Title: Holography and its implications for quantum gravity - VIRTUAL

Speakers: Johanna Erdmenger

Collection: Puzzles in the Quantum Gravity Landscape: viewpoints from different approaches

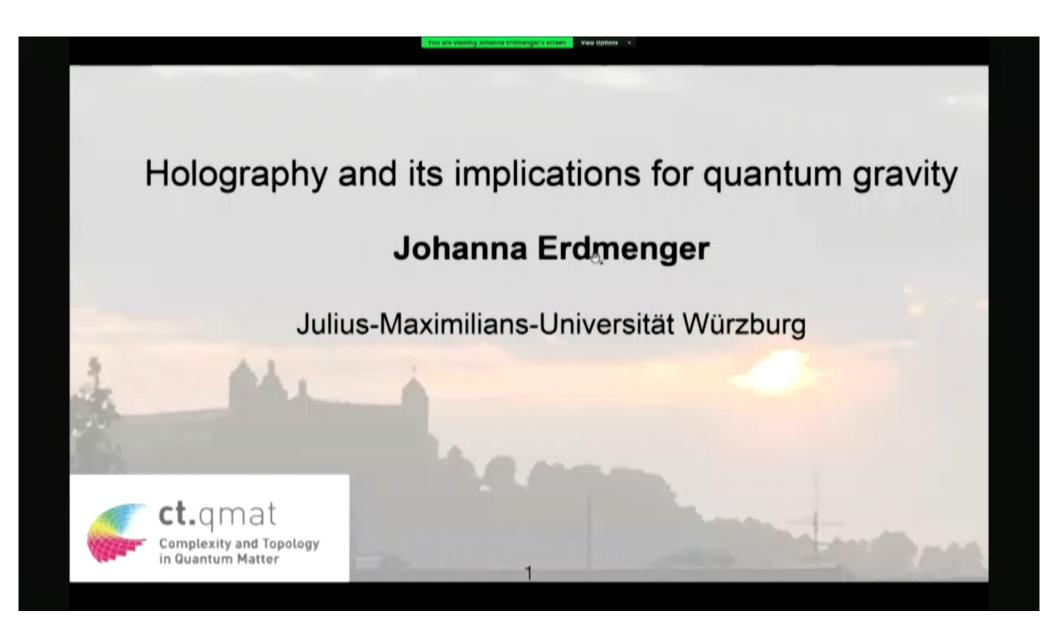
Date: October 23, 2023 - 10:20 AM

URL: https://pirsa.org/23100060

Abstract: I will review implications of the AdS/CFT correspondence for quantum gravity, paying particular attention to holographic RG flows and the processing of information. Deformed hyperbolic geometries correspond to boundary field theory RG flows. Conversely, as I will exemplify using the recent approach of discrete holography, boundary RG flows may be used to reconstruct the bulk spacetime, for instance via tensor networks. I will discuss recent approaches towards including gravity dynamics into this reconstruction process.

The presenter will be joining via Zoom for this talk.

