

Title: Twisted Holography Mini-Course - Lecture 20231109

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Collection: Twisted Holography Mini-Course

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X CY3

$$\mu \in \Omega^{0,1}(X, TX) = \Omega^{0,1}(X, \wedge^2 T^*X) = \Omega^{2,1}(X)$$

$$\mu_j^i dz_i \otimes \frac{\partial}{\partial z_j}$$

Constraint: $\partial \mu = 0$

Lagrangian:

$$\frac{1}{2} \int \bar{\partial} \mu \partial^{-1} \mu + \frac{1}{6} \int \mu^3$$

$$\partial \mu \in \Omega^{3,1}(X) = \Omega^{0,1}(X)$$

$$\mu^3 \in \Omega^{0,3}(X, \wedge^3 T^*X) = \Omega^{0,3}(X) = \Omega^{3,3}(X)$$

$$\partial_{z_i} \vee d^3z = \epsilon_{ijk} dz_j dz_k$$

KS theory has 2 other fields, not as important:

$$\alpha_i \in \bar{\Omega}^{0,1}(X) \quad i=1,2 \quad \text{fermionic } (0,1) \text{ forms}$$

$$\int \epsilon^{ij} \alpha_i \bar{\partial}_\mu \alpha_j \quad X, CY3, \text{ spinors on } X$$

$$S_+ = \Omega^{0, \text{ev}} \quad S_- = \Omega^{0, \text{odd}}$$

Spinors in physical theory, $\not{D} = \not{\partial} + \not{\partial}^*$

Gauge Symmetry

$$V \in \Omega^{0,0}(X, TX)$$

$$V = V_i \frac{\partial}{\partial z^i}$$

vector field pointing in $\frac{\partial}{\partial z}$ directions

Lie derivative

$$\delta \mu = [V, \mu] \circ \bar{\partial} V = [V, \bar{\partial} \mu]$$

generates gauge transformations.

Ex: $X = (\mathbb{C}^x)^3$ where $\mathbb{C}^x = \mathbb{C}/\mathbb{Z} = S^1 \times \mathbb{R}$

Then KK reduction to \mathbb{R}^3 along $(S^1)^3$

Result is 3d CS theory for group
of volume-preserving diffeos of $(S^1)^3$

$$\mu = d\bar{z}_j \wedge \bar{\mu}_j; d z_j$$

$d\bar{z}_j \mapsto 3$ components of 3d gauge field

$V \mapsto$ CS gauge transformations

Branes

$X \subset \mathbb{C}P^3$

"Branes are coherent sheaves on X "

Concretely, brane is

$Y \subseteq X$, Y holomorphic submanifold

E a vector bundle on Y (holomorphic)

Physical String

$$\mathbb{R}_{\varepsilon}^2 \times \mathbb{R}_{-\varepsilon}^2 \times \underline{\mathbb{C}} \times \underline{\mathbb{C}} \times \underline{\mathbb{C}}$$

D7

D5

D3

D1

—

—

—

— —

—

Top' string

Space filling

D5

D3 wraps \mathbb{C}^2

D1 wraps \mathbb{C}

D-1 wraps \emptyset

Branes have gauge fields
also couple to supergravity

D3 in physical string

$F_5 \in \Omega^5(\mathbb{R}^{10})$ 5-form field strength

$$\star F_5 = F_5, \quad dF_5 = 0$$

$$\int_{D3} d^4 F_5 + \dots \quad (\text{determined by SUSY})$$

In \mathbb{R}^3 , F_5 and g are related by SUSY

Can think of $\mu \in \Omega^{2,1}$ as coming from F_5 instead

Roughly, $\mu = \int_{\mathbb{R}^2_\varepsilon} F_5$

Propose (can derive from worldsheet) μ couples to D1 wrapping $C \subseteq X$ $\dim_{\mathbb{C}} C = 1$

