

Title: Horizons and Null Infinity: A Fugue in 4 voices

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Abstract: Horizons of black holes in equilibrium and null infinity of asymptotically flat space-times are null 3-manifolds but have very different physical connotations. We first show that they share a large number of geometric properties, making them both weakly isolated horizons. We then use this new unified perspective to unravel the origin of the drastic differences in the physics they contain. Interestingly, the themes are woven together in a manner reminiscent of voices in a fugue. The talk is based on joint work with Simone Speziale (arXiv:2401.15618 and arXiv 2401:????).

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Zoom link

Walk with Simone Speziale

# HORIZONS & NULL INFINITY: A fugue in 4 voices

- 0 Introduction
- 1 Geometric WIHs
- 2 BH/Cosmological WIHs
- 3 Asymptotic WIHs:  $f^+$
- 4 Symmetries & associated fluxes and charges

0: GR, Null Surfaces — Horizons

Common: Intrinsic derivative operator  $\bar{D}$

Geometrical WIHs:  $(\bar{q}_{ab}, \bar{k}^a)$  } Geo  
} WIHs

NEH (Non-expanding horizons):  $\bar{D}_a \bar{k}^a = 0$  (Expansion-free)

Null, 3-d surface  $\mathbb{R}^2 \times \mathbb{R}$

shear variables  $\bar{D}_a \bar{q}_{bc} = 0$

ruled null normal  $\bar{k}^a$

$\bar{D}_a \bar{q}_{bc} = 0$  ;  $\bar{D}_a \bar{k}^b = \bar{\omega}_a \bar{k}^b$

$\bar{D}_a \bar{k}^a = 0$  ;  $\bar{D}_a \bar{k}^b = \bar{\omega}_a \bar{k}^b$

1-form  $(\bar{D}, \bar{k})$

$\bar{k}^a \rightarrow \bar{k}^b = f \bar{k}^a$

$\bar{\omega}_a \rightarrow \bar{\omega}_a + \bar{D}_a \ln f$

Restriction:

$\bar{k}^a \bar{D}_a \bar{k}^b = 0$  ;  $F \equiv f(\omega, \theta)$

$\bar{\omega}_a$ : divergence, then essential the rescaling freedom

$\bar{k}^a \rightarrow \bar{k}^a \hat{=} c \bar{k}^a$

$\rightarrow$  choose  $[\bar{k}^a]$ ,  $\bar{k}^a \hat{=} c \bar{k}^a$

st.  $\bar{\omega}_a$  is divergence

$\Rightarrow \mathcal{L}_{\bar{k}} \bar{\omega}_a = 0$

Restrictions:

$$\rightarrow \bar{k}^a \bar{D}_a \bar{k}^b = 0; \quad F \equiv f(\rho, \theta)$$

$\bar{\omega}_a$ : divergence, then essentially eliminates the rescaling freedom.

$$\bar{k}^a \rightarrow \bar{k}^a \hat{=} c \bar{k}^a$$

$\rightarrow$  choose  $[\bar{k}^a]$ ,  $\bar{k}^a \hat{=} c \bar{k}^a$ ;  $c > 0$   
 s.t.  $\bar{\omega}_a$  is divergence.

$$\Rightarrow \mathcal{L}_{\bar{k}} \bar{\omega}_a = 0$$

NEH  $\int \mathcal{H}$  equipped with  $[\bar{k}^a]$  s.t.  
 $\mathcal{L}_{\bar{k}} \bar{\omega}_a = 0$  is a WH.

