

Title: Quantum Information Lecture

Speakers: Eduardo Martin-Martinez

Collection: Quantum Information 2023/24

Date: March 15, 2024 - 9:00 AM

URL: <https://pirsa.org/24030057>

Def. (Purification of a state) A state $\hat{\rho}_A$ is said to be pur

Stinespring dilation theorem: Given a CPTP map mapping $\mathcal{L}(\mathcal{H})$ to $\mathcal{L}(\mathcal{H}_B)$,
there exists a Hilbert space \mathcal{H}_{AB} and a state $|\Omega_{AB}\rangle$ for every $\rho \in \mathcal{L}(\mathcal{H})$
If $\dim(\mathcal{H})$ is finite then $\mathcal{H}_{AB} \geq (\dim \mathcal{H})^2$

Def. (Purification of a state) A state $\hat{\rho}_A$ is said to be

Stinespring dilation theorem: Given a CPTP map mapping
there exists a Hilbert space \mathcal{H}_{AB} and a state $|\Omega_{AB}\rangle$ for every
If $\dim(\mathcal{H})$ is finite then $\mathcal{H}_{AB} \geq (\dim \mathcal{H})^2$

to be purified by $|\psi_{AB}\rangle$ if $\hat{\rho}_A = \text{tr}_B(|\psi_{AB}\rangle\langle\psi_{AB}|)$ for $\hat{H}|\psi\rangle = 0$

mapping density operators to density operators $T: \hat{\rho} \in \mathcal{O}(\mathcal{H}) \rightarrow \hat{\rho}' \in \mathcal{O}(\mathcal{H})$
 for every $\hat{\rho}_A \in \mathcal{O}(\mathcal{H}_A)$, for every T , $T(\hat{\rho}_A) = \text{tr}_B [\hat{U}_{AB} |0_{AB}\rangle\langle 0_{AB}| \hat{U}_{AB}^\dagger]$.

$$|0_{AB}\rangle \neq |0_A\rangle \otimes |0_B\rangle$$

$$\sum_i \frac{\hat{p}_i}{i} + m \omega^2 x_i^2 + \lambda (x_1 - x_2)^2$$

Def: Quantum measurement: A measurement of a quantum system accompanied by an update rule for the post-measurement state

—————

Postulate 4: Measurements of a quantum system are represented by one operator per each possible outcome of the measurement that sum to the identity.

If the system is in a state $|\psi\rangle$, the probability of getting outcome n is

is updated as $|\psi'\rangle = \frac{\hat{M}_n |\psi\rangle}{\sqrt{\langle \psi | \hat{M}_n^\dagger \hat{M}_n | \psi \rangle}}$

at of a quantum system is a map from a quantum state $|\psi\rangle$ to a set of real numbers
post-measurement state.

system are represented by a set of operators $\{\hat{M}_n\}$ acting on the Hilbert space of the system
measurement that satisfies $\sum_n \hat{M}_n^\dagger \hat{M}_n = \mathbb{1}$
probability of getting outcome n is $P(n) = \langle \psi | \hat{M}_n^\dagger \hat{M}_n | \psi \rangle$ and the state after the measurement

We call the operators $\hat{E}_n := \hat{M}_n^\dagger \hat{M}_n \Rightarrow \sum_n \hat{E}_n = \mathbb{1}$ POVM elements

(Positive operator-valued measures)

defined by a set of operators $\{\hat{M}_n\}$ acting on the Hilbert space of the system
satisfy $\sum_n \hat{M}_n^\dagger \hat{M}_n = \mathbb{1}$
outcome n is $P(n) = \langle \psi | \hat{M}_n^\dagger \hat{M}_n | \psi \rangle$ and the state after the measurement
operators $\hat{E}_n := \hat{M}_n^\dagger \hat{M}_n \Rightarrow \sum_n \hat{E}_n = \mathbb{1}$ POVM elements
(operator-valued measures)

Scenario 2: PDM:

$$\hat{E}_1 = \frac{\sqrt{2}}{1+\sqrt{2}} |1\rangle\langle 1|$$

$$\hat{E}_2 = \frac{\sqrt{2}}{1+\sqrt{2}} \frac{(|0\rangle - |1\rangle)(\langle 0| - \langle 1|)}{2}$$

$$\hat{E}_3 = \mathbb{1} - \hat{E}_1 - \hat{E}_2$$

$$P_1(|\psi_1\rangle) = \langle \psi_1 | \hat{E}_1 | \psi_1 \rangle = 0$$

$$P_1(|\psi_2\rangle) = \langle \psi_2 | \hat{E}_1 | \psi_2 \rangle = \frac{1}{\sqrt{2}+1}$$

$$P_2(|\psi_1\rangle) = \langle \psi_1 | \hat{E}_2 | \psi_1 \rangle = 1 - \frac{1}{\sqrt{2}}$$

$$P_2(|\psi_2\rangle) = \langle \psi_2 | \hat{E}_2 | \psi_2 \rangle = 0$$

$$P_3(|\psi_1\rangle) = \langle \psi_1 | \hat{E}_3 | \psi_1 \rangle = \frac{1}{\sqrt{2}}$$

$$P_3(|\psi_2\rangle) = \langle \psi_2 | \hat{E}_3 | \psi_2 \rangle = 1 - \frac{1}{\sqrt{2}+1}$$

<11)

$$P_1(|\psi_1\rangle) = \langle \psi_1 | \hat{E}_1 | \psi_1 \rangle = 0 \quad \text{Outcome 1}$$

$$P_1(|\psi_2\rangle) = \langle \psi_2 | \hat{E}_1 | \psi_2 \rangle = \frac{1}{\sqrt{2}+1}$$

$$P_2(|\psi_1\rangle) = \langle \psi_1 | \hat{E}_2 | \psi_1 \rangle = 1 - \frac{1}{\sqrt{2}} \quad \text{Outcome 2}$$

$$P_2(|\psi_2\rangle) = \langle \psi_2 | \hat{E}_2 | \psi_2 \rangle = 0$$

$$P_3(|\psi_1\rangle) = \langle \psi_1 | \hat{E}_3 | \psi_1 \rangle = \frac{1}{\sqrt{2}} \quad \text{Outcome 3}$$

$$P_3(|\psi_2\rangle) = \langle \psi_2 | \hat{E}_3 | \psi_2 \rangle = 1 - \frac{1}{\sqrt{2}+1}$$

$$\langle \psi | M_n^\dagger M_n | \psi \rangle$$

(Positive Operator-Value)

(Corollary of) Naimark's dilation theorem: Consider the a detector, the environment and a target system. Then the act the target system

Operator-valued measures)

Consider the unitary evolution of a coupled system consisting of
when the action of a PVM on the detector always yields a POVM on