Title: Experimental Cosmology

Speakers: Matthew Johnson

Collection: Quantum Spacetime in the Cosmos: From Conception to Reality

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Experimental Cosmology



Matthew C. Johnson York University Perimeter Institute

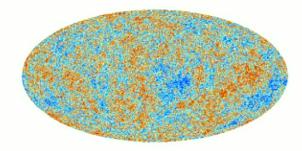


Collaborators: Jonathan Braden, Alex Jenkins, Hiranya Peiris, Dalila Pirvu, Andrew Pontzen, and Silke Weinfurtner.

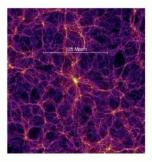
arXiv: 2204.11867, 2109.04496, 1904.07873, 1806.06069, 1712.02356

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- Cosmology ~ archaeology. We get limited and imperfect records of the past, and have to reconstruct what happened.
- Luckily, cosmology is simple enough that we can build theories and calculate or simulate.



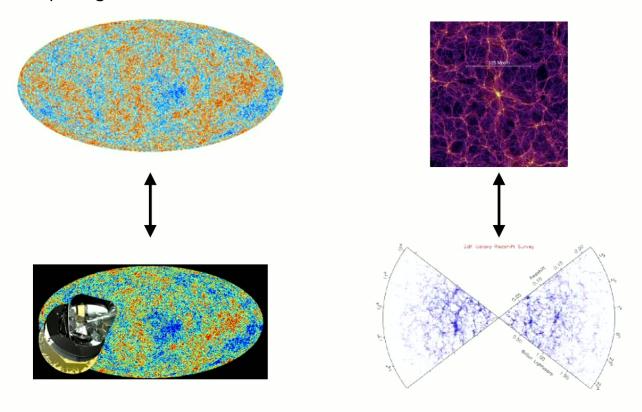
Linear theory



N-body

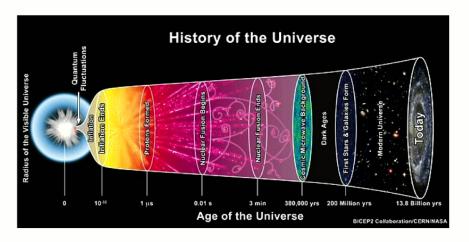
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Comparing with what we observe, we can iterate and learn more.



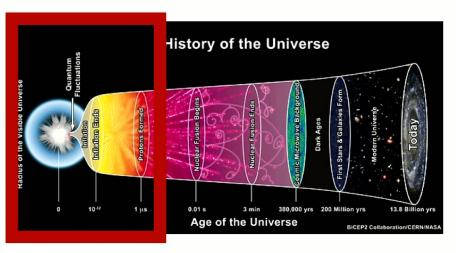
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• Early Universe: little information left over, and unknown/new physics at play.



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• Early Universe: little information left over, and unknown/new physics at play.

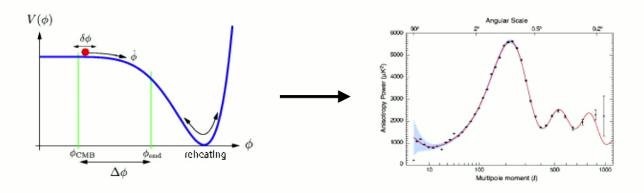


Quantum effects are important.

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Quantum in Cosmology

 Easy quantum: inflationary fluctuations sourced by a field with almost no self-interactions. This is (almost) as easy as the harmonic oscillator problem in QM. Perturbation theory is a useful tool.



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Quantum in Cosmology

 Hard quantum: non-linear and non-perturbative phenomena, strong coupling, etc. — limited tools and limited understanding.

A few things from my corner of the Universe I'd like to understand better:

The 'theory' of de Sitter space.

(I have little to say)

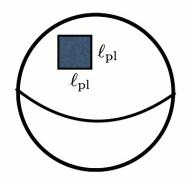
Decay of metastable (false) vacua.

(I have lots to say)

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de Sitter Holography

Does the 'theory' of dS have a finite or infinite number of DOF?



$$S \sim \#$$
 of ℓ_p^2 pixels

$$\sim \#$$
 of bits

$$\sim 100 \text{ Tb}$$

(Inflation consistent with bound on r)

What's the Hamiltonian?

