

Title: Learning Feynman Diagrams with Tensor Trains

Speakers: Xavier Waintal

Collection: Machine Learning for Quantum Many-Body Systems

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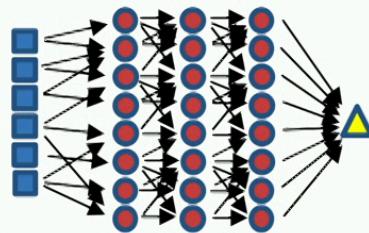
# CORRELATIONS AND COMPUTATIONAL QUANTUM TRANSPORT: LEARNING FEYNMAN DIAGRAMS WITH TENSOR TRAINS

Xavier WAINTEL,  
CEA Grenoble, Phelqs IRIG  
With YURIEL NUNEZ-FERNANDEZ, THOMAS KLOSS, MATTHIEU  
JEANNIN, CORENTIN BERTRAND, PHILIPP DUMITRESCU, Olivier  
Parcollet, Jason Kaye and Serge Florens.

# Machine learning as seen by a not so expert

0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
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3	3	3	3	3	3	3	3	3	3	3	3	3	3	3	3	3	3	3	3	3	3
4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4
5	5	5	5	5	5	5	5	5	5	5	5	5	5	5	5	5	5	5	5	5	5
6	6	6	6	6	6	6	6	6	6	6	6	6	6	6	6	6	6	6	6	6	6
7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7
8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8
9	9	9	9	9	9	9	9	9	9	9	9	9	9	9	9	9	9	9	9	9	9

1) Data



2) Model

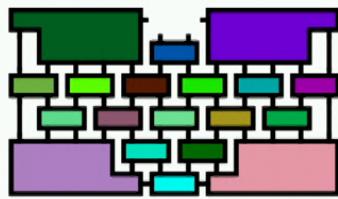
$$C(\mu) = \frac{1}{N} \sum_{i=1}^N (y_i - f(\mu, x_i))^2$$

3) Cost function

4) Minimize Cost function → Enjoy

# Machine learning as seen by a not so expert

0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	
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7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7	7
8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8	8
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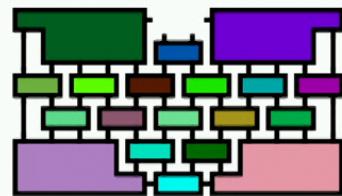


$$C(\mu) = \frac{1}{N} \sum_{i=1}^N (y_i - f(\mu, x_i))^2$$

- 1) Data
- 2) A highly structured model
- 3) Cost function
- 4) Minimize Cost function → Enjoy

# Machine learning as seen by a not so expert

$\tilde{Q}_n(v_1, \dots, v_n)$ .



$$C(\mu) = \frac{1}{N} \sum_{i=1}^N (y_i - f(\mu, x_i))^2$$

1) Data  
(active learning)

2) A highly  
structured model

3) Cost function

4) Minimize Cost function → Enjoy

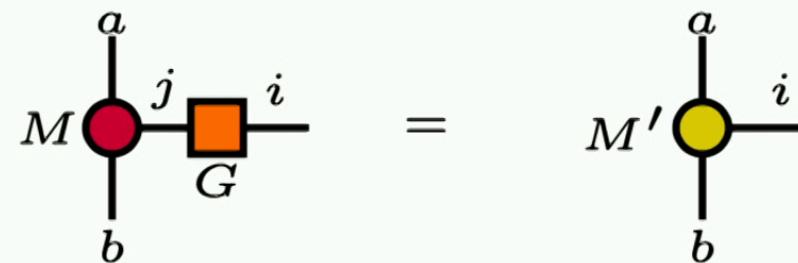
# Today:

- 1) Tensor Trains & the Tensor Cross Interpolation (TCI) algorithm.
- 2) Application to the calculation of Feynman diagrams.
- 3) Some Kondo physics.

Oseledet  
Dolgov & Savostyanov

# TENSOR NETWORKS 101

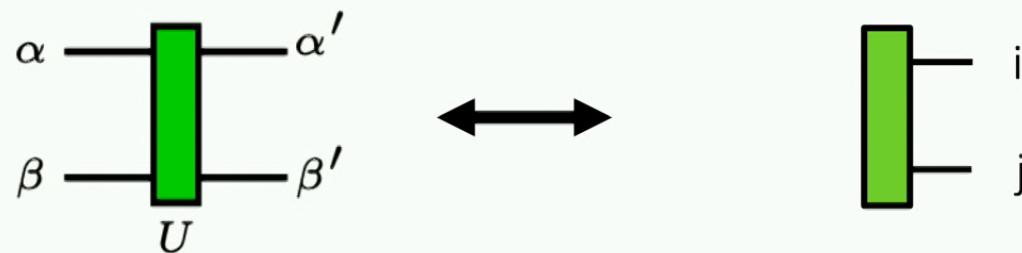
Contractions



$$M'_{abi} = \sum_j G_{ij} M_{abj}.$$

# TENSOR NETWORKS 101

## Grouping and degrouping of indices

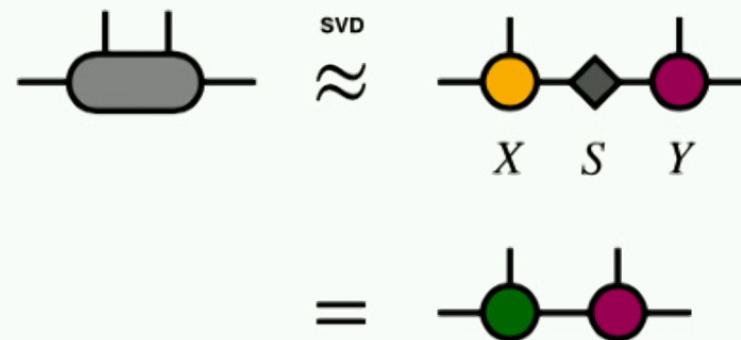


$$i = \alpha + N_\alpha \alpha' \text{ and } j = \beta + N_\beta \beta'$$

$$\hat{U}_{ij} \equiv U_{\alpha(i)\beta(j)\alpha'(i)\beta'(j)}.$$

# TENSOR NETWORKS 101

## Compressing (grouping, svd, degrouping)

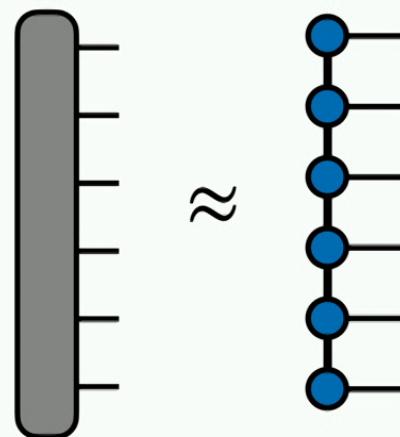


Truncated singular value decomposition provides the optimum low rank approximation of a matrix

By far the most popular tensor network: Matrix Product State (MPS)  
or Tensor Train (TT)

$$|\Psi\rangle = \sum_{i_1 \dots i_N} \sum_{\mu_1 \dots \mu_{N-1}} M_{\mu_1}^{i_1}(1) M_{\mu_1 \mu_2}^{i_2}(2) M_{\mu_2 \mu_3}^{i_3}(3) \dots M_{\mu_{N-1}}^{i_N}(N) |i_1 i_2 i_3 \dots i_N\rangle$$

(Can be seen as a variational ansatz in e.g. DMRG)



A large tool box for MPO/MPS:

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

$A \ x = b$

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = E \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

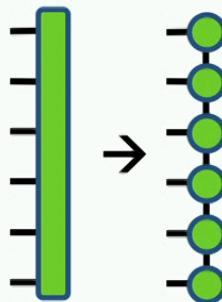
$H \ x = E \ x$

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

$b^T \ x$

(checkout ITensor  
from M. Stoudenmire  
In the room,  
<https://itensor.org>)

Tensor Cross Interpolation allows one to map many things onto this toolbox:



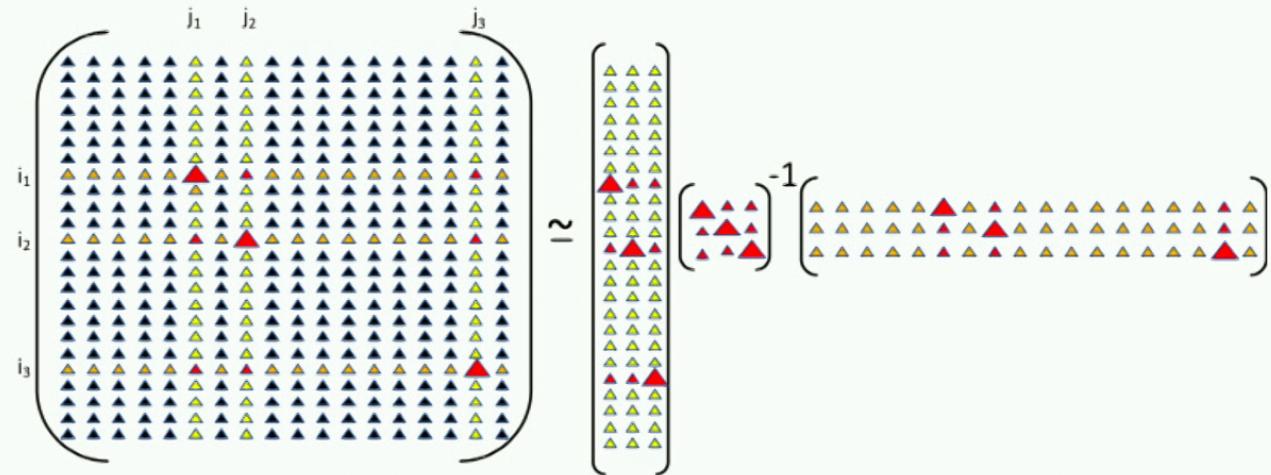
# CROSS INTERPOLATION 101

- Does not necessarily require the knowledge of the full matrix.
- Approximation given as slices of the original matrices.
- In its optimum version CI is to SVD what the infinite norm is to the  $L_2$  norm.

$$A = A \ A^{-1} \ A$$



$$A = A(\mathbb{I}, \mathbb{J}) \approx A(\mathbb{I}, \mathcal{J})A(\mathcal{I}, \mathcal{J})^{-1}A(\mathcal{I}, \mathbb{J})$$

