

Title: QFT1 Lecture - 103023

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Interaction (generic) (LSZ ... Wick

scalar
step 1 $\langle f | S | i \rangle \quad a_i = a_{k_i}$

$$|i\rangle = a_1^\dagger(-\infty) a_2^\dagger(-\infty) |\Omega\rangle$$
$$\langle \Omega | a_4(+\infty) a_3(+\infty) a_1^\dagger(-\infty) a_2^\dagger(-\infty) |\Omega\rangle$$

... Wick)
fermion (2 particle to 2 particle)

$$|i\rangle = b_1^+(-\infty) b_2^+(-\infty) |\Omega\rangle$$

$$b_1 \equiv b_{\vec{k}}$$

→ → change a to bs

Step 2

$$\begin{aligned} & a_1(+\infty) - a_1(-\infty) \\ \equiv & \underline{I}_1 = \int_{-\infty}^{\infty} dt \partial_0 a_1(t) \\ \equiv & -i \int d^4x e^{-ik_1 x} \underline{(\partial^2 + m^2)} \varphi(x) \end{aligned}$$

$$I_1^+ = i \int d^4x e^{+ik_1 x} (\partial^2 + m^2) \varphi(x).$$

$$\begin{aligned} I_{1f} & \equiv b_1(+\infty) - b_1(-\infty) \\ & \equiv \int_{-\infty}^{\infty} dt \partial_0 b_1(t) \end{aligned}$$

$$I_{1,f} \equiv \frac{b_1(+\infty) - b_1(-\infty)}{\int_{-\infty}^{\infty} dt \partial_0 b_1(t)}$$

$$= \# \int d^4x e^{-ik_1 x} \frac{u_1 (i\not{\partial} - m) \psi(x)}{u_1 \gamma^0 (i\not{\partial} \gamma^0 - m \gamma^0)}$$

$$I_{1,f} = \# \int d^4x e^{+ik_1 x} (\partial_{\mu\nu}^{\dagger} \gamma^{\mu} \gamma^{\nu} (-i) \gamma^0 u_1 - m \gamma^0 u_1)$$

$$= \# \int d^4x e^{+ik_1 x} (\partial_{\mu\nu} \gamma^{\mu} u_1 (-i) - m \gamma^0 u_1)$$

$$= \# \int d^4x \frac{\not{\partial} (i\not{\partial} + m) u_1 e^{+ik_1 x}}{\not{\partial} (i\not{\partial} + m) u_1 e^{+ik_1 x}}$$

scalar

$$\langle f | S | i \rangle = \frac{1}{i!} \int d^4x_j e^{-i\lambda_j x_j} (\partial_j^2 + m^2)$$
$$\langle \Omega | T \varphi_1 \varphi_2 \varphi_3 \varphi_4 | \Omega \rangle$$

scalar

$$\langle f | S | i \rangle = \frac{4}{i\pi} \int d^4 x_j e^{-i\lambda_j x_j} (\partial_j^2 + m^2)$$

$$\langle \Omega | T \varphi_1 \varphi_2 \varphi_3 \varphi_4 | \Omega \rangle$$

step 3

$$\langle \Omega | T \varphi_1 \varphi_2 \varphi_3 \varphi_4 | \Omega \rangle = \frac{\langle 0 | T \varphi_1 \varphi_2 \varphi_3 \varphi_4 e^{i\lambda} | 0 \rangle}{\langle 0 | e^{i\lambda} | 0 \rangle}$$

fermion case

$$= \frac{4}{i\pi} \int d^4 x_j e^{-i\lambda_j x_j}$$

$$\langle \Omega | T \varphi_1 \varphi_2 \varphi_3 \varphi_4 | \Omega \rangle$$

$$\frac{1}{i\pi} \int d^4 x_j$$

$$e^{-i\lambda_i \not{x}^i} (\not{\partial}_j^2 + m^2)$$

$$\varphi_3 \varphi_4 | \Omega \rangle$$

$$\varphi_3 \varphi_4 | \Omega \rangle = \frac{\langle 0 | T \varphi_1 \varphi_2 \varphi_3 \varphi_4 e^{i\lambda} | 0 \rangle}{\langle 0 | e^{i\lambda} | 0 \rangle}$$

fermion case

$$= \frac{1}{\prod_{j=3}^4} \int d^4 x_j e^{-i k_j x_j} \underline{u}_j (i \not{\partial} - m)$$

$$\langle \Omega | T \psi \psi \bar{\psi} \bar{\psi} | \Omega \rangle$$

$$\frac{1}{\prod_{j=1}^2} \int d^4 x_j (i \not{\partial} + m) \underline{u}_j e^{+i k_j x_j}$$

$$\langle \Omega | I_4 I_3 I_1^+ I_2^+ | \Omega \rangle$$

fermion case

$$= \frac{1}{\prod_{j=3}^4} \int d^4 x_j e^{-i k_j \cdot x_j} \underline{u_j} (i \not{\partial} - m)$$

$$\langle \Omega | T \psi \psi \bar{\psi} \bar{\psi} | \Omega \rangle$$

$$\frac{1}{\prod_{j=1}^2} \int d^4 x_j (i \not{\partial} + m) \underline{u_j} e^{+i k_j \cdot x_j}$$

