

Title: Quantum Theory Lecture - 100223

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Collection: Quantum Theory 2023/24

Date: October 02, 2023 - 10:45 AM

URL: <https://pirsa.org/23100034>

Interviews: No notes

20 minutes prep

-questions by email

5-7 minute response preparation

+follow-up

Renormalization warm-up quiz
for tomorrow!

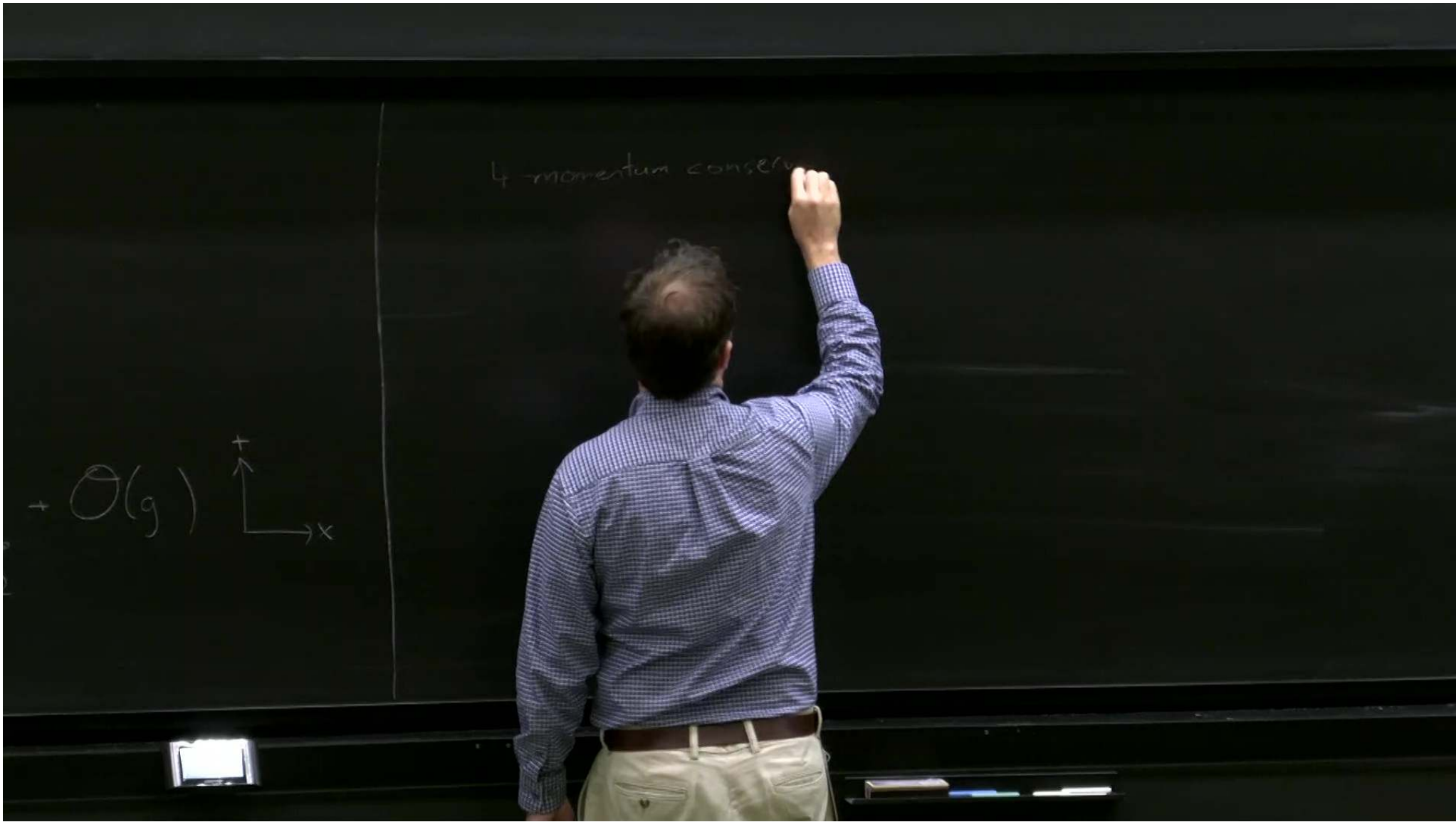
Feynman rules for iM

$$\langle f|S|i\rangle = (2\pi)^4 \delta^4(\sum_{\text{final}} p - \sum_{\text{initial}} p) iM_{i \rightarrow f}$$

