

Title: Testing quantum states

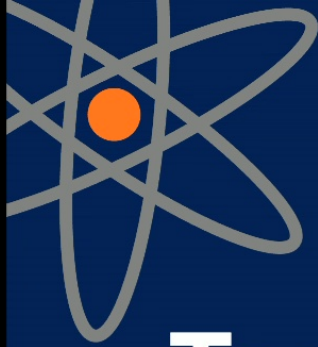
Speakers: Mehdi Soleimanifar

Date: May 24, 2023 - 11:00 AM

URL: <https://pirsa.org/23050151>

Abstract: In this talk, I will present three algorithms that address distinct variants of the problem of testing quantum states. First, I will discuss the problem of statistically testing whether an unknown quantum state is a matrix product state of certain bond dimension or it is far from all such states. Next, I will demonstrate a method for testing whether a bipartite quantum state, shared between two parties, corresponds to the ground state of a given gapped local Hamiltonian. Finally, I will present a scheme for verifying that a machine learning model of an unknown quantum state has high overlap with the actual state.

Zoom Link: <https://pitp.zoom.us/j/99250127489?pwd=UCtXUi9zMzJZamppT29DbWtJcWU3Zz09>

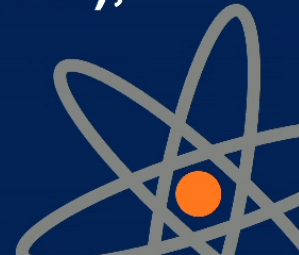


Testing quantum states

Mehdi Soleimanifar (Caltech)

Based on joint works with

**Anurag Anshu (Harvard) , Aram Harrow (MIT),
and Robert Huang (Caltech)**



1. Testing matrix product states

2. Testing machine learning models

3. Two-party testing of ground states



State of n qudits

is a vector in tensor product space $(\mathbb{C}^d)^{\otimes n}$

$$|\psi\rangle = \sum_{i_1, \dots, i_n=1}^d a_{i_1 \dots i_n} |i_1\rangle \cdots |i_n\rangle$$

$$a_{i_1 \dots i_n} \in \mathbb{C} \quad \|\psi\rangle\|_2 = 1$$

$|\psi\rangle$ can be product or entangled:

Product state: $|\psi\rangle = |\psi_1\rangle \otimes |\psi_2\rangle \cdots \otimes |\psi_n\rangle$



Questions we can ask:

Is a state $|\psi\rangle$ entangled or product?

How entangled is a state $|\psi\rangle$?

Long history in quantum information:

Bell test or quantum games

Quantum cryptography

Tensor optimization

Hamiltonian complexity

Quantum many-body physics



Questions we can ask:

Is a state $|\psi\rangle$ entangled or product?

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This talk:

Statistical theory of testing many-body entanglement

Property testing model

Entanglement tester is an algorithm \mathcal{A} such that

1. If $|\psi\rangle$ has at most certain amount of entanglement

$$\Pr[\mathcal{A} \text{ accepts given } |\psi\rangle^{\otimes m}] \geq 2/3 \quad \text{Completeness}$$

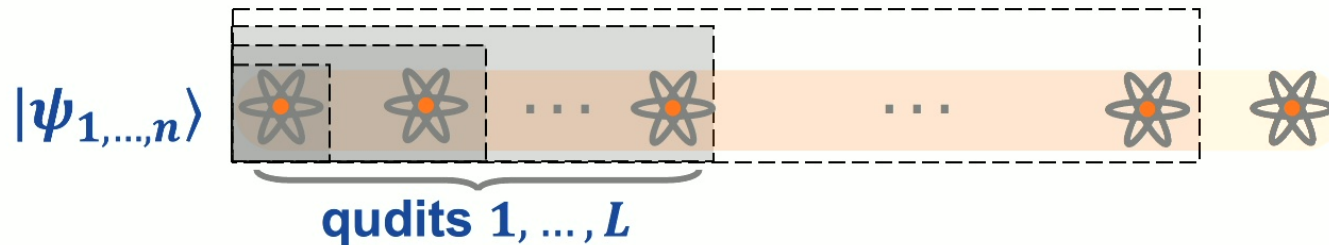
2. If $|\psi\rangle$ is far from states
with at most certain amount of entanglement

$$\Pr[\mathcal{A} \text{ accepts given } |\psi\rangle^{\otimes m}] \leq 1/3 \quad \text{Soundness}$$

***What is the fewest number of copies m
needed for entanglement testing?***

MPS(r): Matrix product states with bond dimension r

Reduced state $\rho_{1,\dots,L} = \text{tr}_{L+1,\dots,n} |\psi_{1,\dots,n}\rangle\langle\psi_{1,\dots,n}|$



$$\text{rank}(\rho_{1,\dots,L}) \leq r \quad \text{for} \quad 1 \leq L \leq n$$

If $r \sim d^n$, any state $|\psi\rangle$ can be written as an MPS
Bond dim limits the amount of entanglement

Let's go back to testing entanglement

Property testing model

MPS tester is an algorithm \mathcal{A} such that

1. If $|\psi\rangle \in \text{MPS}(r)$ then

$$\Pr[\mathcal{A} \text{ accepts given } |\psi\rangle^{\otimes m}] \geq 2/3 \quad \text{Completeness}$$

2. If $\text{Dist}_r(|\psi\rangle) \geq \delta$ then

$$\Pr[\mathcal{A} \text{ accepts given } |\psi\rangle^{\otimes m}] \leq 1/3 \quad \text{Soundness}$$

What does it mean for $|\psi\rangle$ to be far from $\text{MPS}(r)$?

$$\text{Overlap}_r(|\psi\rangle) = \max_{|\phi\rangle \in \text{MPS}(r)} |\langle \psi | \phi \rangle|^2$$

$$\text{Dist}_r(|\psi\rangle) = \min_{|\phi\rangle \in \text{MPS}(r)} D_{\text{trace}}(\psi, \phi) = \min_{|\phi\rangle \in \text{MPS}(r)} \sqrt{1 - |\langle \psi | \phi \rangle|^2}$$

$$\text{Dist}_r(|\psi\rangle) = \sqrt{1 - \text{Overlap}_r(|\psi\rangle)}$$

Property testing model

MPS tester is an algorithm \mathcal{A} such that

1. If $|\psi\rangle \in \text{MPS}(r)$ then

$$\Pr[\mathcal{A} \text{ accepts given } |\psi\rangle^{\otimes m}] \geq 2/3 \quad \text{Completeness}$$

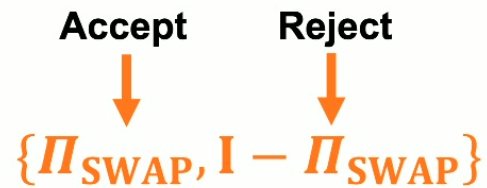
2. If $\text{Dist}_r(|\psi\rangle) \geq \delta$ then

$$\Pr[\mathcal{A} \text{ accepts given } |\psi\rangle^{\otimes m}] \leq 1/3 \quad \text{Soundness}$$

Goal: Finding the smallest **number of copies m**

for a given $\left\{ \begin{array}{l} \text{number of qudits } n \\ \text{bond dimension } r \\ \text{distance } \delta \end{array} \right.$

Product test (testing MPS(r) with $r = 1$)



[Mintert, Kuś, Buchleitner]
[Harrow and Montanaro]



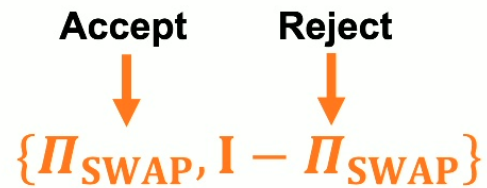
$$\Pi_{\text{SWAP}} = \frac{I + \text{SWAP}}{2}, \quad I - \Pi_{\text{SWAP}} = \frac{I - \text{SWAP}}{2}$$

$$\Pr[\text{SWAP test accepts } |\psi\rangle \otimes |\phi\rangle] = \frac{1}{2} + \frac{1}{2} |\langle \psi | \phi \rangle|^2$$