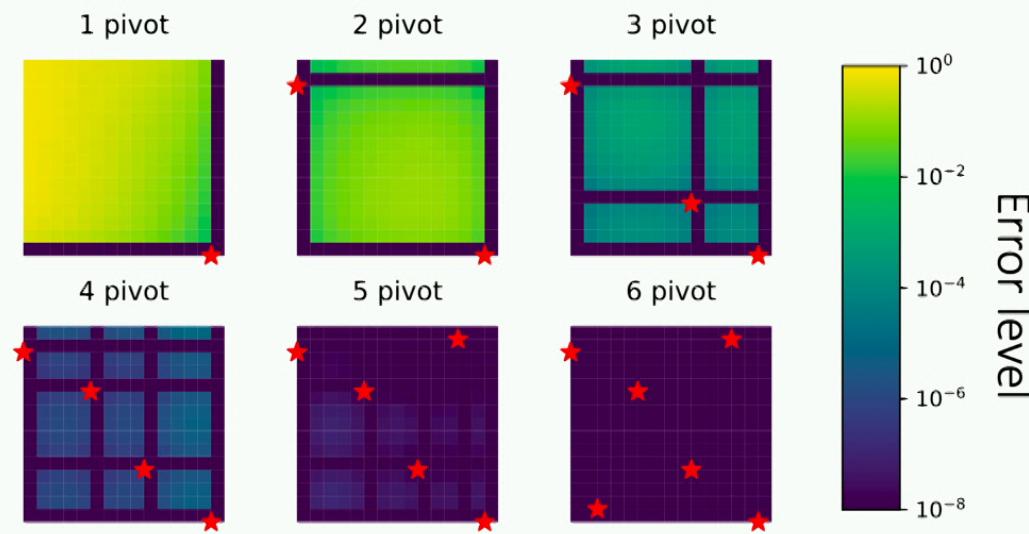


$$\begin{aligned}\mathcal{J} &= \{j_1, j_2, \dots, j_\chi\} \\ \mathbb{J} &= \{1, 2, \dots, N\}\end{aligned}$$

$$A = A(\mathbb{I}, \mathbb{J}) \approx A(\mathbb{I}, \mathcal{J})A(\mathcal{I}, \mathcal{J})^{-1}A(\mathcal{I}, \mathbb{J})$$

- (P1) It is an interpolation, i.e. it is exact for any  $i \in \mathcal{I}$  or  $j \in \mathcal{J}$ . This can be straightforwardly checked from the definition as e.g.  $A(\mathcal{I}, \mathcal{J})A(\mathcal{I}, \mathcal{J})^{-1}A(\mathcal{I}, \mathbb{J}) = A(\mathcal{I}, \mathbb{J})$ .
- (P2) It is exact if the matrix  $A$  has rank  $\chi$  (cf. Appendix B for a simple proof).

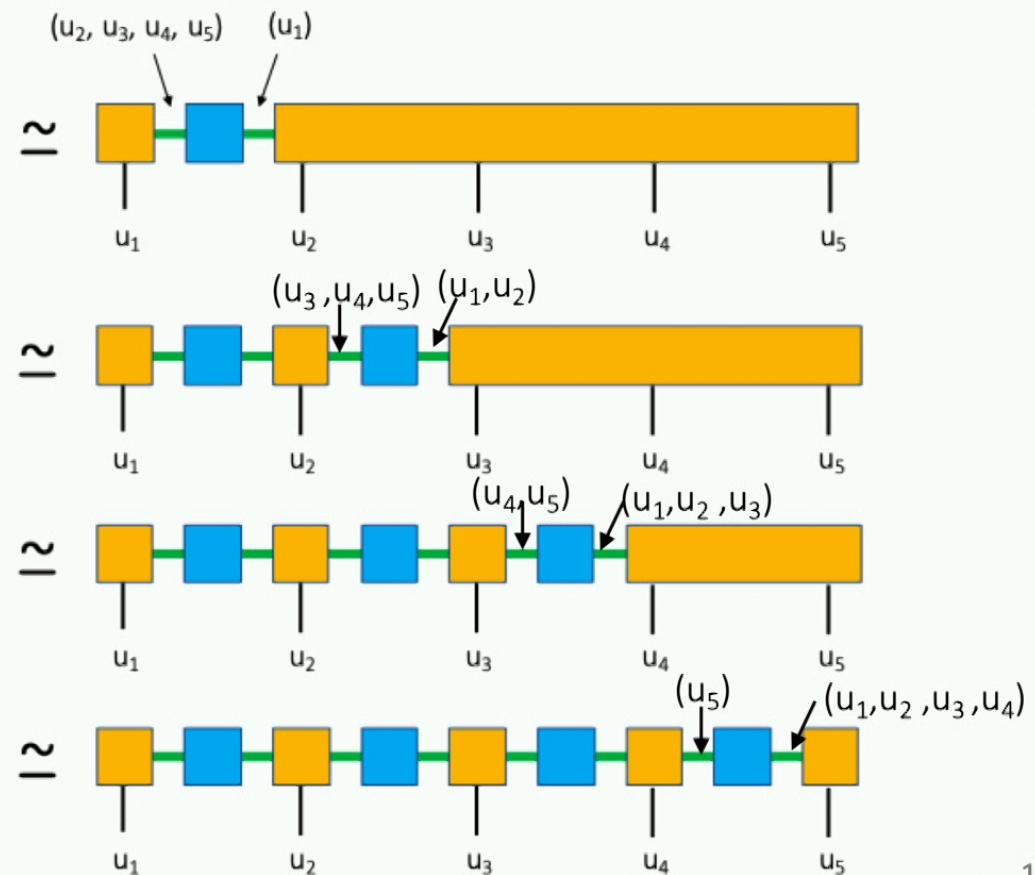
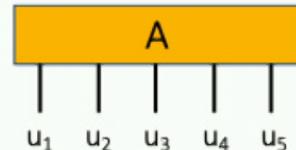
# An example of Cross Interpolation



$$M(i,j) = 1 + (i/N)^2 \cos\left(\frac{ij}{N^2}\right) + (j/N)^2 \sin\left(\frac{ij}{N^2}\right) + \frac{i}{j+1}$$

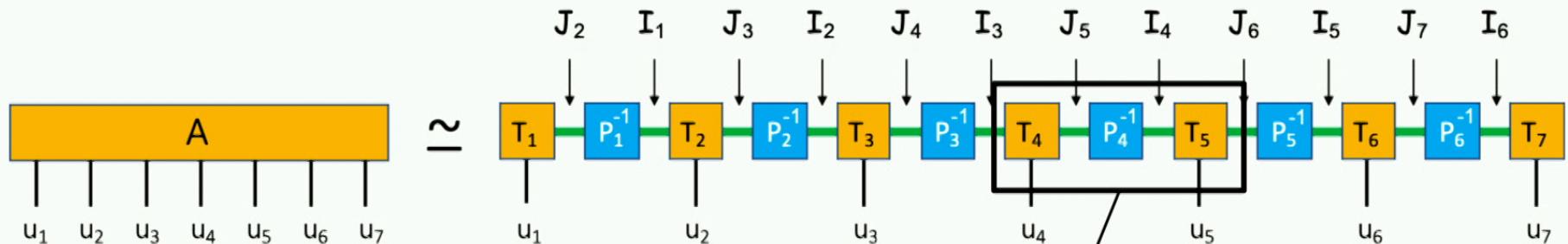
# TENSOR CROSS INTERPOLATION

- toy (exponential) algorithm
- approximation entirely made out of slices of A



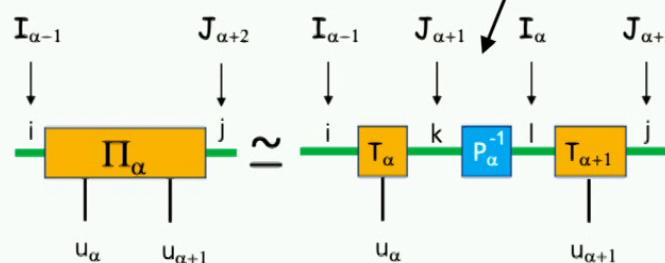
15

# THE ACTUAL ALGORITHM (OLD VERSION)



- Initialize with a rank 1 approximation (a single call to the integrand)
- Loop over pairs of « T » to improve the cross approximation (add pivots)

$$(u_1, \dots, u_n) = (u_1, \dots, u_\alpha) \oplus (u_{\alpha+1}, \dots, u_n)$$

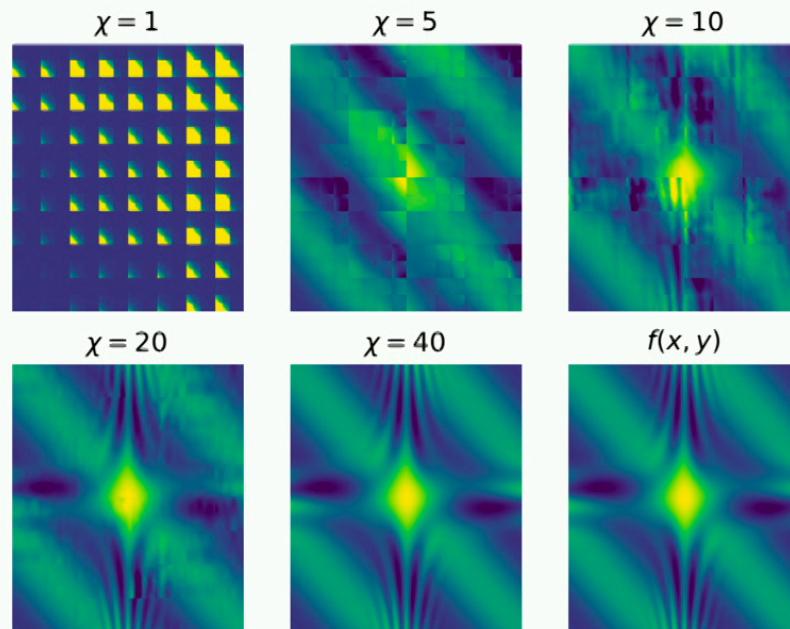


S. Dolgov and D. Savostyanov, Parallel cross interpolation for high-precision calculation of high-dimensional integrals, Computer Physics Communications **246**, 106869 (2020).