Example: $2 \rightarrow 2$ scattering in ϕ^4 theory

$$\langle S|T \phi_1 \phi_2 \phi_3 \phi_4 |S\rangle =$$

1, 2 initial
3, 4 final



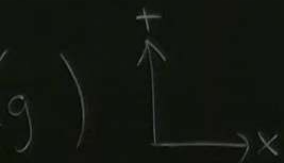
4-momentum conserving delta per connected component!
(last time missed $(\partial^2 + m^2)$)
in $\langle f|S|i \rangle$

$$\langle f|S|i \rangle_1 \propto \delta^4(p_1 - p_3) \delta^4(p_2 - p_4)$$

no scattering!
not interesting

$$\langle f|S|i \rangle_2 \propto \delta^4(p_1 + p_2) \delta^{(4)}(p_3 + p_4)$$

if $m > 0$



4-momentum conserving delta per connected component!
(last time missed $(\partial^2 + m^2)$)
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$$\langle f|S|i \rangle_1 \propto \delta^4(p_1 - p_3) \delta^4(p_2 - p_4) \quad \text{no scattering!}$$

not interesting

$$\langle f|S|i \rangle_2 \propto \delta^4(p_1 + p_2) \delta^{(4)}(p_3 + p_4) = 0$$

if $m > 0$, $p_1^0 + p_2^0 > 0$



Disconnected diagrams do interfere
with fully connected diagrams due
to extra deltas:

Either: arguments of extra deltas are zero \rightarrow infinitely larger
" " " " non-zero \rightarrow terms = 0

we will not consider further.

$\langle \psi^2 \rangle$

$$\langle \Omega | T \phi_1 \phi_2 \phi_3 \phi_4 | \Omega \rangle_{\text{fully corrected}} = \begin{array}{c} 3 \rightarrow \\ \diagdown \quad \diagup \\ \times \\ \diagup \quad \diagdown \\ 1 \quad 2 \end{array} + \mathcal{O}(g^2)$$

ager
= 0

$$\underbrace{\text{LSZ}}_{\text{LSZ}} = -ig \int d^4x \Delta_{1x} \Delta_{2x} \Delta_{3x} \Delta_{4x} \quad \text{plug into LSZ}$$

$$(\partial_1^2 + m_{\text{ph}}^2) \Delta_{1x} = -i \delta^4(x_1 - x) \quad \text{holds if } m_{\text{ph}}^2 = m^2$$

↑
physical mass

$= 0$

$\underbrace{L_S Z}_{\text{LSZ}} = -ig \int d^4x \Delta_{1x} \Delta_{2x} \Delta_{3x} \Delta_{4x}$ plug into LSZ

$(\partial_1^2 + m_{ph}^2) \Delta_{1x} = -i \delta^4(x_1 - x)$ holds if $m_{ph}^2 \stackrel{?}{=} m^2$

↑ physical mass

$\langle f | S | i \rangle_x = -ig \int d^4x \int d^4x_j \left[e^{i(x_j P_j - x_j)} (-i) \delta^4(x_j - x) \right]$



$$\langle f | S | i \rangle_x = -ig \int d^4x e^{i(p_1 + p_2 - p_3 - p_4) \cdot x}$$