SWAP test

[Buhrman, Cleve, Watrous, de Wolf]

Product test (testing MPS(r) with $r = 1$)



[Mintert, Kuś, Buchleitner]
[Harrow and Montanaro]



$$\Pi_{\text{SWAP}} = \frac{I + \text{SWAP}}{2}, \quad I - \Pi_{\text{SWAP}} = \frac{I - \text{SWAP}}{2}$$

$$\Pr[\text{SWAP test accepts } \rho \otimes \rho] = \frac{1}{2} + \frac{1}{2} \text{tr}[\rho^2] \leftarrow \text{Purity}$$

SWAP test

[Buhrman, Cleve, Watrous, de Wolf]

Product test (testing MPS(r) with $r = 1$)

[HM13]:

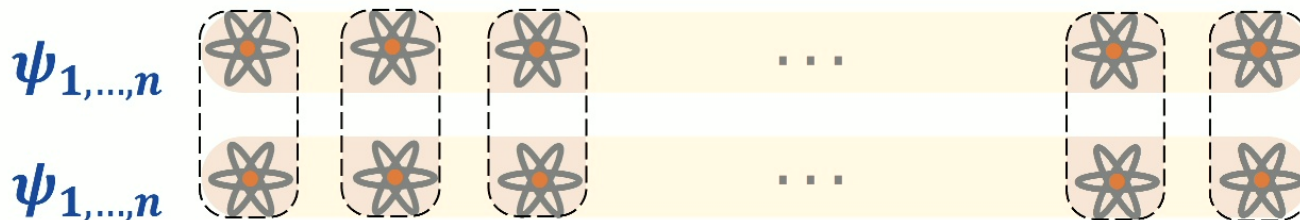
Product states pass this test
with probability 1

States δ -far from product fail this test
with probability $\Omega(\delta^2)$

Why?

Entangled $|\psi_{1,\dots,n}\rangle$ means some mixed subsystems with $\text{tr}[\rho^2] < 1$

$$\Pr[\text{SWAP test accepts } \rho \otimes \rho] = \frac{1}{2} + \frac{1}{2} \text{tr}[\rho^2] < 1$$



Product test (testing $MPS(r)$ with $r = 1$)

[HM13]:

**Product states pass this test
with probability 1**

**States δ -far from product fail this test
with probability $\Omega(\delta^2)$**

Rejection probability can be boosted to $2/3$

by repeating on $m = O\left(\frac{1}{\delta^2}\right)$ pairs

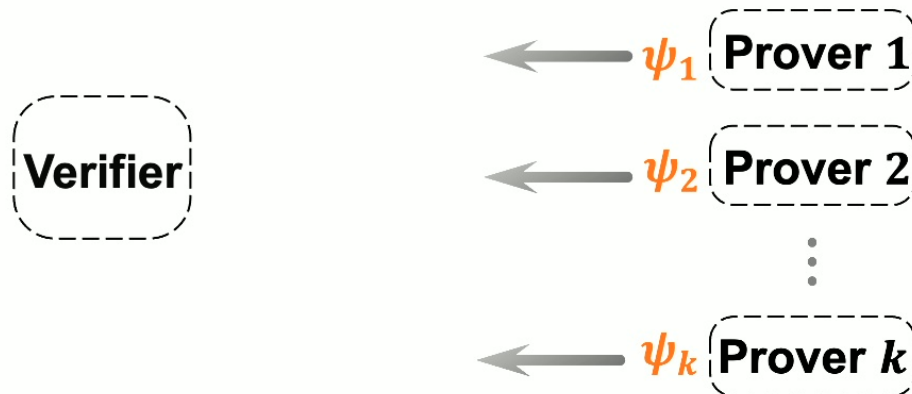
Product test (testing $\text{MPS}(r)$ with $r = 1$)

[HM13]:

**Product states pass this test
with probability 1**

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with probability $\Omega(\delta^2)$**

This implies $\text{QMA}(k) = \text{QMA}(2)$ for $k \geq 2$ [HM13]



Product test (testing MPS(r) with $r = 1$)

[HM13]:

**Product states pass this test
with probability 1**

**States δ -far from product fail this test
with probability $\Omega(\delta^2)$**

This implies $\text{QMA}(k) = \text{QMA}(2)$ for $k \geq 2$ [HM13]

With applications in hardness of tensor optimization problems

Result 1

**Improved and simple analysis of
product test**

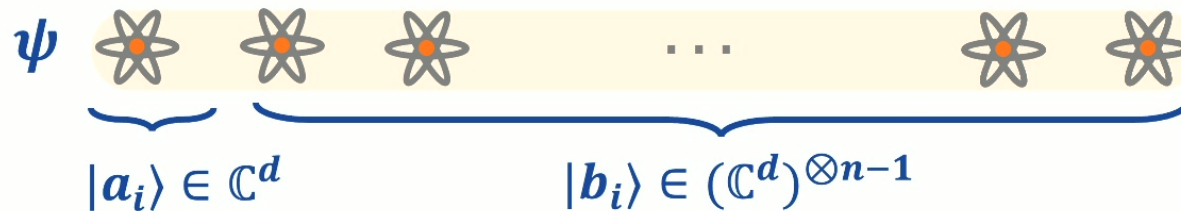
Proof sketch of product test

Schmidt decomposition

$$|\psi\rangle = \sqrt{\lambda_1} |a_1\rangle|b_1\rangle + \sqrt{\lambda_2} |a_2\rangle|b_2\rangle + \cdots + \sqrt{\lambda_d} |a_d\rangle|b_d\rangle$$

$$\begin{aligned} \lambda_1 &\geq \lambda_2 \geq \cdots \geq \lambda_d \geq 0 \\ \lambda_1 + \lambda_2 + \cdots + \lambda_d &= 1 \end{aligned}$$

$$|a_i\rangle \in \mathbb{C}^d, |b_i\rangle \in (\mathbb{C}^d)^{\otimes n-1}$$



Proof sketch of product test

Schmidt decomposition

$$|\psi\rangle = \sqrt{\lambda_1} |a_1\rangle|b_1\rangle + \sqrt{\lambda_2} |a_2\rangle|b_2\rangle + \cdots + \sqrt{\lambda_d} |a_d\rangle|b_d\rangle$$

$$\begin{aligned} \lambda_1 \geq \lambda_2 \geq \cdots \geq \lambda_d \geq 0 & \quad |a_i\rangle \in \mathbb{C}^d, |b_i\rangle \in (\mathbb{C}^d)^{\otimes n-1} \\ \lambda_1 + \lambda_2 + \cdots + \lambda_d = 1 & \end{aligned}$$

Suppose $|\psi\rangle$ is far from product.

If λ_1 is large:

$|\psi\rangle \approx |a_1\rangle \otimes |b_1\rangle$ and **1st SWAP test accepts**

But for $|\psi\rangle$ to be far from product

$|b_1\rangle$ has to be far from product

***Remaining SWAP tests reject with high prob.
(by induction)***

Proof sketch of product test

Schmidt decomposition

$$|\psi\rangle = \sqrt{\lambda_1} |a_1\rangle|b_1\rangle + \sqrt{\lambda_2} |a_2\rangle|b_2\rangle + \cdots + \sqrt{\lambda_d} |a_d\rangle|b_d\rangle$$

$$\begin{aligned} \lambda_1 \geq \lambda_2 \geq \cdots \geq \lambda_d \geq 0 & \quad |a_i\rangle \in \mathbb{C}^d, |b_i\rangle \in (\mathbb{C}^d)^{\otimes n-1} \\ \lambda_1 + \lambda_2 + \cdots + \lambda_d = 1 & \end{aligned}$$

