

Title: Quantum Information Lecture

Speakers: Eduardo Martin-Martinez

Collection: Quantum Information 2023/24

Date: March 08, 2024 - 9:00 AM

URL: <https://pirsa.org/24030054>

Partial transpose

$$\hat{\rho}_{AB} = \sum_{i,j,k,l} p_{ijkl} |i\rangle_A |j\rangle_B \langle k|_A \langle l|_B$$

Question:

Is $\hat{\rho}_{AB}$ separable or entangled?

$$\hat{\rho}_{AB}^{T_B} = \sum_{i,j,k,l} p_{ijkl} |i\rangle_A |l\rangle_B \langle k|_A \langle j|_B$$

$$\hat{\rho}_{AB}^{T_A} = \sum_{i,j,k,l} p_{ijkl} |k\rangle_A |j\rangle_B \langle i|_A \langle l|_B$$

Distillable entanglement: Number of maximally entangled states that can be distilled from N copies of $\hat{\rho}_{AB}$ and LOCC.

Peres criterion

If the partial transpose $\hat{\rho}^{T_{A/B}}$ of a bipartite density matrix $\hat{\rho}_{AB}$ has at least

maximally entangled states of lower dimension (typically Bell pairs) that one can obtain

Peres criterion

If the partial transpose $\hat{\rho}_{AB}^{T_{AB}}$ of a bipartite density matrix $\hat{\rho}_{AB}$ has at least one negative eigenvalue $\hat{\rho}_{AB}$ is entangled.

Locally entangled states of lower dimension (typically Bell pairs) that one can obtain

Partial transpose $\hat{\rho}_{AB} = \sum_{i,j,k,l} \rho_{ijkl} |i\rangle_A |j\rangle_B \langle k|_A \langle l|_B$

Question:
Is $\hat{\rho}_{AB}$ separable or entangled?

$$\hat{\rho}_{AB}^{TA} = \sum_{i,j,k,l} \rho_{ijkl} |i\rangle_A |l\rangle_B \langle k|_A \langle j|_B$$

$$\hat{\rho}_{AB}^{TB} = \sum_{i,j,k,l} \rho_{ijkl} |k\rangle_A |j\rangle_B \langle i|_A \langle l|_B$$

Distillable entanglement: Number of maximally entangled states of lower dimension (typically Bell pairs) that one can obtain from N copies of $\hat{\rho}_{AB}$ and LOCC

Perez criterion (sufficient condition for entanglement)

If the partial transpose $\hat{\rho}_{AB}^{TA}$ of a bipartite density matrix $\hat{\rho}_{AB}$ has at least one negative eigenvalue $\hat{\rho}_{AB}$ is entangled.

if $\hat{\rho}_{AB}$ is 2×2 or 3×2 then the condition is also necessary.

Peres criterion (sufficient condition for entanglement)

If the partial transpose $\hat{\rho}_{AB}^{T_{A/B}}$ of a bipartite density matrix $\hat{\rho}_{AB}$ has at least one negative eigenvalue $\hat{\rho}_{AB}$ is entangled.

if $\hat{\rho}_{AB}$ is 2×2 or 3×2 then the condition is also necessary.

of maximally entangled states of lower dimension (typically Bell pairs) that one can

$\hat{\rho}_{AB}^{T_{A/B}}$ has negative eigenvalues $\Leftrightarrow \hat{\rho}_{AB}$ has distillable entanglement

$$\text{I} \int \hat{\rho}_{AB} = \sum C_i \hat{\rho}_A^i \otimes \hat{\rho}_B^i \Rightarrow \hat{\rho}_{AB}^{T_B} = \sum C_i \hat{\rho}_A^i \otimes \hat{\rho}_B^{iT}$$

$(\hat{\rho}_{AB} \text{ is separable}) \Rightarrow (\hat{\rho}_{AB}^{T_B} \text{ is still a positive operator})$

LOCC: Local operations and Classical Communication

States with bound entanglement but PPT (positive partial transpose) cannot be

Entanglement measures

Def. An entanglement measure is a map $E: \hat{\rho}_{AB} \rightarrow \mathbb{R}$ that

1. Must be maximum iff the state is maximally entangled
2. Must be zero iff the state is separable.
- 3.

Entanglement measures

Def. An entanglement measure is a map $E: \hat{\rho}_{AB} \rightarrow \mathbb{R}$ that

1. Must be maximum iff the state is maximally entangled
2. Must be zero iff the state is separable.
3. Must not increase under LOCC.

Entanglement entropy

The Entanglement entropy of a pure bipartite state $\hat{\rho}_{AB} = |\psi_{AB}\rangle\langle\psi_{AB}|$ is

$$S_E(|\psi_{AB}\rangle) := S[\text{tr}_A(|\psi_{AB}\rangle\langle\psi_{AB}|)] = S[\text{tr}_B(|\psi_{AB}\rangle\langle\psi_{AB}|)]$$

Problem: It dramatically fails to measure entanglement for non-pure states

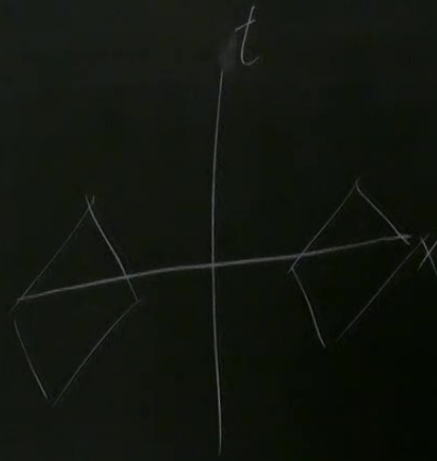
$$\text{Ex } \hat{\rho}_{AB} = \mathbb{1}_A \otimes \mathbb{1}_B \Rightarrow \hat{\rho}_A = \text{tr}_B(\hat{\rho}_{AB}) = \mathbb{1}_A \quad S(\text{tr}_+ \hat{\rho}_{AB}) = \max S$$

Let $\hat{\rho}_{AB} = |\Psi_{AB}\rangle\langle\Psi_{AB}|$ is the Von Neumann entropy of either of the partial subsystems

$$(|\Psi_{AB}\rangle\langle\Psi_{AB}|)$$

system for non-pure state

$$S(\text{tr}_A \hat{\rho}_{AB}) = \max S$$



$\langle \psi_{AB} |$ is the Von Neumann entropy of either of the partial subsystems

$$\hat{\rho}_{AB} = \sum_n C_{ij} |\psi_{AB}^i\rangle \langle \psi_{AB}^j|$$



and state \rangle

$$\rangle = \max S$$

Problem: It dramatically fails to measure entanglement for non-pure states

$$\text{Ex } \hat{\rho}_{AB} = \mathbb{1}_A \otimes \mathbb{1}_B \Rightarrow \hat{\rho}_A = \text{tr}_B(\hat{\rho}_{AB}) = \mathbb{1}_A \quad S(\text{tr}_B \hat{\rho}_{AB}) = \max S$$

Entanglement of Formation (EoF)

$$\hat{\rho}_{AB} = \sum_i P_i |\Psi_{AB}^i\rangle \langle \Psi_{AB}^i| \quad \langle S_E \rangle(\hat{\rho}_{AB}) := \sum_i P_i S_E(|\Psi_{AB}^i\rangle)$$

Def: $EoF(\hat{\rho}_{AB}) := \inf_{\{P_i, |\Psi_{AB}^i\rangle\}} \sum_i P_i EoF(|\Psi_{AB}^i\rangle)$

$$EoF(|\Psi_{AB}^i\rangle) = S_E(|\Psi_{AB}^i\rangle)$$