

Title: Cosmology Lecture

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LAST TIME

MASSLESS SCALAR IN NONDYNAMICAL FLRW

$$S = \frac{1}{2} \int dt d^3x a(t)^2 \left[(\partial_t \phi)^2 - (\partial_i \phi)^2 \right]$$

IF $a(t)$ IS "QUASI DE SITTER" ($w \approx -1$)

THEN ϕ GETS NEARLY SCALE INVARIANT FLUCTUATIONS

$$\langle \phi(k) \phi(k')^* \rangle = \frac{H_*^2}{2k^3} (2\pi)^3 \delta^3(k-k')$$

WELL AFTER HORIZON
CROSSING

$(\cdot)_*$ = "EVALUATED AT HORIZON CROSSING $k=aH$ "

SINGLE FIELD SLOW ROLL INFLATION

$$S = \frac{1}{2} \int dt d^3x \left(\frac{\dot{\phi}^2 a^2}{H^2} \right) \left[(\partial_t \xi)^2 + (\partial_i \xi)^2 \right] \\ + \frac{M_{\text{pl}}^2}{8} \int dt d^3x a(t)^2 \left[(\partial_t \gamma_{ij})^2 - (\partial_k \gamma_{ij})^2 \right]$$

SLOW ROLL BACKGROUND: $\epsilon, \eta \ll 1$

$$\left\{ \begin{array}{l} \dot{\phi} \approx (2\epsilon)^{1/2} H M_{\text{pl}} \\ \dot{H} \approx -\epsilon H \\ \dot{\epsilon} \approx 2\epsilon(2\epsilon - \eta)H \end{array} \right.$$

GRAVITATIONAL WAVES FIRST

$$\gamma_{ij} = \gamma_{ji} \quad \gamma_{ii} = 0 \quad k_i \gamma_{ij}(k) = 0$$

E.G. IF $\vec{k} = k\hat{z}$

$$\gamma_{ij}(\vec{k}) = h_1 \underbrace{\begin{pmatrix} \frac{1}{2} & & \\ & -\frac{1}{2} & \\ & & 0 \end{pmatrix}}_{\epsilon_{ij}^{(1)}} + h_2 \underbrace{\begin{pmatrix} 0 & \frac{1}{2} & \\ \frac{1}{2} & 0 & \\ & & 0 \end{pmatrix}}_{\epsilon_{ij}^{(2)}}$$

GENERALIZE: FOR ANY \vec{k} DEFINE BASIS TENSORS $\epsilon_{ij}^{(s)}(\vec{k})$ SO THAT

$$\epsilon_{ij}^{(s)} \epsilon_{i'j'}^{(s')} = \delta_{ss'}$$

$$\epsilon_{ii}^{(s)} = 0 \quad \epsilon_{ij}^{(s)} = \epsilon_{ji}^{(s)} \quad k_i \epsilon_{ij}^{(s)} = 0$$

$s \in \{1, 2\}$

THEN

$$\gamma_{ij}(\vec{k}) = \sum_{s=1}^2 h_s(\vec{k}) \Sigma_{ij}^{(s)}(\vec{k})$$

IN TERMS OF $h_s(\vec{k})$ THE GW ACTION (S)

$$S_{\text{gw}} = \frac{M_{\text{pl}}^2}{16} \sum_s \int d\tau \frac{d^3k}{(2\pi)^3} a(\tau)^2 \left[\frac{dh_s(\vec{k})}{d\tau} \frac{dh_s(\vec{k})^*}{d\tau} - k^2 h_s(\vec{k})^* h_s(\vec{k}) \right]$$

SAME AS SCALAR IN NONDYNAMICAL FLRW $\phi = \frac{M_{\text{pl}}}{\sqrt{2}} h_s$ ($s \in \{1, 2\}$)

\Rightarrow LATE TIME TWO-POINT FUNCTION

$$\begin{aligned} \langle h_s(\vec{k}) h_s(\vec{k}')^* \rangle &= \frac{8}{M_{\text{pl}}^2} \langle \psi(\vec{k}) \psi(\vec{k}')^* \rangle \delta_{ss'} \\ &= \frac{8}{M_{\text{pl}}^2} \frac{H^2}{2k^3} (2\pi)^3 \delta^3(\vec{k}-\vec{k}') \delta_{ss'} \end{aligned}$$

$$\epsilon_{ij}^{(s)} = 0 \quad \epsilon_{ij}^{(h)} = \epsilon_{ij}^{(s)} \quad k_i \epsilon_{ij}^{(h)} = 0$$