Example: « quantics » , representation of low dimension function with exponential ( $2^n$ ) resolution

$$f(x, y) = e^{-0.4(x^2+y^2)} + 1 + \sin(xy)e^{-x^2} + \cos(3xy)e^{-y^2} + \cos(x+y)$$

$$x = \frac{x_1}{2} + \frac{x_2}{2^2} + \frac{x_3}{2^3} + \dots + \frac{x_n}{2^n} \quad x_i \in \{0, 1\}$$



17

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# TOOLS FOR QUANTUM TRANSPORT

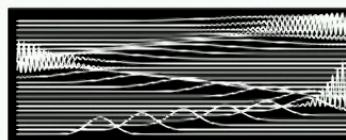
$$\hat{H} = \sum_{ij} H_{ij} c_i^\dagger c_j$$

$$\hat{H}(t) = \sum_{ij} H_{ij}(t) c_i^\dagger c_j$$

$$\hat{\mathbf{H}}(t) = \hat{\mathbf{H}}_0(t) + U\hat{\mathbf{H}}_{\text{int}}(t)$$

$$\hat{\mathbf{H}}_{\text{int}}(t) = \sum_{ijkl} \mathbf{V}_{ijkl}(t) \hat{\mathbf{c}}_i^\dagger \hat{\mathbf{c}}_j^\dagger \hat{\mathbf{c}}_k \hat{\mathbf{c}}_l.$$

**kwant**



[HTTP://KWANT-PROJECT.ORG](http://KWANT-PROJECT.ORG)

With TU Delft (Akhmerov, Wimmer).

Main developer: C. Groth

[HTTP://TKWANT.KWANT-PROJECT.ORG](http://TKWANT.KWANT-PROJECT.ORG)

Main developer: T. Kloss

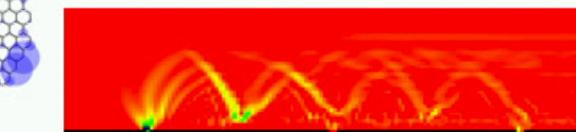
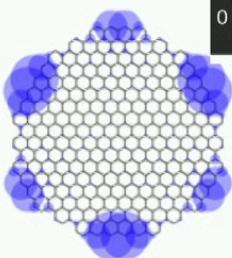
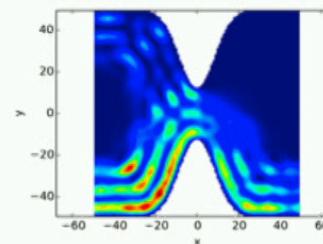
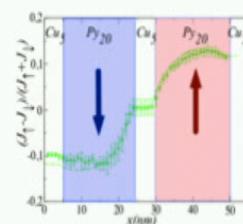
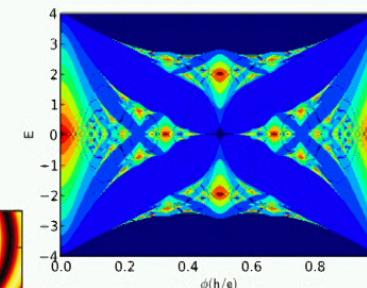
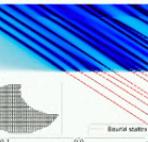
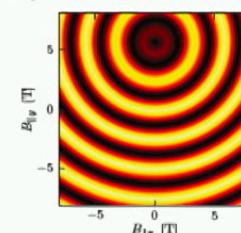
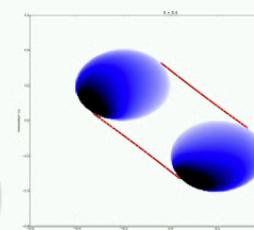
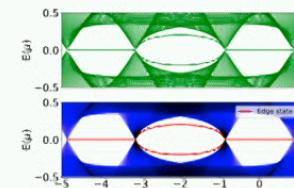
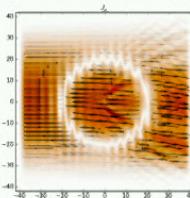
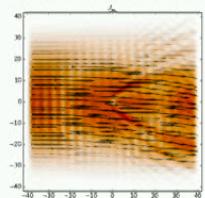
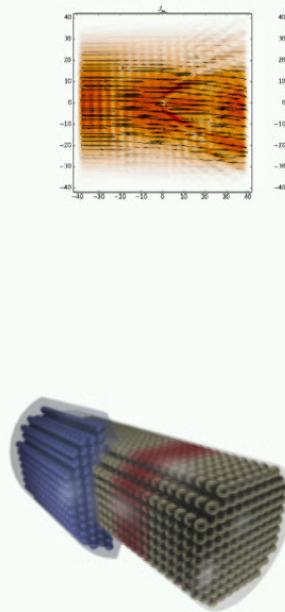
- Kondo physics
- Coulomb blockade, Fermi edge singularity
- 0.7 anomaly
- FQHE
- Quantum computers

...

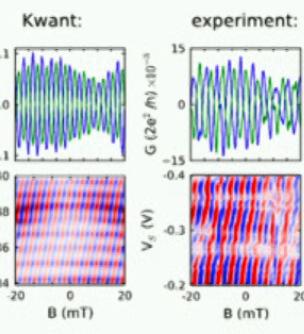
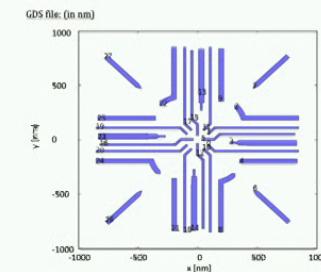
19

# KWANT GALLERY

With TU Delft (Akhmerov, Wimmer et al.)

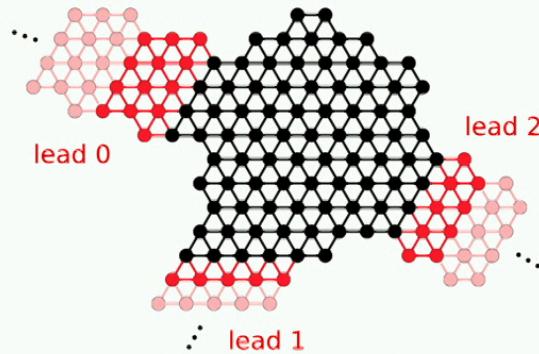


[HTTP://KWANT-PROJECT.ORG](http://kwant-project.org)

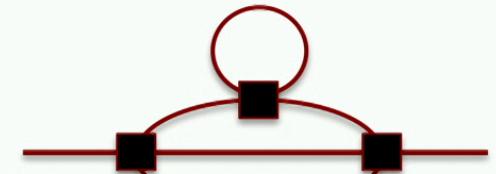


$$\hat{\mathbf{H}}(t) = \hat{\mathbf{H}}_0(t) + U\hat{\mathbf{H}}_{\text{int}}(t)$$

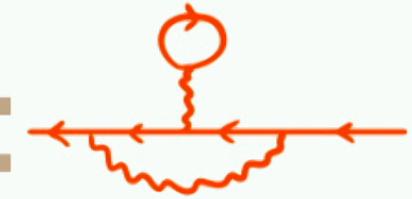
$$\hat{\mathbf{H}}_{\text{int}}(t) = \sum_{ijkl} \mathbf{V}_{ijkl}(t) \hat{\mathbf{c}}_i^\dagger \hat{\mathbf{c}}_j^\dagger \hat{\mathbf{c}}_k \hat{\mathbf{c}}_l.$$



$$Q(U) = \sum_{n=0}^{+\infty} Q_n U^n$$



# A VERY BRAVE/STUPID APPROACH:



→ CALCULATE ALL THE FEYNMAN DIAGRAMS UP TO A GIVEN (LARGE) ORDER

- PROBLEM #1: There are  $n!$  diagrams.

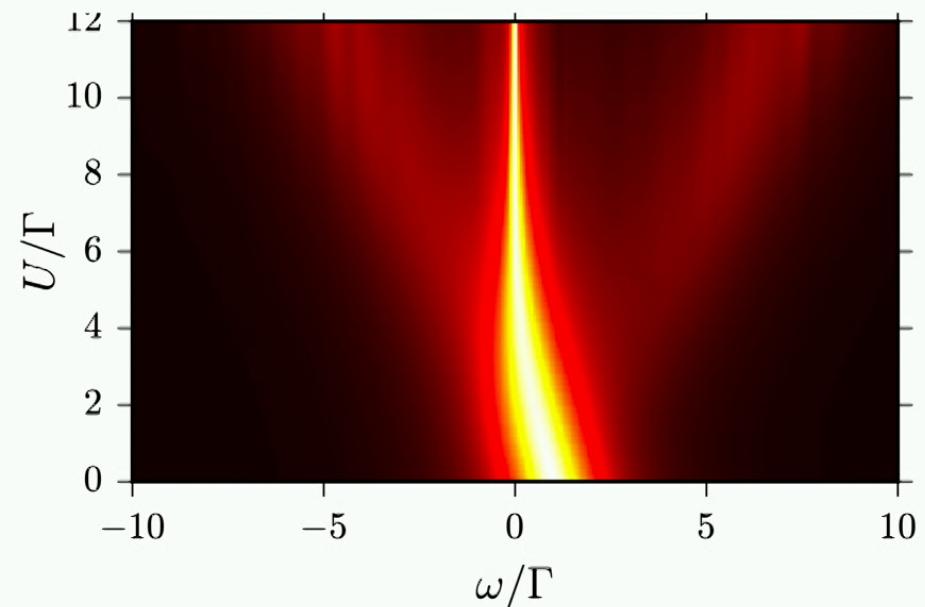
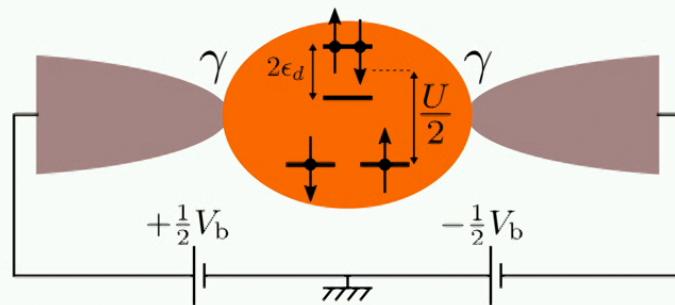
$$F(U) = \sum_n F_n U^n$$

- PROBLEM #2 How to calculate  $n$  dimensional integrals
- PROBLEM #3 How to reconstruct  $F(U)$  from the  $F_n$ .

