

Title: Strong Gravity Lecture

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3+1 Decomposition of Spacetime

References: GR, Wald Chap 10.2
3+1 Formalism,ourgoulhon arXiv:gr-qc/0703035

Hypersurfaces of Spacetime

4-dim manifold \mathcal{M}

3-dim manifold $\hat{\Sigma}$

Embedding of $\hat{\Sigma}$ in \mathcal{M}

$$\bar{\Phi}: \hat{\Sigma} \rightarrow \mathcal{M}$$

$\bar{\Phi}^{-1}$, $\bar{\Phi}$, $\bar{\Phi}^{-1}$ continuous

Then $\Sigma = \Phi^{-1}(\hat{\Sigma})$ is called a hypersurface

Level set: $\Sigma = \{x^q \in M \mid F(x^q) = 0\}$

$\partial_q F$ is normal to Σ

$\Leftrightarrow \partial_q F v^q = 0$ for any vector tangent to Σ

Φ^{-1} continuous

$\partial_a F$ is timelike, call Σ spacelike

$\partial_a F$ is spacelike, call Σ timelike

$\partial_a F$ is null, call Σ null

Restrict to Σ spacelike

Choose coordinates $x^s = (t, x^i)$ with x^i coordinate
3-index Σ

Then $\Sigma = \Phi^{-1}(\hat{\Sigma})$ is called a hypersurface

Level set: $\Sigma = \{x^a \in M \mid F(x^a) = 0\}$

$\partial_a F$ is normal to Σ

$\Leftrightarrow \partial_a F v^a = 0$ for any vector tangent to Σ

inuous

$$\Sigma = \{ x^a \in M, \mid t=0 \}$$

$$n^a = (1, \vec{0})$$

Embedding $\Phi: \hat{\Sigma} \rightarrow M$

$$x' \rightarrow (0, x')$$

Unit normal to Σ :

$$n_a = \frac{-\nabla_a t}{\sqrt{-\nabla_a t \nabla^a t}}, \quad n_a n^a = 1$$

$$\} \rightarrow \mathcal{M}$$

$$\rightarrow (0, x')$$

$$n^a = (1, \vec{0})$$

$$\frac{-\nabla_a t}{\sqrt{-\nabla_a t \nabla^a t}}, \quad n_a n^a = -1$$

$$\Phi: \Sigma \rightarrow \mathcal{M} \quad \text{push forward mapping}$$

$$v^i = (0, V^i)$$

$$\Phi^* \mathcal{M} \rightarrow \Sigma$$

$$(u_+, u_i) \rightarrow (u_i) \quad \text{pull back}$$

$$\int_{\Sigma} \vec{T} \cdot \vec{n} \quad M \quad \int_{\Sigma} \vec{T} \cdot \vec{n} \quad \vec{v} \cdot \vec{u}_i$$

$$(0, v^i) \quad (u_+, u_-) = v^i u_i$$

Caution M
 $(v^+, v^-) \not\rightarrow v^i$
 Not true in general

For gas in M , induce

For gas in M , induced metric γ_{ij} on Σ

$$\gamma_{ij} = g_{ij} \quad (\text{in adapted coordinates})$$

Diff hats $\Sigma \stackrel{\wedge}{=} \hat{\Sigma}$

Not true in general

$$V^a = \underbrace{-V^b n_b n^a}_{\text{orthogonal}} + \underbrace{(V^a + V^b n_b n^a)}_{\text{tangent}} = \delta_b^a V^b$$

$$\delta_b^a = \delta_b^a + n_b n^a$$

$$n^a \delta_a^b V_b = n^a V_a - n_b V^b = 0$$

vectors

$$M \rightarrow \Sigma$$
$$V^a \rightarrow \delta_b^a V^b$$

$$\Sigma \rightarrow M$$

$$U_a = \delta_a^i U_i$$

$$\delta_{ab} = \delta_a^i \delta_b^j \delta_{ij} = g_{ab} + n_a n_b$$

= 0

3-index Σ

Intrinsic curvature same $(g_{ab}, M) \rightarrow (\gamma_{ij}, \Sigma)$

Introduce covariant deriv. $D_i \gamma_{jk} = 0$ (also assume torsion free)

defined ${}^{(3)}M^i_{jk}$

Also introduce 3-Dim Riemann tensor

$$(D_i D_j - D_j D_i) v^k = {}^{(3)}R^k_{lij} v^l$$

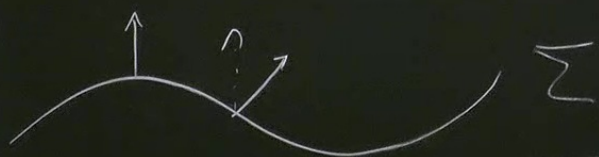
$$3D \text{ Ricci tensor } {}^{(3)}R_{ij} = {}^{(3)}R^k{}_{ikj}$$

$$\text{scalar } {}^{(3)}R = \delta^{ij} {}^{(3)}R_{ij}$$

$$D_a T^{b_1 b_2}{}_{c_1 c_2} = \delta_{d_1}^{b_1} \delta_{c_1}^{e_1} \delta_a^f \nabla_f T^{d_1 d_2}{}_{e_1 e_2} \dots$$

$$D_a \delta_{bc} = \delta_b^d \delta_c^e \delta_a^f \nabla_f (\overset{\circ}{g}_{de} + \overset{\circ}{\eta}_{dne}) = 0$$