Pirsa: 23100060 Page 1/39



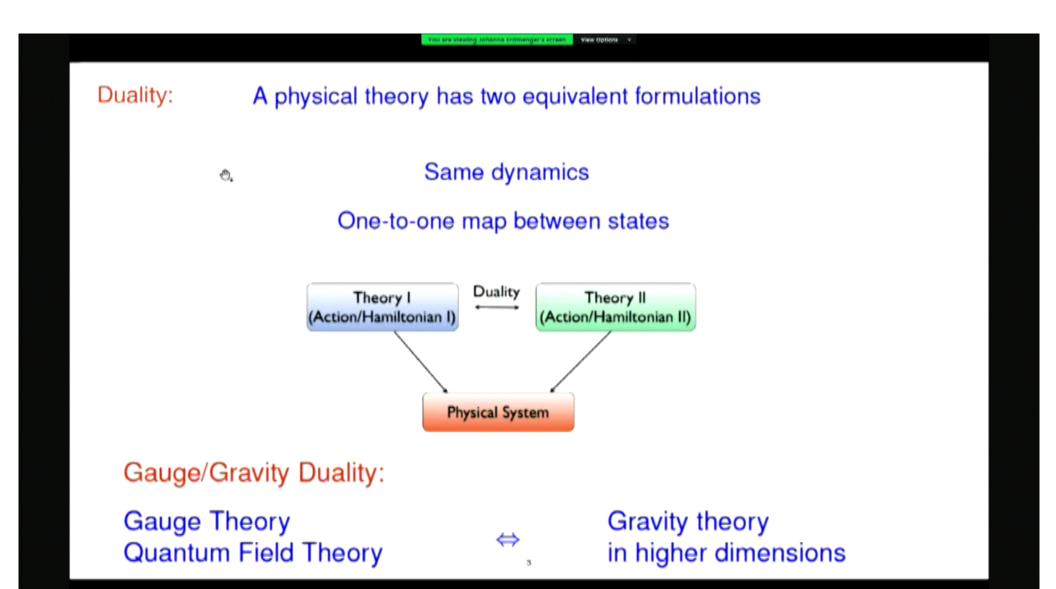
Pirsa: 23100060 Page 2/39

Overview

- I. AdS/CFT correspondence and renormalization group
- II. Information and gravity
- III. Discrete holography

.

Pirsa: 23100060 Page 3/39



Pirsa: 23100060 Page 4/39

Gauge/gravity duality

- Conjecture which follows from a low-energy limit of string theory
- Duality:

Quantum field theory at strong coupling

⇔ Theory of gravitation at weak coupling

₾.

Holography:

Quantum field theory in d dimensions

 \Leftrightarrow Gravitational theory in d+1 dimensions

4

Pirsa: 23100060 Page 5/39

AdS/CFT correspondence

AdS: Anti-de Sitter space, CFT: Conformal field theory

 $\mathcal{N}=4$ SU(N) Super Yang-Mills theory

Anti-de Sitter space:

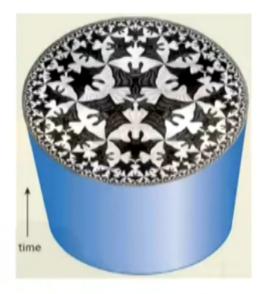
Space of constant negative curvature, has a boundary

In total: $AdS_5 \times S^5$

Isometries of AdS_{d+1} and symmetry of operators in a CFT_d :

SO(d+1,1)

(Euclidean signature)

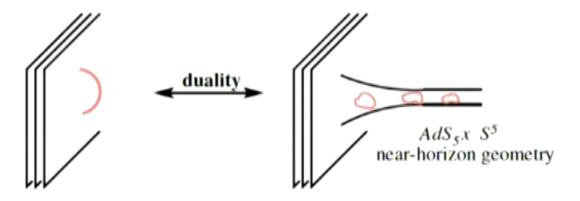


Quelle: Institute of Physics, Copyright: C. Escher

String theory origin of the AdS/CFT correspondence

D3 branes in 10d





↓ Low energy limit

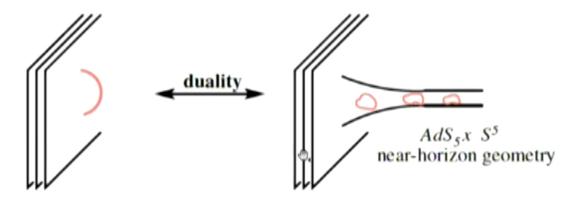
Supersymmetric SU(N) gauge theory in four dimensions $(N \to \infty)$

Supergravity on the space $AdS_5 \times S^5$

.

String theory origin of the AdS/CFT correspondence

D3 branes in 10d



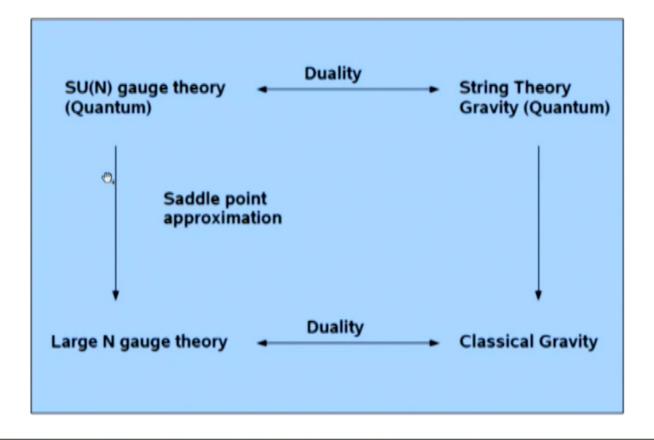
↓ Low energy limit

Supersymmetric SU(N) gauge theory in four dimensions $(N \to \infty)$

Supergravity on the space $AdS_5 \times S^5$

.

