

Title: Relativity Lecture - 110723

Speakers: David Kubiznak

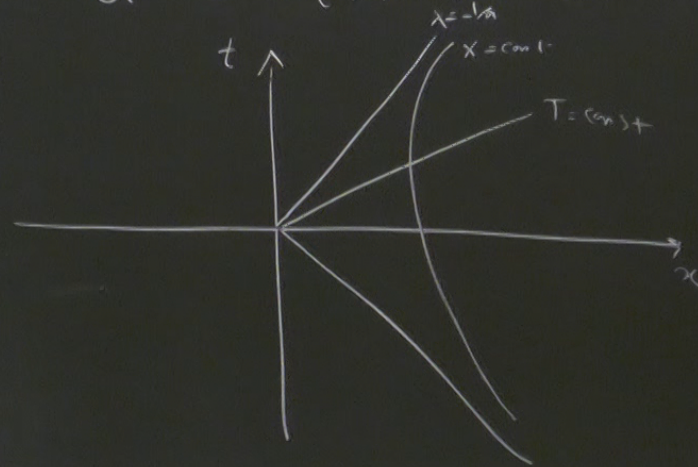
Collection: Relativity 2023/24

Date: November 07, 2023 - 9:00 AM

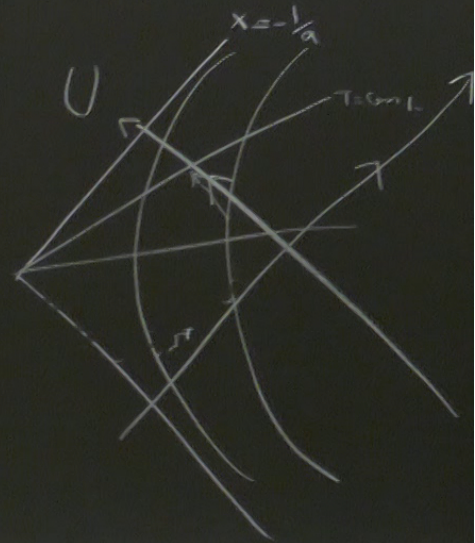
URL: <https://pirsa.org/23110013>

Rindler space-time:

$$ds^2 = -(1+ax)^2 dt^2 + dx^2$$



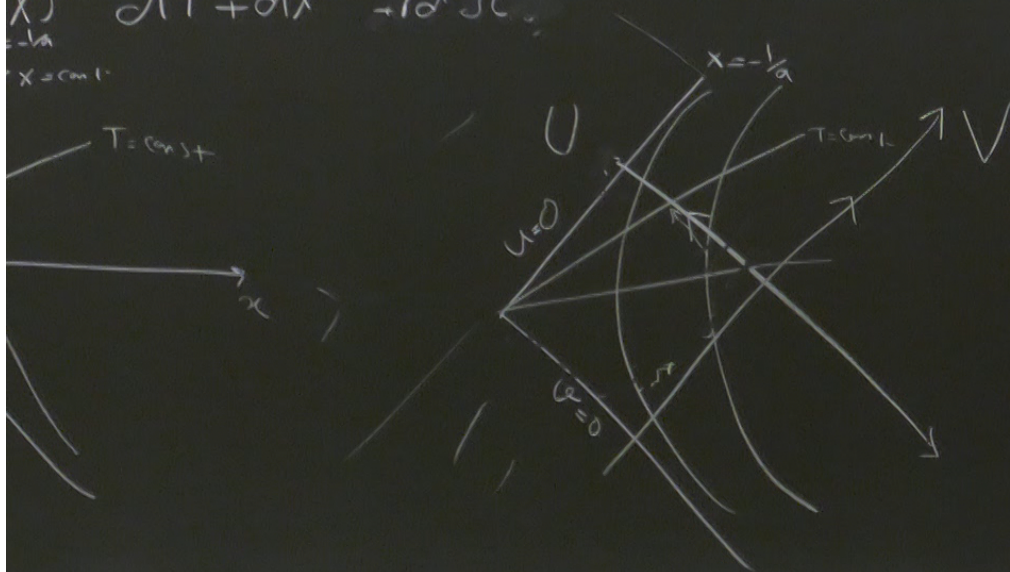
$$x \in \left(-\frac{1}{a}, +\infty\right)$$