$$q_n(t = \infty) \sim \frac{1}{r^n} \quad q_n(t) \sim \frac{1}{n!} t^n$$



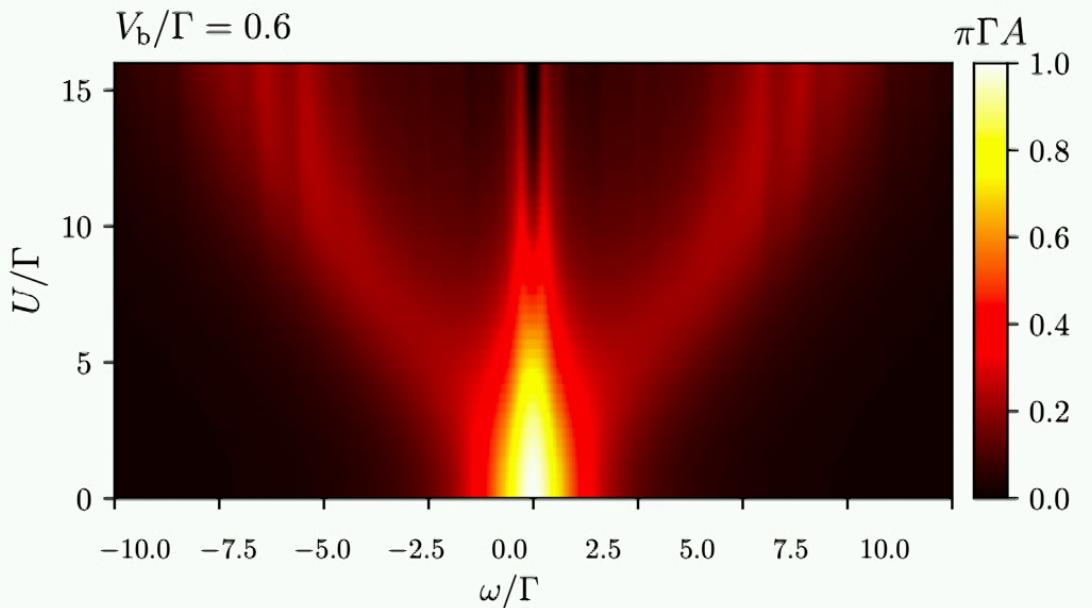
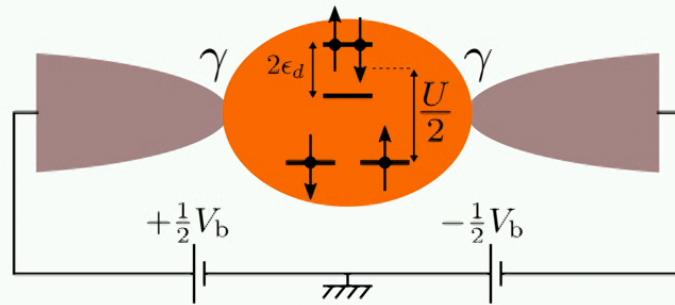
# TEASER: AT EQUILIBRIUM



$$\hat{\mathbf{H}} = \sum_{i=-\infty}^{+\infty} \sum_{\sigma} \gamma_i \hat{\mathbf{c}}_{i,\sigma}^\dagger \hat{\mathbf{c}}_{i+1,\sigma} + h.c. + \epsilon_d (\hat{\mathbf{n}}_\uparrow + \hat{\mathbf{n}}_\downarrow) + U \theta(t) \left( \hat{\mathbf{n}}_\uparrow - \frac{1}{2} \right) \left( \hat{\mathbf{n}}_\downarrow - \frac{1}{2} \right).$$

24

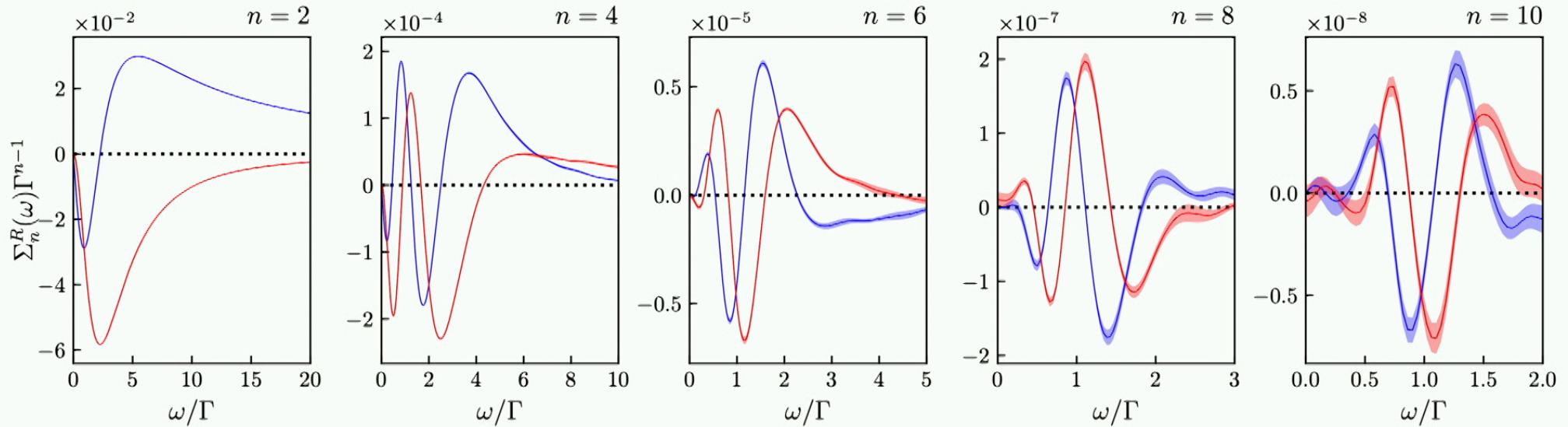
# TEASER: OUT-OF-EQUILIBRIUM



$$\hat{H} = \sum_{i=-\infty}^{+\infty} \sum_{\sigma} \gamma_i \hat{c}_{i,\sigma}^\dagger \hat{c}_{i+1,\sigma} + h.c. + \epsilon_d (\hat{n}_\uparrow + \hat{n}_\downarrow) + U \theta(t) \left( \hat{n}_\uparrow - \frac{1}{2} \right) \left( \hat{n}_\downarrow - \frac{1}{2} \right).$$

25

# PROBLEM #1: THE N! DIAGRAMS

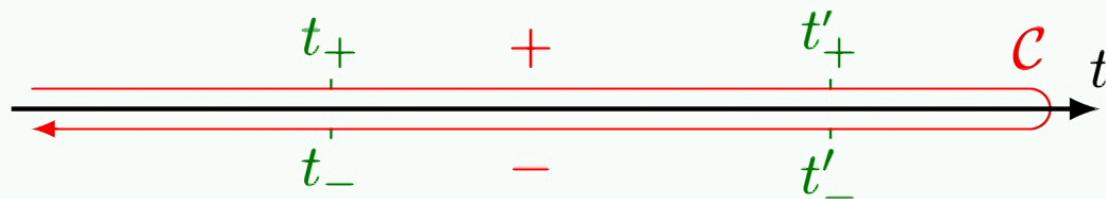


(« old » data, using diagrammatic Monte-Carlo)

# KELDYSH FORMALISM IN A NUTSHELL

$$\langle \mathcal{O}(t) \rangle = \langle \mathcal{U}^\dagger(t) \mathcal{O} \mathcal{U}(t) \rangle$$

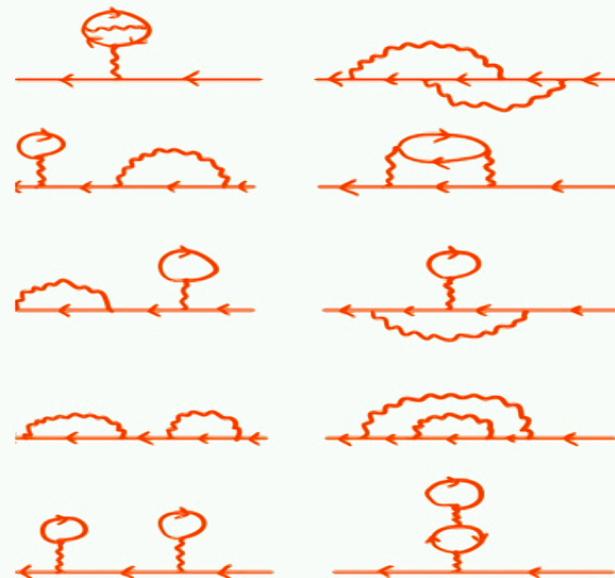
$$\mathcal{U}(t) = T \exp \left( -i \int_0^t \hat{H}_{\text{int}}(u) du \right)$$



$$\langle \mathcal{O}(t) \rangle = \left\langle \mathcal{T}_C \hat{\mathcal{O}}(t) \exp \left( -i \int_C \hat{H}_{\text{int}}(u) du \right) \right\rangle$$

# WICK DETERMINANTS

$$\left\langle c_1^+ c_1 c_2^+ c_2 c_3^+ c_3 c_4^+ c_4 c_5^+ c_5 \right\rangle = \sum_P (-1)^{|P|} \left\langle c_1^+ c_{P(1)} \right\rangle \left\langle c_2^+ c_{P(2)} \right\rangle \left\langle c_3^+ c_{P(3)} \right\rangle \left\langle c_4^+ c_{P(4)} \right\rangle \left\langle c_5^+ c_{P(5)} \right\rangle$$



# WICK DETERMINANTS

$$\langle c_1^+ c_1 c_2^+ c_2 c_3^+ c_3 c_4^+ c_4 c_5^+ c_5 \rangle = \sum_P (-1)^{|P|} \langle c_1^+ c_{P(1)} \rangle \langle c_2^+ c_{P(2)} \rangle \langle c_3^+ c_{P(3)} \rangle \langle c_4^+ c_{P(4)} \rangle \langle c_5^+ c_{P(5)} \rangle$$

$$\langle c_1^+ c_1 c_2^+ c_2 c_3^+ c_3 c_4^+ c_4 c_5^+ c_5 \rangle = \det \langle c_i^+ c_j \rangle$$

# A « VERY SIMPLE » FORMULA

$$G_{ij}^c(\bar{t}, \bar{t}') = \sum_{n=0}^{+\infty} \frac{i^n}{n!} U^n \sum_{\{a_i\}} (-1)^{\sum_i a_i} \int du_1 du_2 \dots du_n \sum_{i_1 j_1 k_1 l_1} V_{i_1 j_1 k_1 l_1}(u_1) \dots \sum_{i_n j_n k_n l_n} V_{i_n j_n k_n l_n}(u_n) \det \mathbf{M}_n$$