Given $|\psi\rangle$ that is δ -far from product states,

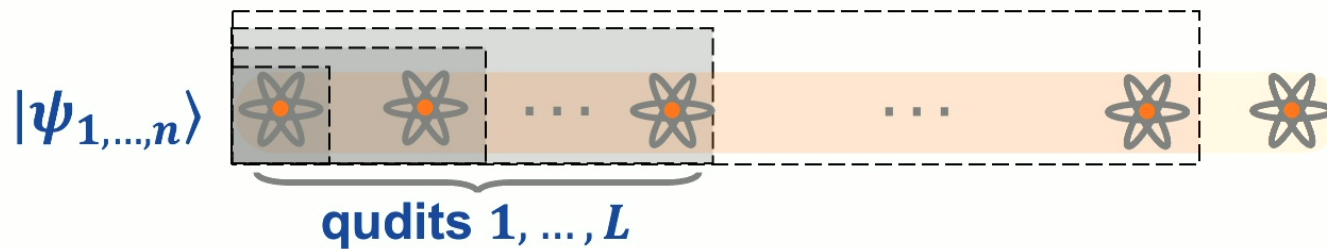
$$\Pr[\text{Product test rejects } |\psi\rangle^{\otimes 2}] \geq \begin{cases} \delta^2 - \delta^4 & \delta \leq \sqrt{1/2} \\ \frac{2}{3}\delta^2 - \frac{1}{3}\delta^4 & \text{otherwise} \end{cases}$$

Our bound is tight for $n \geq 2$, $\delta \leq \sqrt{1/2}$ as shown by

$$|\psi\rangle = \sqrt{1 - \delta^2} |1\rangle|1\rangle + \delta |2\rangle|2\rangle$$

MPS(r): Matrix product states with bond dimension r

Reduced state $\rho_{1,\dots,L} = \text{tr}_{L+1,\dots,n} |\psi_{1,\dots,n}\rangle\langle\psi_{1,\dots,n}|$



$$\text{rank}(\rho_{1,\dots,L}) \leq r \quad \text{for} \quad 1 \leq L \leq n$$

[O'Donnell and Wright'15]:

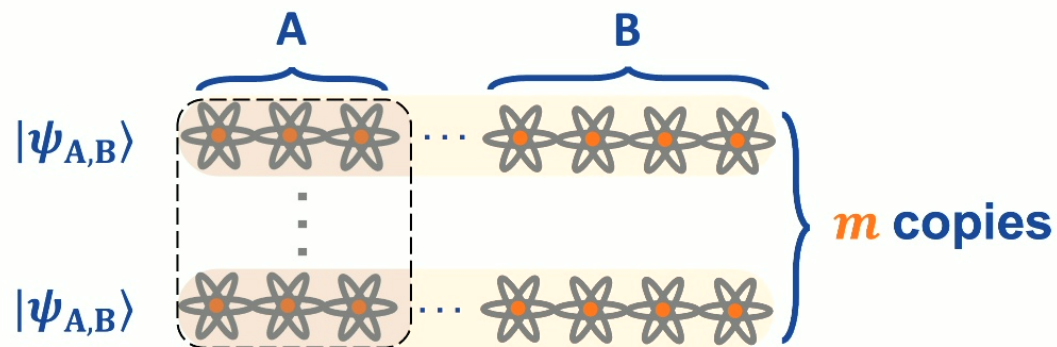
An algorithm to test if

$\text{rank}(\rho) \leq r$ or ρ is ϵ -far from rank- r states
using $m = \Theta(r^2/\epsilon)$ copies

Testing MPS(r) with $r \geq 2$

Upper bound: Our MPS tester requires $m = O(nr^2/\delta^2)$
Proof relies on

1) \exists Cut (A, B) where ρ_A is $\Omega(1/n)$ -far from being rank r



2) The rank tester projectors **mutually commute**

Testing MPS(r) with $r \geq 2$

Upper bound: Our MPS tester requires $m = O(nr^2/\delta^2)$

Proof relies on

Can this analysis be improved

to show $m = O(1)$ copies are sufficient?

Can't be done for general states!

Lower bound: Any MPS tester requires $m = \Omega(\sqrt{n}/\delta^2)$

The hard example: $|\psi\rangle$ and its random local rotations
where $|\phi\rangle$ is $1/\sqrt{n}$ -far from MPS(r)

$$|\psi\rangle = |\phi\rangle^{\otimes n/2} = \overbrace{\begin{array}{ccc} \text{atom} & \text{atom} & \dots & \text{atom} & \text{atom} \\ |\phi\rangle & |\phi\rangle & & |\phi\rangle & |\phi\rangle \end{array}}^{n/2 \text{ copies}}$$

Summary

Developed algorithms for **testing matrix product states**

- 1) **Simple** and improved analysis of the **product test**
- 2) **Upper bound of $O(n)$** for MPS testing with **bond dim ≥ 2**
- 3) **Lower bound of $\Omega(\sqrt{n})$** for MPS testing with **bond dim ≥ 2**

Setup

Suppose we have

1. Representation $R: \{0, 1\}^n \rightarrow \mathbb{C}$ for state

$$|\psi\rangle = \sum_{x_1, \dots, x_n \in \{0, 1\}^n} \psi(x_1, \dots, x_n) |x_1\rangle \cdots |x_n\rangle$$

Such that

$$R(x_1, \dots, x_n) \propto \psi(x_1, \dots, x_n)$$

Setup

Suppose we have

1. Representation $R: \{0, 1\}^n \rightarrow \mathbb{C}$ for state

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Such that

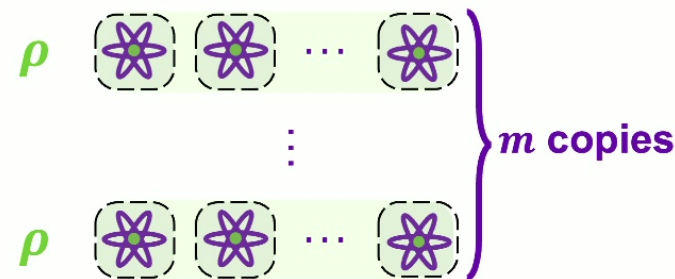
$$R(x_1, \dots, x_n) = c \cdot \psi(x_1, \dots, x_n)$$

$$\frac{R(x_1, \dots, x_n)}{R(y_1, \dots, y_n)} = \frac{\psi(x_1, \dots, x_n)}{\psi(y_1, \dots, y_n)}$$

Setup

Goal:

1. Perform few **single-qubit** measurements on ρ
2. Query representation R of $|\psi\rangle$ a few times



Some applications:

- **Benchmarking** quantum devices

ρ is prepared by a quantum device

$|\psi\rangle$ is the **ideal state** given e.g. by classical simulation

- **Verifying ML models** of quantum states

ρ is an unknown quantum state

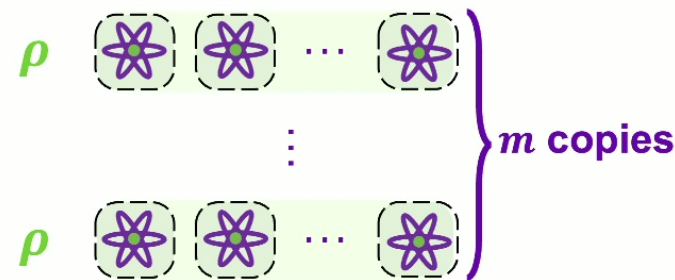
$|\psi_\theta\rangle$ is an **ML model** with **parameters θ**
that allows query access to $\psi_\theta(x_1, \dots, x_n)$

Setup

Goal:

