

Title: QFT III Lecture

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Collection: QFT III 2023/24

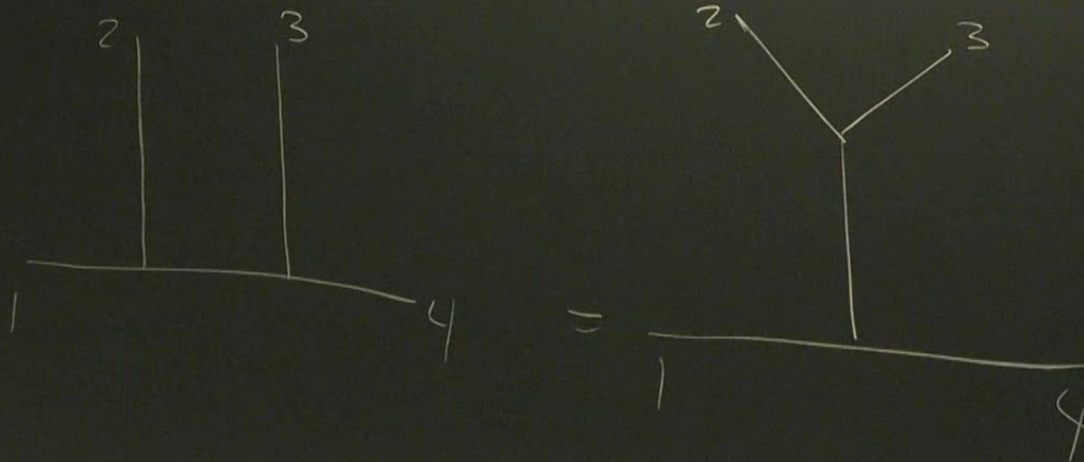
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Recap:

- CFT data : $\{ (\Delta_I, S_I); C_{IJK} \}$
 2 + 3pt functions

- 4pt. function. Associativity of OPE data is constrained



$$g(u, v) = \left(\frac{u}{v} \right)^{\Delta_\phi} g(u, v)$$

$$g(u, v) = \sum_{\mathcal{O}} C_{\phi\phi\mathcal{O}}^2 g(u, v)$$

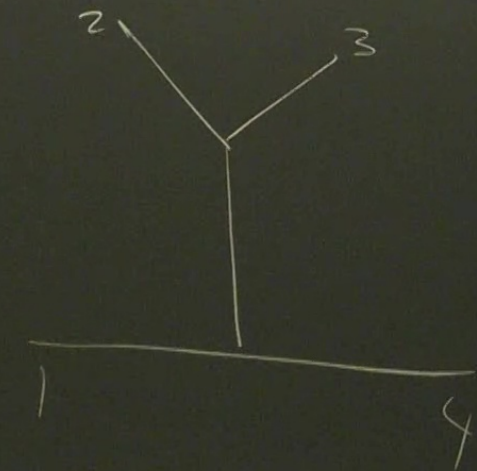
$$\{ (\Delta_I, S_I), C_{IJK} \}$$

associativity of OPE data is constrained $\langle \phi(x_1) \phi(x_2) \phi(x_3) \phi(x_4) \rangle$

$$g(u, v) - \left(\frac{u}{v}\right)^{\Delta_\phi} g(v, u) = 0$$

$$g(u, v) = \sum_{\mathcal{O}} C_{\phi\phi\mathcal{O}}^2 g_{\Delta_{\mathcal{O}}, l_{\mathcal{O}}}(u, v)$$

↓
conformal block.

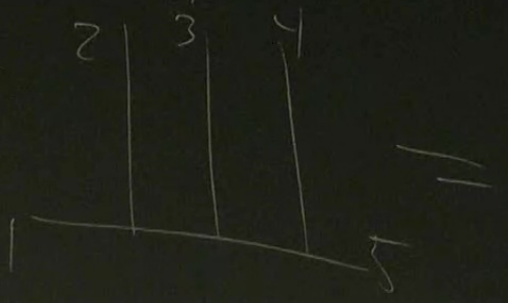


↓
conformal block.

$$g_{\Delta_0, \ell_0}(u, w) = X_{12}^{2\Delta_\phi} X_{34}^{2\Delta_\phi} C_a(X_{12}, \partial_2) C_b(X_{34}, \partial_4) \frac{T^{ab}(X_{34})}{X_{34}^{2\Delta_\sigma}}$$

- Even spacetime dimensions explicitly known.