Is  $\bar{D}$  time dependent?

$$(\mathcal{L}_{\bar{k}} \bar{D}_a - \bar{D}_a \bar{k}^b) \bar{k}^c = \dot{\bar{D}}_a \bar{k}^c$$

If  $\bar{h}_a \bar{k}^a \hat{=} 0$  then  $\dot{\bar{D}}_a \bar{h}^a = 0$

~~$(\mathcal{L}_{\bar{k}} \bar{D}_a - \bar{D}_a \bar{k}^b) \bar{k}^c = 0$~~   
 ~~$\forall \bar{k}^b$  s.t.  $\bar{h}^b$~~   
 Isolated Horizons

suffices to calculate

$$\underbrace{\dot{\bar{D}}_a \bar{J}_b}_{\text{Time-dependence}}$$

$$\bar{J}_b \bar{K}^b = -1$$

$$\begin{aligned} \dot{\bar{D}}_a \bar{J}_b &= \bar{D}_a \bar{\omega}_b + \bar{\omega}_a \bar{\omega}_b \\ &+ \bar{k}^c \bar{C}_{cab} \bar{J}_d + \frac{1}{2} (\bar{S}_{ab} + \bar{\kappa} \bar{g}_{ab}) \bar{J}_d \\ &\quad \underbrace{\bar{k}^c \bar{R}_{cab} \bar{J}_d}_{\bar{R}_{ab} - \frac{1}{6} \bar{R} \bar{g}_{ab}} \end{aligned}$$

Restrictions:

$$\rightarrow \bar{K}^a \bar{D}_a \bar{K}^b = 0; \quad F \equiv f(\rho, \theta)$$

$\bar{\omega}_a$ : divergence, then essentially eliminates the rescaling freedom.

$$\bar{K}^a \rightarrow \bar{K}'^a = c \bar{K}^a$$

$\rightarrow$  choose  $[\bar{K}^a]$ ,  $\bar{K}^a \approx c \bar{K}'^a$ ;  $c > 0$   
 st.  $\bar{\omega}_a$  is divergence.

$$\Rightarrow \bar{D}_a \bar{\omega}^a = 0$$

NEH  $\int h$  equipped with  $[\bar{K}^a]$  st.  $\bar{D}_a \bar{\omega}^a = 0$  is a WIH.

Is  $\bar{D}$  time dependent?

$$(\bar{D}_a \bar{D}_a - \bar{D}_a \bar{D}_a) \bar{K}^a = \bar{D}^a \bar{K}^a$$

If  $\bar{h}_a \bar{K}^a \neq 0$  then  $\bar{D}_a \bar{h}^a = 0$

~~$(\bar{D}_a \bar{D}_a - \bar{D}_a \bar{D}_a) \bar{K}^a = 0$~~   
 ~~$\bar{K}^a$  not to  $\bar{h}$~~   
 Isolated Horizons

suffices to calculate

$$\underbrace{\bar{D}_a \bar{J}_b}_{\text{Time-dependence}}$$

$$\bar{J}_b \bar{K}^b = -1$$

$$\begin{aligned} \bar{D}_a \bar{J}_b &= \bar{D}_a \bar{\omega}_b + \bar{\omega}_a \bar{\omega}_b \\ &+ \bar{K}^c \bar{C}_{cab} \bar{J}_d + \frac{1}{2} (\bar{K}^c \bar{S}_{ab} + \bar{K}^c \bar{R}_{ab} - \frac{1}{6} \bar{R} \bar{g}_{ab}) \bar{J}_d \end{aligned}$$

with radiation  
 $\lambda = 0$

2. BH (cosmological horizon)

- $h \rightarrow \Delta$ : WIH for BH or cosmo
- $\bar{K}^a \rightarrow \bar{K}^a$ : normal
- $\bar{D} \rightarrow D$

calculate

$\bar{J}_b$  dependence

$\bar{J}_b \bar{K}^b = -1$

$\bar{D}_a \bar{w}_b + \bar{w}_a \bar{w}_b$

$+ \bar{K}^c \bar{C}_{cab} \bar{J}_d + \frac{1}{2} \bar{K}^c \bar{g}_{ab} + \dots$

$\bar{K}^c \bar{R}_{cab} \bar{J}_d$

$\Sigma$ : BH (cosmological horizon)

$h \rightarrow \Delta$ : WH for BH or cosmo

$\bar{K}^a \rightarrow \ell^a$ : normal

$\bar{D} \rightarrow D$

$\bar{J}_b \rightarrow \bar{K}_a, \bar{K}_a \bar{K}^a = -1$

$\bar{D}_a \bar{K}_b = \bar{D}_a \bar{w}_b + \bar{w}_a \bar{w}_b + \underbrace{\bar{K}^c \bar{C}_{cab} \bar{J}_d}_{-\frac{1}{2} R \bar{g}_{ab} - D_{ca} \bar{w}_b}$