1. Perform few **single-qubit** measurements on ρ
2. Query representation R of $|\psi\rangle$ a few times

And verify that $\langle \psi | \rho | \psi \rangle$ is **large**



Our test

Part 1

Measuring ρ

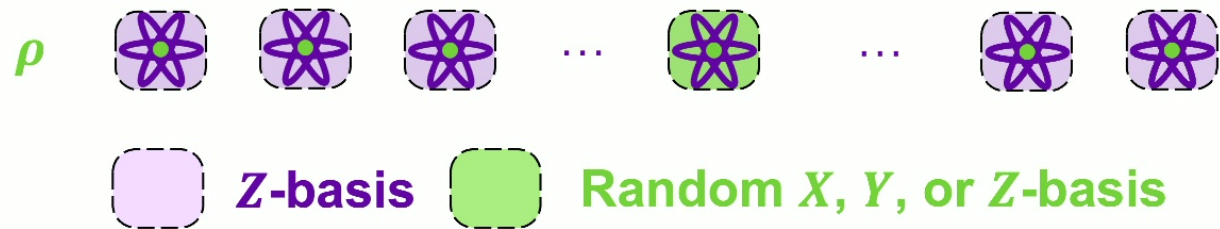
Our test

Choose $i \in \{1, \dots, n\}$ uniformly at random

Perform these measurements:

- Z-basis on all qubits but qubit i

- Random X, Y, or Z-basis on qubit i



Our test

Part 2

Querying representation R

Our test

Query representation R twice

$$|v\rangle = \frac{R(z^{(0)}) |0\rangle + R(z^{(1)}) |1\rangle}{\left(|R(z^{(0)})|^2 + |R(z^{(1)})|^2\right)^{1/2}}$$

$z^{(0)}$ = string $z \in \{0, 1\}^{n-1}$ padded with **0** on i 'th bit

$z^{(1)}$ = string $z \in \{0, 1\}^{n-1}$ padded with **1** on i 'th bit

Our test

Query representation R twice

$$|v\rangle = \frac{R(z^{(0)}) |0\rangle + R(z^{(1)}) |1\rangle}{\left(|R(z^{(0)})|^2 + |R(z^{(1)})|^2\right)^{1/2}}$$

$z^{(0)}$ = string $z \in \{0, 1\}^{n-1}$ padded with 0 on i 'th bit

$z^{(1)}$ = string $z \in \{0, 1\}^{n-1}$ padded with 1 on i 'th bit

Compute and return overlap

$$\langle v | (3|s\rangle\langle s| - I) |v\rangle$$

Claim

$$\mathbb{E} [\langle v | (3|s\rangle\langle s| - I) | v \rangle] \geq 1 - \delta$$



$$\langle \psi | \rho | \psi \rangle \geq 1 - \delta \cdot T$$

$T \sim$ mixing time of certain random walk

Claim

$$\mathbb{E} [\langle v | (3|s\rangle\langle s| - I) | v \rangle] \geq 1 - \delta$$



$$\langle \psi | \rho | \psi \rangle \geq 1 - \delta \cdot T$$

$T \sim$ mixing time of certain random walk

This random walk

- is on Boolean hypercube $\{0, 1\}^n$
- transitions defined by $\frac{\psi(x)}{\psi(y)}$

Claim

$$\mathbb{E} [\langle v | (3|s\rangle\langle s| - I) | v \rangle] \geq 1 - \delta$$



$$\langle \psi | \rho | \psi \rangle \geq 1 - \delta \cdot T$$

$T \sim$ mixing time of certain random walk

This random walk

- is on Boolean hypercube $\{0, 1\}^n$
- transitions defined by $\frac{\psi(x)}{\psi(y)}$
- stationary distribution is $|\psi(x)|^2$

Claim

$$\mathbb{E} [\langle v | (3|s\rangle\langle s| - I) |v\rangle] \geq 1 - \delta$$



$$\langle \psi | \rho | \psi \rangle \geq 1 - \delta \cdot T$$

$T \sim$ mixing time of certain random walk

When $|\psi\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} e^{i\alpha(x)} |x\rangle$ $T = O(n)$

When $|\psi\rangle$ is a Haar-random state $T = \text{poly}(n)$

Claim

$$\mathbb{E} [\langle v | (3|s\rangle\langle s| - I) |v\rangle] \geq 1 - \delta$$



$$\langle \psi | \rho | \psi \rangle \geq 1 - \delta \cdot T$$

$T \sim$ mixing time of certain random walk

When $|\psi\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} e^{i\alpha(x)} |x\rangle$ $T = O(n)$

When $|\psi\rangle$ is a Haar-random state $T = \text{poly}(n)$

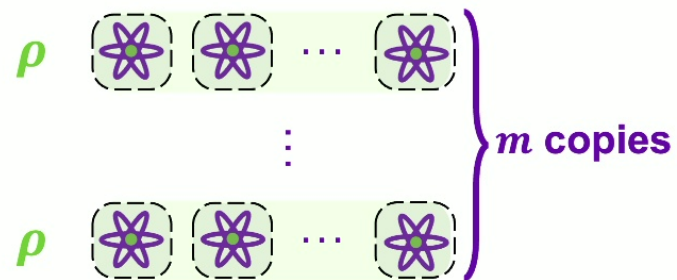
$\text{poly}(n)$ sample complexity as long as $T = \text{poly}(n)$

We can achieve

Goal:

1. Perform few **single-qubit** measurements on ρ
2. Query representation R a few times

And verify that $\langle \psi | \rho | \psi \rangle$ is large



2. Two-party testing of ground states

**Joint work with Anurag Anshu (Harvard)
and Aram Harrow (MIT)**

**Nature Physics 2022
arxiv: 2004.15009**

Testing bipartite states

Alice

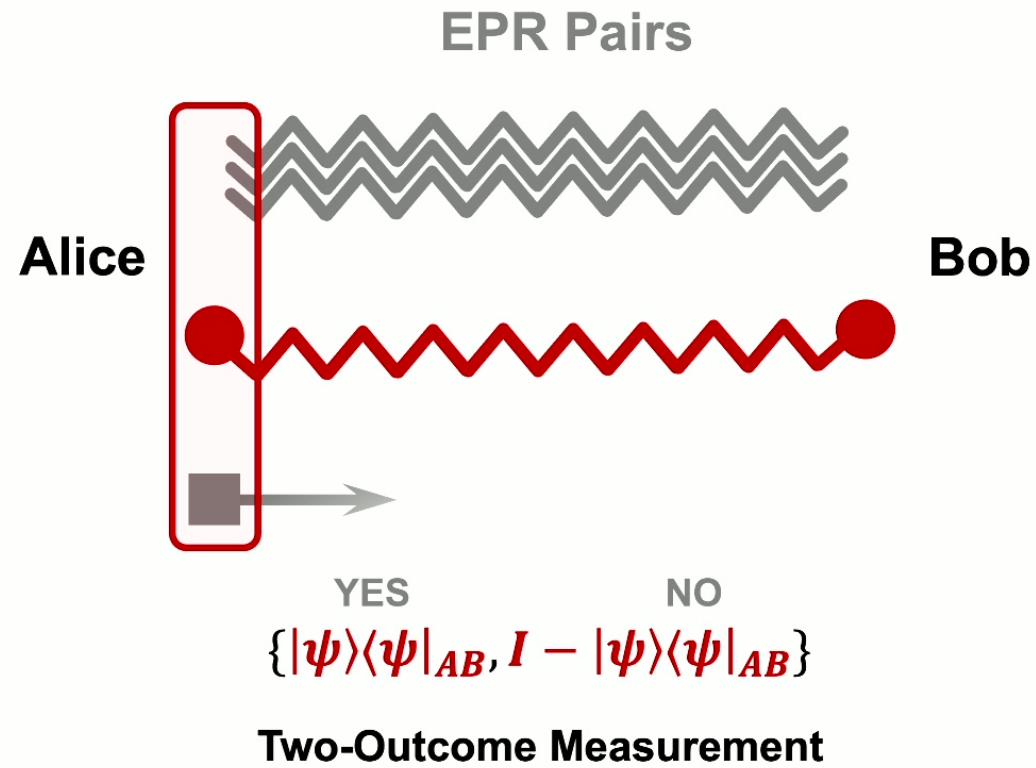


Bob

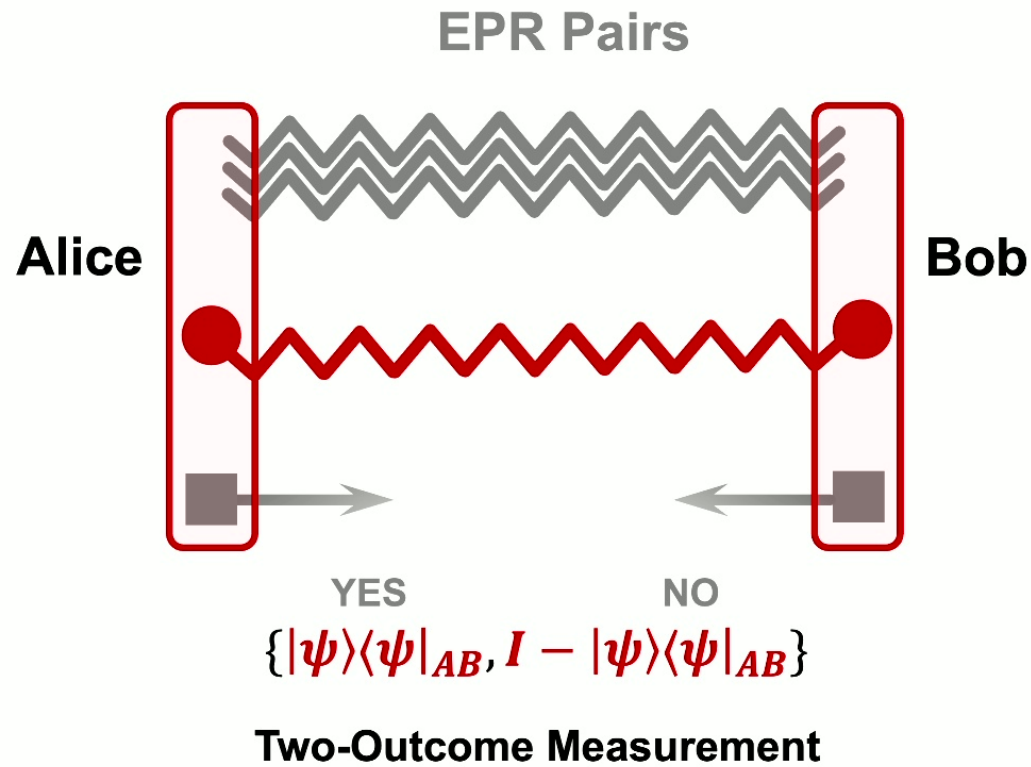
$$\begin{array}{cc} \text{YES} & \text{NO} \\ \{ |\psi\rangle\langle\psi|_{AB}, I - |\psi\rangle\langle\psi|_{AB} \} & \end{array}$$

Two-Outcome Measurement

Testing bipartite states

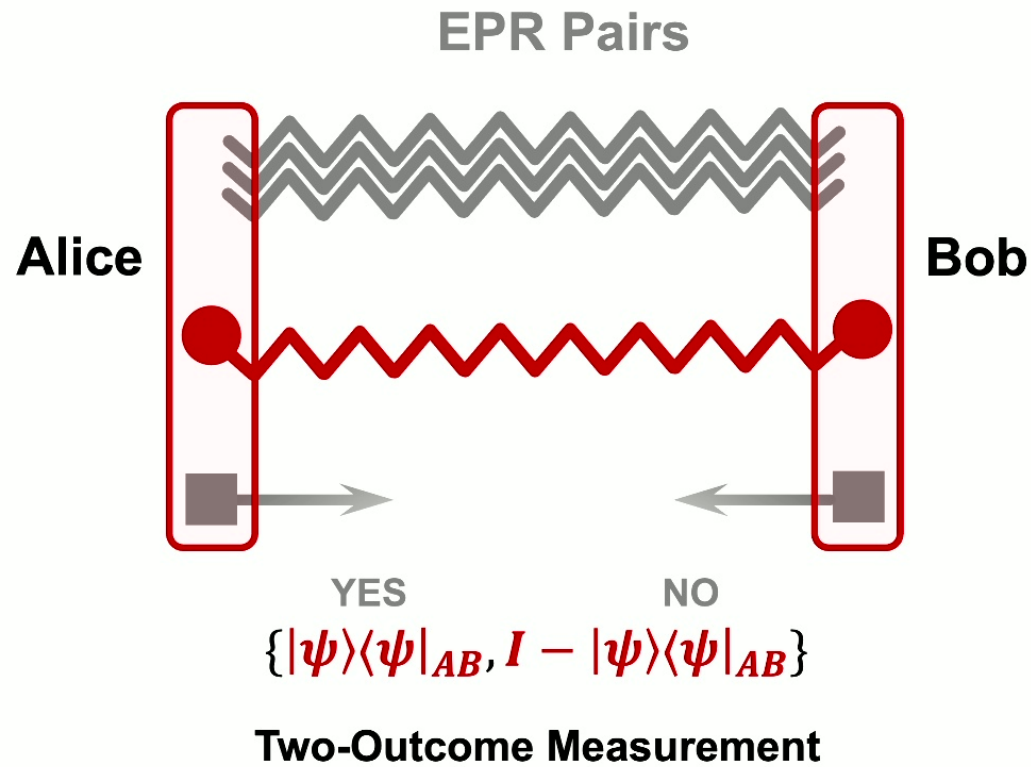


Testing bipartite states



$C(\psi_{AB})$ = Minimum # of exchanged **qubits**
to perform $\{|\psi\rangle\langle\psi|_{AB}, I - |\psi\rangle\langle\psi|_{AB}\}$

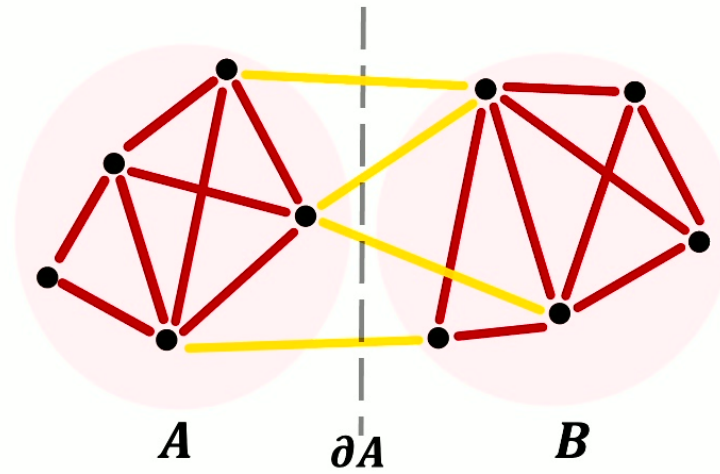
Testing bipartite states



$C_\varepsilon(\psi_{AB}) =$ Minimum # of exchanged **qubits**
to perform **ε approximation** of $\{|\psi\rangle\langle\psi|_{AB}, I - |\psi\rangle\langle\psi|_{AB}\}$

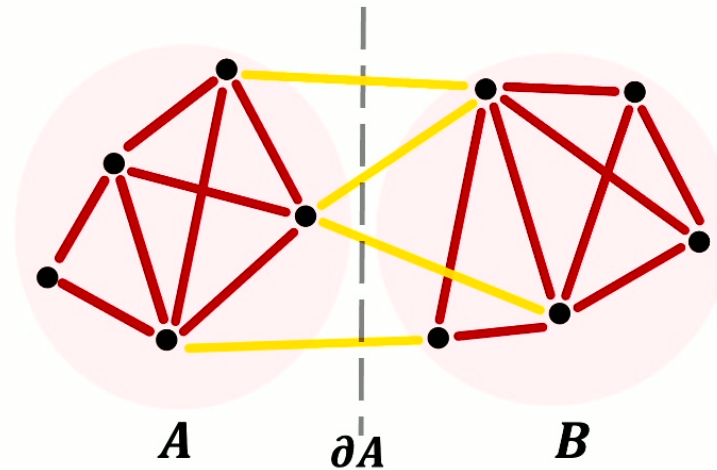
Testing *gapped* ground states

Testing ground states



$$H_{AB} = H_A + H_{\partial A} + H_B$$

Testing ground states



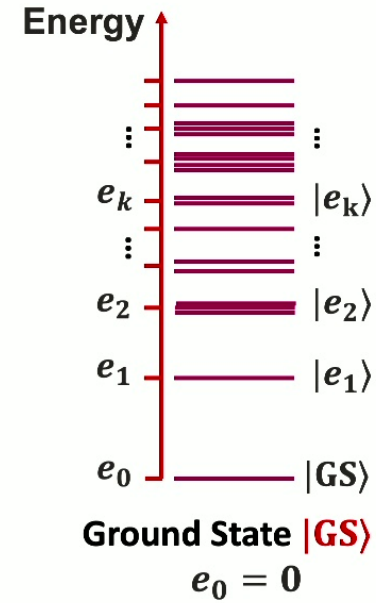
$$H_{AB} = H_A + H_{\partial A} + H_B$$

What is the communication complexity of testing the ground state?