Gauge/Gravity Duality: Low-energy limit



Pirsa: 23100060 Page 9/39

Different levels of the AdS/CFT correspondence

$$g_{\mathrm{YM}}^2 = 2\pi g_s$$

$$\lambda = g_{
m YM}^2 N$$

$$2g_{\mathrm{YM}}^2 N = L^4/\alpha'^2$$

't Hooft coupling

	$\mathcal{N}=4$ SYM theory	IIB theory on $AdS_5 \times S^5$
Strongest form	any N and λ	Quantum string theory, $g_s \neq 0$, $\alpha' \neq 0$
Strong form	$N \to \infty$, λ fixed but arbitrary	Classical string theory, $g_s \to 0$, $\alpha' \neq 0$
Weak form	$N \to \infty$, λ large	Classical supergravity, $g_s \to 0, \alpha' \to 0$

8

Pirsa: 23100060 Page 10/39

Different levels of the AdS/CFT correspondence

$$g_{\mathrm{YM}}^2 = 2\pi g_s$$

$$\lambda = g_{
m YM}^2 N$$

$$2g_{\mathrm{YM}}^2 N = L^4/\alpha'^2$$

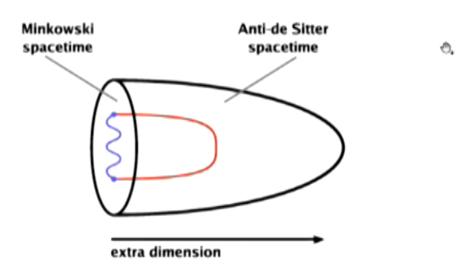
't Hooft coupling

	$\mathcal{N}=4$ SYM theory	IIB theory on $AdS_5 \times S^5$
Strongest form	any N and λ	Quantum string theory, $g_s \neq 0$, $\alpha' \neq 0$
Strong form	$N \to \infty$, λ fixed but arbitrary	Classical string theory, $g_s \to 0$, $\alpha' \neq 0$
Weak form	$N \to \infty$, λ large	Classical supergravity, $g_s \to 0, \alpha' \to 0$

в

Pirsa: 23100060 Page 11/39

Gauge/Gravity Duality



'Dictionary' Gauge invariant field theory operators

⇔ Classical fields in gravity theory

Symmetry properties coincide, generating functionals are identified

Test: (e.g.) Calculation of correlation functions

9

Pirsa: 23100060 Page 12/39

Different levels of the AdS/CFT correspondence

$$g_{\mathrm{YM}}^2 = 2\pi g_s$$

$$g_{
m YM}^2 = 2\pi g_s$$
 $\lambda = g_{
m YM}^2 N$

$$2g_{\mathrm{YM}}^{2}N=L^4/lpha'^2$$

't Hooft coupling

	$\mathcal{N}=4$ SYM theory	IIB theory on $AdS_5 \times S^5$
Strongest form	any N and λ	Quantum string theory, $g_s \neq 0$, $\alpha' \neq 0$
Strong form	$N \to \infty$, λ fixed but arbitrary	Classical string theory, $g_s \to 0$, $\alpha' \neq 0$
Weak form	$N \to \infty$, λ large	Classical supergravity, $g_s \to 0, \alpha' \to 0$

Page 13/39 Pirsa: 23100060

AdS/CFT correspondence

Field-operator correspondence:

$$\langle e^{\int d^d x \, \phi_0(\vec{x}) \mathcal{O}(\vec{x})} \rangle_{CFT} = Z_{sugra} \Big|_{\phi(0,\vec{x}) = \phi_0(\vec{x})}$$