$$\langle \Omega | I_4 I_3 I_1^+ I_2^+ | \Omega \rangle$$

$$p_3 p_4 e^{i k \cdot p}$$

$$| e^{i k \cdot p} \rangle$$

→ no change

Wick theorem

$$T = : : + \square$$

$$T(\varphi(x)\varphi(y)) = : \varphi(x)\varphi(y) : + \overline{\varphi(x)\varphi(y)}$$

$$\langle 0 | T(\varphi_1(x)\varphi_2(y)) | 0 \rangle = \langle 0 | \overline{\varphi(x)\varphi(y)} | 0 \rangle$$

$$\Delta_F(x-y)$$

$$= \int d^4p \frac{i}{p^2 - m^2 + i\epsilon} e^{-ip(x-y)}$$

- ④ assemble with step function
- ① define time ordering.
 - ② split $\varphi \rightarrow \varphi_+$ and φ_-
 - ③ compute $[\varphi_-, \varphi_+]$

starting point

$$T(\psi_a(x)\bar{\psi}_b(y)) = \psi_a(x)\bar{\psi}_b(y) + \psi_a(x)\bar{\psi}_b(y) \leftarrow \text{goal}$$

step 1 time ordering

$$T(\psi(x)\bar{\psi}(y)) = \begin{cases} \psi(x)\bar{\psi}(y) & x^0 > y^0 \\ \bar{\psi}(y)\psi(x) & y^0 > x^0 \end{cases}$$

suppose $x_0 > y_0$

$$T\psi_a(x)\bar{\psi}_b(y) = \psi_a(x)\bar{\psi}_b(y)$$

$$\psi = \int \frac{d^3p}{(2\pi)^3} u(p)e^{-ikx} + v(p)e^{+ikx}$$

$$\psi_a(x)\bar{\psi}_b(y) = -\bar{\psi}_b(y)\psi_a(x) = 0 \dots$$

first $X_0 > y_0$

choose $T(\psi(x)\bar{\psi}(y)) = \psi(x)\bar{\psi}(y)$

$$= \psi_+(x)\bar{\psi}_+(y) + \psi_-(x)\bar{\psi}_-(y)$$

problem $\rightarrow +\psi_-(x)\bar{\psi}_+(y) + \psi_-(x)\bar{\psi}_-(y)$

$$:c c^\dagger := -c^\dagger c \quad \left(-\bar{\psi}_+(y)\psi_-(x) + \bar{\psi}_+(y)\psi_-(x) \right)$$

$$= : \psi(x)\bar{\psi}(y) : + \{ \psi_-(x), \bar{\psi}_+(y) \}$$

$$\psi = \int dV_k u b_k e^{-ikx} + v c_k^+ e^{+ikx}$$

$$\psi^+(y) = \int dV_p u^+ b_p^+ e^{+ipx} + v^+ c_p e^{-ipx}$$

$$\{\psi_-(x), \overline{\psi}_+(y)\} = \left\{ \int dV_k u b_k e^{-ikx}, \int dV_q \bar{u} b_q^+ e^{+iq \cdot y} \right\}$$

$$= \int dV_k u e^{-ikx} \int dV_q \bar{u} e^{+iqy} \underbrace{\{b_k, b_q^+\}}_{N_D \delta(k-q)}$$

$$= \int dV_k u \bar{u} e^{-ik(x-y)}$$

$$= \int dV_k (\not{k} + m) e^{-ik(x-y)}$$

$$= (i \not{\partial}_x + m) \int dV_k e^{-ik(x-y)}$$

$$e^{+i\mathbf{q}\cdot\mathbf{y}}$$

$$\underbrace{\{b_{\mathbf{k}}, b_{\mathbf{q}}^{\dagger}\}}_{N_0 \delta(\mathbf{k}-\mathbf{q})}$$

$$\begin{aligned} S_F(x-y) &= (i\not{\partial}_x + m) \Delta_F(x-y) \\ &= (i\not{\partial}_x + m) \int d^4p \frac{i}{p^2 - m^2 + i\epsilon} e^{-ip(x-y)} \\ &= \int d^4p \frac{i(\not{p} + m)}{p^2 - m^2 + i\epsilon} e^{-ip(x-y)} \end{aligned}$$

Ⓟ as
with st
fur

$$\underline{(\partial^2 + m^2) \Delta_F(x-y) = i \delta(x-y)}$$

$$(i\not{\partial} + m)(i\not{\partial} - m)$$

$$(i\not{\partial} - m)(i\not{\partial} + m) \Delta_F(x-y) = i \delta(x-y)$$

$$\underline{(i\not{\partial} - m) S_F(x-y) = i \delta(x-y)}$$

Wick contraction for many fermion field
 all contractions

$$\begin{array}{c} \square \\ \dots \psi_1 \psi_2 \psi_3 \dots \\ \square \\ \dots \psi_1 \psi_2 \bar{\psi}_3 \dots = - \psi_1 \bar{\psi}_3 (\dots \psi_2 \dots) \end{array}$$

- ① bring together $(-)^s$
- ② \square
 $\psi \bar{\psi}$