$-\frac{1}{2} R \bar{g}_{ab} - D_{ca} \bar{w}_b$

Ricci scalar of  $\bar{g}_{ab}$   
2-d object

$\bar{D}_e R = 0, \bar{D}_e \bar{w}_a = 0, \bar{w}_a \bar{K}^a = 0$

$\bar{w}_a$ : horizontal 1-form

RHS is time independent

$\Rightarrow \bar{D}_a$ : completely determined by data  $(\bar{q}_{ab}, \bar{w}_a)$

$\Rightarrow$  NO 3-d DOF: NO Radiation!!

$$\dot{D}_a h_b = \dot{D}_a \bar{w}_b + \bar{w}_a \bar{w}_b + \underbrace{\frac{1}{2} C_{0ab}^d h_d}_{\substack{\text{Ricci scalar of } g_{ab} \\ \text{2-d object}}}$$

$$-\frac{1}{4} R g_{ab} - D_{[a} \omega_{b]}$$

$$\dot{\Delta}_e R = 0, \quad \dot{\Delta}_e \omega_a = 0, \quad \omega_a \dot{\Delta}^a = 0$$

$\omega_a$ : horizontal 1-form

RHS is time independent

$\Rightarrow D_a$ : completely determined by data  $(\tilde{q}_{ab}, \tilde{\omega}_a)$   
 $\Rightarrow$  No 3-d DoF: NO Radiation!!

### 3. Asymptotic WIH: $\mathcal{I}^+$

$(M, g_{ab})$ : Physical:  $AF @ \mathcal{I}^+$ ,  $\exists \Omega > 0$  on  $M$  st.  $\hat{g}_{ab} = \Omega^2 g_{ab}$   
 is smooth @  $\Omega = 0$  ( $\mathcal{I}^+$ ), st.  $\hat{\nabla}_a \Omega \neq 0$ ,  $\Omega^{-1} T_{ab}$  has a limit

Freedom:  $\Omega \rightarrow \Omega' = \omega \Omega$ , w. no where zero, use it with field eqns (with  $T_{ab}$ )  
 (DIV free  $\mathcal{I}^+$ )  $\hat{\nabla}_a \hat{h}^a = 0$ , field eqn  $\Rightarrow \hat{\nabla}_a \hat{h}_b = 0$

surfaces to calculate

$$\hat{D}_a \bar{J}_b$$

Time-dependence

$$\bar{J}_b \bar{K}^b = -1$$

$$\hat{D}_a \bar{J}_b = \bar{D}_a \bar{W}_b + \bar{w}_a \bar{w}_b + \bar{K}^c \bar{C}_{cab} \bar{J}_d + \frac{1}{2} (\bar{s}_{ab} + d \bar{q}_{ab})$$

with radiation  
 $\lambda = 0$

BH (cosmological horizon)

- $\hat{h} \rightarrow \Delta$  : WIH for BH or cosmo
- $\bar{K}^a \rightarrow \bar{K}^a$  : normal
- $\bar{D} \rightarrow D$
- $\bar{J}_b \rightarrow \bar{K}^a_{;a}, \bar{K}^a_{;a} = -1$

3 contd:

$$\hat{D}_a \hat{h}^b \leq 0 \Rightarrow \hat{\Theta}_{\hat{h}} = 0$$

$$\bar{D}_a \hat{h}^b \leq 0 \Rightarrow \hat{W}_a \leq 0$$

$$\hat{K}^a \hat{W}_a = 0$$

$$\Rightarrow (\hat{J}, \hat{h}^a, \hat{q}_{ab}, \hat{D})_{WIH}$$

$$\hat{K}^a \hat{q}_{ab} = 0$$

$$\hat{h}_b \hat{h}^b = 0$$

$$\hat{D}_a \hat{h}_b = 0$$

Physical information

$$\hat{D}_a \hat{h}_b = \frac{1}{2} (\hat{s}_{ab} + d \hat{q}_{ab})$$

Bondi conf frame:  $TF \hat{S}_{ab} = \hat{N}_{ab}$

$$TF \hat{D}_a \hat{h}_b = \hat{N}_{ab}$$

3-d local DOF  
 Grav. radiation!