Generating functional for correlation functions of particular composite operators in the quantum field theory

coincides with

Classical tree diagram generating functional in supergravity

10

Holographic Renormalization Group Flows

Original AdS/CFT correspondence involves $\mathcal{N}=4$ SU(N) Super Yang-Mills theory

Field content: One gauge field, four fermions, six real scalars

Beta function vanishes identically to all orders in perturbation theory

Theory at UV fixed point

Perturb by relevant operator preserving $\mathcal{N} = 1$ supersymmetry

$$W_{\rm LS} \equiv h \, {\rm Tr} \, (\Phi_3 \, [\Phi_1 \, , \, \Phi_2]) + \frac{m}{2} \, {\rm Tr} \, \Phi_3^2$$
 Leigh-Strassler flow

IR fixed point:
$$\beta(g) \sim 2N(\gamma_1 + \gamma_2 + \gamma_3)$$
 $\beta_h = \gamma_1 + \gamma_2 + \gamma_3, \quad \beta_m = 1 - 2\gamma_3$

$$\beta_h = \gamma_1 + \gamma_2 + \gamma_3,$$

$$\beta_m = 1 - 2 \gamma_3$$

$$\gamma_1 = \gamma_2 = -\frac{\gamma_3}{2} = -\frac{1}{4}$$

Holographic renormalization group flows: field-theory side

Leigh-Strassler flow
$$W_{\rm LS} \equiv h \, {\rm Tr} \, (\Phi_3 \, [\Phi_1 \, , \, \Phi_2]) + \frac{m}{2} \, {\rm Tr} \, \Phi_3^2$$

C theorem:

٥.

$$\frac{a_{\rm IR}}{a_{\rm UV}} = \frac{c_{\rm IR}}{c_{\rm UV}} = \frac{27}{32}$$

For coefficients of conformal anomaly in conformal Ward identity

$$\langle T_{\mu}{}^{\mu} \rangle = \frac{c}{8\pi^2} \left(R^{\mu\nu} R_{\mu\nu} - \frac{1}{3} R^2 \right), \qquad c = \frac{N^2}{4}$$

Holographic Renormalization Group Flows

Gravity side of correspondence: Dual of Leigh-Strassler flow

Look for a solution of supergravity preserving the same global symmetries

Domain wall flow:

$$S = \int d^{d+1}x \sqrt{-g} \left(\frac{R}{16\pi G} - \frac{1}{2} \partial_m \phi \partial^m \phi - V(\phi) \right)$$

$$ds^2 = e^{2A(r)} \eta_{\mu\nu} dx^{\mu} dx^{\nu} + dr^2$$

Conformal anomaly at IR fixed point exactly reproduced

$$\langle T_{\mu}{}^{\mu} \rangle = \frac{k_{\bullet}^{3}}{64\pi G_{5}} \left(R^{\mu\nu} R_{\mu\nu} - \frac{1}{3} R^{2} \right)$$

First order RG equation from gradient flow

$$V(\phi) = \frac{1}{2} \left(\frac{\mathrm{d}W}{\mathrm{d}r} \right)^2 - \frac{d}{d-1} W^2$$

$$V(\phi) = \frac{1}{2} \left(\frac{\mathrm{d}W}{\mathrm{d}r} \right)^2 - \frac{d}{d-1} W^2 \qquad \sqrt{8\pi G} \frac{\mathrm{d}\phi}{\mathrm{d}r} = \frac{\mathrm{d}W}{\mathrm{d}\phi}, \quad A' = -\frac{\sqrt{8\pi G}}{(d-1)} W$$

Freedman, Gubser, Pilch, Warner 1999

Pirsa: 23100060

A(r)

Different levels of the AdS/CFT correspondence

$$g_{\mathrm{YM}}^2 = 2\pi g_s$$

$$\lambda = g_{
m YM}^2 N$$

$$2g_{\mathrm{YM}}^2 N = L^4/\alpha'^2$$

't Hooft coupling

	$\mathcal{N}=4$ SYM theory	IIB theory on $AdS_5 \times S^5$
Strongest form	any N and λ	Quantum string theory, $g_s \neq 0$, $\alpha' \neq 0$
Strong form	$N \to \infty$, λ fixed but arbitrary	Classical string theory, $g_s \to 0$, $\alpha' \neq 0$
Weak form	$N \to \infty$, λ large	Classical supergravity, $g_s \to 0, \alpha' \to 0$