STANDARD NOTATION

$$\langle h_s(\vec{k}) h_s(\vec{k}')^* \rangle = \frac{2\pi^2}{k^3} \underbrace{\Delta_h^2(k)}_{\text{"DIMENSIONLESS GRAVITIC WAVE POWER SPECTRUM"}} \delta_{ss'} (2\pi)^3 \delta^3(\vec{k}-\vec{k}')$$

$$\Delta_h^2(k) = \frac{2H_*^2}{\pi^2 M_{pl}^2}$$

⇒ LATE TIME TWO-POINT FUNCTION

$$\langle h_s(k) h_s(k') \rangle = \frac{\delta}{M_{\text{pl}}^2} \langle \phi(k) \phi(k') \rangle \delta_{\mathcal{S}}^3$$

$$= \frac{\delta}{M_{\text{pl}}^2} \frac{H_e^2}{2k^3} (2\pi)^3 \delta(k-k') \delta_{\mathcal{S}}^3 \lesssim 10^{-10}$$

SPECTRAL INDEX OR "nLT"

$$n_s = \frac{d \log \Delta_h^2(k)}{d \log k}$$

$$= 2 \frac{d \log H_e}{d \log k}$$

GENERAL RULE (WHEN X_* DEPENDS ON k VIA $k = aH$)

$$\begin{aligned}\frac{dX_*}{d \log k} &= \frac{dX}{d \log aH} \\ &= \left(\frac{d \log aH}{dt} \right)^{-1} \frac{dX}{dt} \\ &\approx \left(\underbrace{H}_{\frac{d \log a}{dt}} - \underbrace{\epsilon H}_{\frac{d \log H}{dt}} \right)^{-1} \frac{dX}{dt} \\ &= \frac{1}{H} \frac{dX}{dt}\end{aligned}$$

⇒ LATE TIME TWO-POINT FUNCTION

$$\langle h_s(k) h_s(k') \rangle = \frac{\delta}{M_{pl}^2} \langle \phi(k) \phi(k') \rangle \delta_{\mathcal{S}} \\ = \frac{\delta}{M_{pl}^2} \frac{H^2}{2k^3} (2\pi)^3 \delta(k-k') \delta_{\mathcal{S}} \approx 10^{-10}$$

SPECTRAL INDEX OR "nLT"

$$n_s = \frac{d \log \Delta_h^2(k)}{d \log k}$$

$$= 2 \frac{d \log H_e}{d \log k}$$

$$\approx 2 \frac{1}{H} \frac{d \log H}{dt}$$

$$\approx -2\varepsilon$$

$$\varepsilon \equiv \frac{M_{pl}^2}{2} \left(\frac{v'}{v} \right)^2$$

MOVING ON TO SCALARS (ξ)

$$S = \frac{1}{2} \int dt d^3x \underbrace{\frac{\dot{\phi}^2 a^2}{H^2}} \left[(\partial_\tau \xi)^2 - (\partial_i \xi)^2 \right]$$

SAME AS MASSLESS SCALAR IN FAKE EXPANSION HISTORY $z(\tau) = \frac{\dot{\phi} a}{H}$

TWO HUBBLES

$$H_a = \frac{\dot{a}}{a} \quad H_z = \frac{\dot{z}}{z}$$

\Rightarrow LATE TIME TWO POINT FUNCTION

$$\langle \xi(k) \xi(k')^\dagger \rangle = \left(\frac{H_z^2}{2k^3} \right)_* (2\pi)^3 \delta^3(k-k')$$

COMPUTE H_z IN SLOW ROLL APPROX

$$k = aH_a$$

$$Z(\tau) \equiv \frac{\dot{\phi} a}{H} \approx (2\epsilon)^{1/2} M_{pl} a(\tau)$$

$$H_z \equiv \frac{\partial_\tau Z}{Z^2} = \frac{1}{Z} \partial_\tau \log Z$$

$$\approx \left[(2\epsilon)^{1/2} M_{pl} a \right]^{-1} \left[\frac{1}{2} \frac{\partial_\tau \epsilon}{\epsilon} + \frac{\partial_\tau a}{a} \right]$$

$$\approx \left[(2\epsilon)^{1/2} M_{pl} a \right]^{-1} \left[(2\epsilon - \eta) \frac{aH_a}{a} + aH_a \right]$$

$$\approx \frac{H_a}{(2\epsilon)^{1/2} M_{pl}}$$

NOTE

$k = aH_a = zH_z$ so $(\cdot)_*$ IS UNAMBIGUOUS

$$\Rightarrow \langle \xi(k) \xi(k')^* \rangle = \left(\frac{H_a^2}{4\epsilon M_{pl}^2 k^3} \right)_* (2\pi)^3 \delta^3(k-k')$$