by $\int_C \bar{\partial} \mu$

$\bar{\partial} \mu \in \Omega^{1,1}(X)$

Physical String

Backreact by look at field sourced by D3

EOM in presence of N D3's

$$N \int_{D3} d^4 F_5 + \int F_5^2$$

EOM are

$x_0 - x_3$ D3

$y_4 - y_9$

$$dF_5 = N \delta_{D3}$$

$$\int_{\mathbb{R}^4} F_5 = N$$

\rightsquigarrow $AdS_5 \times S^5$
5-form F
 $\int_{S^5} F = N$

Prop Energy morphism in Bord, can be built from

Topl String N D1 branes then the eqn for μ is

$$\bar{\partial}\mu = \delta_{D1} N$$

$Z \rightsquigarrow$ D1 worldvolume
 $w_1, w_2 \rightsquigarrow$ normal.

Solution to this eqn is Böchner-Martinelli kernel.

$$\mu = N \frac{\epsilon^{ij} \bar{w}_i d\bar{w}_j dw_1 dw_2}{\|w\|^4}$$

Check 1) $\bar{\partial}\mu = 0$ except at $w_i = 0$

$$2) \int_{\|w\|=1} \mu = N(\pi^2)$$

CAUTION

NO SMOKING OR OPEN FLAMES
NO DRINKS OR FOODS
NO BOTTLES OR GLASSES
NO SHARP OBJECTS
NO WEAPONS
NO FIREARMS
NO LASERS
NO TOBACCO
NO ALCOHOL
NO DRUGS
NO RECREATIONAL DRUGS
NO WEAPONS
NO FIREARMS
NO LASERS
NO TOBACCO
NO ALCOHOL
NO DRUGS
NO RECREATIONAL DRUGS

How does geometry change?

$$\mu = N \frac{\varepsilon^{ij} \bar{\omega}_i d\bar{\omega}_j}{\|w\|^4} \quad \partial_z$$

Hol. fns in deformed geometry satisfy

$$\bar{\partial}F + N \frac{\varepsilon^{ij} \bar{\omega}_i d\bar{\omega}_j}{\|w\|^4} \frac{\partial F}{\partial z} = 0$$

$$\Rightarrow \begin{cases} \frac{\partial F}{\partial \bar{\omega}_2} + N \frac{\bar{\omega}_1}{\|w\|^4} \frac{\partial F}{\partial z} = 0 \\ \frac{\partial F}{\partial \bar{\omega}_1} - \frac{N \bar{\omega}_2}{\|w\|^4} \frac{\partial F}{\partial z} = 0 \end{cases}$$

CAUTION

DO NOT USE THIS BOARD FOR
RECORDING IN THE OFFICE OR THE
LABORATORY OR OTHER
HIGH-SENSITIVITY AREAS.
PLEASE REPORT ANY
DAMAGE IMMEDIATELY.

2 solⁿs are $F = \omega_1$, $F = \omega_2$

2 more are

$$u_1 = \omega_1 z + \frac{N \bar{\omega}_2}{\|w\|^2}$$

$$u_2 = \omega_2 z - \frac{N \bar{\omega}_1}{\|w\|^2}$$

4 coords on a 3d geometry. 1 relation

$$\omega_1 u_2 - \omega_2 u_1 = -N$$

= 2x2 matrices $\begin{pmatrix} \omega_1 & u_1 \\ \omega_2 & u_2 \end{pmatrix}$ of determinant $-N$

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EOM are

$$dF_5 = N \delta_{D3}$$

$x_0 - x_3$ D3

$y_4 - y_9$

$$\int_{||y||=1} F_5 = N$$

$$\int_{||y||=1} F = \int_{||y|| \leq 1} dF = N$$

\rightsquigarrow AdS₅ × S⁵
5-form F
 $\int_{S^5} F = N$

Backreaction gives $SL_2(\mathbb{C}) = SO(3,1)$

$$= \mathbb{H}^3 \times S^3 \simeq \text{AdS}_3 \times S^3$$

as $SO(3,1) = \text{isometries of } \mathbb{H}^3$

S^3 bundle over \mathbb{H}^3

$$\text{AdS}_3 \times S^3 \subseteq \text{AdS}_5 \times S^5$$

\mathbb{B} -model

localities $\mathbb{I}\mathbb{B}$