2<sup>n</sup> sum cancels  
disconnected  
diagrams

$$\mathbf{M}_n = \begin{pmatrix} g_{k_1 i_1}^<(\bar{u}_1, \bar{u}_1) & g_{k_1 j_1}^<(\bar{u}_1, \bar{u}_1) & g_{k_1 i_2}^c(\bar{u}_1, \bar{u}_2) & \dots & g_{k_1 j}^c(\bar{u}_1, \bar{t}') \\ g_{l_1 i_1}^<(\bar{u}_1, \bar{u}_1) & g_{l_1 j_1}^<(\bar{u}_1, \bar{u}_1) & g_{l_1 i_2}^c(\bar{u}_1, \bar{u}_2) & \dots & g_{l_1 j}^c(\bar{u}_1, \bar{t}') \\ g_{k_2 i_1}^c(\bar{u}_2, \bar{u}_1) & g_{k_2 j_1}^c(\bar{u}_2, \bar{u}_1) & g_{k_2 i_2}^<(\bar{u}_2, \bar{u}_2) & \dots & g_{k_2 j}^c(\bar{u}_2, \bar{t}') \\ \dots & \dots & \dots & \dots & \dots \\ g_{k_n i_1}^c(\bar{u}_n, \bar{u}_1) & g_{k_n j_1}^c(\bar{u}_n, \bar{u}_1) & g_{k_n i_2}^c(\bar{u}_n, \bar{u}_2) & \dots & g_{k_n j}^c(\bar{u}_n, \bar{t}') \\ g_{l_n i_1}^c(\bar{u}_n, \bar{u}_1) & g_{l_n j_1}^c(\bar{u}_n, \bar{u}_1) & g_{l_n i_2}^c(\bar{u}_n, \bar{u}_2) & \dots & g_{l_n j}^c(\bar{u}_n, \bar{t}') \\ g_{ii_1}^c(\bar{t}, \bar{u}_1) & g_{ij_1}^c(\bar{t}, \bar{u}_1) & g_{ii_2}^c(\bar{t}, \bar{u}_2) & \dots & g_{ij}^c(\bar{t}, \bar{t}') \end{pmatrix}$$

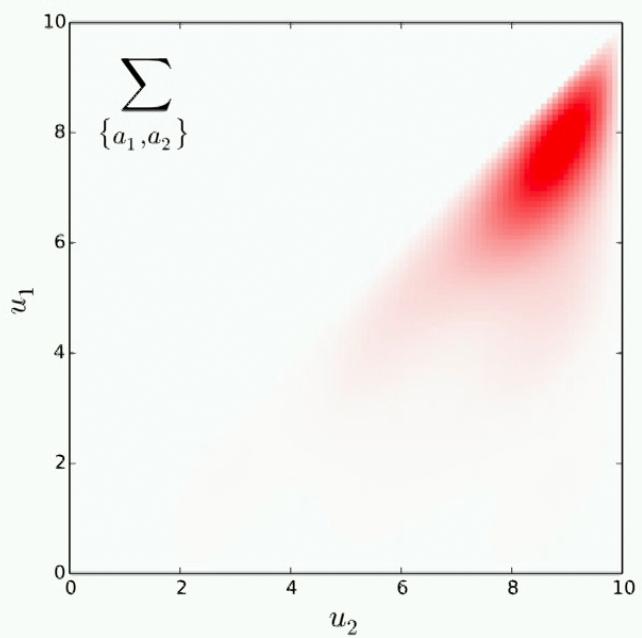
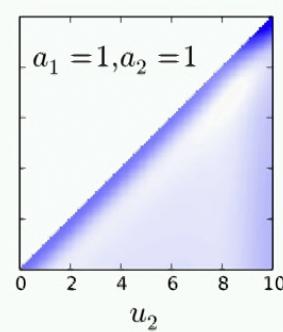
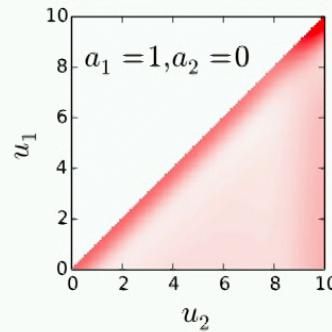
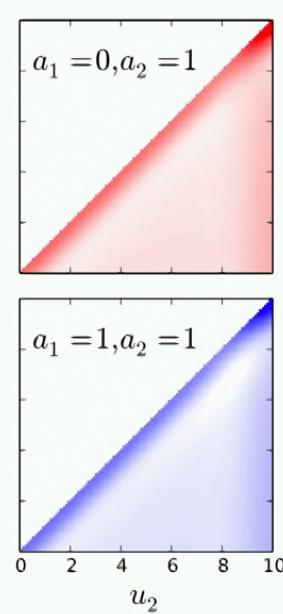
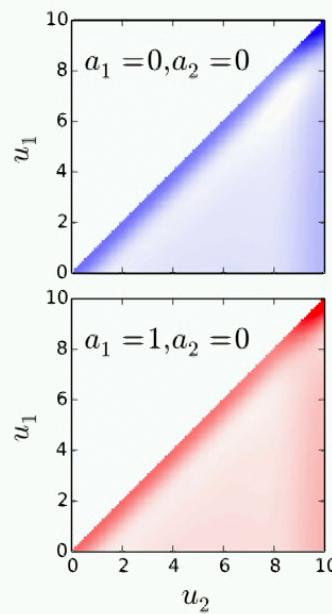


$$g_{ij}^<(t, t') = i \sum_{\alpha} \int \frac{dE}{2\pi} f_{\alpha}(E) \Psi_{\alpha E}(t, i) \Psi_{\alpha E}^*(t', j) + i \sum_n f(E_n) \Psi_n(t, i) \Psi_n^*(t', j)$$

**Known non-interacting functions.**

$$G_{ij}^c(\bar{t}, \bar{t}') = \sum_{n=0}^{+\infty} \frac{i^n}{n!} U^n \sum_{\{a_i\}} (-1)^{\sum_i a_i} \int du_1 du_2 \dots du_n \sum_{i_1 j_1 k_1 l_1} V_{i_1 j_1 k_1 l_1}(u_1) \dots \sum_{i_n j_n k_n l_n} V_{i_n j_n k_n l_n}(u_n) \det \mathbf{M}_n$$

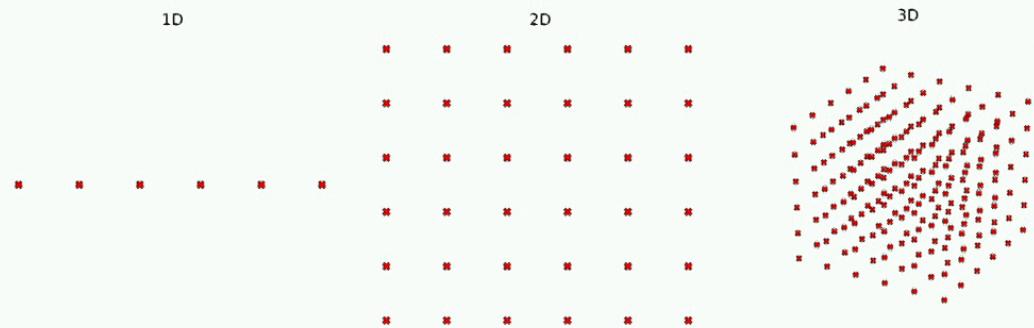
**$2^n$  sum cancels disconnected diagrams**



# PROBLEM #2 THE N DIMENSIONAL INTEGRAL

→ The dimensionality curse,  
 $L^d$  points.

$$Q(U) = \sum_n Q_n U^n$$



$$Q_n = \int dv_1 \dots dv_n \tilde{Q}_n(v_1, \dots, v_n).$$

The standard approach: Metropolis Monte-Carlo.

- Intrinsic very slow convergence  $N^{-1/2}$  (as opposed to  $1/N^{15}$  or even exponential in 1D)
- We do not build any knowledge of  $f(u)$
- The infamous sign problem.

# MACHINE LEARNING THE INTEGRAND

What if the integrand factorized?

$$\tilde{Q}_n(v_1, \dots, v_n) \approx M_1(v_1) \cdots M_n(v_n)$$

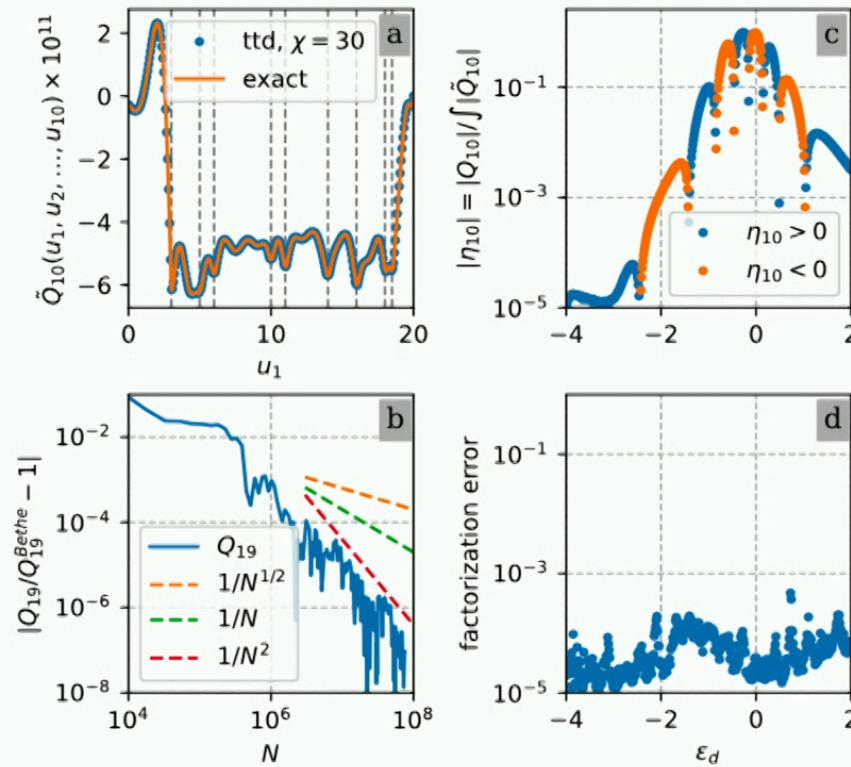
Then the integration would be trivial

$$Q_n \approx \left( \int dv_1 M_1(v_1) \right) \cdots \left( \int dv_n M_n(v_n) \right)$$