Testing **gapped** ground states

Measure energy $\langle \psi | H | \psi \rangle$

- Yes: $\langle \psi | H | \psi \rangle \leq \mathbf{gap}/2$
- No: $\langle \psi | H | \psi \rangle > \mathbf{gap}/2$

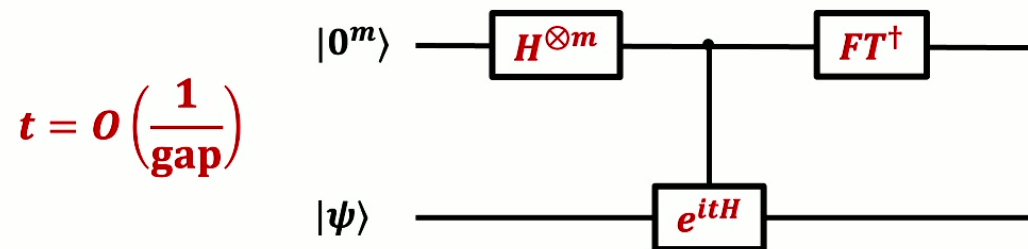


Testing **gapped** ground states

Measure energy $\langle \psi | H | \psi \rangle$

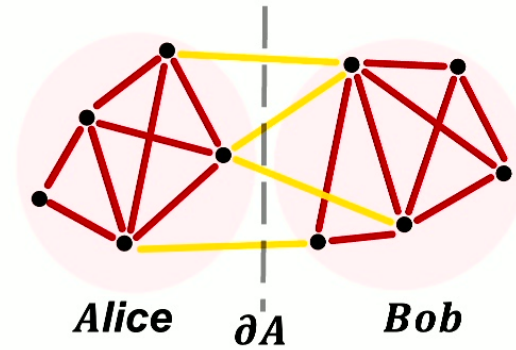
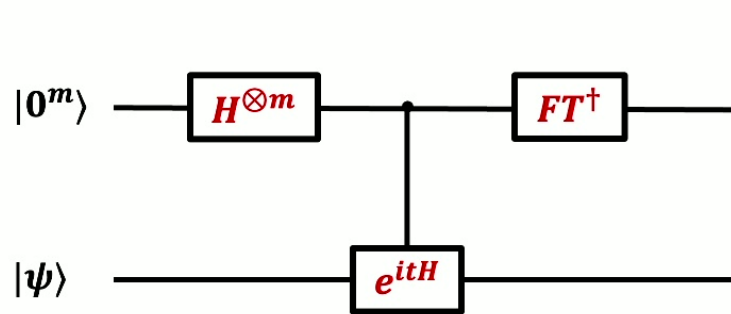
- Yes: $\langle \psi | H | \psi \rangle \leq \text{gap}/2$
- No: $\langle \psi | H | \psi \rangle > \text{gap}/2$

Quantum Phase Estimation



Testing *gapped* ground states

Communication Protocol

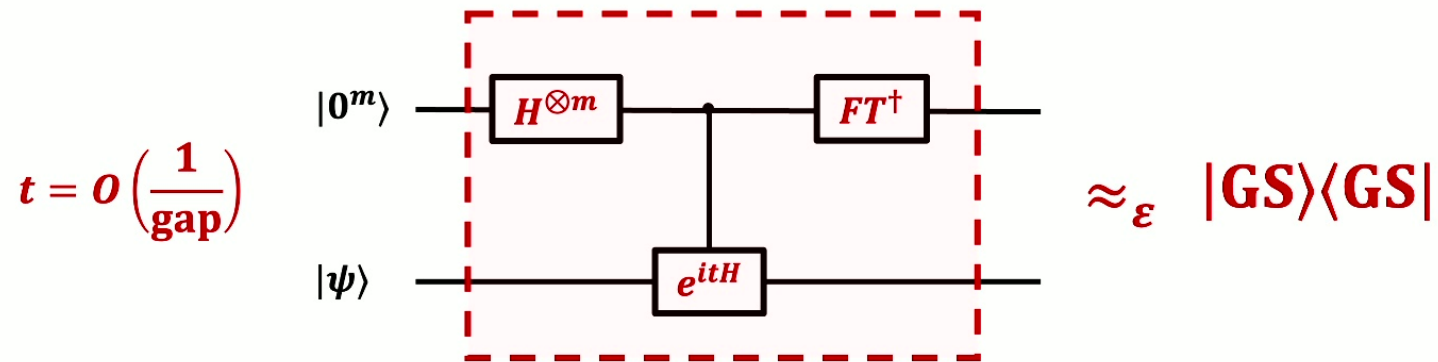


Testing **gapped** ground states

Measure energy $\langle \psi | H | \psi \rangle$

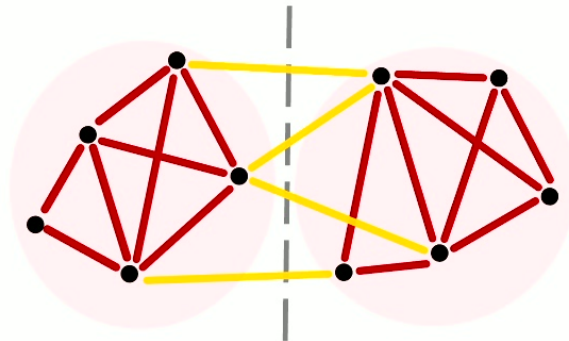
- Yes: $\langle \psi | H | \psi \rangle \leq \text{gap}/2$
- No: $\langle \psi | H | \psi \rangle > \text{gap}/2$

Quantum Phase Estimation



Repeat for $O\left(\log\frac{1}{\epsilon}\right)$ to get ϵ approximation

Hamiltonian Simulation (Performing $e^{itH_{AB}}$)

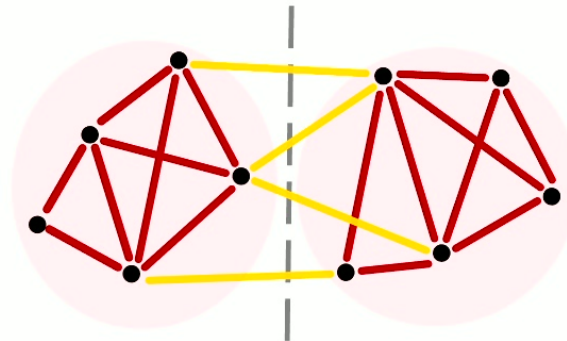


$$H_{AB} = H_A + H_{\partial A} + H_B$$

Depth of Hamiltonian simulation algorithms is $O(t\|H_{AB}\|)$

Communication cost of $e^{itH_{AB}}$ is $O(t\|H_{AB}\|)$

Hamiltonian Simulation (Performing $e^{itH_{AB}}$)