USUAL NOTATION IS

$$\langle \xi(k) \xi(k')^* \rangle = \frac{2\pi^2}{k^3} \underbrace{\Delta_s^2(k)}_{\text{"SONAR POWER SPECTRUM"}} (2\pi)^3 \delta^3(k-k')$$

(L2) M_{pl}

THEORY

$$\Delta_s(k)^2 = \left(\frac{H^2}{8\pi\epsilon M_{pl}^2} \right)^*$$

$$n_s = 1 - 6\epsilon + 2\eta$$

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HISTORICAL

$$r = 16\epsilon$$

$$n_t = -2\epsilon$$

OBSERVATION

$$\left[\equiv 1 + \frac{d \log \Delta_s^2(k)}{d \log k} \right]$$

$$\left[\equiv \frac{\Delta_h^2(k)}{\Delta_s^2(k)} \right]$$

$$\left[\equiv \frac{d \log D_h^2(k)}{d \log k} \right]$$

OBSERVATIONS

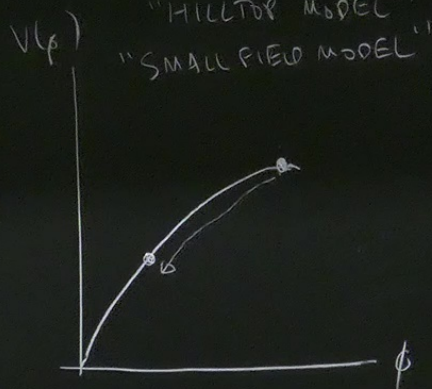
$$n_s^2 = (2.11 \pm 0.03) \times 10^{-9}$$

$$n_s = 0.9649 \pm 0.0042$$

$$r \lesssim 0.036 \quad \text{AT 95\% CL}$$

n_t NO CONSTRAINT

"HILLTOP MODEL"
"SMALL FIELD MODEL"



$$\epsilon \ll \eta \quad \eta \sim -0.018$$

$$H \lesssim \text{FEW} \times 10^{13} \text{ GeV}$$

"LARGE FIELD MODEL"



$m^2 \phi^2 \quad \epsilon = \eta$

DETECTABLE $r = 16\epsilon$

$$r \text{ DETECTABLE} \Leftrightarrow \Delta_h^2 \sim \frac{H^2}{M_{pl}^2} > 10^{-11}$$

$$\Leftrightarrow H \gtrsim (\text{FEW} \times 10^5) M_{pl}$$

$$\sim \text{FEW} \times 10^{13} \text{ GeV}$$

⇒ LATE TIME TWO-POINT FUNCTION $\langle h_s(k) h_s(k') \rangle = \frac{\delta}{M_{pl}^2} \langle \phi(k) \phi(k') \rangle \delta_{\mathcal{S}} \delta_{\mathcal{S}'}$
 $= \frac{\delta}{M_{pl}^2} \frac{H^2}{2k^3} (2\pi)^3 \delta^3(k-k') \delta_{\mathcal{S}} \delta_{\mathcal{S}'} \lesssim 10^{-10}$

$$E_{inf} \sim \rho^{1/4}$$

$$= (3M_{pl}^2 H^2)^{1/4}$$

$$= \left(\frac{3\pi^2 \Delta_s^2 M_{pl}^4}{2} \right)^{1/4} \quad \checkmark^{1/4}$$

$$= (3.23 \times 10^{16} \text{ GeV}) \times r^{1/4}$$

DETECTABLE $r \Leftrightarrow E_{inf} \gtrsim 10^{16} \text{ GeV}$ [GUT SCALE]

$$\epsilon_{ij}^{(s)} = 0 \quad \epsilon_{ij}^{(v)} = \epsilon_{ij}^{(g)} \quad k_{\mu} \epsilon_{ij}^{(s)} = 0$$

INFLATION "FIELD EXCURSION"

$$\begin{aligned} \frac{d\phi}{d \log a} &= \left(\frac{d \log a}{dt} \right)^{-1} \frac{d\phi}{dt} \\ &\approx H^{-1} \left[(2\epsilon)^{1/2} H M_{pl} \right] \\ &= (2\epsilon)^{1/2} M_{pl} \quad \text{LYTH "BOUND"} \end{aligned}$$