p(\xi)} \underline{e} + d\underline{r}(\xi)$$

$$\bullet \underline{J}(\xi) \mapsto \underline{J}(\xi) + (\text{no } \underline{e}) + d(\mathbb{I}_{p(\xi)} \underline{e} - \underline{r}(\xi))$$

bulk unchanged

\approx
 dj

$$\bullet \underline{j} \mapsto \underline{j} + (\mathbb{I}_{p(\xi)} \underline{e} - \underline{r}(\xi)) + d\underline{k}(\xi)$$

$$d\underline{f} = \underline{E} - d\underline{e}$$

$$\mathbb{L}_{p(\xi)} \underline{f} = d\underline{R}(\xi)$$

$$\bullet \underline{\Omega} \mapsto \underline{\Omega} + \cancel{d\underline{t}^2} - d\underline{d}\underline{t}$$

bulk unch

$$\Rightarrow \text{if } \partial \Sigma = \emptyset \quad \mathcal{Q}_\Sigma$$

$$d\tau + d\underline{\tau}(\bar{z})$$

$$d(\underbrace{\rho(\bar{z})}_{\text{bulk uncharged}} - \underline{\tau}(\bar{z}))$$

$$+ d\underline{k}(\bar{z})$$

$$d\underline{I} = \underline{E} - d\underline{Q}$$

$$\underbrace{\rho(\bar{z})}_{\text{bulk uncharged}} \underline{I} = d\underline{R}(\bar{z})$$

$$\underline{\Omega} \mapsto \underline{\Omega} + \cancel{d\underline{I}^2} - d\underline{d\underline{Q}}$$

bulk uncharged

$$\Rightarrow \text{if } \partial\Sigma = \emptyset \quad (\Rightarrow \mathcal{Q}_\Sigma \text{ covered, } \mathcal{Q}_\Sigma \text{ covered})$$

$$\mathcal{Q}_\Sigma = \int_\Sigma \underline{I} \quad \& \quad \mathcal{Q}_\Sigma \text{ are man bi-prod}$$

$$\rightarrow \text{but if } \partial\Sigma \neq \emptyset, \text{ it's the wild west.}$$

→ if all goes for the best", The constraints generates gauge transf. (off shell)

GENERAL RELATIVITY

$$\mathcal{F} = \{ g_{ab}(x) \} = \text{Riem}_{d,1}^{(M)}$$

$$S_{\text{EH}} = \frac{\hbar}{l_{\text{pl}}^2} \int_M \sqrt{g} \left(\frac{1}{2} R - \Lambda \right) + \underbrace{\frac{\hbar}{l_{\text{pl}}^2} \int_{\partial M} \sqrt{h} K}_{\text{YGH}}$$

$$l_{\text{pl}}^2 = 8\pi G\hbar$$

Rmks 1) non polynomial in g_{ab}

(fixable using Palatini-Carter, Ashtekar)
= tetrad & (spin)conn or indep fields

2) $\partial^2 g \Rightarrow$ pbs with action principle $\ominus \alpha S_g$
fixable by adding YGH term

→ if all goes for the "best", the constraints generates gauge transf. (off shell)

$$\underline{j} \mapsto \underline{j} + \left(\frac{\delta}{\delta \underline{j}} \right)^2$$

GENERAL RELATIVITY

$$\mathcal{F} = \{ g_{ab}(x) \} = \text{Riem}_{d,1}(M)$$

$$S_{EH} = \frac{\hbar}{l_{pl}^2} \int_M \sqrt{g} \left(\frac{1}{2} R - \Lambda \right) + \underbrace{\frac{\hbar}{l_{pl}^2} \int_{\partial M} \sqrt{h} K}_{YGH}$$

$$l_{pl}^2 = 8\pi G \hbar$$

- Rmks
- 1) non polynomial in g_{ab}
(fixable using Palatini-Carter, Ashtekar)
= tetrad & (spin) conn or indep fields
 - 2) $\partial^2 g \Rightarrow$ pbs with action principle \ominus or δg
fixable by adding YGH term

$$\Downarrow R = \nabla \Downarrow \Gamma' = \nabla(\nabla dg)$$



$$\Downarrow F = D \Downarrow A$$

$$[\nabla, \nabla] = \text{Riem}$$

$$[D, D] = F$$

$$\Gamma = g^{-1} \partial g + \dots$$

$$d\Gamma = \nabla dg$$

$$(3) \psi - \psi(\xi) + d^4 e(\xi)$$

g)

$$dL_{EH} = dg_{ab} \underline{E}^{ab} - d\underline{\Theta}$$

$$\cdot \underline{E}^{ab} = -\frac{1}{2} (G^{ab} + \Lambda g^{ab}) \underline{\epsilon}$$

$$\cdot \underline{\Theta} = \frac{1}{2} (\nabla_b dg^{ab} - \nabla^a dg) \underline{\epsilon}$$

Sym $\xi(x) \in \text{diff}(M) = \mathcal{X}^1(M)$

$$L_{\rho(\xi)} g_{ab} = L_{\xi} g_{ab} = \nabla_a \xi_b + \nabla_b \xi_a$$

$$[\rho(\xi), \rho(\eta)] = -\rho([\xi, \eta])$$

↑
antihomom.

Background indep
(General covariance)

$$L_{\rho(\xi)} \underline{L} = L_{\xi} \underline{L}$$