$$= -ig (2\pi)^4 \delta^4(p_1 + p_2 - p_3 - p_4)$$

$$\overline{iM}_{i \rightarrow f}$$

$$iM = -ig + \mathcal{O}(g^2)$$

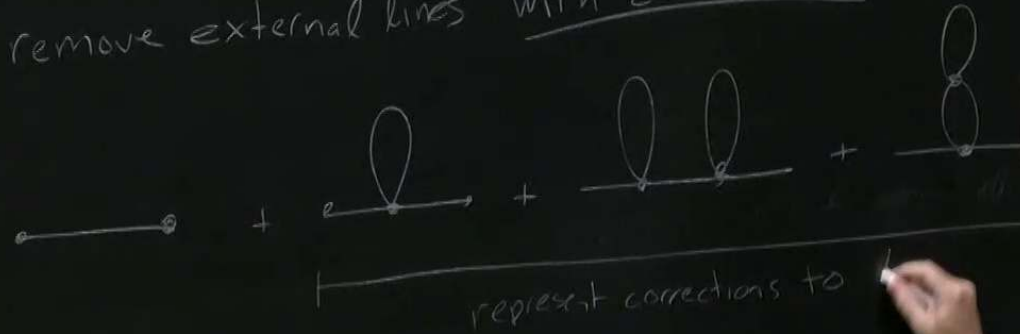
no propagators for external lines!

Amputation

Claim: $(\partial^2 + m_{ph}^2)$

remove external lines with corrections!

$$\langle \Omega | T \phi_1 \phi_2 | \Omega \rangle =$$



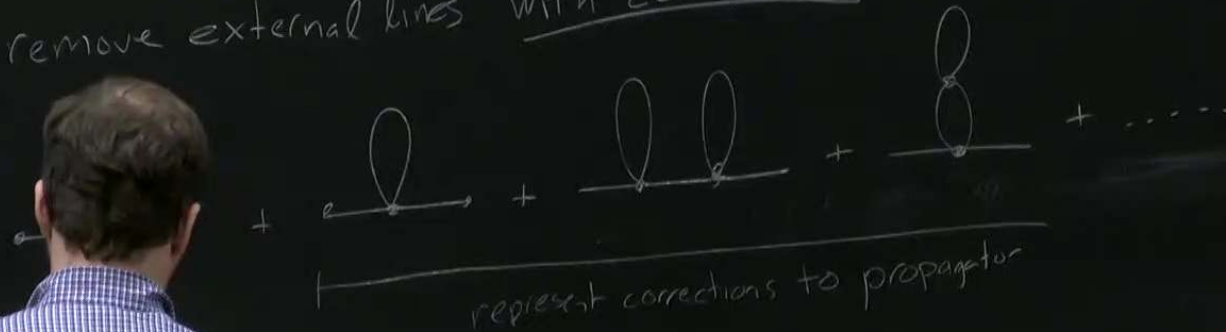
Amputation

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k_0

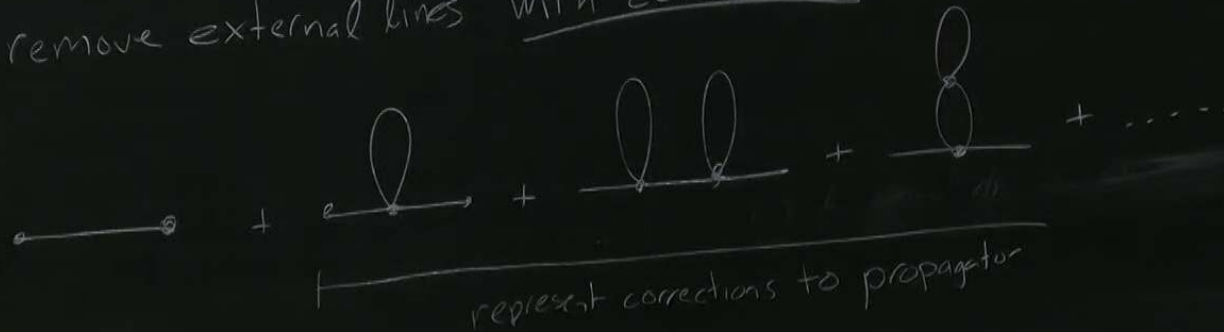


Amputation

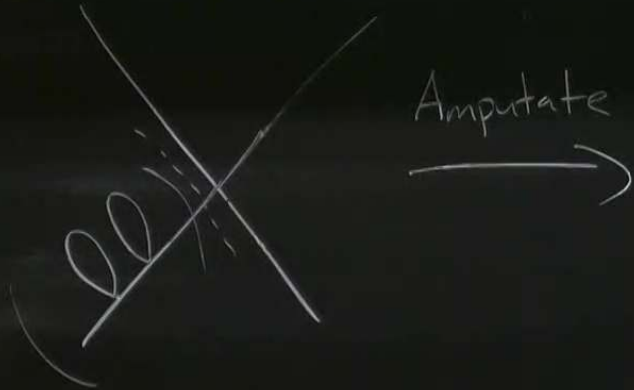
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$\langle T \phi_1 \phi_2 | \Omega \rangle =$