Extrinsic Curvature

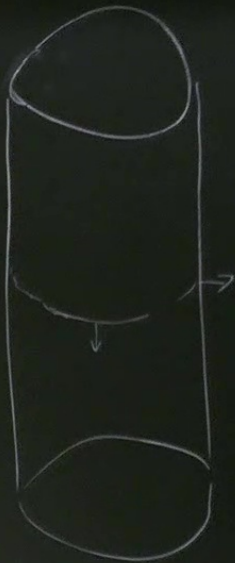


$$k_{ij} = -\nabla_j n_i = \underbrace{(-\nabla_i n_j)}_{\text{symmetric}}$$

Example. Cylinder as hypersurface in \mathbb{R}^3
 $\Sigma = \{(x, y, z) \in \mathbb{R}^3 \mid \rho^2 = x^2 + y^2 = 1\}$

$$n_i = \frac{(-\nabla_i n)}{\text{symmetrisch}}$$

$$\mathbb{R}^3 \left. \begin{array}{l} \\ \\ \end{array} \right\} \rho^2 = x^2 + y^2 = 1$$



$$n^a = \left(\frac{x}{\rho}, \frac{y}{\rho}, 0 \right)$$

$$\nabla_b n_a = \begin{pmatrix} \frac{y}{\rho^2} & -\frac{x}{\rho^2} & 0 \\ -\frac{x}{\rho^2} & \frac{y}{\rho^2} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$i = \{\phi, z\}$$

$$a = \{x, y, z\}$$

$$\tan \phi = \frac{y}{x}$$

$$J_i^a = \frac{dx^a}{dx^i}$$

$$K_{ij} = (-\nabla_b n_a) J_i^a J_j^b$$

$$K_{\phi\phi} = -\frac{(x^2 + y^2)}{\rho^3} = -\rho \Big|_{\rho=1} = -1$$

Close coordinates

$\Lambda = (t, x^i)$ with x^i coordinate
3-index \sum

$$K_{ab} = \delta_a^i \delta_b^j K_{ij} = -\nabla_b n_a - \underbrace{(n^c \nabla_c n_a)}_{=: a^c \text{ acceleration}} n_b \quad (**)$$

$$n_b v^b = 0$$

$$K_{ab} v^b = -v^b \nabla_b n_a$$

$$v^b = \lambda n^b$$

$$\Rightarrow v^b K_{ab} = 0$$

Foliation

M is globally

$$M = \left\{ \Sigma_t \right\}_{t \in \mathbb{R}}$$

$$\Sigma_t = \left\{ x^a \in M \cdot f(x^a) = t \right\}$$

↑ slice or leaf

f smooth and $\nabla_a f \neq 0$

Foliation

$$\mathcal{M} = \left\{ \Sigma_t \right\}_{t \in \mathbb{R}}$$

$$\Sigma_t = \left\{ x^a \in \mathcal{M} \cdot \uparrow(x^a) = t \right\}$$

↑ slice or leaf

↑ smooth and $\nabla_a \uparrow \neq 0$

\mathcal{M} is globally hyperbolic

$$n_a = -K \nabla_c t, \quad n_a n^a = -1$$

$$n_a = (-\alpha, \vec{0})$$

↑ lapse

$$-1 = g^{++} (n_+)^2 = g^{++} \alpha^2$$

$$\alpha = \sqrt{-1/g^{++}}$$

hyperbolic

$$n^a n_a = -1$$

$$n^a = \frac{1}{\alpha} (1, -\beta^i)$$

↑ shift

$$n_{+g^{t_i}} = -\frac{\beta^i}{\alpha} \Rightarrow \beta^i = \alpha^2 g^{t_i}$$

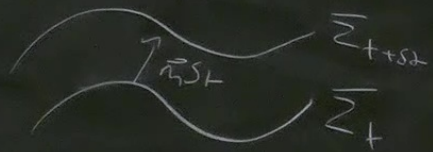
$$\alpha = \frac{d\tau}{dt}, \quad -\beta^i = \frac{dx^i}{dt}$$

Evolut.

Evolution operator:

$$m_a = \alpha \Lambda_a$$

$$m^a \nabla_a \dagger = 1$$



$$\dagger(p + \delta t \vec{m}) = \dagger(p) + \delta t$$

$$\begin{aligned}
\mathcal{L}_{\vec{n}} \gamma_{ab} &= n^c \nabla_c \gamma_{ab} + \gamma_{cb} \nabla_c n^a + \gamma_{ac} \nabla_b n^c \\
&= n^c \nabla_c (n_a n_b) - \gamma_{cb} (K_n^c + n^d (\nabla_d n^c) n_a) - \gamma_{ac} (K_b^c + n^d (\nabla_d n^c) n_b) \\
&= -2K_{ab}
\end{aligned}$$

$+n^d(\nabla_{dn^c})n_b$

$$n^a = \frac{1}{\alpha} (1, -\beta^i)$$

↑ shift

$$1 + g^{ti} = \frac{-\beta^i}{\alpha} \Rightarrow \beta^i = \alpha^2 g^{ti}$$

$$\alpha = \frac{dz}{dt}, \quad -\beta^i = \frac{dx^i}{dt}$$

Evolut.