в

Pirsa: 23100060

Holographic Renormalization Group Flows

Original AdS/CFT correspondence involves $\mathcal{N}=4$ SU(N) Super Yang-Mills theory

Field content: One gauge field, four fermions, six real scalars

Beta function vanishes identically to all orders in perturbation theory

Theory at UV fixed point

Perturb by relevant operator preserving $\mathcal{N} = 1$ supersymmetry

$$W_{\rm LS} \equiv h \, {\rm Tr} \, (\Phi_3 \, [\Phi_1 \, , \, \Phi_2]) \, + \, \frac{m}{2} \, {\rm Tr} \, \Phi_3^2$$
 Leigh-Strassler flow

IR fixed point:
$$\beta(g) \sim 2N(\gamma_1 + \gamma_2 + \gamma_3) \stackrel{\mathfrak{G}}{\longrightarrow} \beta_h = \gamma_1 + \gamma_2 + \gamma_3, \quad \beta_m = 1 - 2\gamma_3$$

$$\gamma_1 = \gamma_2 = -\frac{\gamma_3}{2} = -\frac{1}{4}$$

11

Holographic Renormalization Group Flows

Gravity side of correspondence: Dual of Leigh-Strassler flow

Look for a solution of supergravity preserving the same global symmetries

Domain wall flow:

$$S = \int d^{d+1}x \sqrt{-g} \left(\frac{R}{16\pi G} - \frac{1}{2} \partial_m \phi \partial^m \phi - V(\phi) \right)$$

$$ds^2 = e^{2A(r)} \eta_{\mu\nu} dx^{\mu} dx^{\nu} + dr^2$$

Conformal anomaly at IR fixed point exactly reproduced

$$\langle T_{\mu}{}^{\mu}\rangle = \frac{L^3}{64\pi G_5} \left(R^{\mu\nu} R_{\mu\nu} - \frac{1}{3} R^2 \right)$$

$$\bigoplus_{\bullet}$$

First order RG equation from gradient flow

$$V(\phi) = \frac{1}{2} \left(\frac{\mathrm{d}W}{\mathrm{d}r}\right)^2 - \frac{d}{d-1} W^2$$

$$V(\phi) = \frac{1}{2} \left(\frac{\mathrm{d}W}{\mathrm{d}r} \right)^2 - \frac{d}{d-1} W^2 \qquad \sqrt{8\pi G} \frac{\mathrm{d}\phi}{\mathrm{d}r} = \frac{\mathrm{d}W}{\mathrm{d}\phi}, \quad A' = -\frac{\sqrt{8\pi G}}{(d-1)} W$$

Freedman, Gubser, Pilch, Warner 1999

A(r)

Relation to quantum information

0,

Entanglement entropy

Consider product Hilbert space $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$

Reduced density matrix

$$\rho_A = \text{Tr}_B \rho_{\text{tot}}$$

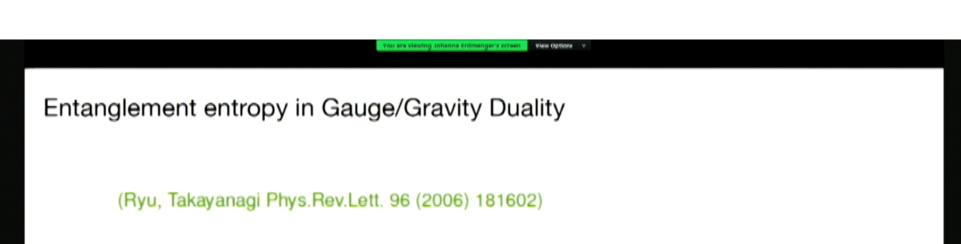
Entanglement entropy

$$S_A = -\mathrm{Tr}_A \rho_A \ln \rho_A$$

Entanglement for space regions

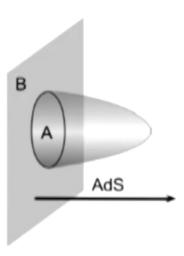
В

Pirsa: 23100060



Leading term in entanglement entropy given by area of minimal surface in holographic dimension