→ A job for Tensor Cross Interpolation

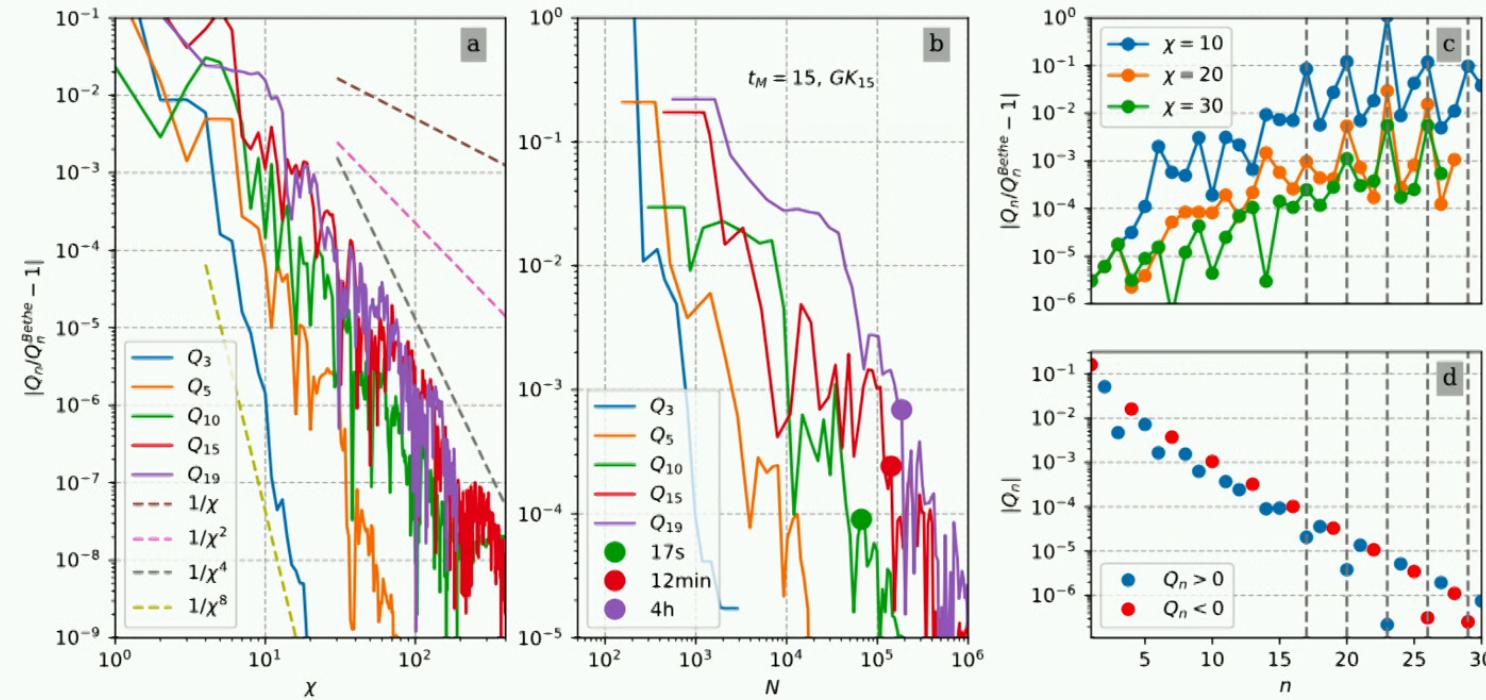
# BENCHMARKING THE RESULTS

Unprecedented convergence →

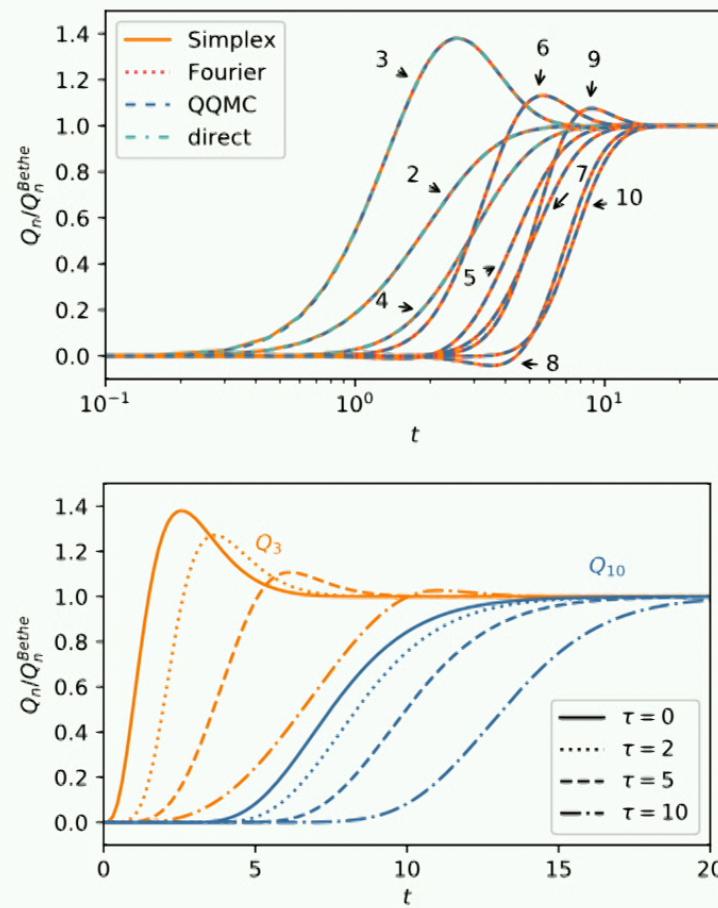


← Agnostic to sign problem

# BENCHMARKING THE RESULTS

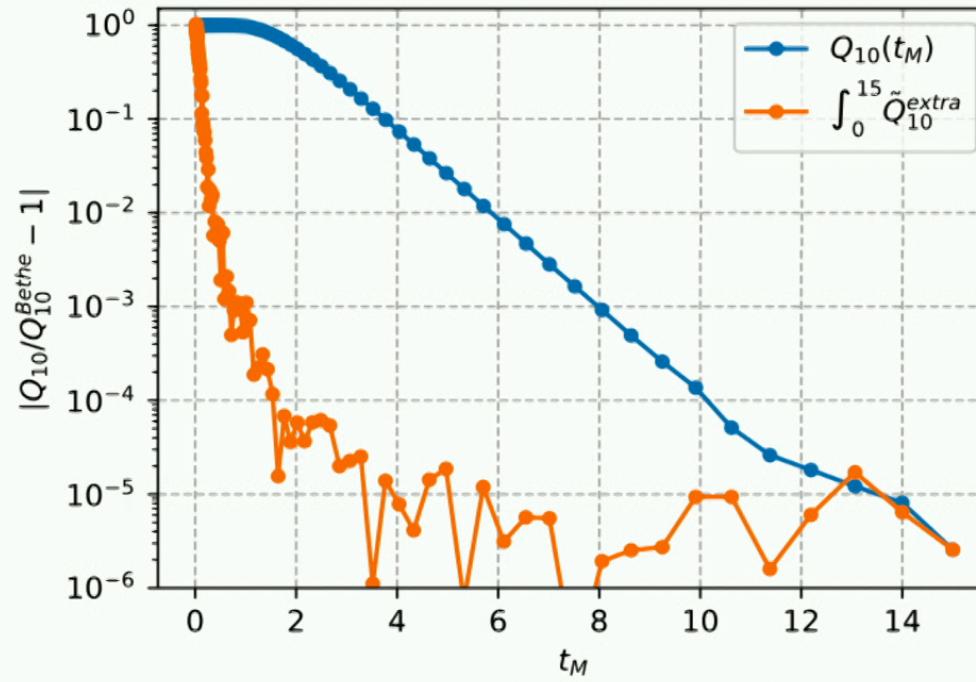


# INTEGRATION IS A COSTLESS POST PROCESSING STEP



36

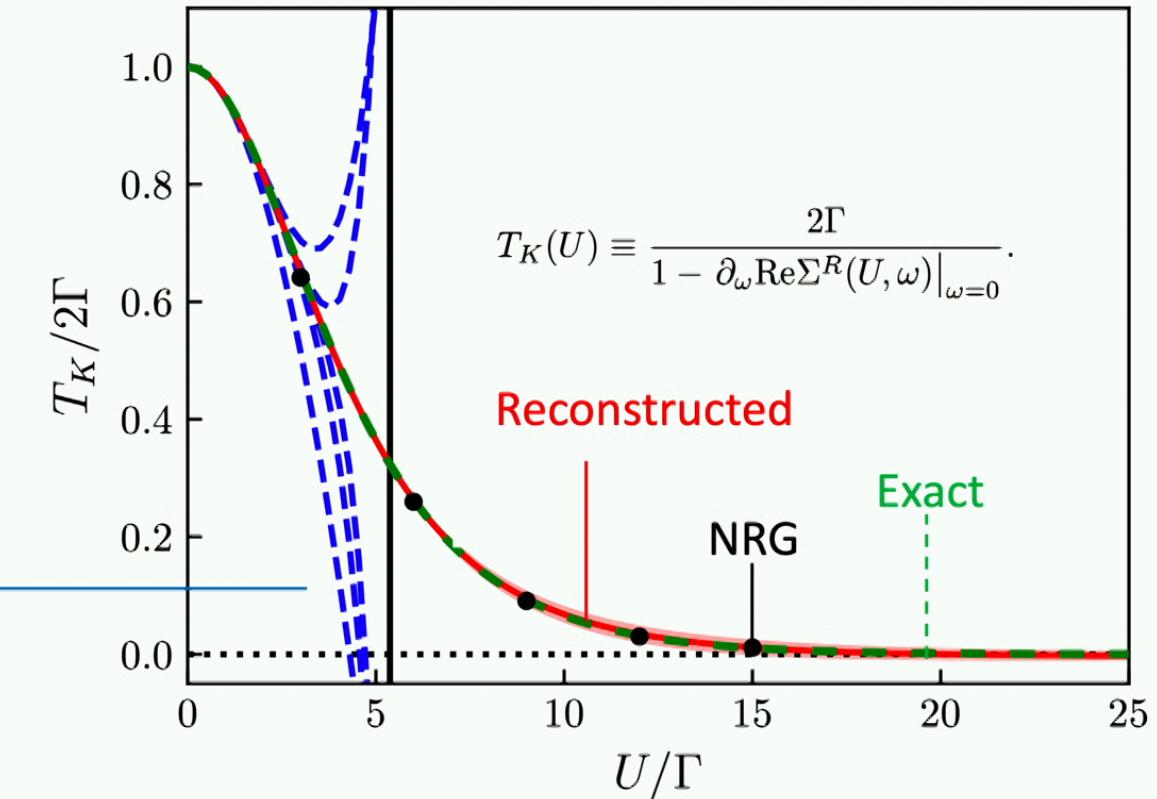
# THIS IS MORE THAN INTERPOLATION



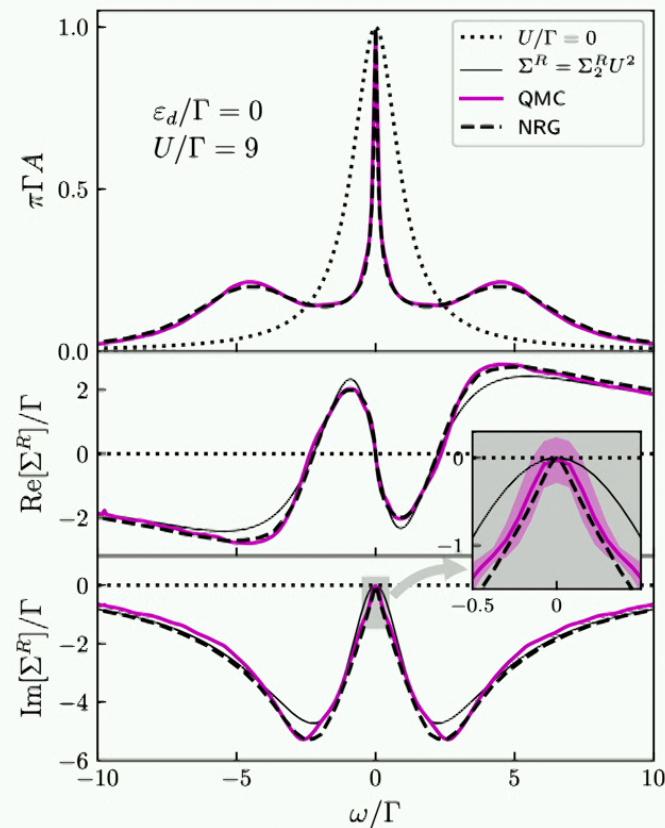
# PROBLEM #3: RECONSTRUCTING $F(U)$ FROM $F_N$