space-time:

$$ds^2 = dt^2 + dx^2 + r^2 d\Omega^2$$

$$x \in (-1/a, +\infty)$$



$$ds^2 = -e^{a(v-u)} du dv + r^2 d\Omega^2$$

$$u = -\frac{1}{a} e^{-av} \quad v = \frac{1}{a} e^{au}$$

$$ds^2 = -du dv + r^2 d\Omega^2$$

$$u \in (0, +\infty) \rightarrow u \in (-\infty, +\infty)$$

$$v \in (-\infty, +\infty)$$

Schwarzschild Black hole:

$$ds^2 = -f_{(r)} dt^2 + f_{(r)}^{-1} dr^2 + r^2 d\Omega^2$$

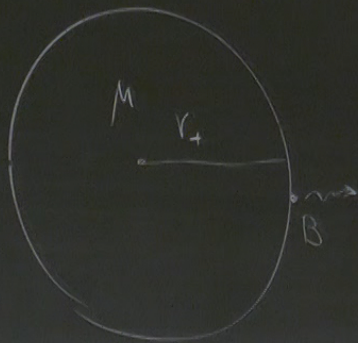
$$f(r) = 1 - \frac{2M}{r}$$

$$R_{\oplus} \sim 6400 \text{ km}$$

$$r_+ \sim 9 \text{ mm}$$

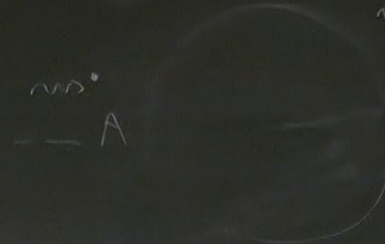
$$R_{\odot} \sim 8 \times 10^5 \text{ km}$$