Thermodynamics of horizon -> density fluctuations?

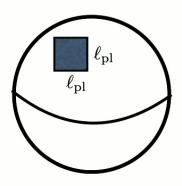
Lots more DOF today. What about then?

Complex mapping?

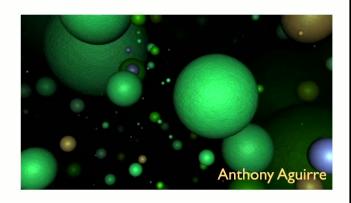
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de Sitter Holography

• Does the 'theory' of dS have a finite or infinite number of DOF?



?OR?

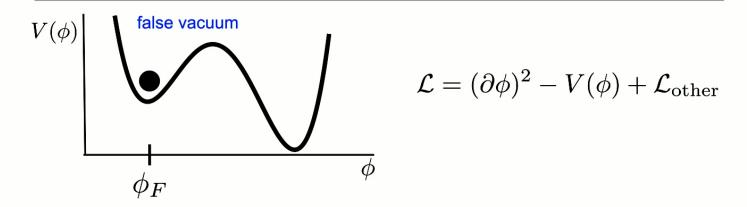


 ∞ ?

Among the most important questions: can computation/simulation help?

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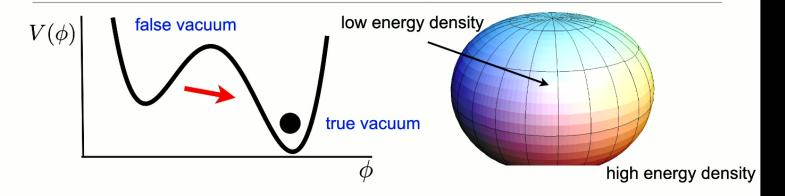
Metastable (false) vacuum decay



Classically, if the field is everywhere in the false vacuum, it will remain there.

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Metastable (false) vacuum decay

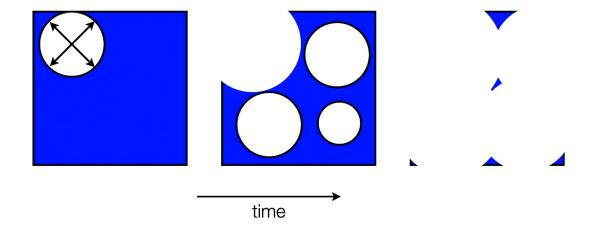


Quantum mechanically, the field can transition to the true vacuum!

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Metastable (false) vacuum decay

 Many nucleation events eventually (in flat space) complete the phase transition.



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Applications of vacuum decay

- Vacuum decay is ubiquitous in (B)SM physics and Early Universe Cosmology.
 - Higgs stability.
 - Electroweak phase transition.
 - Baryogensis.
 - Inflation/eternal Inflation.
 - SUSY/GUTs/string theory.
 - More.....

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Fundamentals of vacuum decay

 Vacuum decay is a non-perturbative feature of non-linear quantum field theories.

We might learn something new!

- There are approximate techniques for calculating the rate of individual nucleation events, but still lots left to learn:
 - Real-time description.
 - Many-bubbles/fast decays.
 - Correlations between nucleation events.
 - Non-trivial initial states.
 - Time-dependent backgrounds.

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Experimental cosmology

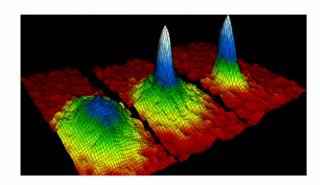
• Wouldn't it be nice if we could make a Universe in the lab to learn more about these and other questions?

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Experimental cosmology

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Quantum simulators



Lab systems like a BEC which can be used to simulate the early Universe.



Like an analog quantum computer - or actual quantum computer.

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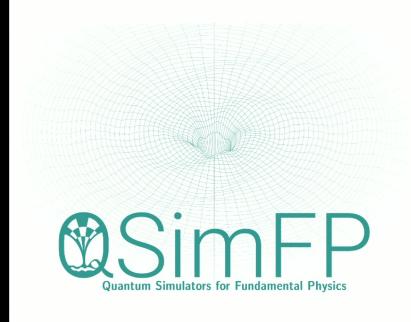
Experimental cosmology

40+ years after Unruh suggested black hole evaporation could be studied with sound horizons and 20 years after BECs were proposed as analogue systems for quantum fields in an expanding Universe...