$$F(U) = \sum_n F_n U^n$$

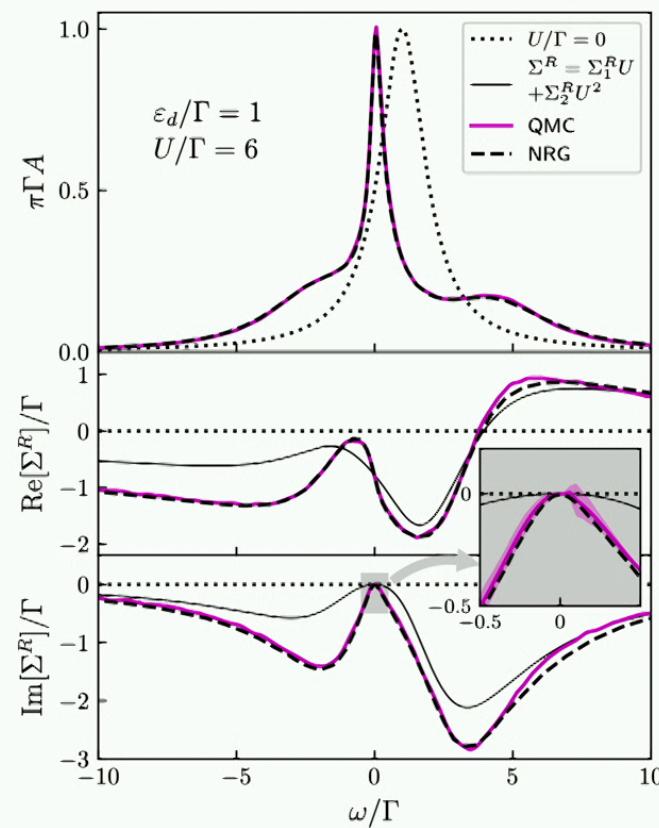
Truncated sum  
(Monte-Carlo data)



# SOME EQUILIBRIUM BENCHMARKS

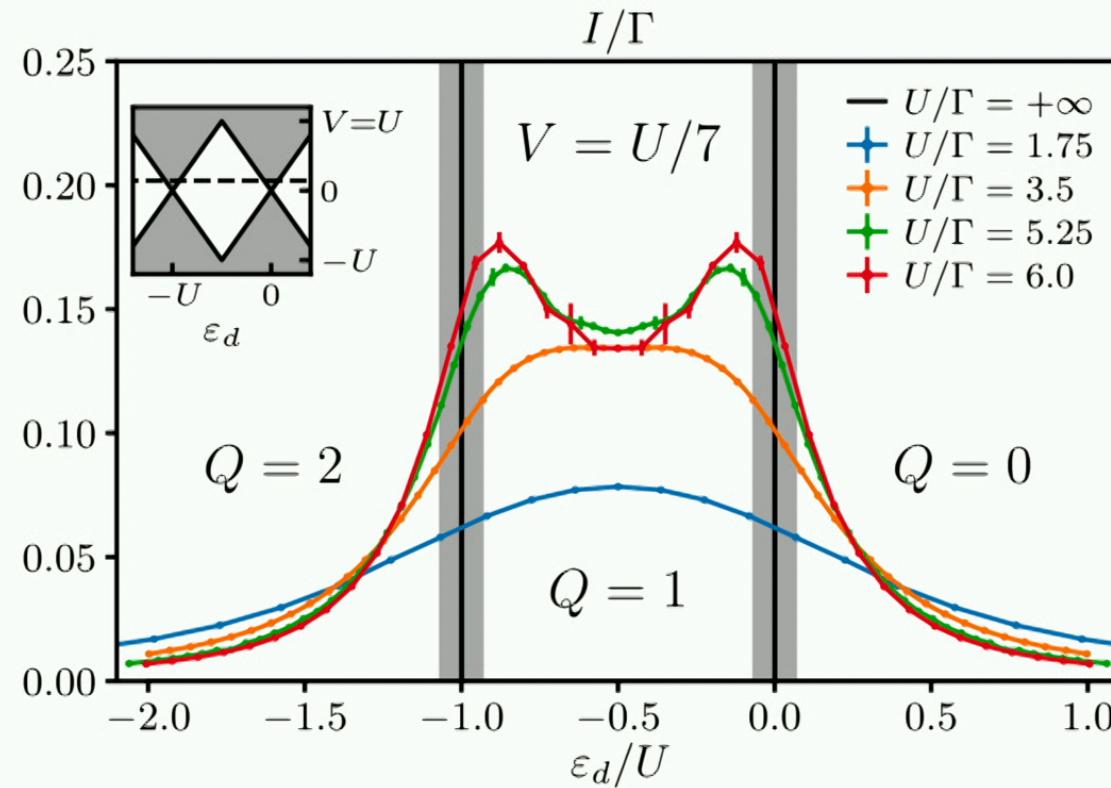


(Monte-Carlo data)



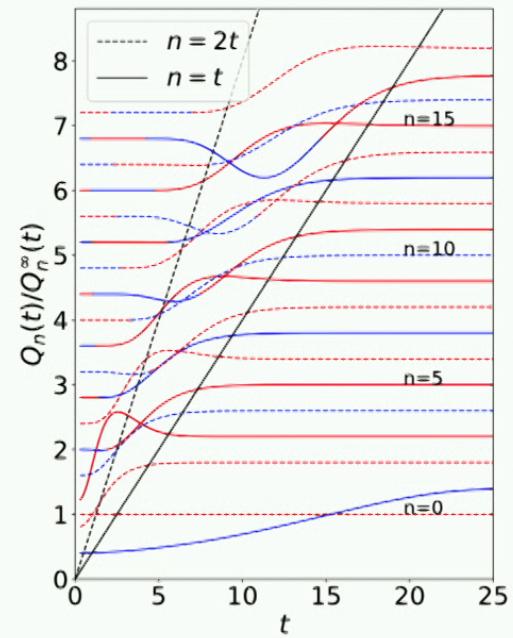
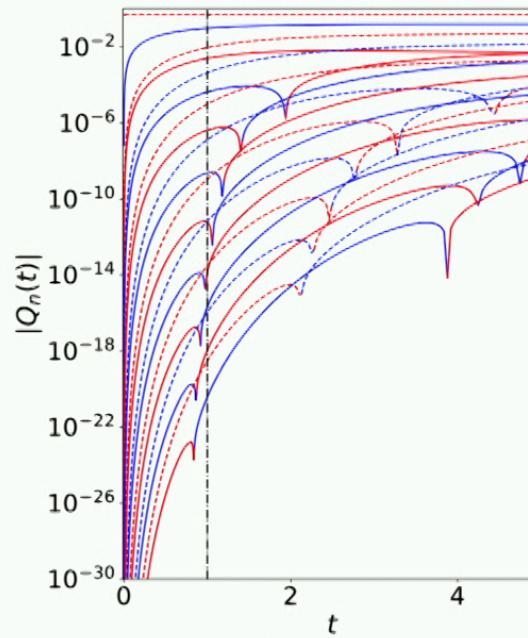
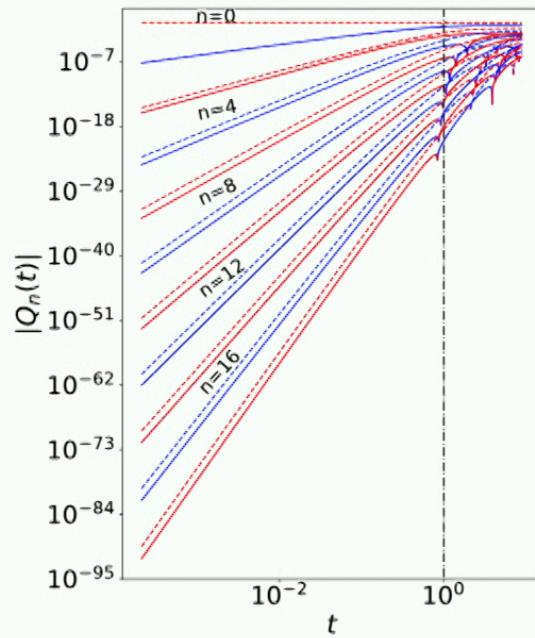
39