Pirsa: 23100060 Page 22/39

Entanglement entropy: Quantum field theory

Conformal field theory in 1+1 dimensions (Cardy, Calabrese, J.Stat.Mech. 0406 (2004) P06002):

 $S = \frac{c}{3} \ln(\ell \Lambda)$

Reproduced by Ryu-Takayanagi result

 $\Lambda \propto 1/\epsilon, \, \epsilon$ boundary cut-off in radial direction

 $c = 3L/(2G_3)$ central charge of CFT

Finite temperature (at small ℓ):

$$S(\ell) = \frac{c}{3} \ln \left(\frac{1}{\pi \epsilon T} \sinh(2\pi \ell T) \right)$$

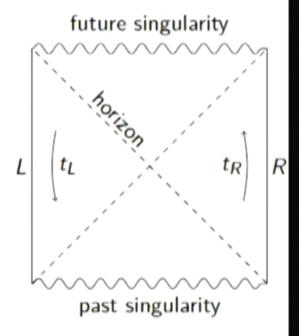
Pirsa: 23100060 Page 23/39

Von Neumann algebras, quantum gravity and AdS/CFT

Factorization puzzle Maldacena+Maoz '13; Harlow '16

• The two CFTs have disjoint Hilbert spaces since there is no interaction between them $~\mathcal{H}_L\otimes\mathcal{H}_\mathcal{R}$

- The wormhole Hilbert space does not factorize
- Apparent contradiction?



Eternal AdS black hole

٠

17

Pirsa: 23100060 Page 24/39

View Options

von Neumann algebras

von Neumann 1930

Jefferson; Liu, Leutheusser; Witten; Chandrasekaran, Penington, Witten

Concept of algebraic QFT

for classifying operator algebras w. r. t. entanglement properties

Type I - density matrix and trace (as in quantum mechanics), admits irreducible representations

Type II - trace prescription, but does not act irreducibly

Type III - no trace prescription (eg. free QFTs)

18

Pirsa: 23100060 Page 25/39

View Options

Type II vs. type III von Neumann algebra for eternal black hole

Including 1/N corrections gives rise to type II algebra with vanishing center

Allows for well-defined trace functional for a particular state Witten '21

Geometric phase approach: Non-exact symplectic form leads to Berry phase

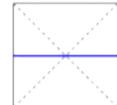
Trace functional proportional to Berry phase

Banerjee, Dorband, J.E., Weigel arXiv:2306.00055

Non-factorization

-> non-zero geometric phase -> no trace definition -> type III vN algebra

Maximally entangled state -> geometric phase vanishes -> type II vN algebra



19

Pirsa: 23100060 Page 26/39

von Neumann algebras

von Neumann 1930

Jefferson; Liu, Leutheusser; Witten; Chandrasekaran, Penington, Witten

Concept of algebraic QFT

for classifying operator algebras w. r. t. entanglement properties

0,

Type I - density matrix and trace (as in quantum mechanics), admits irreducible representations

Type II - trace prescription, but does not act irreducibly

Type III - no trace prescription (eg. free QFTs)

18

Pirsa: 23100060 Page 27/39

View Options V

Type II vs. type III von Neumann algebra for eternal black hole

Including 1/N corrections gives rise to type II algebra with vanishing center

Allows for well-defined trace functional for a particular state Witten '21

Geometric phase approach: Non-exact symplectic form leads to Berry phase

Trace functional proportional to Berry phase

Banerjee, Dorband, J.E., Weigel arXiv:2306.00055

Non-factorization

-> non-zero geometric phase -> no trace definition -> type III vN algebra

Maximally entangled state -> geometric phase vanishes -> type II vN algebra

0,

19

Pirsa: 23100060 Page 28/39

von Neumann algebras

von Neumann 1930

Jefferson; Liu, Leutheusser; Witten; Chandrasekaran, Penington, Witten

Concept of algebraic QFT

for classifying operator algebras w. r. t. entanglement properties

Type I - density matrix and trace (as in quantum mechanics), admits irreducible representations