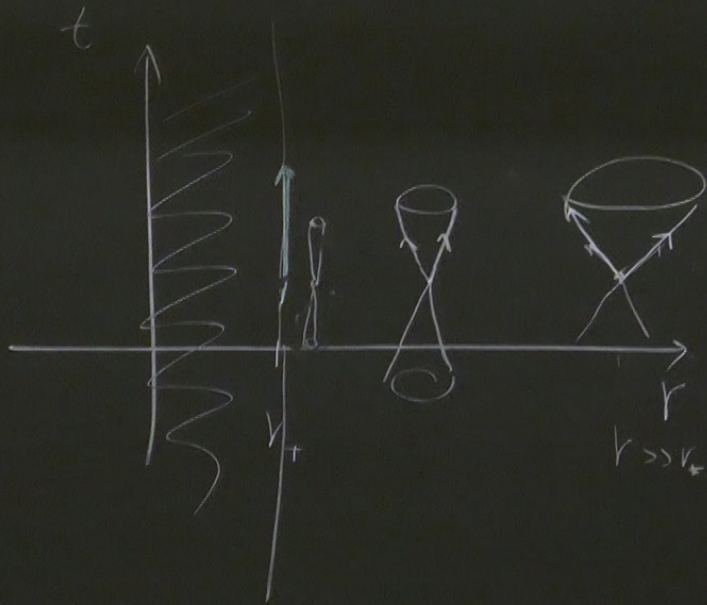
$$r_+ \sim 3 \text{ km}$$



$$r_+ = 2M$$

HBH





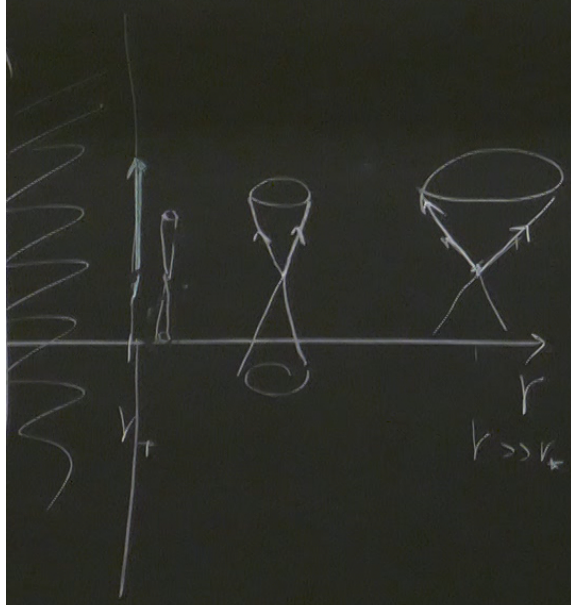
$$ds^2 = 0 \quad (ds^2 = 0)$$

$$dt^2 = \frac{1}{f(r)}$$

$$\frac{dt}{dr} = \pm \frac{1}{f(r)} = \pm \frac{1}{1 - \frac{2M}{r}}$$

$$R = g^{ab} R_{ab} = 0$$

$$R_{ab} R^{ab} = 0$$



$$ds^2 = 0 \quad (ds^2 = 0)$$

$$dt^2 = \frac{1}{f(r)} dr^2 \Rightarrow dt = \frac{1}{\sqrt{f(r)}} dr$$

$$\frac{dt}{dr} = \frac{1}{\sqrt{f(r)}} = \frac{1}{\sqrt{1 - \frac{2M}{r}}}$$

$$R = g^{ab} R_{ab} = 0$$

$$R_{ab} R^{ab} = 0$$

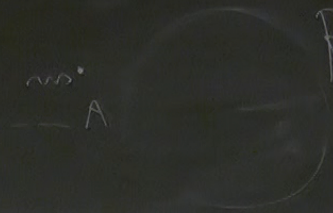
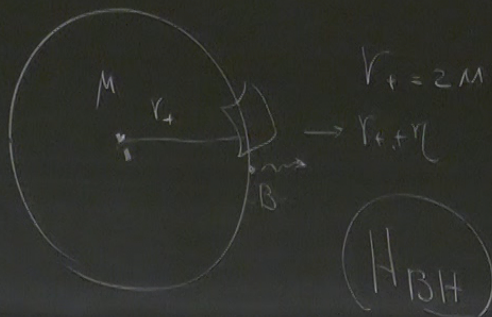
$$R_{\alpha\beta\gamma\delta} R^{\alpha\beta\gamma\delta} \rightarrow +\infty$$

$$r \rightarrow 0$$

Schwarzschild Black hole:

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

$$f(r) = 1 - \frac{2M}{r}$$

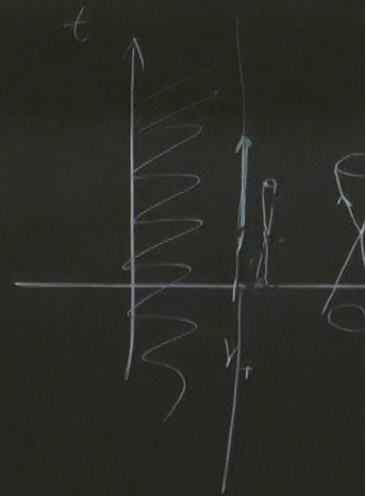


$$R_+ \sim 6400 \text{ km}$$

$$r_+ \sim 9 \text{ mm}$$

$$R_+ \sim 8 \times 10^5 \text{ km}$$

$$r_+ \sim 3 \text{ km}$$



$$R = g^{ab} R_{ab}$$

$$R_{mn} R^{mn} = 0$$

$$ds^2 = 0 \quad (d\Omega = 0)$$

$$dt^2 = \frac{1}{f^2} dr^2 \Rightarrow dt = \pm \frac{1}{f(r)} dr$$

$$\frac{dt}{dr} = \pm \frac{1}{f(r)} = \pm \frac{1}{1 - \frac{2M}{r}}$$



Kruskal Extension, 1960

$$\mathbb{R}_{\text{ext}} \times \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R} \times \mathbb{R}$$