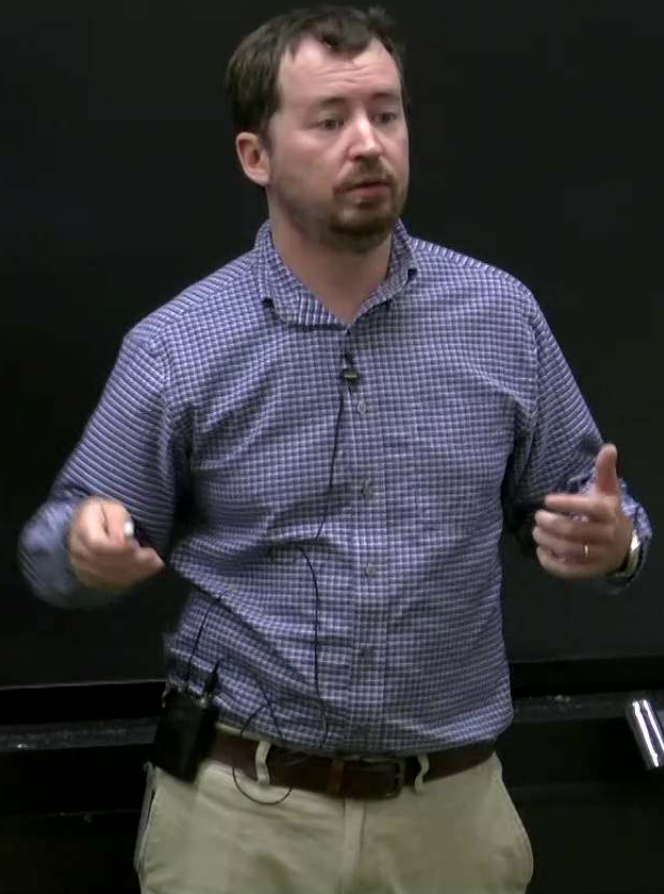
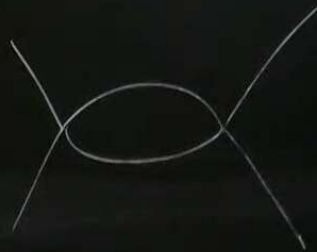
Källén-Lehmann tomorrow

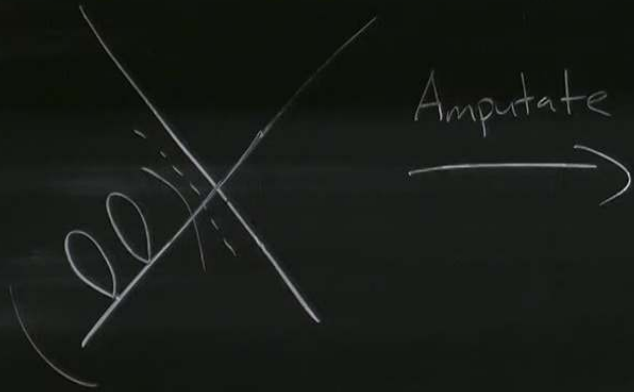


Amputate

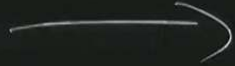


represents corrections to propagator
for an external particle

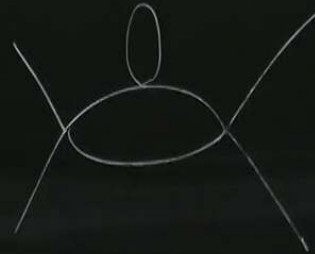




Amputate



represents corrections to propagator
for an external particle



Momentum conservation

higher order diagrams have more { vertices
propagators

$$\left. \begin{array}{l} \int d^4 x \\ \int d^4 k e^{ik \cdot x} \end{array} \right\} \text{lots of } \delta^4(p) \int d^4 p$$

Example: $2 \rightarrow 2$ scattering in φ^4 theory (completely connected)