- Viermann et. al. Nature (2022) Massless scalar in AdS from phonons in a single-component trapped BEC.
- Eckel et. al. PRX (2018) Scalar in 1+1 dS from phonons in expanding BEC ring.
- Wittemer et. al. PRL (2020) Pair creation in expanding Universe from ion trap.
- Schmitt et. al. Nature (2023) Schwinger pair production in graphene.

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QSimFP



https://qsimfp.org

Consortium of theorists/
experimentalists in UK/Canada
funded by UKRI to develop
quantum simulators of
fundamental physics in
cosmology and black holes.

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 Two-species Bose-Einstein Condensate systems can serve as quantum simulators of false vacuum decay:

$$\begin{aligned} \text{GPEs:} & i\hbar\dot{\psi}_1 = -\frac{\hbar^2\nabla^2}{2m}\psi_1 + \underbrace{\left(g_{\mathrm{s}}\left|\psi_1\right|^2\psi_1 + g_{\mathrm{c}}\left|\psi_2\right|^2\psi_1}_{\left(g_{\mathrm{s}}\left|\psi_2\right|^2\psi_2 + g_{\mathrm{c}}\left|\psi_1\right|^2\psi_2\right)} - \underbrace{\left(\nu\psi_2\right)^2\psi_1}_{\left(\psi_1\right)} \end{aligned} \\ & \text{Son \& Stephanov; Fialkov et. al.} \\ & \text{scattering} \qquad \text{mixing/coupling} \end{aligned}$$

• Change basis $\psi_i = \sqrt{\rho_i} e^{i\phi_i}$ and integrate out density and total phase:

$$\mathcal{L}_{\text{eff}} = \frac{\dot{\varphi}^2}{2} - \frac{(\nabla \varphi)^2}{2} + \frac{4\nu \bar{n}(g_s - g_c)}{\hbar^2} \cos \varphi \qquad \qquad \varphi \equiv \phi_1 - \phi_2$$

See 1712.02356 for details.

• We can create a false vacuum by introducing a time-dependent coupling:

$$\nu(t) = \nu_0 + \delta\hbar\omega\cos(\omega t)$$

Fialkov et. al. Braden, MCJ, et. al.



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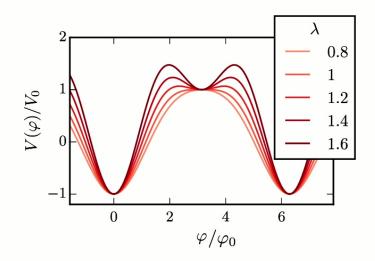
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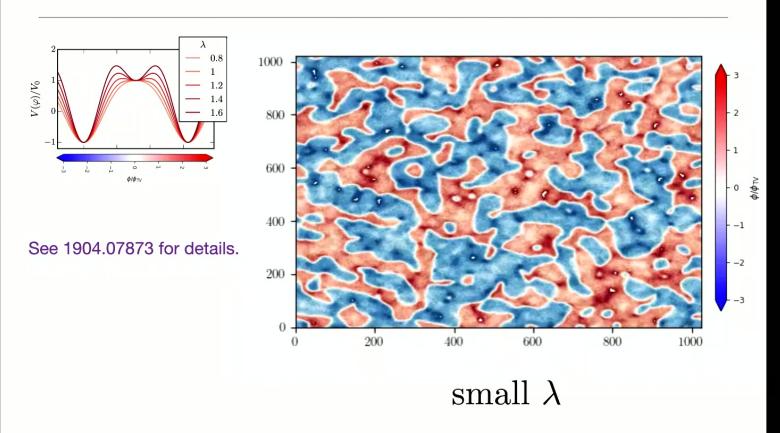
See 1712.02356 for details.