# AND BACK TO NON-EQUILIBRIUM

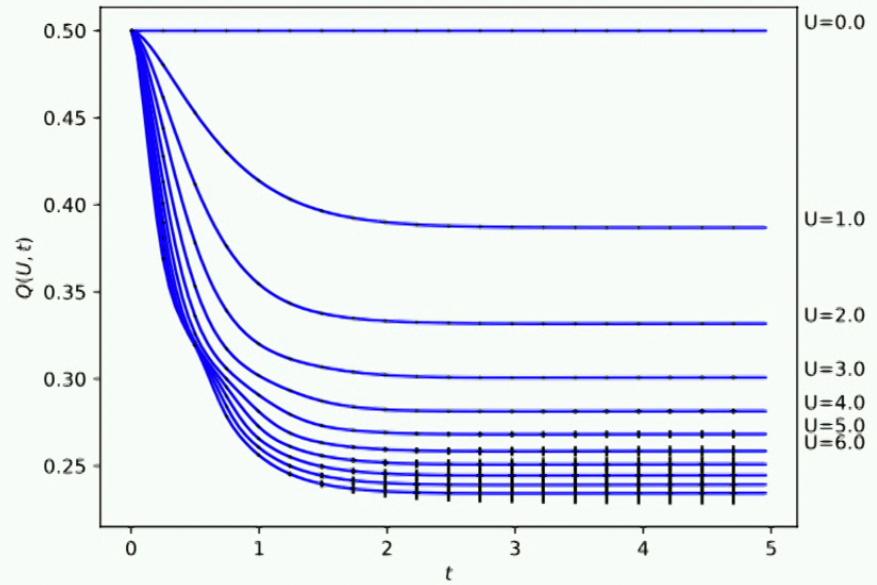
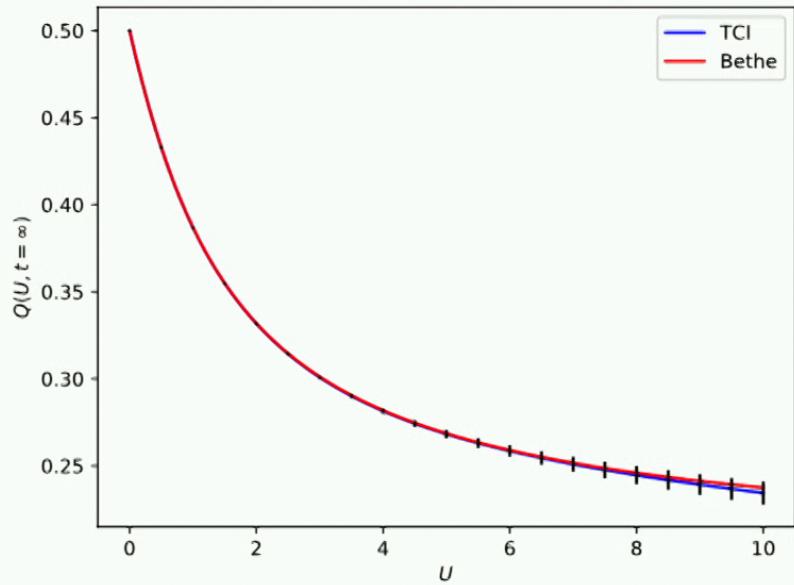


(low discrepancy  
sequence integratic  
data

# Data with the new tensor train technique



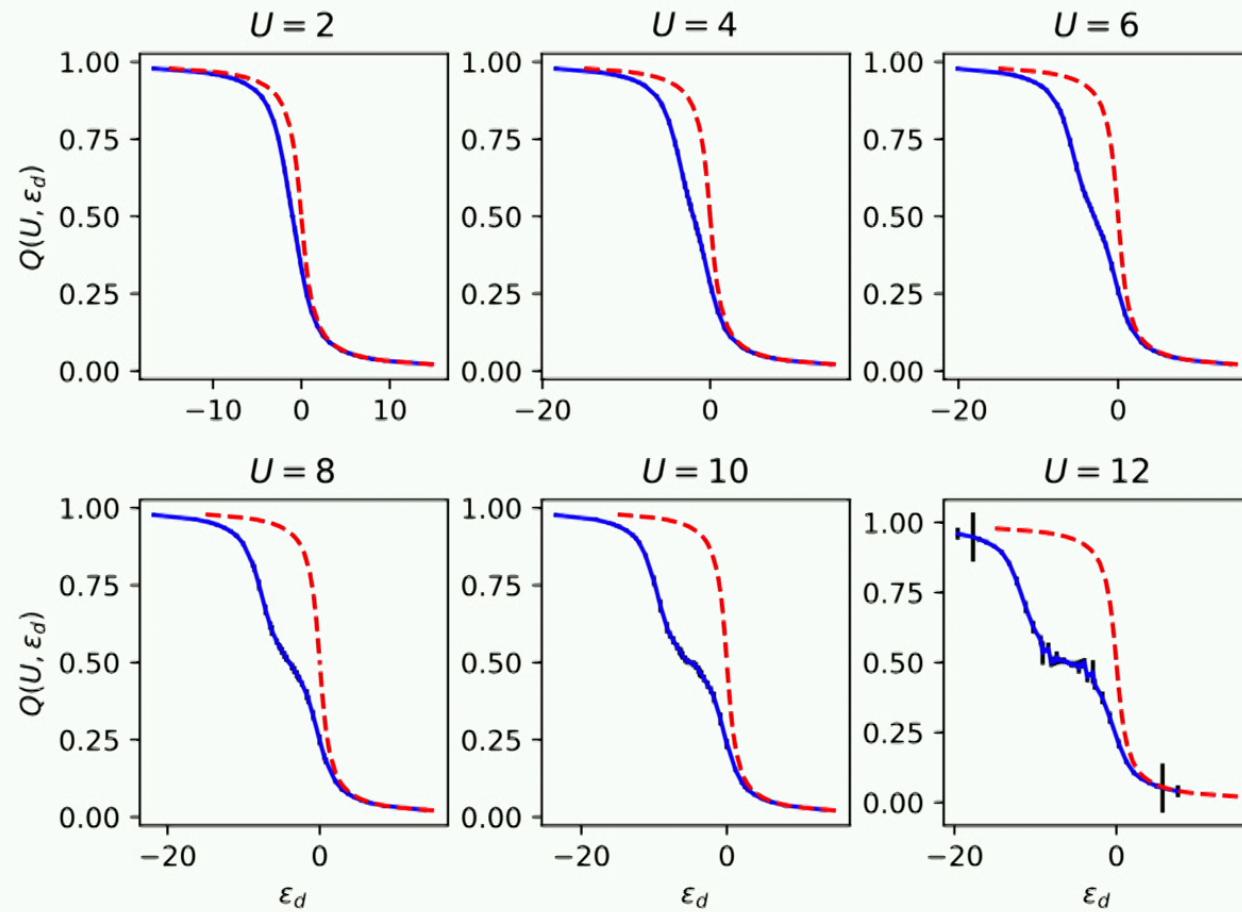
# Data with the new tensor train technique



$$Q(U, t) \approx \sum_{n=0}^N q_n(t) U^n$$

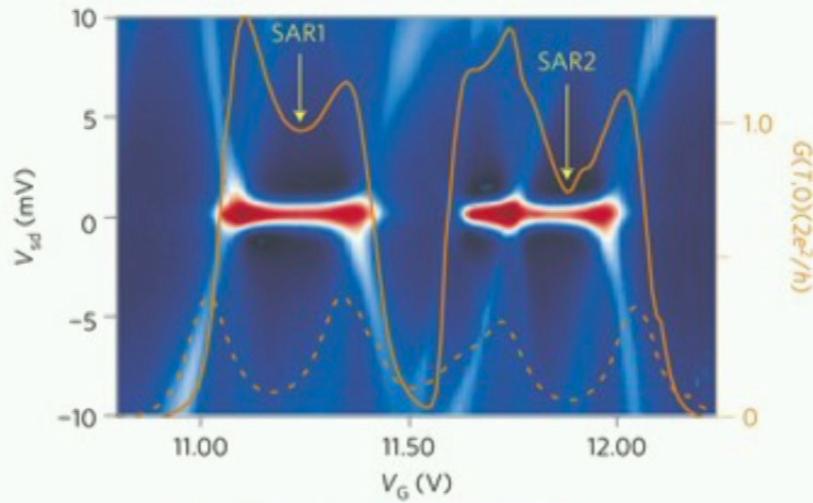
42

## Data with the new tensor train technique

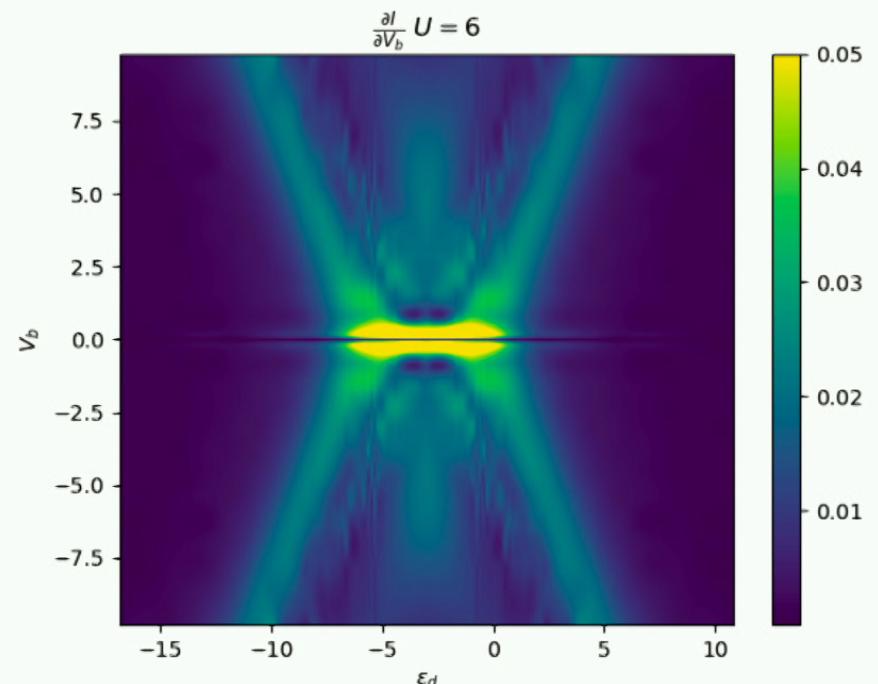


43

# Data with the new tensor train technique



Nature Physics 5, 208 (2009)



Caution: fresh from the oven)

# CONCLUSION

- In principle, tensor cross-interpolation can be used (tried) wherever Monte-Carlo has been.  
→ that's a lot of things to explore (imaginary time, quantum field approaches, classical problems...) → Got started on hybridization expansion [arXiv:2303.11199](https://arxiv.org/abs/2303.11199)
- It can also be used for other problems (quantics/superfast FFT, learning an experiment, ising models) [arXiv:2303.11819](https://arxiv.org/abs/2303.11819)
- We do not know yet how far we will be able to go with our tensor train diagrammatic approach  
→ Hubbard models, DMFT, 0.7 anomaly, few qubits and their baths
  - Open source software « xfac » coming up (with a better TCI algorithm).
  - PEPS...

PHYSICAL REVIEW LETTERS **125**, 047702 (2020)

PHYSICAL REVIEW X **9**, 041008 (2019)

PHYSICAL REVIEW B **100**, 125129 (2019)

PHYSICAL REVIEW B **91**, 245154 (2015)

<https://arxiv.org/abs/2207.06135>

Phys. Rev. X **12**, 041018 (2022).

I am hiring! (postdocs/PhD)

45