V vertices

E edges or lines

momentum integrals \leftarrow one delta pulled out of iM

$$\underbrace{E - 4}_{\text{internal lines}} - (V - 1) = E - V - 3$$

1, 2 initial

$$\text{Euler: } \chi = -E + V + L$$

Euler characteristic

loops

constant for all 2→2 connected diagrams

$$\text{for } X \quad \chi = -4 + 1 + 0 = -3$$

$$L = E - V - 3$$

in general

momentum integrals = loops



L =

1, 2 initial

$$\chi = -E + V + L$$

↑
characteristic

↑
loops

constant for all 2→2 connected diagrams

$$X \quad \chi = -4 + 1 + 0 = -3$$

$$L = E - V - 3$$

general

momentum integrals = loops



$$L = 1$$
$$E = 6$$
$$V = 2$$

$$\chi = -6 + 1 + 2 = -3$$



$$L = 2$$
$$E = 8$$
$$V = 3$$

$$\chi = -8 + 2 + 3 = -3$$

1, 2 initial
3, 4 final

Feynman rules for ϕ^4

$\Gamma =$ sum of all completely connected, amputated diagrams


(interesting part)

internal line


$$\begin{array}{c} \text{---} \\ \nearrow \vec{p} \\ \text{arrow next to} \\ \text{line} \end{array} = \frac{i}{p^2 - m^2 + i\epsilon}$$

external lines

$$\begin{array}{c} \bullet \text{---} \\ \rightarrow \vec{p} \end{array} = 1$$

3.  = $-ig$ (depends on theory)

4. Impose (by hand) momentum conservation at each vertex

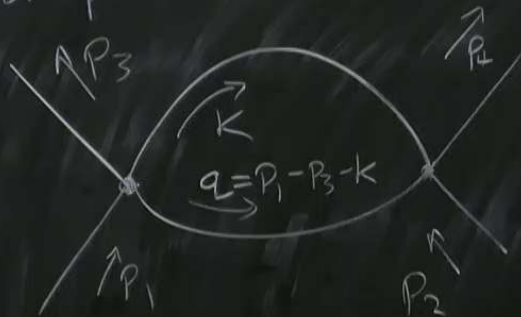
3.  = $-ig$ (depends on theory)

4. Impose (by hand) momentum conservation at each vertex (No δ !)

5. Integrate over undetermined momenta $\int \frac{d^4 k}{(2\pi)^4}$ (1 per loop)

6. Divide by symmetry factor

Example:



=

$$\frac{i}{k^2 - m^2 + i\epsilon}$$

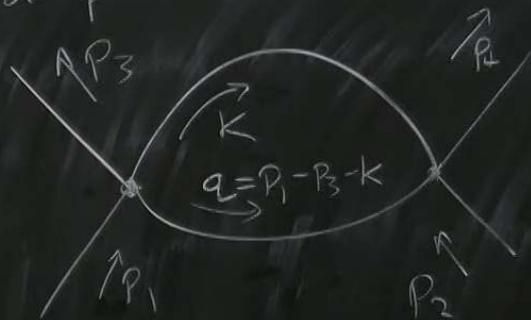
$$\frac{i}{(p_1 - p_3 - k)^2 - m^2 + i\epsilon}$$

$$p_1 = p_3 + k + q$$

$$p_2 + k + p_1, p_3 - k = p_4$$



Example



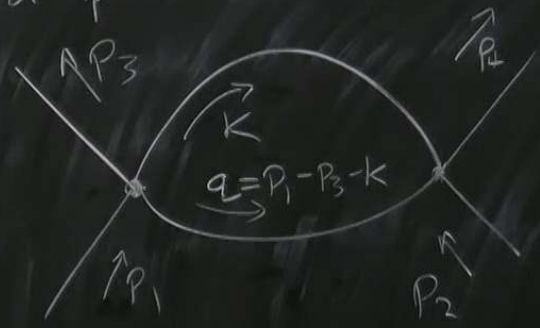
$$= (-ig)^2 \int \frac{d^4 k}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon}$$

$$\frac{i}{(p_1 - p_3 - k)^2 - m^2 + i\epsilon}$$

$$p_1 = p_3 + k + q$$

$$p_2 + k + p_1, p_3 - k = p_4$$

Example:

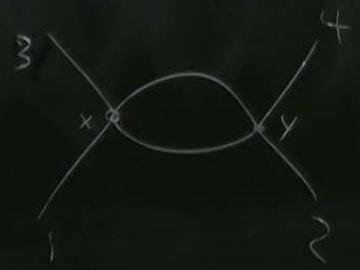


$$= \left(\frac{-ig}{2}\right)^2 \int \frac{d^4k}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon}$$

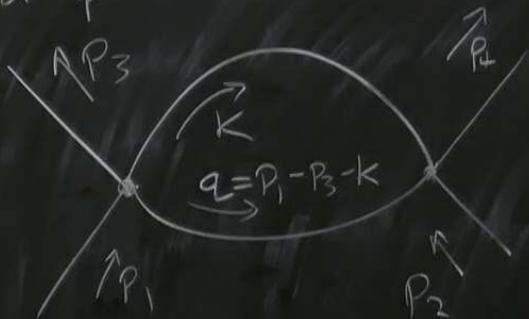
$$\frac{i}{(p_1 - p_3 - k)^2 - m^2 + i\epsilon}$$

$$p_1 = p_3 + k + q$$

$$p_2 + k + p_1 = p_3 - k = p_4$$



Example

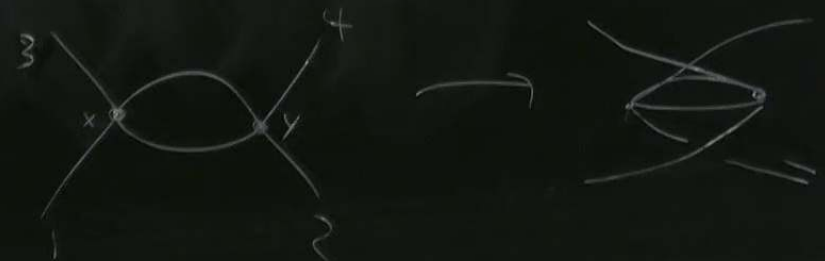


$$= \left(\frac{-ig}{2}\right)^2 \int \frac{d^4k}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon}$$

$$\frac{i}{(p_1 - p_3 - k)^2 - m^2 + i\epsilon}$$

$$p_1 = p_3 + k + q$$

$$p_2 + k + p_1, p_3 - k = p_4$$



Renormalization

Separation of scales : we can predict how an apple falls
with knowing cosmology or atomic physics

$$-k^2 - m^2 + i\epsilon$$

Goal: predict near a physical scale M
 $O(p)$ for $|p| \approx M$



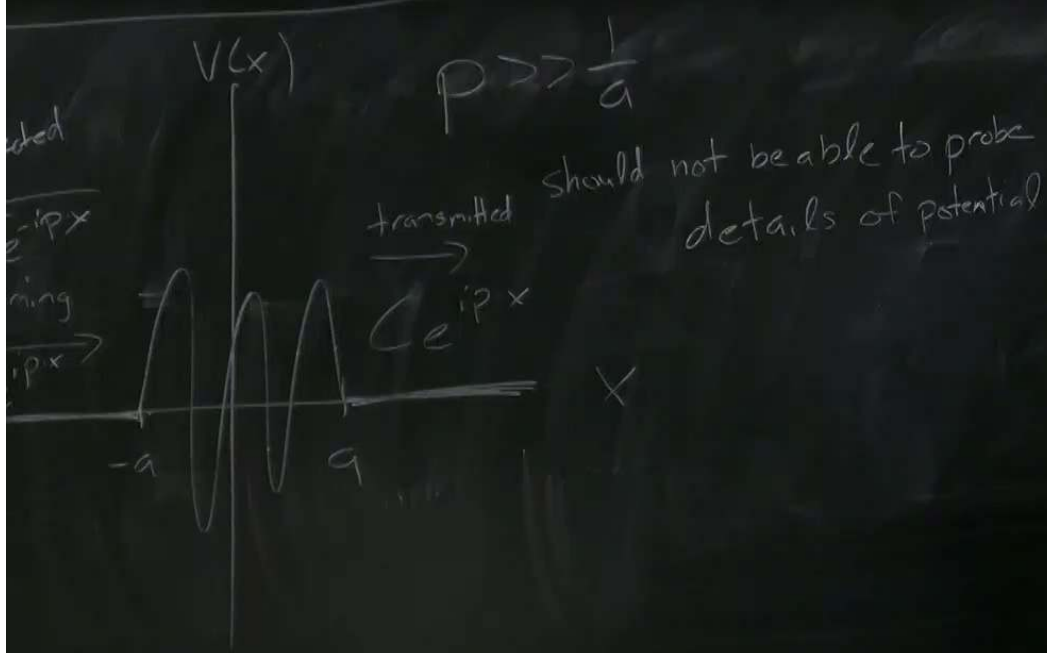
introduce unobservable quantity coupling constants