• Time-averaging, we obtain:

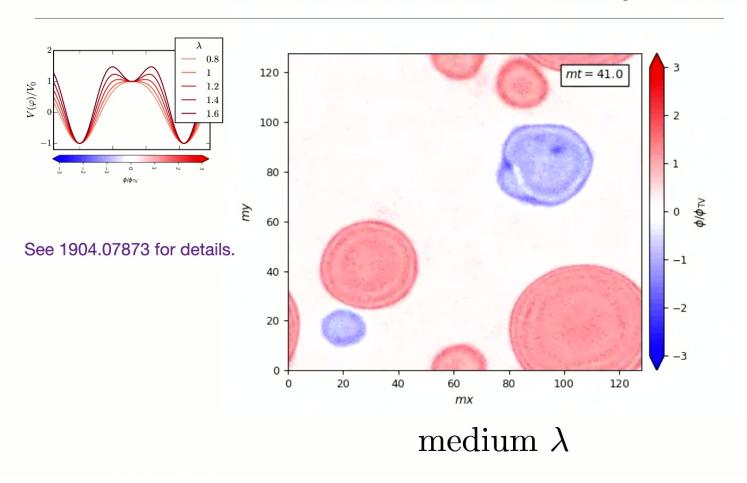
$$\mathcal{L}_{ ext{eff}} = rac{\dot{arphi}^2}{2} - rac{(
abla arphi)^2}{2} - V_0 \left(-\cosarphi + rac{\lambda^2}{2}\sin^2arphi
ight)$$

Tuneable potential: $\lambda^2 \propto \delta^2/\nu_0$

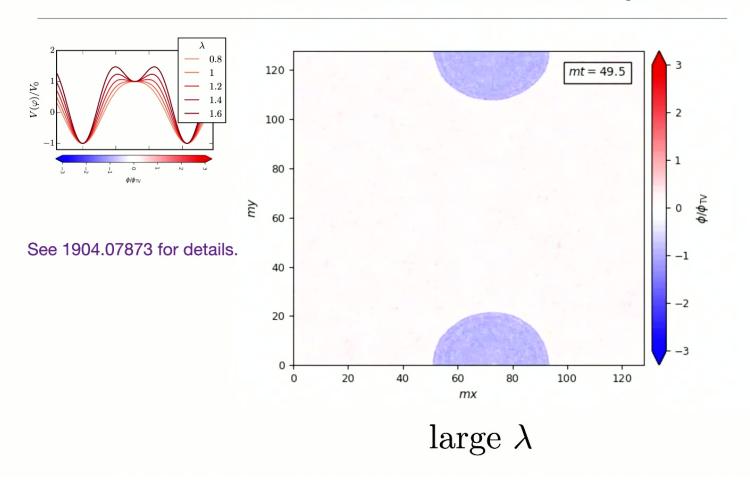




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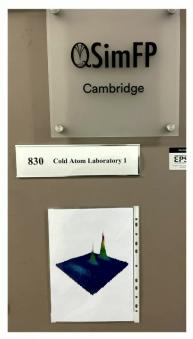


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Can we do it?



Hadzibabic lab Cambridge

- Many in-practice challenges (I'm a theorist:).
- In-principle challenges & questions:
 - Time-dependent couplings drive parametric resonance: do timescales work?
 - Expanding bubble walls length-contract: UV sensitivity (what happens sub-healing length?).
 - How faithful is the analogue?
 - What should we be measuring?
 - How can we compare to theory?
 - Quantum vs thermal?

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A great re-think

- Even if the experiments are years off, already lots of new insights into vacuum decay:
 - New semi-classical real-time description of vacuum decay: 1806.06069
 - Bubble nucleation sites are correlated: 2109.04496
 - Bubbles and RG: 2204.11867
 - Bubble nucleation pre-cursors: Coming soon!
 - Bubbles don't nucleate at rest: Coming soon!

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From quantum mechanics:

$$V \propto x^2$$
 $\langle x(t) \rangle = x_{cl}(t), \langle p(t) \rangle = p_{cl}(t)$

• For non-linear potentials, this is approximately true as long as

$$(\Delta x)^2 \ll V'/V'''$$

 Can extend to statements about time evolution of a quantum state using the Wigner Function:

$$W(x,p) = \frac{1}{\pi\hbar} \int \langle x+y|\hat{\rho}|x-y\rangle e^{-2ipy/\hbar}dy$$

Quasi-probability distribution (e.g. not positive definite) on phase space.

 For a quadratic potential, the classical time-evolution of the Wigner function is equivalent to the Schrodinger equation:

$$\partial_t W = \{W, H\} \iff i\hbar \partial_t \psi = \hat{H}\psi$$

- A recipe for evolving the Wigner function:
 - 1. Create an ensemble of systems with initial conditions fairly drawn from Wigner function (works when positive definite: ground state).
 - 2. Time-evolve with the classical Hamiltonian.
 - 3. Re-calculate the distribution on phase space at a later time.