$r \rightarrow 0$

External $M_{\text{Sch}} \rightarrow$

- Following Null geodesics
- Switching coordinates to affine parametrization

(H_{BH})

$V_+ \sim 3 \ln r$

Null ray trajectories:

$$dt = \pm \frac{dr}{f(r)}$$

$$\left\{ \begin{array}{l} t - r_{(r)}^* = U \\ t + r_{(r)}^* = V \end{array} \right. \quad \begin{array}{l} \text{Out-gang rays} \\ \text{in-gang rays} \end{array}$$

$$r^* = \frac{V - U}{2} \rightarrow r(U, V)$$

$$t = \pm \int_a^r \frac{dr}{f(r)} + \text{const} \quad \Rightarrow \quad t = \pm \underbrace{\left(r + 2M \ln \left| \frac{r}{2M} - 1 \right| \right)}_{r_{(r)}^*} + \text{const}$$

$$\dots \rightarrow ds^2 = -f(r) dU dV + r^2 d\Omega^2$$

\downarrow
 $1 - \frac{2M}{r}$

HBH

trajectories: $dt = \pm \frac{dr}{f(r)}$ $t = \pm \int_a^r \frac{dr}{f(r)} + \text{const} \Rightarrow t = \pm$

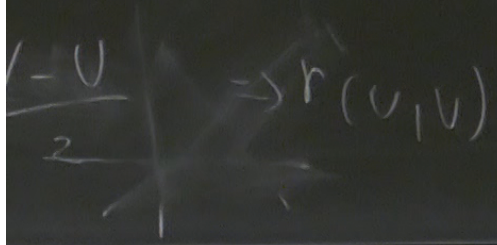
U Out-going rays

$$\rightarrow ds^2 = -f(r) dU dV + r^2 d\Omega^2$$

V in-going rays

$$\downarrow$$

$$1 - \frac{2M}{r}$$



$$r \rightarrow r_+ \Rightarrow \begin{cases} U \rightarrow +\infty \\ V \rightarrow -\infty \end{cases}$$

$$\int_a^r \frac{dr}{f(r)} + \text{const} \Rightarrow t = \underbrace{-t}_{r^*(r)} \left(r + 2M \ln \left| \frac{r}{2M} - 1 \right| \right) + \text{const}$$

$$> ds^2 = -f(r) dU dV + r^2 d\Omega^2$$

$$\downarrow$$

$$1 - \frac{2M}{r}$$

$$(U, 0, 0, 0)$$

$$\ln |x|$$

$$x \rightarrow 0$$

$+\infty$

$-\infty$

Next: affine parametrization.

$$u = -e^{-\frac{v}{4M}} \quad v = e^{\frac{u}{4M}}$$

$$ds^2 = -\frac{32M^3}{r^2} e^{-\frac{r}{2M}} du dv + r^2 d\Omega^2$$

$$u \rightarrow 0, \quad v \rightarrow 0 \\ r \rightarrow r_+, \quad r \rightarrow r_+$$

$$u \in (0, +\infty)$$

$$v \in (0, +\infty)$$

$$\rightarrow \begin{aligned} u &\in (-\infty, +\infty) \\ v &\in (-\infty, +\infty) \end{aligned}$$

$$dt = \pm \frac{dr}{f(r)} \quad t = \pm \int_a^r \frac{dr}{f(r)} + \text{const} \Rightarrow t = \pm \left(r + 2M \ln \left| \frac{r}{2M} \right. \right)$$