$$Q(\Lambda, M, L, P, c)$$

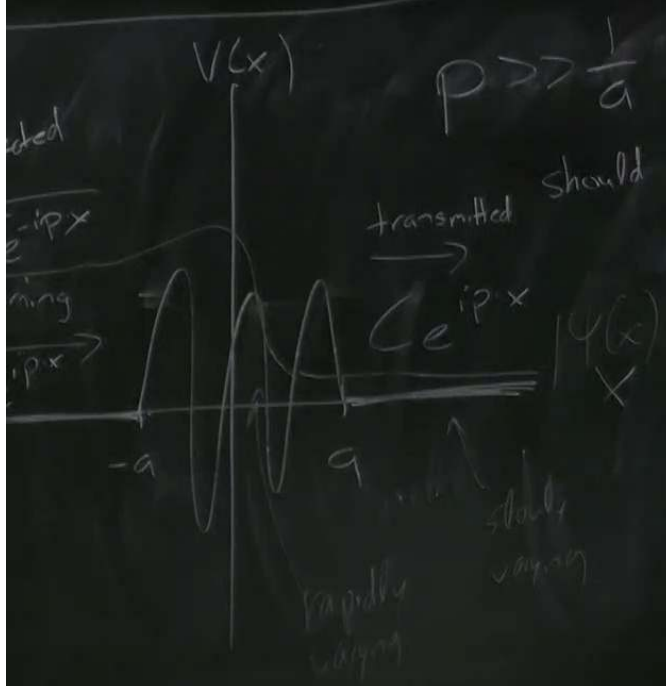
scale at which we are measuring

$$O(p, c)$$

Example: local potential in 1d non-relativistic QM



example: local potential in 1d non-relativistic QM



$$p \gg \frac{1}{a}$$

should not be able to probe details of potential

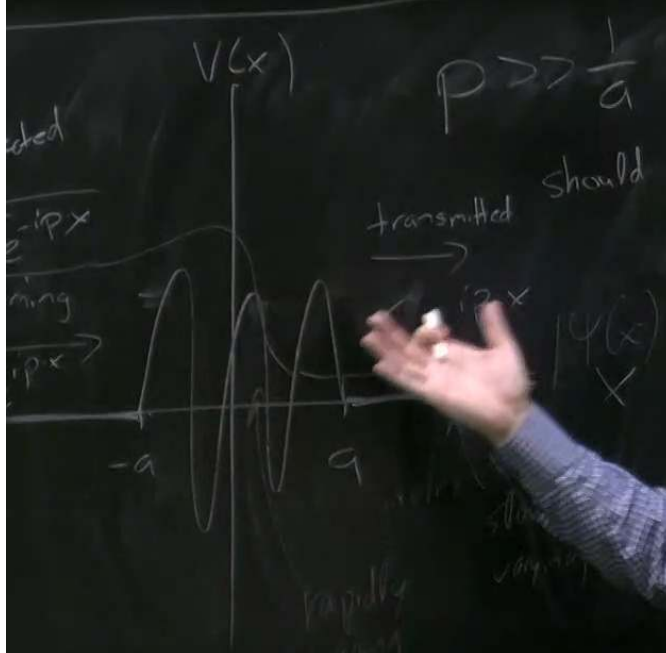
transmission coefficient is observable

$$T = \frac{C}{A}$$

cannot (easily) calculate T without taking into high energy mode



example: local potential in 1d non-relativistic QM



$$p \gg \frac{1}{a}$$

transmitted should not be able to see details of

cannot (easily) calculate T without taking into account high energy modes

how to define coupling constants