$$- g_{\Delta_0}(u, w) = u^{\Delta/2} (1 + \dots)$$



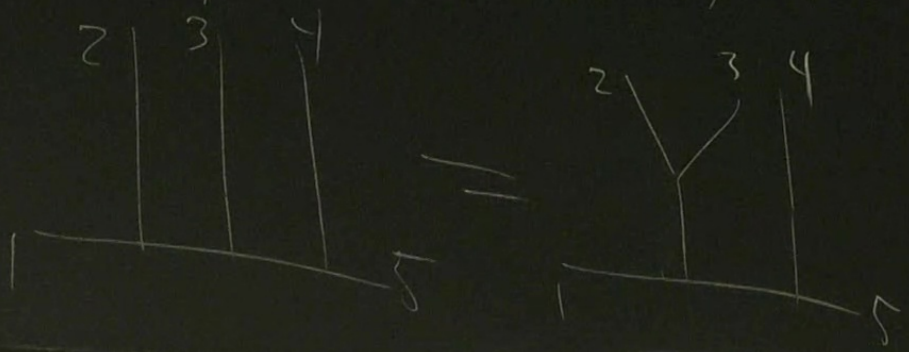
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↓ conformal block

$$g_{\Delta_0, \theta}(\mu, \nu) = X_{12}^{2\Delta_\phi} X_{34}^{2\Delta_\phi} C_a(X_{12}, \partial_2) C_b(X_{34}, \partial_4) \frac{T^{ab}(X_{34})}{X_{34}^{2\Delta_\theta}}$$

- Even spacetime dimensions explicitly known.

$$- g_{\Delta_0}(\mu, \nu) = u^{\Delta/2} (1 + \dots)$$



Anomalies

- Symmetry . if symmetry is continuous, $\partial^\mu j_\mu = 0$ (Noether's theorem).

Quantum mechanically $\partial^\mu j_\mu = 0(\hbar)$

$$Q = \int d^{d-1}x j_0$$

Examples:

- $T_\mu^\mu = 0 = \beta(g) \text{Tr} F_{\mu\nu}^2$ (you were naive).

- 't Hooft anomalies

Anomalies

- Symmetry: if symmetry is continuous, $\partial^\mu j_\mu = 0$ (Noether's theorem).

Quantum mechanically $\partial^\mu j_\mu = \mathcal{O}(\hbar)$ $Q = \int d^{d+1}x j_0$

Examples:

- $T_\mu^\mu = 0 = \beta(g) \text{Tr} F_{\mu\nu}^2$ (you were name).

- 't Hooft anomalies: - couple system to a background gauge field

Anomalies

- Symmetry : if symmetry is continuous, $\partial^\mu j_\mu = 0$ (Noether's theorem).

Quantum mechanically $\partial^\mu j_\mu = \theta(t)$ $Q = \int d^{d-1}x j_0$

Examples:

- $T_\mu^\mu = 0 = \beta(g) \text{Tr} F_{\mu\nu}^2$ (you were naive).

- 't Hooft anomaly: - couple system to a background gauge field

anomaly coefficient

$$\partial^\mu j_\mu = \boxed{K} f(F_{\mu\nu})$$

topological invariant

$$\int dx f_{\mu\nu} j^\mu$$

$D=2 \quad F = dA$
 $D=4 \quad \int_1(F_1 F)$

UV

$$\langle T_{\mu\nu}(p) T_{\rho\sigma}(-p) \rangle = (\eta_{\mu\rho} \eta_{\nu\sigma} + \eta_{\mu\sigma} \eta_{\nu\rho}) F_1(p^2) + \eta_{\mu\nu} \eta_{\rho\sigma} F_2(p^2) + \dots +$$

$$P_\mu P_\nu P_\rho P_\sigma F_5(p^2)$$

- Demand that $P^\mu \langle \quad \rangle = 0$

3 relations between $F_i(p^2)$.

$$\partial^\mu T_{\mu\nu} = 0 \Rightarrow P^\mu T_{\mu\nu} = 0$$

$$\Pi_{\mu\nu} = (P_\mu P_\nu - \eta_{\mu\nu} p^2)$$

$$\langle T_{\mu\nu}(p) T_{\rho\sigma}(-p) \rangle = \left[\Pi_{\mu\rho} \Pi_{\nu\sigma} + \Pi_{\mu\sigma} \Pi_{\nu\rho} - \frac{2}{D-1} \Pi_{\mu\nu} \Pi_{\rho\sigma} \right] F(p^2) + \Pi_{\mu\nu} \Pi_{\rho\sigma} G(p^2)$$

$$\tilde{P}_M = \epsilon_{\mu\nu} P^\nu$$

$$\tilde{\Pi}_{\mu\nu} = \tilde{P}_\mu \tilde{P}_\nu$$

$$\epsilon_{\mu\alpha} \epsilon_{\nu\beta} = \eta_{\mu\nu} \eta_{\alpha\beta} - \eta_{\mu\beta} \eta_{\alpha\nu}$$