Type II - trace prescription, but does not act irreducibly

Type III - no trace prescription (eg. free QFTs)

18

Pirsa: 23100060 Page 29/39

Type II vs. type III von Neumann algebra for eternal black hole

Including 1/N corrections gives rise to type II algebra with vanishing center

Allows for well-defined trace functional for a particular state Witten '21

Geometric phase approach: Non-exact symplectic form leads to Berry phase

Trace functional proportional to Berry phase

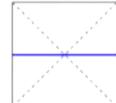
Banerjee, Dorband, J.E., Weigel arXiv:2306.00055

€.

Non-factorization

-> non-zero geometric phase -> no trace definition -> type III vN algebra

Maximally entangled state -> geometric phase vanishes -> type II vN algebra

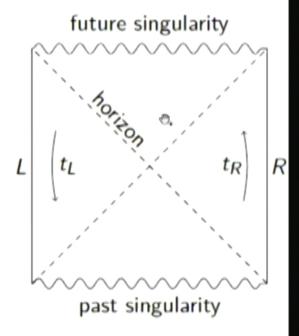


Von Neumann algebras, quantum gravity and AdS/CFT

Factorization puzzle M

Maldacena+Maoz '13; Harlow '16

- The two CFTs have disjoint Hilbert spaces since there is no interaction between them $~\mathcal{H}_L\otimes\mathcal{H}_\mathcal{R}$
- The wormhole Hilbert space does not factorize
- Apparent contradiction?



Eternal AdS black hole

17

Pirsa: 23100060 Page 31/39

Type II vs. type III von Neumann algebra for eternal black hole

Including 1/N corrections gives rise to type II algebra with vanishing center

Allows for well-defined trace functional for a particular state Witten '21

Geometric phase approach: Non-exact symplectic form leads to Berry phase

Trace functional proportional to Berry phase

Banerjee, Dorband, J.E., Weigel arXiv:2306.00055

Non-factorization

-> non-zero geometric phase -> no trace definition -> type III vN algebra

Maximally entangled state -> geometric phase vanishes -> type II vN algebra



Pirsa: 23100060 Page 32/39

View Options V

Type II vs. type III von Neumann algebra for eternal black hole

Including 1/N corrections gives rise to type II algebra with vanishing center

Allows for well-defined trace functional for a particular state Witten '21

Geometric phase approach: Non-exact symplectic form leads to Berry phase

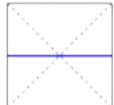
Trace functional proportional to Berry phase

Banerjee, Dorband, J.E., Weigel arXiv:2306.00055

Non-factorization

-> non-zero geometric phase -> no trace definition -> type III vN algebra

Maximally entangled state -> geometric phase vanishes -> type II vN algebra



19

Analogy: Interacting two-spin system: Berry phase and type I Von Neumann algebra

$$H = J\mathbf{S}_1 \cdot \mathbf{S}_2 - 2\mu_B B S_{1z}$$
 $\tan \alpha = 2\mu_B \frac{B}{I}$

$$\tan \alpha = 2\mu_B \frac{B}{J}$$

Entanglement entropy

$$S_{\text{EE}} = -\sum_{i=\uparrow,\downarrow} \kappa_i^2 \ln \kappa_i^2 = \sin \alpha \ln \frac{1 - \sin \alpha}{\cos \alpha} - \ln \frac{\cos \alpha}{2}$$