rays

$$\dots \rightarrow ds^2 = -f(r) dU dV + r^2 d\Omega^2$$

r^*
(r)

ing rays

$$\downarrow$$

$$1 - \frac{2M}{r} \quad (U, 0, 0, 0)$$

$$\left. \begin{array}{l} \leftarrow +\infty \\ \leftarrow 0 \\ \leftarrow -\infty \end{array} \right\} \begin{array}{l} U \rightarrow -\infty \\ r \rightarrow \infty \\ V \rightarrow \end{array}$$

$$\begin{array}{l} r_h \rightarrow -\infty \\ r \rightarrow r_+ \end{array} \Rightarrow \begin{cases} U \rightarrow +\infty \\ V \rightarrow -\infty \end{cases}$$

$$dt = \pm \frac{dr}{f(r)}$$

$$t = \pm \int_a^r \frac{dr}{f(r)} + \text{const} \Rightarrow t = \pm \left(\right)$$

Out-going rays

$$\rightarrow ds^2 = -f(r) dU dV + r^2 d\Omega^2$$

in-going rays

$$r_0 \rightarrow -\infty$$

$$r \rightarrow r_+$$

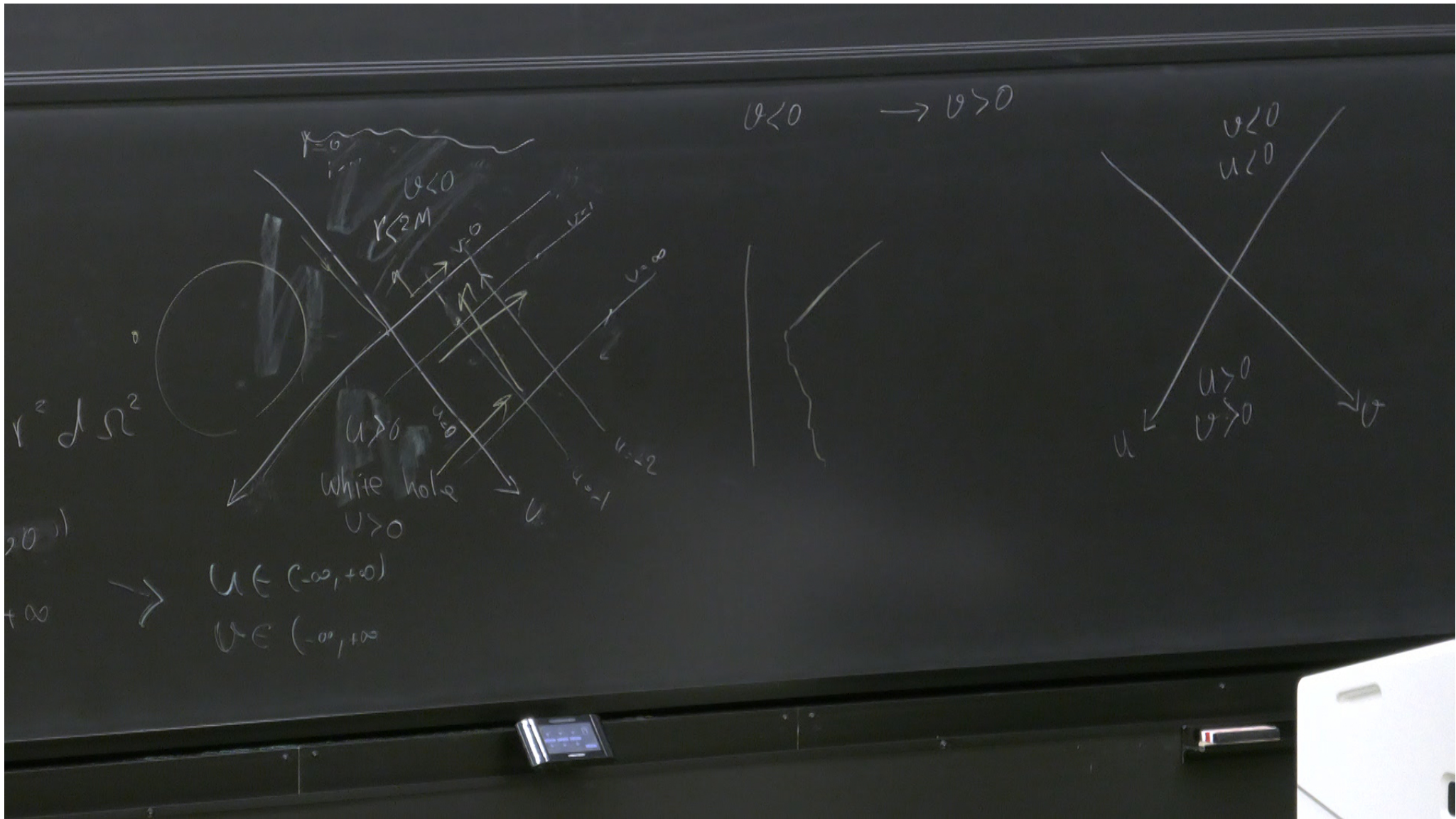
$$\Rightarrow \begin{cases} U \rightarrow +\infty \\ V \rightarrow -\infty \end{cases}$$