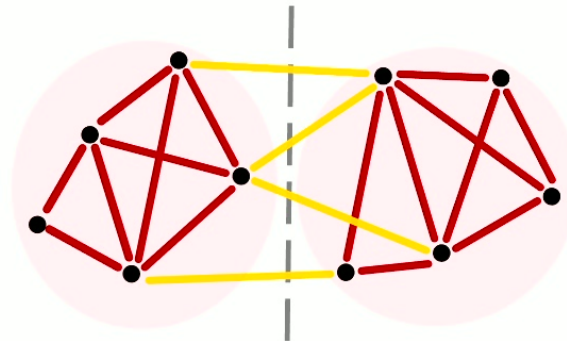


$$H_{AB} = H_A + H_{\partial A} + H_B$$

$$e^{itH_{AB}} = e^{itH_A} \cdot e^{itH_B} \cdot e^{itH_{\partial A}} \quad \text{when } H_A, H_B, H_{\partial A} \text{ Commute}$$

Interaction Picture: Time-dependent Hamiltonian [LW18]

Hamiltonian Simulation (Performing $e^{itH_{AB}}$)



$$H_{AB} = H_A + H_{\partial A} + H_B$$

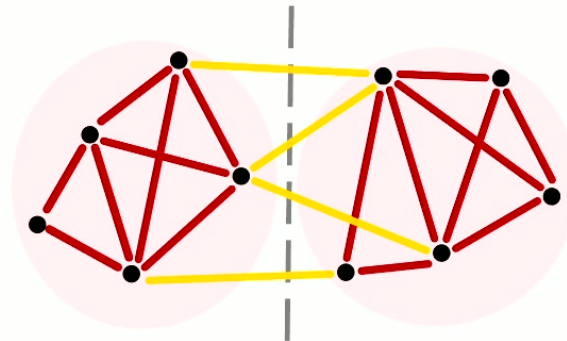
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Interaction Picture: Time-dependent Hamiltonian [LW18]

$$H_I(t) = e^{-it(H_A+H_B)} \cdot H_{\partial A} \cdot e^{it(H_A+H_B)}$$

$$e^{itH_{AB}} = e^{itH_A} \cdot e^{itH_B} \cdot e^{\int_{\tau=0}^t iH_I(\tau) d\tau}$$

Hamiltonian Simulation (Performing $e^{itH_{AB}}$)



$$H_{AB} = H_A + H_{\partial A} + H_B$$

$$e^{itH_{AB}} = e^{itH_A} \cdot e^{itH_B} \cdot e^{itH_{\partial A}} \quad \text{when } H_A, H_B, H_{\partial A} \text{ Commute}$$

Interaction Picture: Time-dependent Hamiltonian [LW18]

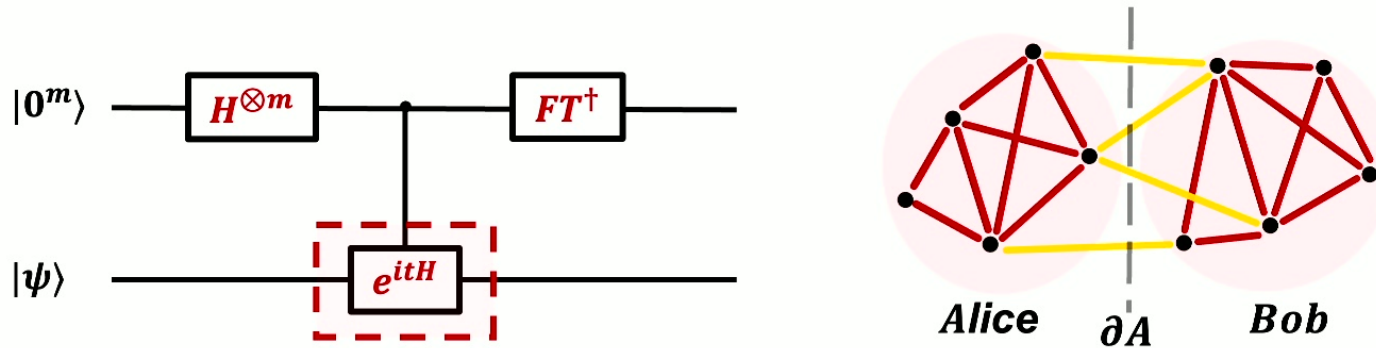
$$H_I(t) = e^{-it(H_A+H_B)} \cdot H_{\partial A} \cdot e^{it(H_A+H_B)}$$

$$e^{itH_{AB}} = e^{itH_A} \cdot e^{itH_B} \cdot e^{\int_{\tau=0}^t iH_I(\tau) d\tau}$$

$$O(t\|H_I\|) = O(t\|H_{\partial A}\|)$$

Testing *gapped* ground states

Communication Protocol



Overall **Communication Cost**: $\tilde{O}(|\partial A|/\text{gap} \cdot \log 1/\varepsilon)$

Our protocol gives an **upper bound** on $C_\varepsilon(\text{GS}_{AB})$

$$\tilde{O}\left(\frac{|\partial A|}{\text{gap}} \cdot \log \frac{1}{\varepsilon}\right) \geq C_\varepsilon(\text{GS}_{AB})$$

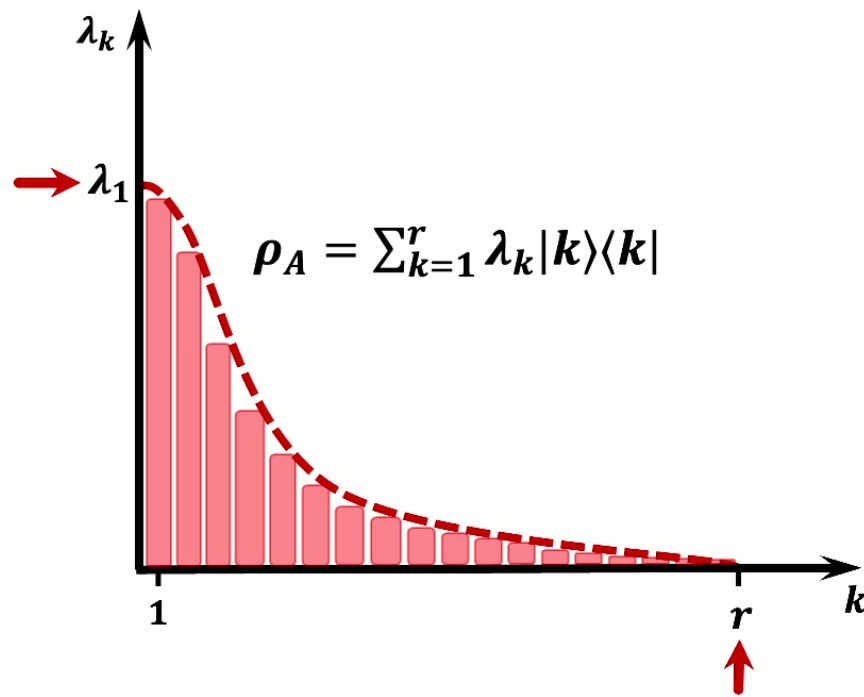
But there is also a known **lower bound** on $C_\varepsilon(\text{GS}_{AB})$:

Communication Complexity \geq Entanglement Spread

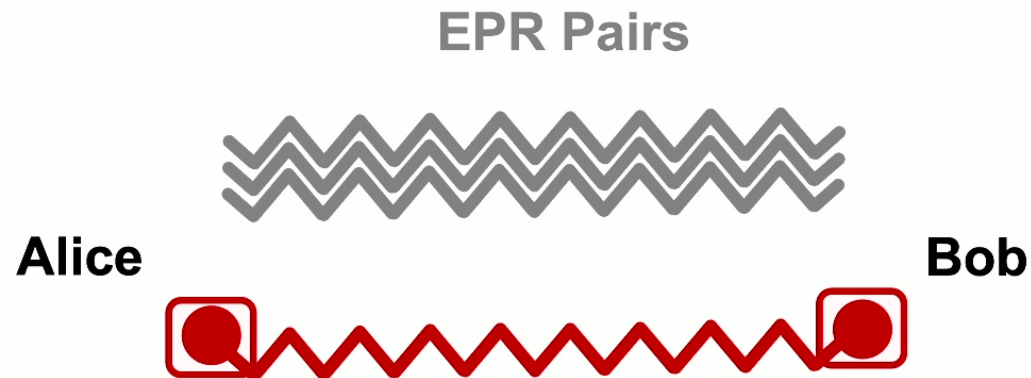
[Hayden, Winter'03]
[Coudron, Harrow'19]
[Harrow, Leung'11]