No entanglement (J=0)

$$S_{EE} = 0$$

 $S_{EE} = 0$ Entanglement orbit $\mathbb{C}P^1 \times \mathbb{C}P^1$

$$\mathbb{C}\mathrm{P}^1 \times \mathbb{C}\mathrm{P}^1$$

Berry phase maximal

Maximal entanglement (J very large)

$$S_{EE} = \ln 2$$

$$\frac{SU(2)}{\mathbb{Z}_2} = \mathbb{R}P^3$$

Berry phase vanishes

Intermediate entanglement

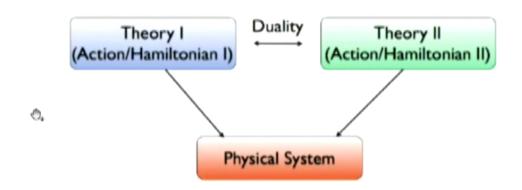
$$\mathbb{C}P^1 \times \mathbb{R}P^3$$

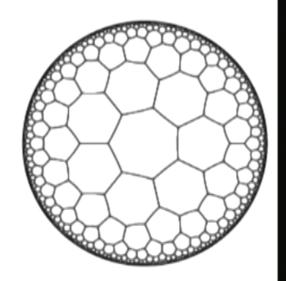
$$\int \Omega = \frac{\sin \alpha}{2} V(S^2) = 2\pi \sin \alpha$$

20

Discrete Holography

Goal: Establish holographic duality on hyperbolic tiling





arXiv:2205.05081, Phys. Rev. Lett. '23 (Basteiro, Dusel, J.E., Hinrichsen, Meyer, Herdt, Schrauth)

arXiv:2205.05693, Sci Post Physics (Basteiro, Di Giulio, J.E., Karl, Meyer, Xian)

arXiv: 2212.11292 (Basteiro, Das, Di Giulio, J.E.)

Pirsa: 23100060 Page 35/39

Hyperbolic tilings through inflation rules

Starting point: central tile = 0-th layer

↓

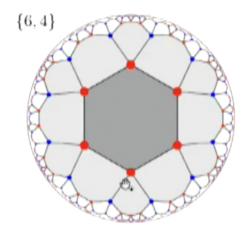
we construct concentric layers of tiles

After a layers; outermost set of edges

After n layers: outermost set of edges and vertices = **boundary of the tiling**

2 internal neighbours = ● ← → a

3 internal neighbours = • ← → b



{4,5}

Inflation rule

$$\sigma_{\{p,q\}}: \begin{cases} a \to w_a(a,b) \\ b \to w_b(a,b) \end{cases}$$

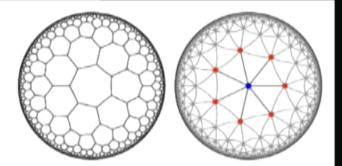
$$\sigma_{\{6,4\}}: \left\{ egin{aligned} a
ightarrow aabaaab \ b
ightarrow aab \end{aligned}
ight.$$

$$\sigma_{\{4,5\}}: \begin{cases} a \to babab \\ b \to bab \end{cases}$$

The boundary of the tiling after a large number of inflation steps is characterised by a long **aperiodic sequence**, whose size grows by a factor $\lambda_+(p,q)$ at each step

Discrete holography





Spin chain at boundary realises bulk tiling construction principle at boundary Aperiodic spin chain

Ground state of spin chain Hamiltonian determined using RG approach

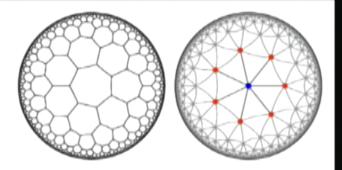
Associated tensor network representing this ground state bulk reconstruction

Gravity dynamics on bulk side — dynamical triangulations in hyperbolic space?

24

Pirsa: 23100060 Page 37/39

Discrete holography



Spin chain at boundary realises bulk tiling construction principle at boundary

Aperiodic spin chain

Ground state of spin chain Hamiltonian determined using RG approach

Associated tensor network representing this ground state - bulk reconstruction

Gravity dynamics on bulk side — dynamical triangulations in hyperbolic space?

24

Pirsa: 23100060 Page 38/39

View Options

Conclusion and Outlook

Gauge/Gravity Duality

Holographic principle

Holographic RG flows implemented from supergravity equations of motion

(Gradient flow) - Constructed to match global symmetries

New relations between quantum information and gravity

Entanglement entropy, mutual information, computational complexity

von Neumann algebras, Berry phases

Reconstrucion of spacetime, discrete holography

23

Pirsa: 23100060 Page 39/39