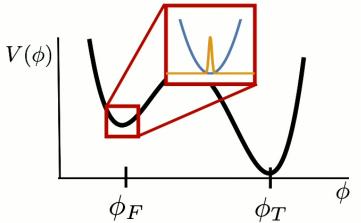
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 Field theory: promote functions to functionals, and proceed as before:

$$\left[\frac{\partial}{\partial t} + \int d^dx \left(\dot{\phi} \frac{\delta}{\delta \phi} + \dot{\Pi} \frac{\delta}{\delta \Pi} \right) + \mathcal{O} \left(\hbar^2 \ V^{\prime\prime\prime}(\phi) \frac{\delta^3}{\delta \Pi^3} \right) \right] W \left[\phi(x), \Pi(x); t \right] = 0$$
 classical evolution quantum noise initial state

- For a massive scalar, "quantum effects" enter only through the initial state.
- Interactions lead to additional effects, which can in principle be included perturbatively as stochastic contributions to classical evolution.

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Sufficiently stable vacua can be approximated by a massive theory.

 Follow the recipe for evolving the Wigner Function(al) by performing ensembles of classical lattice simulations with different initial conditions.

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- Similar stochastic methods have been used to study preheating, sphalerons, cosmic strings, etc.
- This is leading order in hbar through the initial conditions, but no virtual particles - it is a semi-classical approximation.
- Bubbles, and their precursors, are many-particle states. A semiclassical description should be OK.

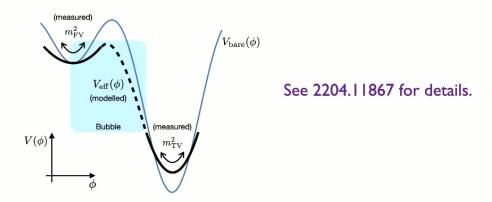
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Comparing decay rates

Coleman prediction for the decay rate:

$$rac{\Gamma}{V} = Ae^{-B}$$
 — Depends on the potential.

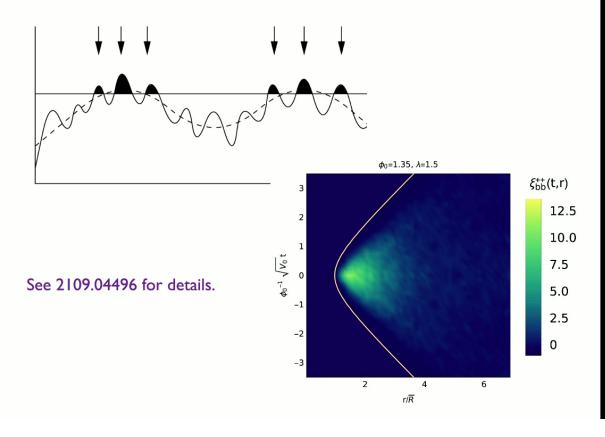
- The potential input to the simulations is not the one seen by the IR modes making up the bubble -> RG effects.
- Around minima, clear what to do. Over the barrier quite subtle....



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Bubble bias

Bubbles are biased tracers of the underlying quantum vacuum fluctuations.



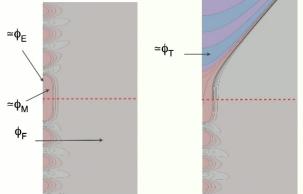
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How do bubbles form?

- Q: What is the most likely time-evolution of large fluctuations away from thermal equilibrium?
- A: The CPT conjugate of the most likely way the system relaxes back to equilibrium.
 See 1108.0417 for details.

 Thermal bubbles: a true vacuum bubble in unstable equilibrium between expansion and collapse.

Oscillon pre-cursor observed in sims.



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Experimental Cosmology

- Can laboratory-based systems, quantum computers, etc. teach us something new about the early Universe?
- Cosmology is dynamics -> interface with quantum dynamics is fertile ground.
- Goal: connect back to observation! What are we not looking for?
- Goal: connect with big questions. Physics of accelerated expansion?
 Quantum Gravity? Non-perturbative effects in QFT?

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