$$\hat{h} \rightarrow \hat{J}$$

$$\bar{K}^a \rightarrow \hat{h}^a$$

$$\bar{D}_a \rightarrow \hat{D}_a$$

$$\bar{J}_a \rightarrow \hat{J}_a$$

At  $\hat{J}$ : due to  $S^2 \times \mathbb{R}$  topology we have  $\hat{C}_{abcd} \leq 0$ .

$$\int_{\hat{J}} \hat{C}_{abcd} = \hat{K}_{abcd}$$

$$\hat{K}_{abcd} = \int_{\hat{J}} \psi_a^+ \psi_b^+ \psi_c^- \psi_d^-$$

calculate

$$\bar{J}_b \bar{K}^b = -1$$

since

$$\bar{D}_a \bar{w}_b + \bar{w}_a \bar{w}_b$$

$$\bar{K}^c \bar{C}_{cab} \bar{J}_d + \frac{1}{2} (\bar{q}_{ab} + \alpha \bar{q}_{ab})$$

$$\bar{K}^c \bar{R}_{cab} \bar{J}_d + \frac{1}{2} \bar{R}_{ab} - \frac{1}{2} \bar{R} \bar{g}_{ab}$$

z: BH (cosmological horizon)

h → Δ: WIH for BH or cosmo

R<sup>a</sup> → t<sup>a</sup>: normal

D → D

J<sub>b</sub> → K<sub>a</sub>, K<sub>a</sub> t<sup>a</sup> = -1

3 contd:

$$\hat{D}_a \hat{n}^b \leq 0 \Rightarrow \hat{O}_a = 0$$

$$\hat{D}_a \hat{n}^b \leq 0 \Rightarrow \hat{w}_a \leq 0$$

$$\hat{D}_a \hat{n}^b = 0$$

$$\Rightarrow (t, \hat{n}^a, \hat{q}_{ab}, \hat{D}) \text{ WIH}$$

$$\hat{D}_a \hat{n}^b = 0$$

$$\hat{D}_a \hat{h}_b = 0$$

Physical information

Bondi conf frame: TF S<sub>ab</sub> = N<sub>ab</sub>

4: symmetry

WIH: universal structure

At J<sub>+</sub>

group B

$$\hat{D}_z \hat{q}_{ab} = 2\beta \hat{q}_{ab}$$

$$\hat{D}_z \hat{n}^a = -\beta \hat{n}^a$$

$$h \rightarrow J$$

$$R^a \rightarrow \hat{n}^a$$

$$\bar{D}_a \rightarrow \hat{D}_a$$

$$\bar{J}_a \rightarrow \hat{t}_a$$

At J<sub>+</sub> due to S<sup>2</sup> x R topology we have C<sub>abcd</sub> = 0

$$\int_{S^2} \hat{C}_{abcd} = \hat{K}_{abcd}$$

$$\hat{D}_a \hat{t}_b = \frac{1}{2} (\hat{q}_{ab} + \alpha \hat{q}_{ab})$$

3d local DOF  
Grav: radiation

group G = B x D

generator

$$\hat{D}_d \hat{q}_{ab} = 0$$

$$\hat{D}_d \hat{K} = c \hat{K}$$

General  $\sum \in G$  Lie's

$$\hat{D}_z \hat{q}_{ab} = (2\beta + \gamma) \hat{q}_{ab}$$

$$\hat{D}_z \hat{K} = (\beta + \epsilon) \hat{K}$$

$$\hat{D}_a \hat{h}_b = \hat{D}_a \hat{h}_b$$

$$\Rightarrow \hat{D}_a \Rightarrow \text{No}$$

3. Asymptotic

(M, g<sub>ab</sub>): P is smooth