Entanglement spread

$$\text{ES}(\rho_A) = \log(r\lambda_1)$$



Testing Bipartite States



$$|\psi\rangle_{AB} = \sum_{k=1}^r \sqrt{\lambda_k} |k\rangle_A |k\rangle_B$$

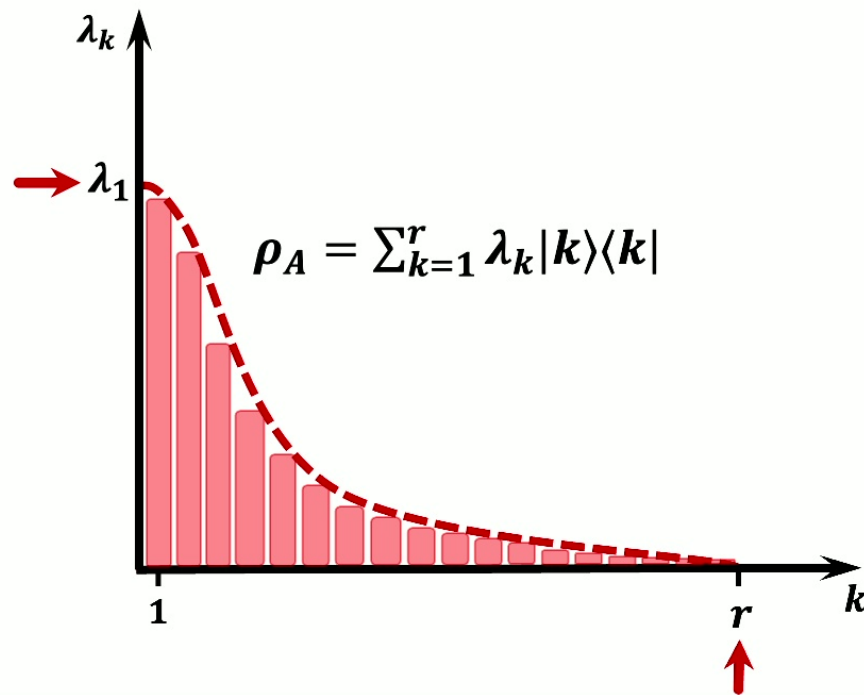
$$\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_r > 0$$

$$\lambda_1 + \lambda_2 + \dots + \lambda_r = 1$$

Schmidt Form

Entanglement spread

$$\begin{aligned}\text{ES}(\rho_A) &= \log(r\lambda_1) \\ &= S_{\max}(\rho_A) - S_{\min}(\rho_A) \\ &\approx \log(\lambda_1/\lambda_r)\end{aligned}$$



Entanglement spread

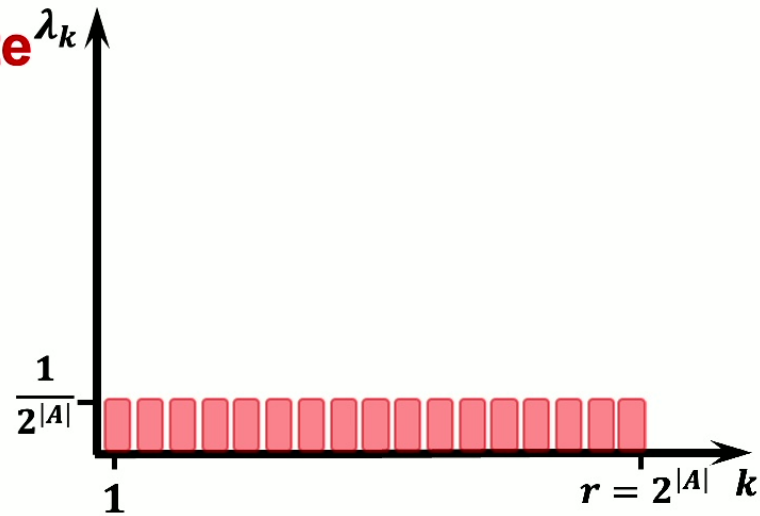
$$\begin{aligned} \text{ES}(\rho_A) &= \log(r\lambda_1) \\ &= S_{\max}(\rho_A) - S_{\min}(\rho_A) \\ &\approx \log(\lambda_1/\lambda_r) \end{aligned}$$

Some examples:

Maximally entangled state λ_k

$$|\phi\rangle_{AB} = \sum_{k=1}^{2^{|A|}} \sqrt{1/2^{|A|}} |k\rangle_A |k\rangle_B$$

$$\text{ES}(\rho_A) = 0$$



Entanglement spread

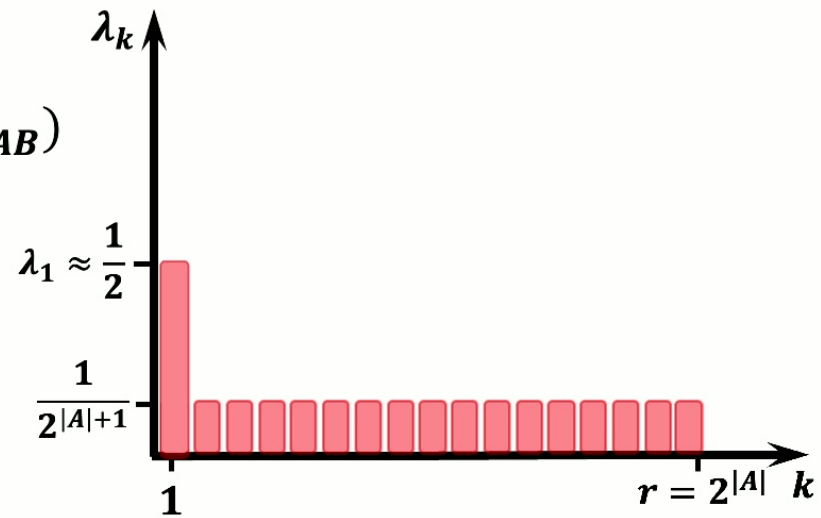
$$\begin{aligned}\text{ES}(\rho_A) &= \log(r\lambda_1) \\ &= S_{\max}(\rho_A) - S_{\min}(\rho_A) \\ &\approx \log(\lambda_1/\lambda_r)\end{aligned}$$

Some examples:

or more exotic

$$|\omega\rangle_{AB} \approx 1/\sqrt{2} (|\psi\rangle_{AB} + |\phi\rangle_{AB})$$

$$\text{ES}(\rho_A) = \mathcal{O}(|A|)$$



Our protocol gives an **upper bound** on $C_\varepsilon(\text{GS}_{AB})$

$$\tilde{O}\left(\frac{|\partial A|}{\text{gap}} \cdot \log \frac{1}{\varepsilon}\right) \geq C_\varepsilon(\text{GS}_{AB})$$

But there is also a known **lower bound** on $C_\varepsilon(\text{GS}_{AB})$:

$$\tilde{O}\left(\frac{|\partial A|}{\text{gap}} \cdot \log \frac{1}{\varepsilon}\right) \geq \text{Entanglement Spread}$$

Our protocol gives an **upper bound** on $C_\varepsilon(\text{GS}_{AB})$

$$\tilde{O}\left(\frac{|\partial A|}{\text{gap}} \cdot \log \frac{1}{\varepsilon}\right) \geq C_\varepsilon(\text{GS}_{AB})$$

But there is also a known **lower bound** on $C_\varepsilon(\text{GS}_{AB})$:

$$\tilde{O}\left(\frac{|\partial A|}{\text{gap}} \cdot \log \frac{1}{\varepsilon}\right) \geq \text{Entanglement Spread}$$

Compared to area law: $O\left(\frac{|\partial A|}{\text{gap}}\right) \geq \text{Entanglement Entropy}$

1. Testing matrix product states

2. Testing machine learning models

3. Two-party testing of ground states