$$\downarrow$$

$$1 - \frac{2M}{r}$$

$$\left(U, V \right)$$

$$\left. \begin{array}{l} r \rightarrow r_+ \\ r \rightarrow \infty \end{array} \right\} \begin{array}{l} U \rightarrow -\infty \\ V \rightarrow +\infty \end{array}$$

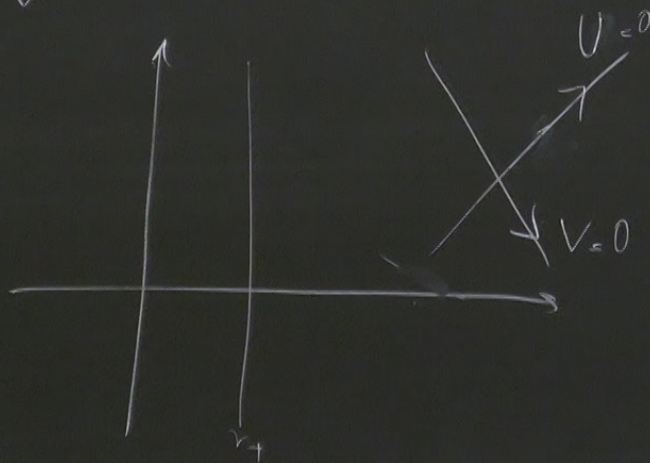


$$t - r^* = U \quad \Rightarrow \quad t = r^* + U = r + 2M \ln \left| \frac{r}{2M} - 1 \right| + U$$

$$t + r^* = V$$

$$r > 2M$$

ln ↗



$$V=0 \Rightarrow t = -r - 2M \ln \left(\frac{r}{2M} - 1 \right)$$



Null ray trajectories:

$$dt = \pm \frac{dr}{f(r)}$$

$$t = \pm \int_a^r \frac{dr}{f(r)} + \dots$$

$$\left\{ \begin{array}{l} t - r^* = U \\ t + r^* = V \end{array} \right.$$

Out-going rays

$$\left\{ \begin{array}{l} t - r^* = U \\ t + r^* = V \end{array} \right.$$

in-going rays

$$\dots \rightarrow ds^2 = -$$

$$r^* = \frac{V - U}{2} \Rightarrow r(U, V)$$

$$\begin{array}{l} r_h \rightarrow -\infty \\ r \rightarrow r_+ \end{array}$$

$$\Rightarrow \begin{cases} U \rightarrow +\infty \\ V \rightarrow -\infty \end{cases}$$