In $D=2$

$$\langle T_{\mu\nu}(p) T_{\rho\sigma}(-p) \rangle = \tilde{P}_\mu \tilde{P}_\nu \tilde{P}_\rho \tilde{P}_\sigma G(p^2) = \epsilon \frac{\tilde{P}_\mu \tilde{P}_\nu \tilde{P}_\rho \tilde{P}_\sigma}{p^2}$$

- Scale symmetry.

$$\tilde{T}_{\mu\nu} = \int d^D x e^{i p \cdot x} T_{\mu\nu}(x)$$

$$\tilde{P}_\mu = \epsilon_{\mu\nu} P^\nu$$

$$\tilde{\Pi}_{\mu\nu} = \tilde{P}_\mu \tilde{P}_\nu$$

$$\epsilon_{\mu\alpha} \epsilon_{\nu\beta} = \eta_{\mu\nu} \eta_{\alpha\beta} - \eta_{\mu\beta} \eta_{\alpha\nu}$$

$$T_{\mu\nu}(p) T_{\rho\sigma}(q)$$

In D=2

$$\langle T_{\mu\nu}(p) T_{\rho\sigma}(-p) \rangle = \tilde{P}_\mu \tilde{P}_\nu \tilde{P}_\rho \tilde{P}_\sigma G(p^2) = \epsilon \frac{\tilde{P}_\mu \tilde{P}_\nu \tilde{P}_\rho \tilde{P}_\sigma}{p^2}$$

- Scale symmetry:

$$\tilde{T}_{\mu\nu} = \int d^D x e^{i p \cdot x} T_{\mu\nu}(x)$$

$$\langle T_{\mu\nu}(p) T_{\rho\sigma}(-p) \rangle = c \frac{\hat{P}_\mu \hat{P}_\nu \hat{P}_\rho \hat{P}_\sigma}{p^2}$$

$$\eta^{\mu\nu} \tilde{P}_\mu \hat{P}_\nu = p^2$$

Naively $T_{\mu}^{\mu} = 0$

$$\langle T_{\mu}^{\mu}(p) T_{\rho\sigma}(-p) \rangle = c \tilde{P}_\rho \tilde{P}_\sigma \quad \text{conformal anomaly}$$

$$\langle T_{\mu\nu}(x) T_{\rho\sigma}(y) \rangle =$$

$$\langle T_{\mu}^{\nu}(0) \rangle_h = \langle T_{\mu}^{\nu}(0) e^{\int h_{\rho\sigma} T^{\rho\sigma} d^2x} \rangle$$

$$= \int d^2x \langle T_{\mu}^{\nu}(0) T_{\rho\sigma}(x) \rangle h^{\rho\sigma}(x)$$

$$R|_{g=\eta+h}$$

$$= c (\partial_{\rho} \partial_{\sigma} - \square \eta_{\rho\sigma}) h^{\rho\sigma} = c R|_h$$

$$T_{\mu}^{\nu} = c R$$

- In odd spacetime dimension, ∇ conformal anomalies

$$T^{\mu}_{\mu} = 0$$

- $D=4$

$$a_W > a_{IR}$$

RG monotonicity.

$$T^{\mu}_{\mu} = a \text{ Euler} + c \text{ Weyl}^2$$

$$\sum_{i=0}^3 (-1)^i b_i = \chi = \int_{M_4} \text{Euler}$$

A congruence analysis

$\mathbb{Q} \subset \mathbb{R}$

$\mathbb{R} \subset \mathbb{C}$ monotonicity.

$\mathbb{C} \xrightarrow{2}$ Weyl

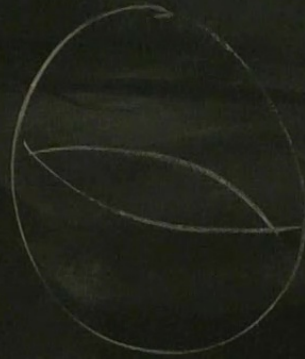
$\dim H^1 = \{ d\alpha = 0 \text{ but not exact } \alpha = d\beta. \}$

$$a_{UV} > a_{IR}$$

$$\chi = a \text{ Euler} + c \text{ Weyl}^2$$

$$\chi = \int_{M_4} \text{Euler}$$

RG monotonicity.



$$b_0 = 1$$

$$b_1 = b_2 = b_3 = 0$$

$$b_4 = 1$$

$\dim H^i = \begin{cases} d\alpha = 0 \text{ but not} \\ \text{exact } \alpha = d\beta. \end{cases}$

$$\chi = 2$$

$$\chi(T^4)$$